BAYESUVIUS QUANTICO

a visual dictionary of Quantum Bayesian Networks



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Bayesuvius Quantico,

a visual dictionary of Quantum Bayesian Networks

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This book is constantly being expanded and improved. To download the latest version, go to

https://github.com/rrtucci/bayes-quantico

Bayes Quantico

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Appendices

Appendix A

Spectral Decomposition and Eigenvalue Projection Operators: COMING SOON

 $M \in \mathbb{C}^{d \times d}$

$$M|v\rangle = \lambda|v\rangle \tag{A.1}$$

If M is Hermitian $(H^{\dagger} = H)$, its eigenvalues are real. $(\lambda = \langle \lambda | M \lambda \rangle \in \mathbb{R})$

$$cp(\lambda) \stackrel{\text{def}}{=} \det(M - \lambda) = 0$$
 (A.2)

If M is a Hermitain matrix, then there exists a unitary matric ($CC^{\dagger}=C^{\dagger}C=1$) such that

$$CMC^{\dagger} = \begin{bmatrix} D_{\lambda_1} & 0 & 0 & 0 \\ 0 & D_{\lambda_2} & 0 & 0 \\ 0 & 0 & \ddots & 0 \\ 0 & 0 & 0 & D_{\lambda} \end{bmatrix}$$
 (A.3)

where

$$D_{\lambda_i} = \operatorname{diag}\underbrace{(\lambda_i, \lambda_i, \dots, \lambda_i)}_{d_i \text{ times}} \tag{A.4}$$

$$d = \sum_{i=1}^{r} d_i \tag{A.5}$$

$$CMC^{\dagger} = \begin{bmatrix} \lambda_1 & 0\\ 0 & \lambda_2 \end{bmatrix} \tag{A.6}$$

$$CP_1C^{\dagger} = \begin{bmatrix} 1 & 0 \\ 0 & 0 \end{bmatrix} = \frac{CMC^{\dagger} - \lambda_2}{\lambda_1 - \lambda_2}$$
 (A.7)

$$CP_2C^{\dagger} = \begin{bmatrix} 0 & 0 \\ 0 & 1 \end{bmatrix} = \frac{CMC^{\dagger} - \lambda_1}{\lambda_2 - \lambda_1}$$
 (A.8)

If $I^{d_i \times d_i}$ is the d_i dimensional unit matrix,

$$P_i = C^{\dagger} diag(0, \dots, 0, I^{d_i \times d_i}, 0, \dots, 0)C \tag{A.9}$$

$$= \prod_{j \neq i} \frac{M - \lambda_j}{\lambda_i - \lambda_j} \tag{A.10}$$

Note that P_i are Hermitian $(P_i^{\dagger} = P_i)$ because M is Hermitian and its eigenvalues are real.)

Note that P_i and M commute

$$[P_i, M] = P_i M - M P_i = 0 (A.11)$$

orthogonal

$$P_i P_i = \delta(i, j) P_i \tag{A.12}$$

complete

$$\sum_{i} P_i = 1 \tag{A.13}$$

$$M = \sum_{i=1}^{r} P_i M P_i \tag{A.14}$$

$$d_i = \operatorname{tr} P_i \tag{A.15}$$

$$CMP_1C^{\dagger} = \begin{bmatrix} \lambda_1 & 0 \\ 0 & \lambda_2 \end{bmatrix} \begin{bmatrix} 1 & 0 \\ 0 & 0 \end{bmatrix}$$
 (A.16)

$$= \lambda_1 \begin{bmatrix} 1 & 0 \\ 0 & 0 \end{bmatrix} \tag{A.17}$$

$$MP_i = \lambda_i P_i \text{ (no } i \text{ sum)}$$
 (A.18)

$$f(M)P_i = f(\lambda_i)P_i \text{ (no } i \text{ sum)}$$
 (A.19)

 $M^{(1)}, M^{(2)}$

$$[M^{(1)}, M^{(2)}] = 0 (A.20)$$

Use $M^{(1)}$ to decompose V into $\bigoplus_i V_i$. Use $M^{(2)}$ to decompose V_i into $\bigoplus_j V_{i,j}$. If $M^{(1)}$ and $M^{(2)}$ don't commute, let $P_i^{(1)}$ be the eigenvalue projection operators of $M^{(1)}$. The replace $M^{(2)}$ by $P_i^{(1)}M^{(2)}P_i^{(1)}$

$$[M^{(1)}, P_i^{(1)}M^{(2)}P_i^{(1)}] = 0 (A.21)$$

Appendix B

Birdtracks: COMING SOON

Cvitanovic Birdtracks book [1]

Elliott-Dawber book [2]

My paper "Quantum Bayesian Nets" [3]

B.1 Classical Bayesian Networks and their Instantiations

TPM (Transition Probability Matrix) $P(y|x) \in [0,1]$ where $x \in val(\underline{x})$ and $y \in val(y)$

$$\sum_{y \in val(\underline{y})} P(y|x) = 1 \tag{B.1}$$

$$C = \underbrace{\frac{b}{\underline{c}}}_{\underline{a}} \tag{B.2}$$

$$C(a,b,c) = P(c|b,a)P(b|a)P(a) = b$$

$$C \leftarrow a$$

$$P(a)$$

$$(B.3)$$

$$a^{:2} = (a_1, a_2)$$

$$C' = \underbrace{\frac{b}{\underline{a}_1}}_{\underline{a}_2 - \underline{a}_2 - \underline{a}_2$$

$$C'(a^{:2}, b, c) = P(c|b, a_2)P(a_2|a^{:2})P(b|a_1)P(a_1|a^{:2})P(a^{:2}) =$$

$$c = a_2$$

$$a^{:2}$$
(B.5)

Marginalizer nodes \underline{a}_1 and \underline{a}_2 have the TPMs

$$P(a_i'|\underline{a}^{:2} = (a_1, a_2)) = \delta(a_i', a_i)$$
(B.6)

for i = 1, 2

B.2 Quantum Bayesian Networks and their Instantiations

TPM (Transition Probability Matrix) $A(y|x) \in \mathbb{C}$ where $x \in val(\underline{x})$ and $y \in val(y)$

$$\sum_{y \in val(\underline{y})} |A(y|x)|^2 = 1 \tag{B.7}$$

$$Q = \underbrace{\frac{b}{c}}_{\underline{a}}$$
 (B.8)

$$Q(a,b,c) = A(c|b,a)A(b|a)A(a) =$$

$$c = a$$
(B.9)

$$a^{:2} = (a_1, a_2)$$

$$Q' = \underbrace{\frac{b}{\underline{a}_1}}_{\underline{a}_2 \underline{\underline{a}_2 \underline{a}_2 \underline{a}$$

$$Q'(a^{:2}, b, c) = A(c|b, a_2)A(a_2|a^{:2})A(b|a_1)A(a_1|a^{:2})A(a^{:2}) = b$$

$$c \stackrel{b}{\longleftarrow} a_1 \qquad A(a^{:2})$$
(B.11)

Marginalizer nodes \underline{a}_1 and \underline{a}_2 have the TAMs

$$A(a_i'|\underline{a}^{:2} = (a_1, a_2)) = \delta(a_i', a_i)$$
(B.12)

for i = 1, 2

B.3 Birdtracks

$$\delta(b, a) = \mathbb{1}(a = b) = \delta_a^b = a - b \tag{B.13}$$

$$\underline{a} = a \longleftarrow X_{\underline{a}\underline{b}}^{\underline{c}\underline{d}}$$

$$\langle a, b | X_{\underline{a}\underline{b}}^{\underline{c}\underline{d}} | c, d \rangle = X_{ab}^{\underline{c}d} = b$$

$$\underline{c} = c$$

$$d = d$$
(B.14)

$$a \longleftarrow X_{\underline{ab}} \stackrel{cd}{=} \qquad a, b \longleftarrow X_{\underline{ab}} \stackrel{cd}{=}$$

$$b \qquad \qquad \rightarrow \qquad a, b \qquad \qquad (B.15)$$

$$c \qquad \qquad c \qquad \qquad d$$

 $X_{\underline{a}\underline{b}} \stackrel{cd}{\in} V^2 \otimes V_2$. Sometimes, we will omit denote this node simply by X. This if okay as long as we are not using, X to also denote a different version of $X_{\underline{a}\underline{b}} \stackrel{cd}{=}$ with some of the indices raised or lowered or their order has been changed. ¹

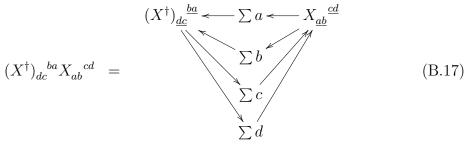
$$(X^{\dagger})_{\underline{dc}} \xrightarrow{\underline{ba}} \underbrace{\underline{a} = a}$$

$$(X^{\dagger})_{\underline{dc}} \xrightarrow{ba} = \underline{\underline{b}} = b$$

$$\underline{\underline{c}} = c$$

$$d = d$$
(B.16)

¹For matrices, $(A^{\dagger})_{i,j} = (A_{j,i})^*$ so taking a Hermitian conjugate involves both taking the complex conjugate of the matrix element and reversing the left-to-right (L2R) order of its indices. This generalizes to $(X^{\dagger})_{dc}^{\ \ ba} = (X_{ab}^{\ \ cd})^*$. Besides raising and lowering indices, we reverse their L2R order.



$$= X^{\dagger} - X$$

$$= (B.18)$$

Birdtracks originated as a graphical way to represent the tensors in General Relativity (Gravitation). In General Relativity, one deals with tensors such as $T_{a\ c}^{\ b}$ which have some indices raised and some lowered. One can use the metric $g^{a,b}$ to raise all the lowered indices to get T^{abc} . If we represent this graphically as a node with incoming arrows a, b, c, we need to follow one of the following 2 conventions: either

- 1. label the arrows as \underline{a} , \underline{b} , \underline{c} , and define the node as $T^{\underline{abc}}$, or
- 2. instead of labelling the arrows explicitly $\underline{a}, \underline{b}, \underline{c}$, indicate in the node where is the first arrow \underline{a} , and draw the arrows $\underline{a}, \underline{b}, \underline{c}$ so that they enter the node in **counterclockwise** (CC) order. The **left-to-right** (L2R) order of the indices on T corresponds the CC order of the arrows.

If we don't do either 1 or 2, we won't be able to distinguish between the graphical representations of $T^{1,2,3}$ and $T^{2,1,3}$, for example. Cvitanovic's Birdtracks book Ref.[1] follows Convention 2, but most of the time, in this book, we will follow Convention 1 ² The reason I chose to do so is for the sake of consistency: Convention 2 is closer to the quantum bnet conventions.

Another issue that arises in using birdtracks is this. When is it permissible to represent a tensor by T_{ab}^{cd} ? If we define T_{ab}^{cd} by

$$T_{ab}^{cd} = T_{ab}^{\quad cd} \tag{B.19}$$

then it's always permissible. Then one can define tensors like $T_a{}^{bcd}$ as

$$T_a^{\ bcd} = g^{bb'} T_{ab'}^{\ cd} = g^{bb'} T_{ab'}^{cd} \tag{B.20}$$

²If we follow Convention 1, we don't need to reverse the L2R order of the indices when taking a Hermitian conjugate. Thus, $(X^{\dagger})^{\underline{ab}}_{\underline{cd}} = X_{\underline{ab}}^{\underline{cd}} = X_{\underline{ba}}^{\underline{dc}}$. As long as $\underline{a}, \underline{b}$ are lower indices and $\underline{c}, \underline{d}$ are upper indices of X, any L2R order of $\underline{a}, \underline{b}, \underline{c}, \underline{d}$ is equivalent under Convention 1.

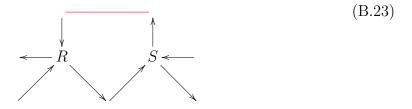
Hence, one drawback of using the notation T_{ab}^{cd} is that if one is interested in using versions of T_{ab}^{cd} with some indices raised or lowered, one has to write down explicitly the metric tensors that do the lowering and raising. Instead of writing $T_a^{\ bcd}$, you'll have to write $g^{bb'}T_{ab'}^{cd}$. This is not very onerous when explaining a topic in which not much lowering and raising of indices is done. But in topics like General Relativity that do use a lot of raising and lowering of indices, it might not be too elegantly concise.

$$a^{:m} \in \mathbb{Z}_+^m$$

$$R_{b_{3}^{im_{3}},a_{2}^{im_{2}}}^{im_{3},b_{2}^{in_{2}}}S_{b_{2}^{in_{2}},a_{1}^{im_{1}}}^{im_{1}} = b_{3}^{in_{3}} \underbrace{-\sum b_{2}^{in_{2}}}_{S^{in_{2}}} S \underbrace{-\sum b_{2}^{in_{2}}}_{S^{in_{1}}} S \underbrace{-\sum b_{2}^{in_{1}}}_{S^{in_{1}}} S \underbrace{-\sum a_{2}^{in_{1}}}_{S^{in_{1}}} S \underbrace{-\sum a_{2}^{in_{1}}$$

$$\operatorname{tr}_{\underline{b}} X_{a\underline{b}}^{\underline{b}\underline{d}} = \sum_{b} X_{ab}^{bd} = \begin{pmatrix} A & A_{\underline{a}\underline{b}}^{\underline{c}\underline{d}} \\ A & A_{\underline{a}\underline{b}}^{\underline{b}\underline{d}} \end{pmatrix}$$

$$(B.22)$$



Casimir Operators: COMING SOON

Clebsch-Gordan Coefficients: COMING SOON

$$\begin{bmatrix} 0 \\ C_{\lambda}^{d_{\lambda} \times d} \\ 0 \end{bmatrix}^{d \times d} = \begin{bmatrix} 0 & 0 & 0 \\ 0 & I^{d_{\lambda} \times d_{\lambda}} & 0 \\ 0 & 0 & 0 \end{bmatrix}^{d \times d} C^{d \times d}$$
 (2.1)

Let $b^{:nb} = (b_1, b_2, \dots, b_{nb})$ where $b_i \in Z_{[0,db_i]}$ and $a \in Z_{[1,d_{\lambda}]}$. Hence,

$$d_{\lambda} = \prod_{i=1}^{:nb} db_i \tag{2.2}$$

$$(C_{\lambda})_{a}^{b:nb} = a \longleftarrow C_{\lambda} \longleftarrow b_{2}$$

$$b_{nb}$$

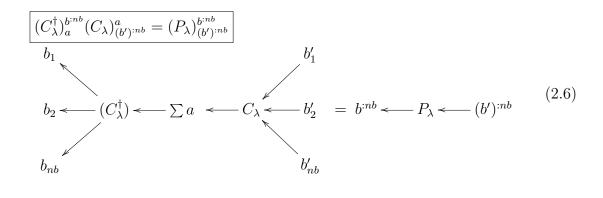
$$(2.3)$$

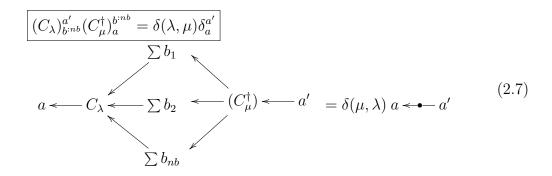
$$\begin{bmatrix} 0 & (C^{\dagger})_{\lambda}^{d \times d_{\lambda}} & 0 \end{bmatrix}^{d \times d} = (C^{\dagger})^{d \times d} \begin{bmatrix} 0 & 0 & 0 \\ 0 & I^{d_{\lambda} \times d_{\lambda}} & 0 \\ 0 & 0 & 0 \end{bmatrix}^{d \times d}$$
(2.4)

$$(C_{\lambda}^{\dagger})_{b:nb}^{a} = b_{2} \longleftarrow (C_{\lambda}^{\dagger}) \longleftarrow a \tag{2.5}$$

$$b_{nb}$$

More generally, some of the b_i indices may lowered and their arrows changed to outgoing instead of ingoing. Each b_i represents a different rep (or irrep)





Determinants: COMING SOON

General Relativity Nets: COMING SOON

Group Integrals: COMING SOON

Invariants: COMING SOON

Levi-Civita Tensor

$$\epsilon^{123...p} = \epsilon_{123...p} = 1 \tag{7.1}$$

$$\epsilon_{rev(a^{:p})} = (-1)^{\binom{p}{2}} \epsilon_{a^{:p}} \tag{7.2}$$

where $rev(a^{p})$ is the reverse of a^{p} . $rev(a_1, a_2, \ldots, a_p) = (a_p, a_{p-1}, \ldots, a_1)$

$$(C_{\mathcal{A}_p})_1^{a^{:p}} = e^{i\phi} \frac{\epsilon^{a^{:p}}}{\sqrt{p!}} = \mathcal{A}_p \leftarrow a_1$$

$$= a_2$$

$$\vdots$$

$$(C_{\mathcal{A}_{p}}^{\dagger})_{a:p}^{1} = e^{-i\phi} \frac{\epsilon_{a:p}}{\sqrt{p!}} = a_{1} \leftarrow \mathcal{A}_{p}$$

$$a_{2} \leftarrow \parallel$$

$$\vdots$$

$$\underbrace{e^{i2\phi} \frac{1}{p!} \epsilon^{a^{:n}} \epsilon_{a^{:n}} = \delta_1^1 = 1}_{\qquad \vdots \qquad \qquad = 1 \qquad (7.6)$$

For Convention 1, we will use $\phi = 0$. For Convention 2, we must choose

$$e^{i2\phi} = (-1)^{\binom{p}{2}} = e^{i\pi\frac{p(p-1)}{2}}$$
 (7.7)

SO

$$\phi = \frac{\pi}{4}p(p-1) \tag{7.8}$$

Lie Algebra Definition: COMING SOON

 $i \in \mathbb{Z}_{[1,N]}, a, b \in \mathbb{Z}_{[1,n]}$

$$(C_{Adj}^{i})_{b}^{a} = \frac{1}{\sqrt{K}} (T^{i})_{b}^{a} = i \sim C_{Adj}^{i}$$

$$\downarrow$$

$$b$$
(8.1)

The matrices T^i are called the generators. It's customary to choose them so that they are Hermitian and $K=\frac{1}{2}$

$$\underbrace{(T^i)_b^a (T^j)_a^b = \operatorname{tr}(T^i T^j) = K\delta(i,j)}_{\sum a} \quad i \leadsto T^j \leadsto j = K \longleftrightarrow (8.2)$$

$$(P_{Adj})_{b,d}^{a,c} = \sum_{i} \frac{1}{K} (T^{i})_{b}^{a} (T^{i})_{d}^{c} = \frac{1}{K} \left| \begin{matrix} a & d \\ \\ \\ b & c \end{matrix} \right|$$
(8.3)

 $H \in V^{\underline{a}} \otimes V_{\underline{a}}$

$$(P_{Adj})_{bd}^{ac}H_c^d = \sum_i (T^i)_b^a \underbrace{\left[\frac{1}{K}(T^i)_d^c H_c^d\right]}_{h_i \in \mathbb{R}}$$
(8.4)

$$G = 1 + iD \in \mathcal{G}$$

$$\epsilon_i \in \mathbb{R}, \ |\epsilon_i| << 1$$

$$D = \sum_{i} \epsilon_i T^i = V^{\underline{a}} \otimes V_{\underline{a}}$$

$$\mathcal{T}^i q = 0 \tag{8.5}$$

$$(\mathcal{T}^i)_{b^{nb}c^{na}}^{a^{na}d^{nb}} = \tag{8.6}$$

Lie Algebra Classification, Dynkin Diagrams: COMING SOON

Orthogonal Groups: COMING SOON

Quantum Shannon Information Theory: COMING SOON

Recoupling Equations: COMING SOON

Reducibility: COMING SOON

Spinors: COMING SOON

Squashed Entanglement: COMING SOON

Symplectic Groups: COMING SOON

Symmetrization and Antisymmetrization: COMING SOON

(1,2) transposition, swaps 1 and 2, $1 \to 2 \to 1$. (3,2,1) means $3 \to 2 \to 1 \to 3$. A reordering of $(1,2,3,\ldots,p)$ is a permutation on p letters. A permutation can be expressed as a product of transpositions (3,2,1)=(3,2)(2,1) is an even permutation because it can be expressed as a product of an even number of transpositions. An odd permutation can be expressed as a product of an odd number of permutations.

17.1 Symmetrization

$$\mathbb{1}_{a_1,a_2}^{b_2,b_1} = \delta_{a_1}^{b_1} \delta_{a_2}^{b_2} = a_1 \leftarrow b_1$$

$$a_2 \leftarrow b_2$$

$$(17.1)$$

$$(\sigma_{(1,2)})_{a_1,a_2}^{b_2,b_1} = \delta_{a_1}^{b_2} \delta_{a_2}^{b_1} = \begin{pmatrix} a_1 \leftarrow \bullet \leftarrow b_1 \\ \downarrow \\ a_2 \leftarrow \bullet \leftarrow b_2 \end{pmatrix}$$
 (17.2)

$$1 = \leftarrow (17.3)$$

$$\sigma_{(1,2)} = \begin{array}{c} & \longleftarrow & \longleftarrow & \longleftarrow \\ & \downarrow & \\ & \downarrow & \\ & \longleftarrow & \\ & (17.4)$$

$$\sigma_{(1,3,2)} = \underbrace{\hspace{1cm}}_{\bullet} \underbrace{\hspace{1cm}}_{\bullet}$$

Claim 1

proof: We only prove it for p = 3.

QED

$$=\frac{n+p-1}{p} \begin{pmatrix} \leftarrow \mathcal{S}_{p-1} \leftarrow \\ \leftarrow \\ \vdots \\ \vdots \\ \leftarrow \end{pmatrix}$$

$$(17.17)$$

$$\operatorname{tr}_{\underline{a}_{1}} \mathcal{S}_{p} = \frac{n+p-1}{p} \mathcal{S}_{p-1}$$
 (17.18)

$$\operatorname{tr}_{\underline{a}_1,\underline{a}_2,\dots,\underline{a}_k} \mathcal{S}_p = \frac{(n+p-1)(n+p-2)\dots(n=p-k)}{p(p-1)\dots(p-k+1)} \mathcal{S}_{p-k}$$
 (17.19)

$$d_{S_p} = \operatorname{tr}_{\underline{a}^p} S_p = \frac{(n+p-1)!}{p!(n-1)!} = \binom{n+p-1}{p}$$
 (17.20)

For p=2,

$$d_{\mathcal{S}_2} = \frac{(n+1)n}{2} \tag{17.21}$$

17.2 Antisymmetrization

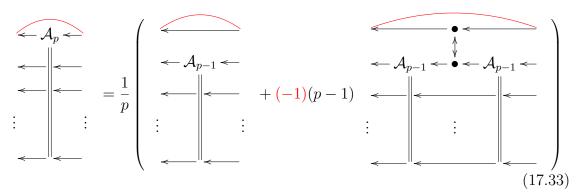
$$\begin{aligned}
S_{p}A_{[1,q]} &= A_{p}S_{[1,q]} = 0 \\
&\leftarrow S_{p} \leftarrow &\leftarrow A_{[1,q]} \leftarrow &\leftarrow &\leftarrow S_{[1,q]} \leftarrow \\
&\leftarrow &\parallel \leftarrow &\leftarrow &\parallel \leftarrow &\leftarrow &\parallel \leftarrow \\
&\leftarrow &\parallel \leftarrow &\leftarrow &\parallel \leftarrow &\leftarrow &\parallel \leftarrow \\
&\leftarrow &\parallel \leftarrow &\leftarrow &\parallel \leftarrow &\leftarrow &\parallel \leftarrow &\leftarrow \\
&\leftarrow &\parallel \leftarrow &\leftarrow &\leftarrow &\parallel \leftarrow &\leftarrow &\parallel \leftarrow &\leftarrow \\
&\vdots &\vdots &\vdots &\vdots &\vdots &\vdots &\vdots
\end{aligned}$$

$$(17.27)$$

Claim 2

proof: We only prove it for p = 3.

QED



$$\operatorname{tr}_{\underline{a}_1} \mathcal{A}_p = \frac{n-p+1}{p} \mathcal{A}_{p-1} \tag{17.35}$$

$$\operatorname{tr}_{\underline{a}_{1},\underline{a}_{2},\dots,\underline{a}_{k}} \mathcal{A}_{p} = \frac{(n-p+1)(n-p+2)\dots(n-p+k)}{p(p-1)\dots(p-k+1)} \mathcal{A}_{p-k}$$
 (17.36)

$$d_{\mathcal{A}_p} = \operatorname{tr}_{\underline{a}^p} \mathcal{A}_p = \frac{\prod_{i=n-p+1}^n i}{p!}$$
(17.37)

$$= \frac{\prod_{i=n}^{n-p+1} i}{p!} \tag{17.38}$$

$$= \begin{cases} \frac{n!}{p!(n-p)!} = \binom{n}{p} & \text{if } p \le n\\ 0 & \text{otherwise} \end{cases}$$
 (17.39)

For $p = 2 \le n$,

$$d_{\mathcal{A}_2} = \binom{n}{2} \tag{17.40}$$

$$\mathcal{A}_p = 0 \text{ if } n$$

For example, for n=2 and p=3

$$\mathcal{A}_{3}|a,a,b\rangle = \frac{1}{6} \begin{pmatrix} |a,a,b\rangle + |a,b,a\rangle + |b,a,a\rangle \\ -|a,b,a\rangle - |a,a,b\rangle - |b,a,a\rangle \end{pmatrix}$$

$$= 0$$

$$(17.43)$$

Unitary Groups: COMING SOON

SU(n) 18.1

$$m(p,q) = \delta_b^a \sum_{a=1}^n (p_a)^* q_a$$
 (18.1)

$$d \leftarrow c$$

$$\mathbb{1}_{d,b}^{a,c} = \delta_b^a \delta_d^c =$$

$$a \rightarrow b$$
(18.2)

$$\uparrow\downarrow_{d,b}^{a,c} = \delta_d^a \delta_b^c = \begin{pmatrix} d & c \\ \uparrow & \downarrow \\ a & b \end{pmatrix}$$
 (18.3)

$$P_i = \sum_{j \neq i} \frac{M - \lambda_j}{\lambda_i - \lambda_j} \tag{18.5}$$

 $\lambda_1 = n$

$$P_{1} = \frac{\uparrow \downarrow -n}{0-n} = 1 - \frac{1}{n} \uparrow \downarrow$$

$$c \qquad d \qquad e \qquad b \qquad a \qquad b \qquad b \qquad d$$

$$c \leftarrow -\frac{1}{n} \uparrow \qquad d$$

$$\lambda_{2} = 0 \qquad (18.6)$$

$$\operatorname{tr} P_{1} = \frac{1}{n} \left(\frac{18.8}{18.9} \right)$$

$$= n^{2} - 1 \qquad (18.9)$$

$$= n^2 - 1 (18.9)$$

$$trP_2 = \frac{1}{n}$$

$$(18.10)$$

$$= 1 \tag{18.11}$$

$$(T_i)_a^b = i \sim T_i$$

$$\downarrow$$

$$a$$
(18.12)

$$T_i^{\dagger} = T_i \tag{18.13}$$

Claim 3

$$C_F \delta_a^b = (T_i T_i)_a^b = \frac{n^2 - 1}{n} \delta_a^b$$
 (18.14)

proof:

$$(T_{i}T_{i})_{a}^{b} = \sum_{i} i \sim T_{i} \qquad T_{i} \sim i$$

$$= \sum_{i} i \sim T_{i} \qquad T_{i} \sim i$$

$$(18.15)$$

$$= \sum_{i} i \sim T_{i} T_{i} \sim i$$
 (18.16)

QED

Wigner Coefficients: COMING SOON

Wigner-Ekart Theorem: COMING SOON

Young Tableau: COMING SOON

Bibliography

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