Chapter 1

The Supersymmetric Extension to the Standard Model

The following introduction to SUSY focuses primarily on the aspects of the formalism that are relevant to phenomenology. In particular, most of the details of SUSY breaking (about which there is little theoretical consensus) are omitted, except where they are relevant to experiment. The notation is similar to that used in refs. **??** and **??**.

1.1 Supermultiplet Representation

The Standard Model is extended to include supersymmetry by the introduction of a supersymmetry transformation that takes fermionic states to bosonic states and vice versa. In analogy with the known symmetries of the Standard Model, the SUSY transformation has associated generators that obey defining commutation relations, and a fundamental representation. All SM particles and their *superpartners* fall into one of two *supermultiplet* representations. Using the property that

$$n_F = n_B, \tag{1.1}$$

where n_F is the number of fermionic degrees of freedom per supermultiplet and n_B is the number of bosonic degrees of freedom, the two types of supermultiplets are

1. *Chiral supermultiplets*: one Weyl fermion (two helicity states $\Rightarrow n_F = 2$) and one complex scalar field (with two real components $\Rightarrow n_B = 2$)

2. *Gauge supermultiplets*: One spin-1 vector boson (two helicity states $\Rightarrow n_B = 2$) and one Weyl fermion (two helicity states $\Rightarrow n_F = 2$)

In the gauge supermultiplet, the vector boson is assumed massless (i.e. before EWSB generates a mass for it). Since the superpartners to the SM particles have not yet been discovered, they must be significantly heavier than their SM counterparts. Unbroken SUSY predicts that the SM particles and their superpartners must have exactly the same mass, so ultimately a mechanism for SUSY breaking must be introduced to generate masses for the superpartners (see Sec. 1.3). Tables 1.1 and 1.2

show the chiral and gauge supermultiplets of the supersymmetric Standard Model, respectively. Note that the scalar partners to the SM fermions are denoted by placing an "s" in front of their names, while the chiral fermion partners to the SM gauge bosons are denoted by appending "ino" to their names.

Table 1.1: Chiral supermultiplets of the supersymmetric Standard Model. Adapted from Table 1.1 of [**?**].

Type of supermultiplet	Notation	Spin-0 component	Spin-1/2 component	Representation under $SU(3)_C \otimes SU(2)_L \otimes U(1)_Y$
Left-handed quark/squark doublet (× 3 families)	Q	$(\widetilde{u}_L\ \widetilde{d}_L)$	$(u_L\ d_L)$	$(\mathbf{3},\ \mathbf{2},\ \frac{1}{6})$
Right-handed up-type quark/squark singlet (× 3 families)	$\overline{u}$	$\widetilde{u}^*_R$	$u^\dagger_R$	$(\overline{\mathbf{3}},\ \mathbf{1},\ -\frac{2}{3})$
Right-handed down-type quark/squark singlet (× 3 families)	$\overline{d}$	$\widetilde{d}^*_R$	$d^\dagger_R$	$(\overline{\mathbf{3}},\ \mathbf{1},\ \frac{1}{3})$
Left-handed lepton/slepton doublet (× 3 families)	L	$(\widetilde{\overline{\nu}}_{eL}\ \widetilde{e}_L)$	$(\overline{\nu}_{eL}\ e_L)$	$(\mathbf{1},\ \mathbf{2},\ -\frac{1}{2})$
Right-handed lepton/slepton singlet (× 3 families)	$\overline{e}$	$\widetilde{e}^*_R$	$e^\dagger_R$	$(\overline{\mathbf{1}},\ \mathbf{1},\ 1)$
Up-type Higgs/Higgsino doublet	H_u	$(H_u^+\ H_u^0)$	$(\widetilde{H}_u^+\ \widetilde{H}_u^0)$	$(\mathbf{1},\ \mathbf{2},\ \frac{1}{2})$
Down-type Higgs/Higgsino doublet	H_d	$(H_d^0\ H_d^-)$	$(\widetilde{H}_d^0\ \widetilde{H}_d^-)$	$(\mathbf{1},\ \mathbf{2},\ -\frac{1}{2})$

Table 1.2: Gauge supermultiplets of the supersymmetric Standard Model. Adapted from Table 1.2 of [?].

Type of supermultiplet	Spin-1/2 component	Spin-1 component	Representation under $SU(3)_C \otimes SU(2)_L \otimes U(1)_Y$
Gluon/gluino	$\widetilde{g}$	g	$(\mathbf{8}, \mathbf{1}, 0)$
W/wino	$\widetilde{W}^{\pm}\ \widetilde{W}^{0}$	$W^{\pm}\ W^{0}$	$(\mathbf{1}, \mathbf{3}, 0)$
B/bino	$\widetilde{B}^{0}$	B^{0}	$(\mathbf{1}, \mathbf{1}, 0)$

1.2 The Unbroken SUSY Lagrangian

The first piece of the full unbroken SUSY Lagrangian density consists of the kinetic and interacting terms related to the chiral supermultiplets. As explained in Sec. 1.1, a chiral supermultiplet consists of a Weyl fermion ψ (the ordinary fermion) and a complex scalar ϕ (the sfermion). For a collection of such chiral supermultiplets, the Lagrangian is

$$\begin{aligned}\mathcal{L}_{chiral} \;=\; & -\partial^{\mu}\phi^{*i}\partial_{\mu}\phi_i \;-\; V_{chiral}(\phi,\phi^*) \;-\; i\psi^{\dagger i}\overline{\sigma}^{\mu}\partial_{\mu}\psi_i \;-\; \frac{1}{2}M^{ij}\psi_i\psi_j \\ & -\; \frac{1}{2}M^*_{ij}\psi^{\dagger i}\psi^{\dagger j} \;-\; \frac{1}{2}y^{ijk}\phi_i\psi_j\psi_k \;-\; \frac{1}{2}y^*_{ijk}\phi^{*i}\psi^{\dagger j}\psi^{\dagger k}\end{aligned} \tag{1.2}$$

where i runs over all supermultiplets in Table **??**, $\overline{\sigma}^{\mu}$ are -1 $\times$ the Pauli matrices (except for $\sigma^0 = \overline{\sigma}^0$), M^{ij} is a mass matrix for the fermions, y^{ijk} are the Yukawa couplings between one scalar and two spinor fields, and $V_{chiral}(\phi, \phi^*$ is the scalar potential

$$\begin{aligned}V_{chiral}(\phi,\phi^*) \;=\; & M^*_{ik}M^{kj}\phi^{*i}\phi_j \;+\; \frac{1}{2}M^{in}y^*_{jkn}\phi_i\phi^{*j}\phi^{*k} \\ & +\; \frac{1}{2}M^*_{in}y^{jkn}\phi^{*i}\phi_j\phi_k \;+\; \frac{1}{4}y^{ijn}y^*_{kln}\phi_i\phi_j\phi^{*k}\phi^{*l}.\end{aligned} \tag{1.3}$$

The Lagrangian can also be written as the kinetic terms plus derivatives of the *superpotential* W:

$$\begin{aligned}\mathcal{L}_{chiral} \;=\; & -\partial^{\mu}\phi^{*i}\partial_{\mu}\phi_i \;-\; i\psi^{\dagger i}\overline{\sigma}^{\mu}\partial_{\mu}\psi_i \\ & -\; \frac{1}{2}\Big(\frac{\delta^2 W}{\delta\phi^i\delta\phi^j}\psi_i\psi_j + \frac{\delta^2 W^*}{\delta\phi_i\delta\phi_j}\psi^{\dagger i}\psi^{\dagger j}\Big) \;-\; \frac{\delta W}{\delta\phi^i}\frac{\delta W^*}{\delta\phi_i}\end{aligned} \tag{1.4}$$

where

$$W \;=\; M^{ij}\phi_i\phi_j \;+\; \frac{1}{6}y^{ijk}\phi_i\phi_j\phi_k. \tag{1.5}$$

The second part of the Lagrangian involves the gauge supermultiplets. In terms of the spin-1 ordinary gauge boson A^a_μ and the spin-1/2 Weyl spinor gaugino λ^a of the gauge supermultiplet, where a runs over the number of generators for the SM subgroup (i.e. 1-8 for $SU(3)_C$, 1-3 for $SU(2)_L$, and 1 for $U(1)_Y$), this part of the Lagrangian is

$$\mathcal{L}_{gauge} = -\frac{1}{4}F^a_{\mu\nu}F^{\mu\nu a} - i\lambda^{\dagger a}\overline{\sigma}^\mu D_\mu\lambda^a + \frac{1}{2}D^a D^a \tag{1.6}$$

where

$$F^a_{\mu\nu} = \partial_\mu A^a_\nu - \partial_\nu A^a_\mu + gf^{abc}A^b_\mu A^c_\nu \tag{1.7}$$

(g is the coupling constant and f^{abc} are the structure constants for the particular SM gauge group),

$$D_\mu\lambda^a = \partial_\mu\lambda^a + gf^{abc}A^b_\mu\lambda^c, \tag{1.8}$$

and D^a is an auxiliary field that does not propagate (in the literature, it is used as a bookkeeping tool and can be removed via its algebraic equation of motion).

To build a fully supersymmetric and gauge-invariant Lagrangian, the ordinary derivatives in $\mathcal{L}_{chiral}$ (Eq. 1.2) must be replaced by covariant derivatives

$$\begin{aligned} D_\mu\phi_i &= \partial_\mu\phi_i - igA^a_\mu(T^a\phi)_i & (1.9)\\ D_\mu\phi^{*i} &= \partial_\mu\phi^{*i} + igA^a_\mu(\phi^*T^a)^i & (1.10)\\ D_\mu\psi_i &= \partial_\mu\psi_i - igA^a_\mu(T^a\psi)_i. & (1.11)\end{aligned}$$

This leads to the full Lagrangian

$$\begin{aligned} \mathcal{L} &= \mathcal{L}_{chiral} + \mathcal{L}_{gauge}\\ &\quad - \sqrt{2}g(\phi^{*i}T^a\psi_i)\lambda^a - \sqrt{2}g\lambda^{\dagger a}(\psi^{\dagger i}T^a\phi_i) + g(\phi^{*i}T^a\phi_i)D^a\\ &= -\partial^\mu\phi^{*i}\partial_\mu\phi_i - i\psi^{\dagger i}\overline{\sigma}^\mu\partial_\mu\psi_i + ig\partial^\mu\phi^{*i}A^a_\mu(T^a\phi)_i - ig\partial_\mu\phi_i A^{\mu a}(\phi^*T^a)^i\\ &\quad - g^2A^{\mu a}(\phi^*T^a)^iA^a_\mu(T^a\phi)_i - g\psi^{\dagger i}\overline{\sigma}^\mu A^a_\mu(T^a\psi)_i - V_{chiral}(\phi,\phi^*)\\ &\quad - \frac{1}{2}M^{ij}\psi_i\psi_j - \frac{1}{2}M^*_{ij}\psi^{\dagger i}\psi^{\dagger j} - \frac{1}{2}y^{ijk}\phi_i\psi_j\psi_k - \frac{1}{2}y^*_{ijk}\phi^{*i}\psi^{\dagger j}\psi^{\dagger k}\\ &\quad - \frac{1}{4}F^a_{\mu\nu}F^{\mu\nu a} - i\lambda^{\dagger a}\overline{\sigma}^\mu\partial_\mu\lambda^a - ig\lambda^{\dagger a}\overline{\sigma}^\mu f^{abc}A^b_\mu\lambda^c + \frac{1}{2}D^aD^a\\ &\quad - \sqrt{2}g(\phi^{*i}T^a\psi_i)\lambda^a - \sqrt{2}g\lambda^{\dagger a}(\psi^{\dagger i}T^a\phi_i) + g(\phi^{*i}T^a\phi_i)D^a. \end{aligned} \tag{1.12}$$

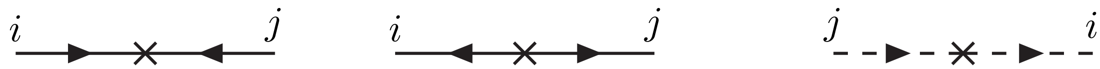
Writing out $F^a_{\mu\nu}$ and $V_{chiral}(\phi,\phi^*)$ explicitly combining the D^a terms using the equation of motion $D^a = -g\phi^*T^a\phi$, and rearranging some terms, the final unbroken SUSY Lagrangian is

$$
\begin{aligned}
\mathcal{L} \;=\; & -\partial^\mu \phi^{*i} \partial_\mu \phi_i \;-\; i\psi^{\dagger i}\overline{\sigma}^\mu \partial_\mu \psi_i \\
& -\frac{1}{4}(\partial_\mu A^a_\nu - \partial_\nu A^a_\mu)(\partial^\mu A^{\nu a} - \partial^\nu A^{\mu a}) \;-\; i\lambda^{\dagger a}\overline{\sigma}^\mu \partial_\mu \lambda^a \\
& - M^*_{ik} M^{kj} \phi^{*i}\phi_j \;-\; \frac{1}{2} M^{ij}\psi_i\psi_j \;-\; \frac{1}{2} M^*_{ij}\psi^{\dagger i}\psi^{\dagger j} \\
& + ig\partial^\mu \phi^{*i} A^a_\mu (T^a\phi)_i \;-\; ig\partial_\mu \phi_i A^{\mu a}(\phi^* T^a)^i \;-\; g\psi^{\dagger i}\overline{\sigma}^\mu A^a_\mu (T^a\psi)_i \\
& - ig\lambda^{\dagger a}\overline{\sigma}^\mu f^{abc} A^b_\mu \lambda^c \\
& - \frac{1}{4} g f^{abc}[(\partial_\mu A^a_\nu - \partial_\nu A^a_\mu) A^{\mu b} A^{\nu c} + A^b_\mu A^c_\nu (\partial^\mu A^{\nu a} - \partial^\nu A^{\mu a})] \\
& - \frac{1}{2} M^*_{in} y^{jkn} \phi^{*i}\phi_j\phi_k \;-\; \frac{1}{2} M^{in} y^*_{jkn} \phi_i \phi^{*j}\phi^{*k} \\
& - \frac{1}{2} y^{ijk}\phi_i\psi_j\psi_k \;-\; \frac{1}{2} y^*_{ijk}\phi^{*i}\psi^{\dagger j}\psi^{\dagger k} \\
& - \sqrt{2} g(\phi^{*i} T^a \psi_i)\lambda^a \;-\; \sqrt{2} g\lambda^{\dagger a}(\psi^{\dagger i} T^a \phi_i) \\
& - g^2 A^{\mu a}(\phi^* T^a)^i A^a_\mu (T^a\phi)_i \;-\; \frac{1}{4} g^2 f^{abc} A^b_\mu A^c_\nu f^{abc} A^{\mu b} A^{\nu c} \\
& - \frac{1}{4} y^{ijn} y^*_{kln} \phi_i\phi_j \phi^{*k}\phi^{*l} \;-\; \frac{1}{2} g^2 (\phi^{*i} T^a \phi_i)^2 .
\end{aligned} \tag{1.13}
$$

The above Lagrangian applies to chiral supermultiplets interacting with one kind of gauge supermultiplet (i.e. one SM gauge group). In the general case, there are additional terms corresponding to interactions with all three SM gauge groups.

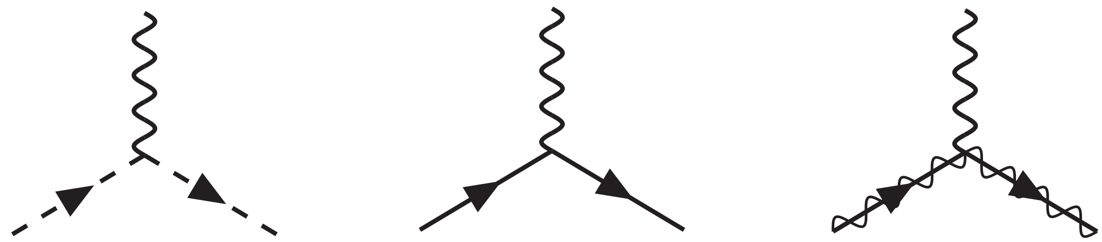
The following list gives a description of the terms in Eq. 1.13:

- First two lines: kinetic terms for the four types of fields ϕ_i, ψ_i, A^a_μ, and λ^a
- Third line: mass terms for the ϕ_i and ψ_i (see Figs. 1.1a and 1.1b)
- Fourth and fifth lines: cubic couplings in which ϕ_i, ψ_i, or λ^a radiates an A^a_μ (see Figs. 1.1c, 1.1d, and 1.1e)
- Sixth line: triple gauge boson couplings (see Fig. 1.1f)
- Seventh line: triple sfermion couplings (see Fig. 1.1g)
- Eighth line: cubic couplings in which ψ_i radiates a ϕ_i (see FIg. 1.1h)
- Ninth line: ϕ_i-ψ_i-λ^a vertices (see FIg. 1.1i)
- 10^{th} line: A^a_μ-A^a_μ-ϕ_i-ϕ_i and quadruple gauge boson couplings (see Figs. 1.1j and 1.1k)
- 11^{th} line: ϕ_i^4 vertices (see Figs. 1.1l and 1.1m)



(a) Fermion mass terms. Reprinted from Figs. 3.2(c) and 3.2(d) of ref. **??**.

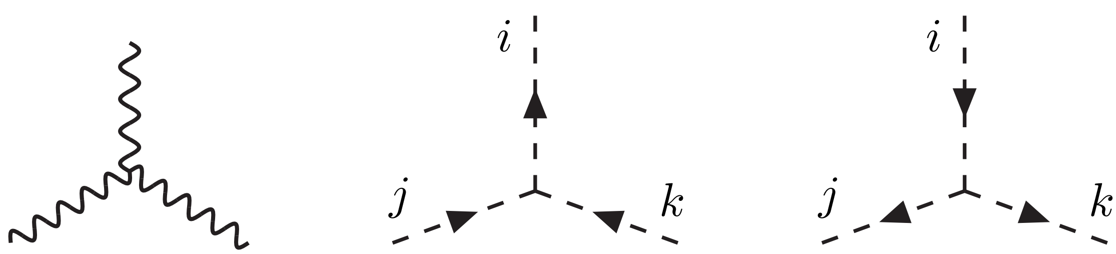
(b) Sfermion mass terms. Reprinted from Fig. 3.2(e) of ref. **??**.



(c) $A\phi\phi$ vertex. Reprinted from Fig. 3.3(e) of ref. **??**.

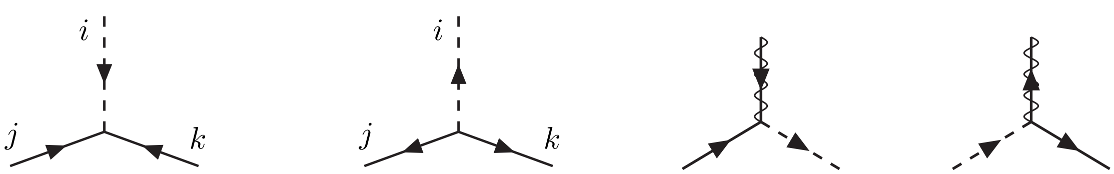
(d) $A\psi\psi$ vertex. Reprinted from Fig. 3.3(f) of ref. **??**.

(e) $A\lambda\lambda$ vertex. Reprinted from Fig. 3.3(c) of ref. **??**.



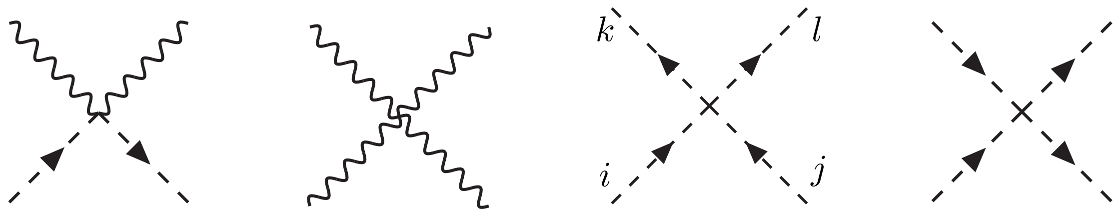
(f) Triple gauge boson vertex. Reprinted from Fig. 3.3(b) of ref. **??**.

(g) Triple sfermion vertices. Reprinted from Figs. 3.2(a) and 3.2(b) of ref. **??**.



(h) Cubic couplings in which the fermion radiates a sfermion. Reprinted from Figs. 3.1(a) and 3.1(b) of ref. **??**.

(i) Sfermion-fermion-gaugino vertices. Reprinted from Figs. 3.3(g) and 3.3(h) of ref. **??**.



(j) $AA\phi\phi$ vertex. Reprinted from Fig. 3.3(d) of ref. **??**.

(k) Quadruple gauge boson vertex. Reprinted from Fig. 3.3(a) of ref. **??**.

(l) Non-gauge ϕ^4 vertex. Reprinted from Fig. 3.1(c) of ref. **??**.

(m) Gauge ϕ^4 vertex. Reprinted from Fig. 3.3(i) of ref. **??**.

Figure 1.1: Interactions in the unbroken SUSY Lagrangian.

1.3 Soft SUSY Breaking

Since quadratic divergences in sfermion masses vanish to all orders in perturbation theory in plain unbroken SUSY[?] due to the presence of gauge and Yukawa interactions with the necessary relationships between coupling constants, it is desirable that the terms that break SUSY not disturb this property. In addition, SUSY should be broken spontaneously, as electroweak symmetry is broken in the Standard Model, so that it is only made manifest at high energies. To satisfy these constraints, SUSY-breaking terms are simply added to the unbroken SUSY Lagrangian in Eq. 1.13 such that $\mathcal{L}_{\rm total} = \mathcal{L}_{\rm unbroken} + \mathcal{L}_{\rm breaking}$. The coefficients of terms in $\mathcal{L}_{\rm breaking}$ must have positive mass dimension in order not to contribute quadratically divergent loop corrections to the scalar masses (like the Higgs mass).[1] This form of SUSY breaking is called *soft*, and all coefficients of soft SUSY breaking terms are expected to of order $\mathrm{m_{soft}}$ or $\mathrm{m^2_{soft}}$.

Soft SUSY breaking terms give masses to the sfermions and gauginos and introduce a cubic sfermion vertex. The soft terms are given by

$$\begin{aligned}
\mathcal{L}_{soft} \;=\; & -\frac{1}{2}(M_3\widetilde{g}^a\widetilde{g}^a + M_2\widetilde{W}^a\widetilde{W}^a + M_1\widetilde{B}\widetilde{B} + \text{h.c.}) \\
& - (a_u^{ij}\widetilde{u}^*_{Ri}\widetilde{Q}_jH_u - a_d^{ij}\widetilde{d}^*_{Ri}\widetilde{Q}_jH_d - a_e^{ij}\widetilde{e}^*_{Ri}\widetilde{L}_jH_d + \text{h.c.}) \\
& - m^2_{\widetilde{Q}ij}\widetilde{Q}_i^\dagger\widetilde{Q}_j \;-\; m^2_{\widetilde{L}ij}\widetilde{L}_i^\dagger\widetilde{L}_j \\
& - m^2_{\widetilde{\overline{u}}ij}\widetilde{u}_{Ri}\widetilde{u}^*_{Rj} \;-\; m^2_{\widetilde{\overline{d}}ij}\widetilde{d}_{Ri}\widetilde{d}^*_{Rj} \;-\; m^2_{\widetilde{\overline{e}}ij}\widetilde{e}_{Ri}\widetilde{e}^*_{Rj} \\
& - m^2_{H_u}H_u^*H_u \;-\; m^2_{H_d}H_d^*H_d \;-\; (bH_uH_d + \text{h.c.})
\end{aligned} \tag{1.14}$$

where a runs from 1-8 for $\widetilde{g}^a$ and from 1-3 for $\widetilde{W}^a$, and i, j run over the three families. The first line of Eq. 1.14 contains the gaugino mass terms. The second line contains cubic scalar couplings that contribute to mixing between the left- and right-handed third generation sfermions (it is assumed in the supersymmetric Standard Model that the a_u^{ij}, a_d^{ij}, and a_e^{ij} are negligible unless $i = j = 3$). The third and fourth lines of Eq. 1.14 contain squark and slepton mass terms, and finally the last line contains the Higgs mass terms.

[1] This point can be argued via dimensional analysis. Radiative corrections take the form $\Delta\mathrm{m_S^2}$, where $\mathrm{m_S}$ is the mass of the scalar particle in question. The dimensions of $\Delta\mathrm{m_S^2}$ are mass^2. $\Delta\mathrm{m_S^2}$ is proportional to some coupling constant or mass coefficient k multiplied by a function of $\Lambda_{\rm UV}$, the high energy cutoff scale. The function of $\Lambda_{\rm UV}$ is determined by a loop integral, and thus typically takes the form $\Lambda^2_{\rm UV}$ (quadratically divergent) or $\ln\frac{\Lambda_{\rm UV}}{\mathrm{m_{low}}}$ (logarithmically divergent, where $\mathrm{m_{low}}$ is some other lower-mass scale in the problem). Now, if k already contributes at least one power of mass to $\Delta\mathrm{m_S^2}$, then the high-energy behavior—the function of $\Lambda_{\rm UV}$—can only contribute at most one power of the dimensionful parameter $\Lambda_{\rm UV}$. However, there are typically no loop integrals that diverge linearly in $\Lambda_{\rm UV}$, so by forcing k to have positive mass dimension, the form of the radiative corrections contributed by SUSY-breaking terms is limited to $\Delta\mathrm{m_S^2} \sim \mathrm{m^2_{low}}\ln\frac{\Lambda_{\rm UV}}{\mathrm{m_{low}}}$. In effect, the possibility of dangerous corrections proportional to $\Lambda^2_{\rm UV}$ is excluded by dimensional analysis if the requirement that k contribute at least one power of mass is enforced.

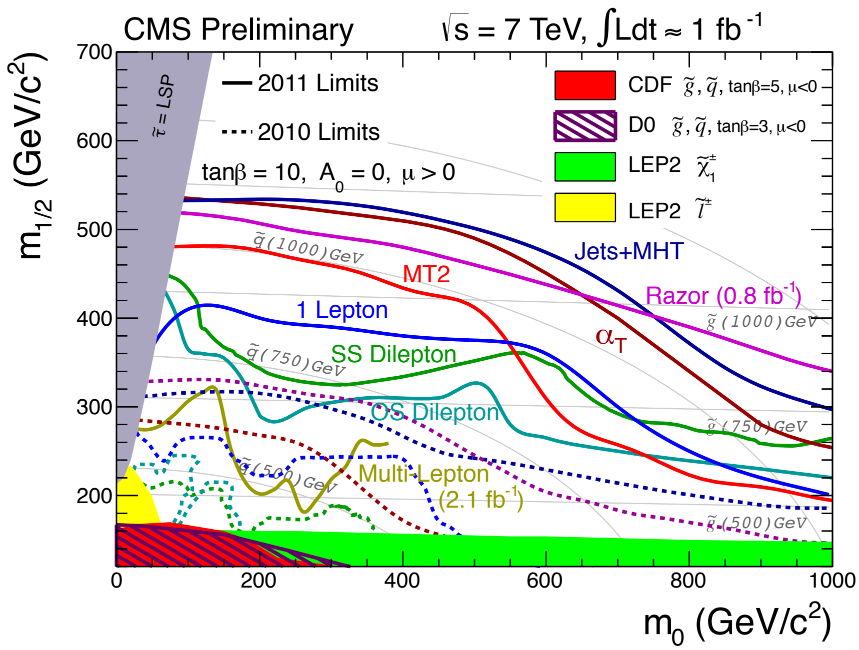


Figure 1.2: CMS limits on mSUGRA with tan $\beta = 10$. The limits set by individual searches are shown as separate colored lines. Solid lines refer to 2011 searches (i.e. using an integrated luminosity of ~1 fb^{-1}), while dashed lines refer to 2010 searches (~36 pb^{-1}). Reprinted from [?].

1.4 Dark Matter and the WIMP Miracle

1.5 Gauge-Mediated SUSY Breaking

1.6 Experimental Status of SUSY

Collider searches for evidence of supersymmetry began in earnest in the 1980s [?] and continue to this day. Most recently, the LHC and Tevatron[2] experiments have set the strictest limits on a variety of SUSY breaking scenarios, including GMSB and mSUGRA (discussed below).

Figure 1.2 shows the current limits set by the CMS experiment on the mSUGRA model (with tan $\beta = 10$) in the m_0-$m_{1/2}$ plane. (Note that although the plot is truncated at $m_0 = 1000$ GeV/c^2, some searches are sensitive out to $m_0 \sim 2000$ GeV/c^2.) Although the LHC has pushed m_0 above ~1 TeV/c^2 for $m_{1/2}$ up to ~400 GeV/c^2, casting some doubt onto the theory's prospects for solving the hierarchy problem, there is still a sizable chunk of mSUGRA parameter space that is not ruled out by collider experiments. Furthermore, parts of the CMS unexplored regions overlap with areas allowed by astrophysics experiments [?].

Figure 1.3 shows the most up-to-date limit (using 1 fb^{-1} of integrated luminosity

[2] Located on the Fermilab site in Batavia, Illinois, the Tevatron was a proton-antiproton collider operating at 1.96 TeV center-of-mass energy. The Tevatron ran from 1987 to 2011 [?].

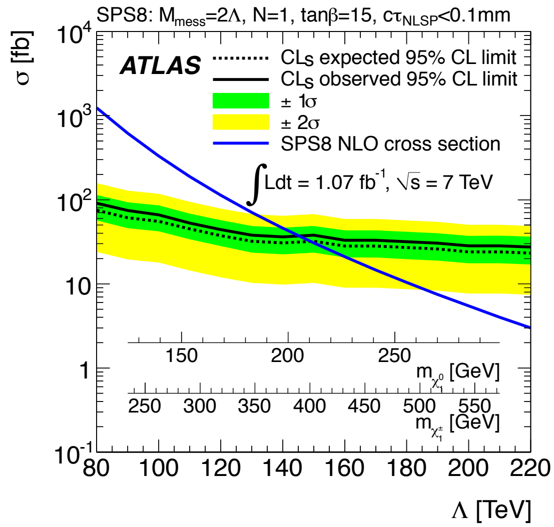


Figure 1.3: ATLAS cross section upper limit on the SPS8 [?] model of mGMSB as a function of SUSY breaking scale Λ, lightest neutralino mass $m_{\tilde{\chi}_1^0}$, or lightest chargino mass $m_{\tilde{\chi}_1^\pm}$. Values of Λ, $m_{\tilde{\chi}_1^0}$, or $m_{\tilde{\chi}_1^\pm}$ below the intersection point between the blue (predicted SPS8 cross section) and black (observed cross section upper limit) curves are excluded. The model parameters listed above the plot are defined in Sec. 1.5. Reprinted from [?].

collected by the ATLAS experiment [?] at the LHC) on the Snowmass Points and Slopes (SPS) model of minimal GMSB (mGMSB), dubbed SPS8 [?]. SPS8 represents the simplest class of GMSB models described in Sec. 1.5. The best limits on a variety of general gauge mediation (GGM) models, from the same ATLAS study, are shown in Figure 1.4. In these models, no assumptions are made about the specific parameters common to many gauge mediation models (e.g. the number of messengers or the relationship between the messenger mass and the SUSY breaking scale). Instead, it is only assumed that the lightest neutralino is light enough to be produced on-shell at the LHC (by setting M_1 and M_2 appropriately, see Sec. ??) and that it decays to a gravitino, that the gravitino is extremely relativistic (mass of order eV-keV), and that the gravitino is stable. The one-dimensional scan over SUSY breaking scales in the SPS8 model (in which the full sparticle spectrum is specified by the model parameters) is replaced by a two-dimensional scan over gluino and lightest neutralino mass in the GGM models (in which all sparticles except the gluino, first- and second-generation squarks, and neutralinos are forced to be at ∼1.5 TeV/c^2, effectively decoupling them from the dynamics that can be probed with 1 fb^{-1} at a 7 TeV/c pp collider).

In general, the lifetime of the lightest neutralino in GMSB models can take on any value between hundreds of nanometers to a few kilometers depending on the mass of the lightest neutralino and the SUSY breaking scale [?]. The search published in [?] (from which Figs. 1.3 and 1.4 are culled) considers only *prompt* neutralino variants, i.e. with neutralino lifetime short enough that the distance traveled by the neutralino before decay cannot be resolved by the detector. The most recent limits on non-

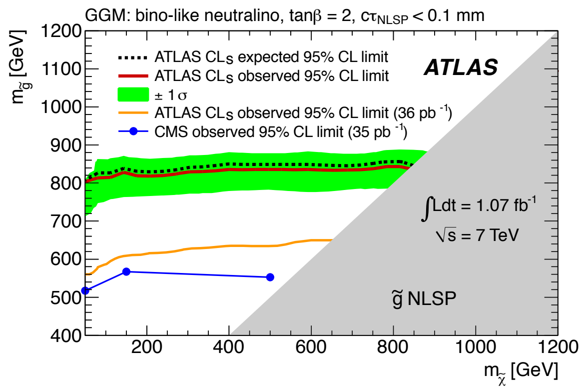


Figure 1.4: ATLAS exclusion contour in the $m_{\tilde{g}}$-$m_{\tilde{\chi}_1^0}$ plane. Values of $m_{\tilde{g}}$-$m_{\tilde{\chi}_1^0}$ below the red curve are excluded. The gray region is theoretically excluded in the GGM models considered. "Bino-like neutralino" means that $M_2 = 1.5$ TeV/c^2. Reprinted from [?].

prompt SPS8-style neutralino models were set by the Collider Detector at Fermilab (CDF) collaboration with 570 pb^{-1}, and are shown in Figure 1.5 [?].

Finally, if the gravitino is to make up some or all of the dark matter, constraints on the form of gauge mediation must come from cosmological considerations and astronomical observations. The gravitino in gauge mediation models is usually very light ($\mathcal{O}$(eV-MeV)) because it is proportional to the SUSY breaking scale divided by the Planck mass, and in GMSB the breaking scale is typically only of order a few hundred TeV ([?] and Sec. **??**). A light, highly relativistic dark matter particle might have been produced, for instance, in the early, radiation-dominated period of the universe [?]. This *warm dark matter* (WDM) may be responsible for all of the dark matter needed to account for galactic structure, or it may share the duties with *cold dark matter* (CDM, the classic WIMPs of Sec. **??**). In any viable model, the predicted relic density of the dark matter species must match the observed value of $\Omega h^2 \sim 0.1$ [?]. For many GMSB models, this measurement constrains the gravitino mass to the keV range [?]. This constraint, however, does not translate into a very strong bound on the lifetime of the lightest neutralino. Using the following equation (taken from [?]):

$$\tau_{\tilde{\chi}_1^0} \sim 130\left(\frac{100\mathrm{GeV}}{m_{\tilde{\chi}_1^0}}\right)^5\left(\frac{\sqrt{F}}{100\mathrm{TeV}}\right)^4\ \mu\mathrm{m} \tag{1.15}$$

where $\sqrt{F}$ is approximately the SUSY breaking scale, and applying the gravitino mass constraint $\sqrt{F} \lesssim 3000$ TeV (cf. Eq. X with $m_{\tilde{G}} \sim$ keV) and $m_{\tilde{\chi}_1^0} = 100$ GeV, the upper bound on the neutralino lifetime is 100 meters. For $\sqrt{F} \sim 100$ TeV, the neutralino lifetime is detectable on collider time scales.

Recently, a lower bound on the WDM particle mass in either pure warm or mixed warm and cold dark matter scenarios was set using observations of the Lyman-α forest. For pure WDM, $m_{WDM} > 8$ keV, while for some mixed WDM-CDM scenarios, $m_{WDM} > 1.1$-1.5 keV [?, ?]. These bounds and others have motivated the development

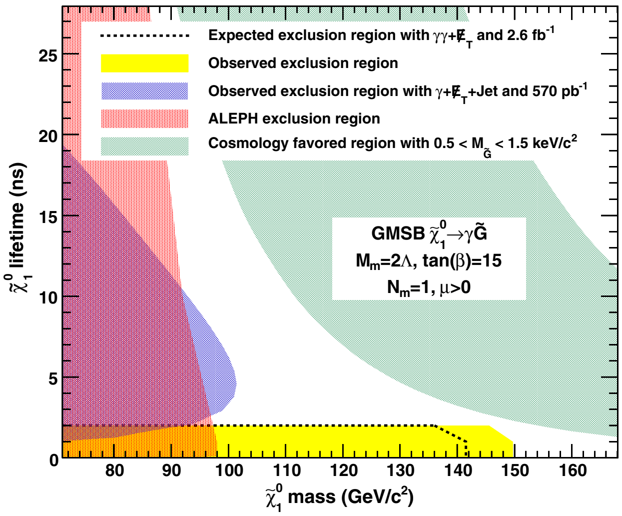


Figure 1.5: CDF exclusion contour in the $\tau_{\tilde{\chi}_1^0}$-$\mathrm{m}_{\tilde{\chi}_1^0}$ plane, where $\tau_{\tilde{\chi}_1^0}$ is the lifetime of the neutralino. Reprinted from [**?**].

of more complicated gauge mediation models [**?**]. However, rather than focus on a specific GMSB model, of which there are many, the search detailed here is interpreted in a minimally model dependent way. With this approach, the results can be applied to many competing models. The remainder of this thesis is devoted to the experimental details of the search, analysis strategy, and presentation of the results.