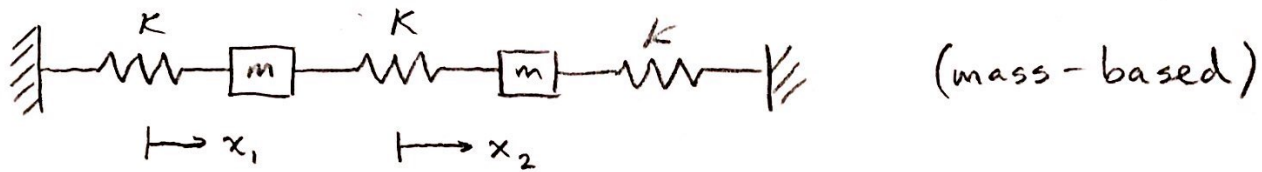


Principal Coordinates



$$\begin{bmatrix} m & 0 \\ 0 & m \end{bmatrix} \begin{bmatrix} \ddot{x}_1 \\ \ddot{x}_2 \end{bmatrix} + \begin{bmatrix} 2K & -K \\ -K & 2K \end{bmatrix} \begin{bmatrix} x_1 \\ x_2 \end{bmatrix} = \begin{bmatrix} 0 \\ 0 \end{bmatrix} \Rightarrow \begin{cases} \omega_1^2 = \frac{K}{m} \\ \omega_2^2 = \frac{3K}{m} \end{cases} \begin{cases} u_1 = 1 \\ u_2 = -1 \end{cases}$$

$$\text{Let } [U] = \begin{bmatrix} 1 & 1 \\ 1 & -1 \end{bmatrix} \quad \text{and} \quad \vec{x} = [U] \vec{p}$$

$$\text{Then } \begin{bmatrix} x_1 \\ x_2 \end{bmatrix} = \begin{bmatrix} 1 & 1 \\ 1 & -1 \end{bmatrix} \begin{bmatrix} p_1 \\ p_2 \end{bmatrix} \Rightarrow \begin{bmatrix} p_1 \\ p_2 \end{bmatrix} = \begin{bmatrix} 1/2 & 1/2 \\ 1/2 & -1/2 \end{bmatrix} \begin{bmatrix} x_1 \\ x_2 \end{bmatrix}$$

$$\text{So } p_1 = \frac{x_1 + x_2}{2} \quad \text{and} \quad p_2 = \frac{x_1 - x_2}{2}$$

$$\text{Let } [M^*] = [U]^T [M] [U]$$

$$[M^*] = \begin{bmatrix} 1 & 1 \\ 1 & -1 \end{bmatrix} \begin{bmatrix} m & 0 \\ 0 & m \end{bmatrix} \begin{bmatrix} 1 & 1 \\ 1 & -1 \end{bmatrix} = \begin{bmatrix} 2m & 0 \\ 0 & 2m \end{bmatrix}$$

$$\text{And } [K^*] = [U]^T [K] [U]$$

$$[K^*] = \begin{bmatrix} 1 & 1 \\ 1 & -1 \end{bmatrix} \begin{bmatrix} 2K & -K \\ -K & 2K \end{bmatrix} \begin{bmatrix} 1 & 1 \\ 1 & -1 \end{bmatrix} = \begin{bmatrix} 2K & 0 \\ 0 & 6K \end{bmatrix}$$

Equation with principal coordinates:

$$[M^*] \ddot{\vec{p}} + [K^*] \vec{p} = \vec{0}$$

$$\begin{bmatrix} 2m & 0 \\ 0 & 2m \end{bmatrix} \begin{bmatrix} \ddot{p}_1 \\ \ddot{p}_2 \end{bmatrix} + \begin{bmatrix} 2k & 0 \\ 0 & 6k \end{bmatrix} \begin{bmatrix} p_1 \\ p_2 \end{bmatrix} = \begin{bmatrix} 0 \\ 0 \end{bmatrix}$$

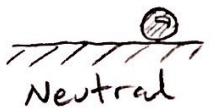
Uncoupled equations:

$$\left. \begin{aligned} 2m\ddot{p}_1 + 2kp_1 &= 0 \Rightarrow \omega_1^2 = \frac{2k}{2m} = \frac{k}{m} \\ 2m\ddot{p}_2 + 6kp_2 &= 0 \Rightarrow \omega_2^2 = \frac{6k}{2m} = \frac{3k}{m} \end{aligned} \right\} \text{As before}$$

Vibration Stability



This system has a higher ω_n than a flatter bowl. ($\omega_n^2 > 0$)

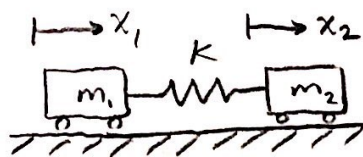


As the bowl becomes flatter $\omega_n^2 \rightarrow 0$.
For perfectly flat plane $\omega_n^2 = 0$.

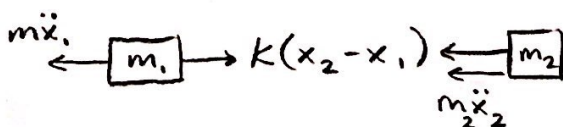


There is no frequency here. The ball rolls off and does return. ($\omega_n^2 < 0$)

Semi-Definite Systems



$$\begin{bmatrix} m_1 & 0 \\ 0 & m_2 \end{bmatrix} \begin{bmatrix} \ddot{x}_1 \\ \ddot{x}_2 \end{bmatrix} + \begin{bmatrix} k & -k \\ -k & k \end{bmatrix} \begin{bmatrix} x_1 \\ x_2 \end{bmatrix} = \begin{bmatrix} 0 \\ 0 \end{bmatrix}$$



Try solution $\begin{bmatrix} x_1 \\ x_2 \end{bmatrix} = \begin{bmatrix} X_1 \\ X_2 \end{bmatrix} \cos(\omega t + \phi)$ ← constants

Equation of motion: $\begin{bmatrix} k - \omega^2 m_1 & -k \\ -k & k - \omega^2 m_2 \end{bmatrix} \begin{bmatrix} X_1 \\ X_2 \end{bmatrix} \cos(\omega t + \phi) = \begin{bmatrix} 0 \\ 0 \end{bmatrix}$

Non-trivial solution: $\det \begin{bmatrix} k - \omega^2 m_1 & -k \\ -k & k - \omega^2 m_2 \end{bmatrix} = 0$

Then $m_1 m_2 \omega^4 - k(m_1 + m_2)\omega^2 + k^2 - k^2 = 0$

$\Rightarrow \omega_1^2 = 0 \quad \omega_2^2 = \frac{k(m_1 + m_2)}{m_1 m_2}$

Mode shapes $\begin{bmatrix} k - \omega^2 m_1 & -k \\ -k & k - \omega^2 m_2 \end{bmatrix} \begin{bmatrix} 1 \\ u \end{bmatrix} = \begin{bmatrix} 0 \\ 0 \end{bmatrix}$

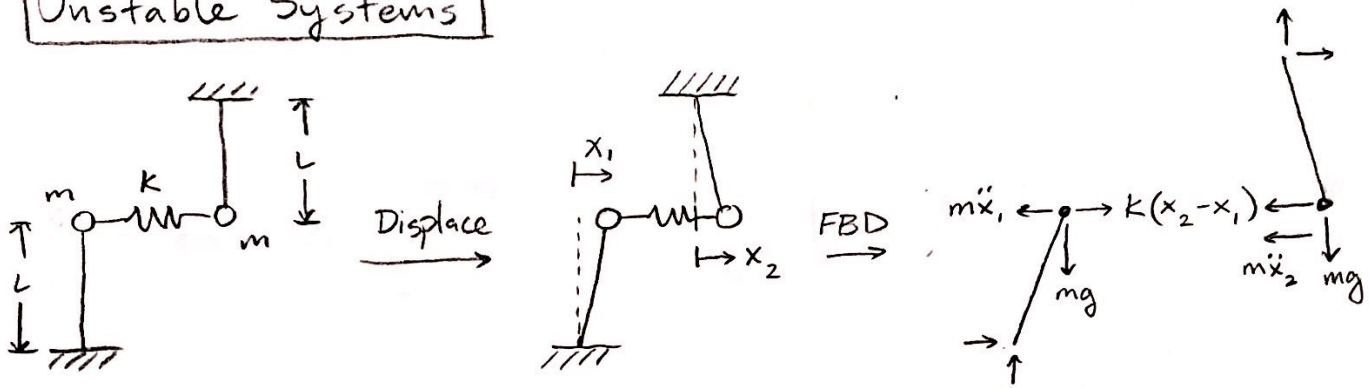
When $\omega_1^2 = 0 \Rightarrow k - k u_1 = 0 \Rightarrow \boxed{u_1 = 1}$ (top line)

$\omega_2^2 = \frac{k(m_1 + m_2)}{m_1 m_2} \Rightarrow -k \frac{m_1}{m_2} - k u_2 = 0 \Rightarrow \boxed{u_2 = -\frac{m_1}{m_2}}$ (top line)

Physically, for $u_1 = 1$, both amplitudes are equal, so the train moves forever. The period of vibration is infinite. Frequency is period⁻¹ $\rightarrow \omega = 0$. This is rigid body translation.

For $u_2 = -\frac{m_1}{m_2}$ the cars vibrate around the center of mass of the whole system.

Unstable Systems



Moments at ends $\sum M = 0 \Rightarrow$ (small angles)

$$\begin{cases} m\ddot{x}_1 l - mgx_1 - k(x_2 - x_1)l = 0 \\ m\ddot{x}_2 l + mgx_2 + k(x_2 - x_1)l = 0 \end{cases}$$

Matrices

$$\begin{bmatrix} m & 0 \\ 0 & m \end{bmatrix} \begin{bmatrix} \ddot{x}_1 \\ \ddot{x}_2 \end{bmatrix} + \begin{bmatrix} k - \frac{mg}{l} & -k \\ -k & k + \frac{mg}{l} \end{bmatrix} \begin{bmatrix} x_1 \\ x_2 \end{bmatrix} = \begin{bmatrix} 0 \\ 0 \end{bmatrix}$$

Require

$$\det \begin{bmatrix} k - \frac{mg}{l} - m\omega^2 & -k \\ -k & k + \frac{mg}{l} - m\omega^2 \end{bmatrix} = 0$$

Solve

$$m^2\omega^4 - 2mk\omega^2 - \frac{m^2g^2}{l^2} = 0 \Rightarrow \omega^2 = \frac{k}{m} \pm \sqrt{\left(\frac{k}{m}\right)^2 + \left(\frac{g}{l}\right)^2}$$

Since $\sqrt{\left(\frac{k}{m}\right)^2 + \left(\frac{g}{l}\right)^2} > \left(\frac{k}{m}\right) \Rightarrow \omega_1^2 < 0$ and $\omega_2^2 > 0$

collapse ↑ Vibration around C.O.M.

Meaning of $\omega^2 < 0$

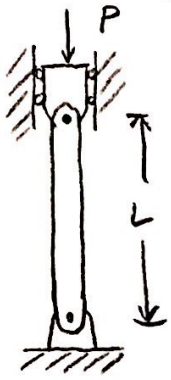
Let $\omega^2 = -\lambda^2 \Rightarrow \lambda = i\omega$. If $\omega^2 < 0 \Rightarrow \lambda^2 > 0$

Then $x = C_1 e^{i\omega t} + C_2 e^{-i\omega t}$

$x = C_1 e^{\lambda t} + C_2 e^{-\lambda t}$

↑
Unstable portion

Buckling



$$\text{Frequency [Hz]} \quad f = \frac{1}{2\pi L} \sqrt{\left(\frac{EI}{\rho A}\right) \left(\frac{\pi^4}{L^2} - \frac{P\pi^2}{EI}\right)}$$

The column buckles when $f = 0$.

$$\Rightarrow \text{when } \frac{\pi^4}{L^2} = \frac{P\pi^2}{EI} \Rightarrow P = \frac{\pi^2 EI}{L^2}$$

Pinned-Pinned
Column

This is the Euler buckling load.