# BAYESUVIUS QUANTICO

a visual dictionary of Quantum Bayesian Networks



ROBERT R. TUCCI

### Bayesuvius Quantico,

a visual dictionary of Quantum Bayesian Networks

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This book is constantly being expanded and improved. To download the latest version, go to

https://github.com/rrtucci/bayes-quantico

#### **Bayes Quantico**

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## Appendix A

# Notational Conventions and Preliminaries

#### A.1 Group

A group  $\mathcal{G}$  is a set of elements with a multiplication map  $\mathcal{G} \times \mathcal{G} \to \mathcal{G}$  such that

1. the multiplication is **associative**; i.e.,

$$(ab)c = a(bc) \tag{A.1}$$

for  $a, b, c \in \mathcal{G}$ .

2. there exists an **identity element**  $e \in \mathcal{G}$  such that

$$ea = ae = a \tag{A.2}$$

for all  $a \in \mathcal{G}$ 

3. for any  $g \in \mathcal{G}$ , there exists an **inverse**  $a^{-1} \in \mathcal{G}$  such that

$$aa^{-1} = a^{-1}a = e (A.3)$$

The number of elements in any set S is denoted by |S|.  $|\mathcal{G}|$  is called the **order** of the group.

If multiplication is **commutative** (i.e., ab = ba for all  $a, b \in \mathcal{G}$ , the group is said to be **abelian**.

A subgroup  $\mathcal{H}$  of  $\mathcal{G}$  is a subset of  $\mathcal{G}$  ( $\mathcal{H} \subset \mathcal{G}$ ) which is also a group. It's easy to show that any  $\mathcal{H} \subset \mathcal{G}$  is a group if it contains the identity and is **closed under multiplication** (i.e.,  $ab \in \mathcal{H}$  for all  $a, b \in \mathcal{H}$ )

#### A.2 Group Representation

A group representation of a group  $\mathcal{G}$  is a map  $\phi: \mathcal{G} \to \mathbb{C}^{n \times n1}$  such that

$$\phi(a)\phi(b) = \phi(ab) \tag{A.4}$$

Such a map is called a **homomorphism**. When a group is defined using matrices, those matrices are called the **defining representation**. The map  $\phi$  partitions  $\mathcal{G}$  into disjoints subsets (equivalence classes), such that all elements of  $\mathcal{G}$  in a disjoint set are represented by the same matrix.

For example, the group of **General Linear Transformations** is defined by

$$GL(n, \mathbb{C}) = \{ M \in \mathbb{C}^{n \times n} : \det M \neq 0 \}$$
 (A.5)

#### A.3 Vector Space and Algebra over a field $\mathbb{F}$

A vector (or linear) space  $\mathcal{V}$  is defined as a set endowed with two operations: vector addition  $+: \mathcal{V} \times \mathcal{V} \to \mathcal{V}$ , and scalar multiplication  $\mathbb{F} \times \mathcal{V} \to \mathcal{V}$ , such that

- $\mathcal{V}$  is an abelian group under + with identity 0 and inverse of  $x \in \mathcal{V}$  equal to  $-x \in \mathcal{V}$
- For  $\alpha, \beta \in \mathbb{F}$  and  $x, y \in \mathcal{V}$

$$\alpha(x+y) = \alpha x + \alpha y \tag{A.6}$$

$$(\alpha + \beta)x = \alpha x + \beta y \tag{A.7}$$

$$\alpha(\beta x) = (\alpha \beta)x \tag{A.8}$$

$$1x = x \tag{A.9}$$

$$0x = 0 (A.10)$$

In this book, we will always use either  $\mathbb{C}$  or  $\mathbb{R}$  for  $\mathbb{F}$ . Both of these fields are infinite but some fields are finite.

An algebra  $\mathcal{A}$  is a vector space which, besides being endowed with vector addition and scalar multiplication with which all vector spaces are, it has a **bilinear vector product**. A bilinear vector product is a product that is linear on both sides; i.e.,

$$(\alpha x + \beta y) \cdot z = \alpha x \cdot z + \beta y \cdot z \tag{A.11}$$

and

$$z \cdot (\alpha x + \beta y) = \alpha z \cdot x + \beta z \cdot y \tag{A.12}$$

<sup>&</sup>lt;sup>1</sup>More generally, the  $\mathbb{C}^{n\times n}$  can be replaced by  $\mathbb{R}^{n\times n}$  or by  $\mathbb{F}^{n\times n}$  for any field  $\mathbb{F}$ 

for  $x, y, z \in \mathcal{A}$  and  $\alpha, \beta \in \mathbb{C}$ . The cross product (but not the dot product) for vectors in  $\mathbb{R}^3$ , the multiplication of 2 complex numbers, and the commutator for square matrices, are all good examples of bilinear vector products.

Let  $B = \{\tau_i : i = 1, 2, ..., r\}$  be a basis for the vector space  $\mathcal{A}$ . Then note that B is closed under vector multiplication.

$$\tau_i \cdot \tau_j = \sum_k c_{ij}^{\ k} \tau_k \tag{A.13}$$

where  $c_{ij}^{\ k} \in \mathbb{C}$ . The  $c_{ij}^{\ k}$  are called **structure constants** of B. An **associative algebra** satisfies  $(x \cdot y) \cdot z = x \cdot (y \cdot z)$  for  $x, y, z \in \mathcal{A}$ .

- Not associative: cross product for vectors in  $\mathbb{R}^3$ .
- Associative: the commutator for square matrices and product of complex numbers

#### Tensors A.4

 $(x_1, x_2, \dots, x_n) = x^{:n} \in V^n = \mathbb{C}^{n \times 1}$ Reverse of vector  $rev(x_1, x_2, \dots, x_n) = (x_n, x_{n-1}, \dots, x_1)$  $y^b = \sum_b g^{ba} x^{:n}$  $(y^1,y^2,\ldots,y^n)=\bar{y}^{:n}\in \bar{V}^n=\mathbb{C}^{n\times 1}.\ V^n\ \mathrm{and}\ \bar{V}^n\ \mathrm{are}\ \mathrm{dual}\ \mathrm{vector}\ \mathrm{spaces}.$   $M_a{}^b\in\mathbb{C}^{n\times n},\ a,b\in\mathbb{Z}_{[1,n]}$ Implicit Summation Convention

$$M_a{}^b x_b = \sum_{b=1}^n M_a{}^b x_b \tag{A.14}$$

$$(M^{\dagger})_b^{\ a} = (M^*)_a^{\ b}$$
 (A.15)

$$= M_b^a \quad \text{(only if } M \text{ is a unitary matrix)} \tag{A.16}$$

For  $x_a \in V^n$ ,

$$(x')_a = M_a{}^b x_b \tag{A.17}$$

For  $x^a \in \bar{V}^n$ ,

$$(x'^*)^a = x^{*b}(M^*)^a_b$$
 (A.18)  
=  $x^{*b}(M^{\dagger})^a_b$  (A.19)

$$= x^{*b} (M^{\dagger})_b^{\ a} \tag{A.19}$$

SO

$$(M^{\dagger})_b^{\ a} = (M^*)_b^a \tag{A.20}$$

If the Hermitian conjugate  $\dagger$  equals \*T where \* is complex conjugation and T is transpose,

$$(M^T)_b^{\ a} = M^a_{\ b}$$
 (A.21)

This corresponds to flipping M along its horizontal.

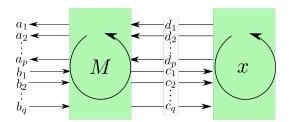


Figure A.1: Index labels for Mx where  $M \in \mathbb{C}^{n^{p+q} \times n^{p+q}}$  and  $x \in V^{n^p} \otimes \bar{V}^{n^q}$ . Note that we list indices in counterclockwise (CC) direction, starting at the top.

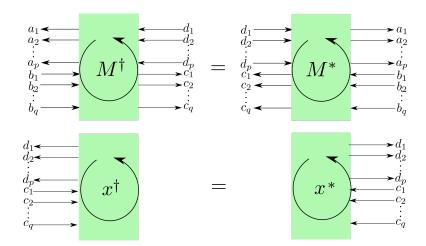


Figure A.2: Index labels for  $M^{\dagger}$  where  $M \in \mathbb{C}^{n^{p+q} \times n^{p+q}}$ . Note that we list indices in counterclockwise (CC) direction, starting at the top.

Suppose  $a_i, b_i, c_i, d_i \in \mathbb{Z}_{[1,n]}$ . From Fig.A.1

$$y_{a^{:p}}^{b^{:q}} = M_{a^{:p}}^{b^{:q}} rev(c^{:q})^{rev(d^{:p})} x_{d^{:p}}^{c^{:q}}$$
(A.22)

$$X_{\alpha} = X_{a^{:p}}^{b^{:q}}, \quad X^{\alpha} = X_{rev(b^{:q})}^{rev(a^{:p})}$$
 (A.23)

$$x_{\alpha} = M_{\alpha}{}^{\beta} x_{\beta} \tag{A.24}$$

Hermitian conjugation (see Fig.A.2)

$$\begin{cases} (M^{\dagger})_a{}^d = (M^d{}_a)^* \\ (M^{\dagger})_{\alpha}{}^{\delta} = (M^{rev(\delta)}{}_{rev(\alpha)})^* \end{cases}$$
 (A.25)

Hermitian matrix

$$M^{\dagger} = M, \quad \left\{ \begin{array}{l} (M_a^d)^* = M_a^{\phantom{a}d} \\ (M^{rev(\delta)}_{\phantom{a}rev(\alpha)})^* = M_{\alpha}^{\phantom{a}\delta} \end{array} \right. \tag{A.26}$$

Note that for  $x \in V^n$ ,  $y \in \bar{V}^n$ , and  $G \in \mathcal{G} \subset GL(n,\mathbb{C})$ ,

$$(x')_a(y')^b = G^b_{\ c} G_a^{\ d} x_d y^c \tag{A.27}$$

If  $x \in V^{n^p} \otimes \bar{V}^{n^q}$ ,  $\mathbb{G} \in \mathcal{G} \subset GL(n^{p+q}, \mathbb{C})$ ,

$$(x')_{a:p}^{b:q} = \mathbb{G}_{a:p}^{b:q} rev(c:q)^{rev(d:p)} x_{d:p}^{c:q}, \quad (x'_{\alpha} = \mathbb{G}_{\alpha}^{\beta} x_{\beta})$$
 (A.28)

where we define

$$\mathbb{G}_{a:p}^{b:q} \stackrel{rev(c:q)}{=} rev(c:q) \stackrel{\text{def}}{=} \prod_{i=1}^{q} G^{b_i}_{c_i} \prod_{i=1}^{p} G_{a_i}^{d_i}$$
(A.29)

An issue that arises with tensors is this: When is it permissible to represent a tensor by  $T_{ab}^{cd}$ ? If we define  $T_{ab}^{cd}$  by

$$T_{ab}^{cd} = T_{ab}^{cd} \tag{A.30}$$

then it's always permissible. Then one can define tensors like  $T_a^{\ bcd}$  as

$$T_a^{bcd} = g^{bb'}T_{ab'}^{cd} = g^{bb'}T_{ab'}^{cd}$$
 (A.31)

Hence, one drawback of using the notation  $T_{ab}^{cd}$  is that if one is interested in using versions of  $T_{ab}^{cd}$  with some indices raised or lowered, one has to write down explicitly the metric tensors that do the lowering and raising. Instead of writing  $T_a^{bcd}$ , you'll have to write  $g^{bb'}T_{ab'}^{cd}$ . This is not very onerous when explaining a topic in which not much lowering and raising of indices is done. But in topics like General Relativity that do use a lot of raising and lowering of indices, it might not be too elegantly concise.

### Appendix B

### **Birdtracks**

Cvitanovic Birdtracks book [1]

Elliott-Dawber book [2]

My paper "Quantum Bayesian Nets" [3]

# B.1 Classical Bayesian Networks and their Instantiations

TPM (Transition Probability Matrix)  $P(y|x) \in [0,1]$  where  $x \in val(\underline{x})$  and  $y \in val(y)$ 

$$\sum_{y \in val(y)} P(y|x) = 1 \tag{B.1}$$

$$C = \underbrace{\frac{b}{\underline{c}}}_{\underline{a}} \tag{B.2}$$

$$C(a,b,c) = P(c|b,a)P(b|a)P(a) = \sum_{c = a}^{b} P(a)$$
 (B.3)

$$a^{:2} = (a_1, a_2)$$

$$C' = \underbrace{\frac{b}{\underline{a}_1}}_{\underline{a}_2 - \underline{a}_2 - \underline{a}_2$$

$$C'(a^{:2}, b, c) = P(c|b, a_2)P(a_2|a^{:2})P(b|a_1)P(a_1|a^{:2})P(a^{:2}) = e^{b} a_1 \qquad P(a^{:2})$$
(B.5)

Marginalizer nodes  $\underline{a}_1$  and  $\underline{a}_2$  have the TPMs

$$P(a_i'|\underline{a}^{:2} = (a_1, a_2)) = \delta(a_i', a_i)$$
(B.6)

for i = 1, 2

# B.2 Quantum Bayesian Networks and their Instantiations

TPM (Transition Probability Matrix)  $A(y|x) \in \mathbb{C}$  where  $x \in val(\underline{x})$  and  $y \in val(y)$ 

$$\sum_{y \in val(\underline{y})} |A(y|x)|^2 = 1 \tag{B.7}$$

$$Q = \underbrace{\frac{b}{c}}_{a}$$
 (B.8)

$$Q(a,b,c) = A(c|b,a)A(b|a)A(a) =$$

$$C = a$$

$$A(a)$$
(B.9)

$$a^{:2} = (a_1, a_2)$$

$$Q' = \underbrace{\frac{b}{\underline{a}_1}}_{\underline{a}_2 - \underline{a}_2 - \underline{a}_2 - \underline{a}_2 - \underline{a}_2}$$
(B.10)

$$Q'(a^{:2}, b, c) = A(c|b, a_2)A(a_2|a^{:2})A(b|a_1)A(a_1|a^{:2})A(a^{:2}) = b$$

$$c \stackrel{b}{\longleftarrow} a_1 \qquad A(a^{:2})$$
(B.11)

Marginalizer nodes  $\underline{a}_1$  and  $\underline{a}_2$  have the TAMs

$$A(a_i'|\underline{a}^{:2} = (a_1, a_2)) = \delta(a_i', a_i)$$
(B.12)

for i = 1, 2

#### B.3 Birdtracks

$$\delta(b, a) = \mathbb{1}(a = b) = \delta_a^b = a - b \tag{B.13}$$

$$\underline{a} = a \longleftarrow X_{\underline{a}\underline{b}}^{\underline{c}\underline{d}}$$

$$\langle a, b | X_{\underline{a}\underline{b}}^{\underline{c}\underline{d}} | c, d \rangle = X_{ab}^{\underline{c}d} = b$$

$$\underline{c} = c$$

$$d = d$$
(B.14)

$$a \longleftarrow X_{\underline{ab}} \stackrel{cd}{=} \qquad a, b \longleftarrow X_{\underline{ab}} \stackrel{cd}{=}$$

$$b \qquad \qquad \rightarrow \qquad a, b \qquad \qquad (B.15)$$

$$c \qquad \qquad c \qquad \qquad d$$

 $X_{\underline{a}\underline{b}} \stackrel{cd}{\in} V^2 \otimes V_2$ . Sometimes, we will omit denote this node simply by X. This if okay as long as we are not using, X to also denote a different version of  $X_{\underline{a}\underline{b}} \stackrel{cd}{=}$  with some of the indices raised or lowered or their order has been changed. <sup>1</sup>

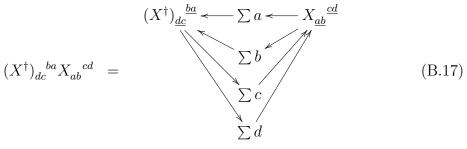
$$(X^{\dagger})_{\underline{dc}} \xrightarrow{\underline{ba}} \underline{a} = a$$

$$(X^{\dagger})_{\underline{dc}} \xrightarrow{ba} \underline{b} = b$$

$$\underline{c} = c$$

$$d = d$$
(B.16)

<sup>&</sup>lt;sup>1</sup>For matrices,  $(A^{\dagger})_{i,j} = (A_{j,i})^*$  so taking a Hermitian conjugate involves both taking the complex conjugate of the matrix element and reversing the left-to-right (L2R) order of its indices. This generalizes to  $(X^{\dagger})_{dc}^{\ \ ba} = (X_{ab}^{\ \ cd})^*$ . Besides raising and lowering indices, we reverse their L2R order.



$$= X^{\dagger} - X$$

$$= (B.18)$$

Birdtracks originated as a graphical way to represent the tensors in General Relativity (Gravitation). In General Relativity, one deals with tensors such as  $T_{a\ c}^{\ b}$  which have some indices raised and some lowered. One can use the metric  $g^{a,b}$  to raise all the lowered indices to get  $T^{abc}$ . If we represent this graphically as a node with incoming arrows a, b, c, we need to follow one of the following 2 conventions: either

- 1. label the arrows as  $\underline{a}$ ,  $\underline{b}$ ,  $\underline{c}$ , and define the node as  $T^{\underline{abc}}$ , or
- 2. instead of labelling the arrows explicitly  $\underline{a}, \underline{b}, \underline{c}$ , indicate in the node where is the first arrow  $\underline{a}$ , and draw the arrows  $\underline{a}, \underline{b}, \underline{c}$  so that they enter the node in **counterclockwise** (CC) order. The **left-to-right** (L2R) order of the indices on T corresponds the CC order of the arrows.

If we don't do either 1 or 2, we won't be able to distinguish between the graphical representations of  $T^{1,2,3}$  and  $T^{2,1,3}$ , for example. Cvitanovic's Birdtracks book Ref.[1] follows Convention 2, but most of the time, in this book, we will follow Convention 1 <sup>2</sup> The reason I chose to do so is for the sake of consistency: Convention 2 is closer to the quantum bnet conventions.

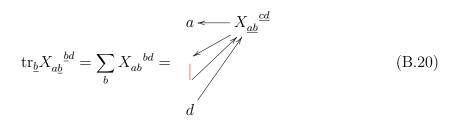
$$a^{:m} \in \mathbb{Z}_+^m$$

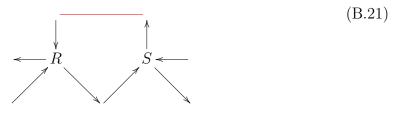
$$R_{b_{3}^{i,m_{3}},a_{2}^{i,m_{2}}}^{i,m_{3},b_{2}^{i,n_{2}}}S_{b_{2}^{i,n_{2}},a_{1}^{i,m_{1}}}^{i,m_{3},b_{2}^{i,m_{2}}}= \begin{pmatrix} b_{3}^{i,n_{3}} & & \sum b_{2}^{i,n_{2}} & & \sum b_{2}^{i,n_{2}} & & \\ & & & & \sum a_{2}^{i,m_{2}} & & a_{1}^{i,m_{1}} \end{pmatrix}$$

$$= \begin{pmatrix} b_{3}^{i,n_{3}} & & & \sum b_{2}^{i,n_{2}} & & & \\ & & & & \sum a_{2}^{i,m_{2}} & & a_{1}^{i,m_{1}} \end{pmatrix}$$

$$= \begin{pmatrix} b_{3}^{i,n_{3}} & & & \sum b_{2}^{i,n_{2}} & & & \\ & & & & \sum a_{2}^{i,m_{2}} & & & \\ & & & & \sum a_{2}^{i,m_{2}} & & \\ & & & & \sum a_{2}^{i,m_{2}} & & \\ & & & & \sum a_{2}^{i,m_{2}} & & \\ & & & & \sum a_{2}^{i,m_{2}} & & \\ & & & & \sum a_{2}^{i,m_{2}} & & \\ & & & & \sum a_{2}^{i,m_{2}} & & \\ & & & & \sum a_{2}^{i,m_{2}} & & \\ & & & & \sum a_{2}^{i,m_{2}} & & \\ & & & & \sum a_{2}^{i,m_{2}} & & \\ & & & & \sum a_{2}^{i,m_{2}} & & \\ & & & & \sum a_{2}^{i,m_{2}} & & \\ & & & & \sum a_{2}^{i,m_{2}} & & \\ & & & & \sum a_{2}^{i,m_{2}} & & \\ & & & & \sum a_{2}^{i,m_{2}} & & \\ & & & & \sum a_{2}^{i,m_{2}} & & \\ & \sum a_{2}^{i,m_{2}}$$

<sup>&</sup>lt;sup>2</sup>If we follow Convention 1, we don't need to reverse the L2R order of the indices when taking a Hermitian conjugate. Thus,  $(X^{\dagger})^{\underline{ab}}_{\underline{cd}} = X_{\underline{ab}}^{\underline{cd}} = X^{\underline{dc}}_{\underline{ba}}$ . As long as  $\underline{a}, \underline{b}$  are lower indices and  $\underline{c}, \underline{d}$  are upper indices of X, any L2R order of  $\underline{a}, \underline{b}, \underline{c}, \underline{d}$  is equivalent under Convention 1.





Casimir Operators: COMING SOON

## Clebsch-Gordan Coefficients

$$\begin{bmatrix} 0 \\ C_{\lambda}^{d_{\lambda} \times d} \\ 0 \end{bmatrix}^{d \times d} = \begin{bmatrix} 0 & 0 & 0 \\ 0 & I^{d_{\lambda} \times d_{\lambda}} & 0 \\ 0 & 0 & 0 \end{bmatrix}^{d \times d} C^{d \times d}$$
 (2.1)

Let  $b^{:nb} = (b_1, b_2, ..., b_{nb})$  where  $b_i \in Z_{[0,db_i]}$  and  $a \in Z_{[1,d_{\lambda}]}$ . Hence,

$$d_{\lambda} = \prod_{i=1}^{:nb} db_i \tag{2.2}$$

$$(C_{\lambda})_{a}^{rev(b:nb)} = a \leftarrow C_{\lambda} \leftarrow b_{2}$$

$$b_{nb}$$

$$(2.3)$$

$$\begin{bmatrix} 0 & (C^{\dagger})_{\lambda}^{d \times d_{\lambda}} & 0 \end{bmatrix}^{d \times d} = (C^{\dagger})^{d \times d} \begin{bmatrix} 0 & 0 & 0 \\ 0 & I^{d_{\lambda} \times d_{\lambda}} & 0 \\ 0 & 0 & 0 \end{bmatrix}^{d \times d}$$
(2.4)

$$(C_{\lambda}^{\dagger})_{b:nb}^{a} = b_{2} \leftarrow (C_{\lambda}^{\dagger}) \leftarrow a$$

$$b_{nb}$$

$$(2.5)$$

More generally, some of the  $b_i$  indices may lowered and their arrows changed to outgoing instead of ingoing. Each  $b_i$  represents a different rep (or irrep)

$$C_{\lambda})_{a}^{rev(b:nb)}(C_{\mu}^{\dagger})_{b:nb}^{a'} = \delta(\lambda,\mu)\delta_{a}^{a'}, \quad C_{\lambda}C_{\mu}^{\dagger} = \delta(\mu,\lambda)$$

$$a \leftarrow C_{\lambda} \leftarrow \sum b_{2} \leftarrow (C_{\mu}^{\dagger}) \leftarrow a' = \delta(\mu,\lambda) \ a \leftarrow a'$$

$$\sum b_{nb}$$
(2.7)

Completeness relation

$$\sum_{\lambda} (C_{\lambda}^{\dagger})_{b:nb}^{a} (C_{\lambda})_{a}^{rev((b'):nb)} = \delta_{b:nb}^{rev((b'):nb)}, \quad \sum_{\lambda} C_{\lambda}^{\dagger} C_{\lambda} = 1$$

$$b_{1} \qquad b_{1} \qquad b_{1} \qquad b_{1} \qquad b'_{1}$$

$$\sum_{\lambda} b_{2} \leftarrow (C_{\lambda}^{\dagger}) \leftarrow \sum_{\alpha} a \leftarrow C_{\lambda} \leftarrow b'_{2} = b_{2} \leftarrow b'_{2}$$

$$b_{nb} \qquad b'_{nb} \qquad b_{nb} \leftarrow b'_{nb}$$

$$(2.8)$$

$$(C_{\lambda})_{a}^{rev((b'):nb)}(P_{\mu})_{(b'):nb}^{rev(b:nb)} = \delta(\mu,\lambda)(C_{\mu})_{a}^{rev(b:nb)}, \quad C_{\lambda}P_{\mu} = \delta(\mu,\lambda)C_{\mu}$$

$$b_{1}$$

$$b_{1}$$

$$b_{2}$$

$$b_{2}$$

$$b_{nb}$$

$$b_{nb}$$

$$b_{nb}$$

$$(2.9)$$

$$(P_{\mu})_{b:nb}^{rev((b'):nb)}(C_{\lambda}^{\dagger})_{(b'):nb}^{a} = \delta(\mu,\lambda)(C_{\mu}^{\dagger})_{b:nb}^{a}, \quad P_{\mu}C_{\lambda}^{\dagger} = \delta(\mu,\lambda)C_{\mu}^{\dagger}$$

$$b_{1} \qquad \qquad b_{1} \qquad \qquad b_{1} \qquad \qquad b_{1} \qquad \qquad b_{1} \qquad \qquad b_{2} \leftarrow P_{\mu} \leftarrow \sum b_{2}' \qquad \leftarrow (C_{\lambda}^{\dagger}) \leftarrow a \qquad = \delta(\mu,\lambda) \quad b_{2} \leftarrow (C_{\lambda}^{\dagger}) \leftarrow a \qquad b_{nb} \qquad (2.10)$$

Determinants: COMING SOON

Dynkin Diagrams: COMING SOON

General Relativity Nets: COMING SOON

Group Integrals: COMING SOON

### **Invariants**

Given a bilinear form

$$m(\bar{x}^{:n}, y^{:n}) = x^a M_a{}^b y_b \qquad M$$

$$a \qquad b$$

$$(7.1)$$

is invariant if

$$m(\bar{x}^{:n}, y^{:n}) = m(\bar{x}^{:n}G^{\dagger}, Gy^{:n})$$
 (7.2)

matrix invariant

$$M_a{}^b = (G^{\dagger})_a{}^{a'}G_{b'}{}^b M_{a'}{}^{b'} \qquad M_b = M_b \qquad (7.3)$$

$$M = G^{\dagger}MG \tag{7.4}$$

$$GM = MG, \quad [G, M] = 0 \tag{7.5}$$

multilinear form

$$h(\bar{w}, \bar{x}, y, z) = h_{ab}^{cd} w^a x^b y_c z_d \qquad \qquad \begin{matrix} h \\ \downarrow \\ a \end{matrix} \qquad \qquad \begin{matrix} h \\ \downarrow \\ b \end{matrix} \qquad \qquad (7.6)$$

is invariant if

$$h(\bar{w}, \bar{x}, y, z) = h(\bar{w}G^{\dagger}, \bar{x}G^{\dagger}, Gy, Gz)$$
(7.7)

tensor invariant (TI)

$$h_{ab}^{cd} = (G^{\dagger})_a^{a'} (G^{\dagger})_b^{b'} h_{a'b'}^{c'd'} G_{c'}^{c} G_{d'}^{d} \qquad h \qquad \qquad h \qquad (7.8)$$

A **composed TI** is a TI that can be written as a product or contraction of TIs.

A tree TI is a composed TIs without any loops.

A **primitive TI** is a TI that can be expressed as a linear combination of a finite number of tree TIs.

The **primitiveness assumption**: All TI are primitive.

Examples. Consider  $\mathbb{R}^3$  vector space.

#### • Primitive TIs

$$length(x) = \delta_{ij}x_ix_i \quad volume(x, y, z) = \epsilon_{ijk}x_iy_jz_k$$
 (7.9)

#### • Tree TIs

$$\delta_{ij}\epsilon_{klm} = \begin{vmatrix} i & & \epsilon \\ & & \\ j & & k \end{vmatrix}$$

$$(7.11)$$

$$\epsilon_{ijm}\delta_{mn}\epsilon_{nkl} = \begin{vmatrix} \epsilon_{ijm} - \epsilon_{nkl} \\ \\ \\ i \end{vmatrix}$$

$$i \qquad k \qquad l$$

$$(7.12)$$

#### • Non-tree TI

$$\epsilon_{ims}\epsilon_{jnm}\epsilon_{krn}\epsilon_{lsr} = 
\begin{bmatrix}
i & ---- \epsilon_{ims} - \sum s - \epsilon_{lsr} & ---- l \\
\sum m & \sum r \\
j & ---- \epsilon_{jnm} - \sum n - \epsilon_{krn} & ---- k
\end{bmatrix} (7.13)$$

#### • Primitiveness Assumption

Suppose  $\mathcal{P} = \{\delta_{ij}, f_{ijk}\}$  where  $f_{ijk}$  is not  $\epsilon_{ijk}$ 

$$- \underbrace{\hspace{1cm}} = A - \underbrace{\hspace{1cm}} (7.15)$$

$$= B \qquad | \qquad (7.16)$$

#### An algebra of invariants

Let  $\mathcal{P} = (p_1, p_2, \dots, p_k)$  be a full set of primitive TIs. By "full", we mean no others exist.

An invariance group  $\mathcal{G}$  is the set of all linear transformation  $G \in \mathcal{G}$  such that

$$p_1(x,\bar{y}) = p_1(Gx,\bar{y}G^{\dagger}) \tag{7.18}$$

$$p_2(w, x, \bar{y}, \bar{z}) = p_2(Gw, Gx, \bar{y}G^{\dagger}, \bar{z}G^{\dagger})$$

$$(7.19)$$

etc. 
$$(7.20)$$

Example

$$p(\bar{x}, y) = \delta_a^b x^a y_b = x^b y_b \tag{7.21}$$

$$(x')^{a}(y')_{a} = x^{b}(G^{\dagger}G)_{b}^{c}y_{c} = x^{b}y_{b}$$
(7.22)

So G must be unitary

$$G^{\dagger}G = 1 \tag{7.23}$$

The group of n dimensional unitary matrices is called U(n)

### Levi-Civita Tensor

$$\epsilon^{123...p} = \epsilon_{123...p} = 1 \tag{8.1}$$

$$\epsilon_{rev(a^{:p})} = (-1)^{\binom{p}{2}} \epsilon_{a^{:p}} \tag{8.2}$$

where  $rev(a^{p})$  is the reverse of  $a^{p}$ .  $rev(a_1, a_2, \ldots, a_p) = (a_p, a_{p-1}, \ldots, a_1)$ 

$$(C_{\mathcal{A}_p}^{\dagger})_{a:p}^{1} = e^{-i\phi} \frac{\epsilon_{a:p}}{\sqrt{p!}} = a_1 \leftarrow \mathcal{A}_p$$

$$a_2 \leftarrow$$

$$\vdots$$

$$\underbrace{e^{i2\phi} \frac{1}{p!} \epsilon^{a^{:n}} \epsilon_{a^{:n}} = \delta_1^1 = 1}_{\qquad \vdots \qquad \qquad } = 1$$
(8.6)

For Convention 1, we will use  $\phi = 0$ . For Convention 2, we must choose

$$e^{i2\phi} = (-1)^{\binom{p}{2}} = e^{i\pi\frac{p(p-1)}{2}}$$
 (8.7)

SO

$$\phi = \frac{\pi}{4}p(p-1) \tag{8.8}$$

### Lie Algebras

 $i \in \mathbb{Z}_{[1,N]}, a,b \in \mathbb{Z}_{[1,n]}$ 

$$(C_{Adj}^{i})_{b}^{a} = \frac{1}{\sqrt{K}} (T^{i})_{b}^{a} = i \sim C_{Adj}^{i}$$

$$\downarrow$$

$$b$$

$$(9.1)$$

Note that we list the indices of  $T^i$  in the counter-clockwise (CC) direction, starting at the i leg. The matrices  $T^i$  are called the generators. It's customary to choose them so that they are Hermitian and  $K = \frac{1}{2}$ .

$$\underbrace{\left[ (T^i)_b{}^a (T^j)_a{}^b = \operatorname{tr}(T^i T^j) = K \delta(i,j) \right]}_{\sum b} i \sim T^i \qquad T^j \sim j = K \leftarrow \bullet \quad (9.2)$$

$$(P_{Adj})_{b,d}^{a,c} = \sum_{i} \frac{1}{K} (T^{i})_{b}^{a} (T^{i})_{d}^{c} = \frac{1}{K} \left| \begin{array}{c} a & d \\ \\ \\ \\ b & c \end{array} \right|$$
(9.3)

 $H\in V^n\otimes \bar V^n$ 

$$(P_{Adj})_{bd}^{ac} H_c^{\ d} = \sum_{i} (T^i)_b^{\ a} \underbrace{\left[\frac{1}{K} (T^i)_d^{\ c} H_c^{\ d}\right]}_{h_i \in \mathbb{R}}$$
(9.4)

For SU(2), it is customary to use  $T^i = \frac{1}{2}\sigma_i$ , where  $\sigma_i$  for i = 1, 2, 3 are the Pauli matrices. For SU(3), it is customary to choose  $T^i = \frac{1}{2}\lambda_i$  where  $\lambda_i$  for  $i = 1, 2, \ldots, 8$  are the Gelll-Mann matrices.

$$G = 1 + iD \in \mathcal{G}$$

$$\epsilon_i \in \mathbb{R}, |\epsilon_i| << 1$$

$$D = \sum_i \epsilon_i T^i = V^n \otimes \bar{V}^n$$
Recall Eq.(A.28). If  $x \in V^{n^p} \otimes \bar{V}^{n^q}$ ,  $\mathbb{G} \in \mathcal{G} \subset GL(n^{p+q}, \mathbb{C})$ ,

$$(x')_{a^{:p}}^{b^{:q}} = \mathbb{G}_{a^{:p}}^{b^{:q}}_{rev(c^{:q})}^{rev(d^{:p})} x_{d^{:p}}^{c^{:q}}, \quad (x'_{\alpha} = \mathbb{G}_{\alpha}^{\beta} x_{\beta})$$
 (9.5)

where we define

$$\mathbb{G}_{\alpha}^{\beta} \stackrel{\text{def}}{=} \prod_{i=1}^{q} G^{b_i}{}_{c_i} \prod_{i=1}^{p} G_{a_i}{}^{d_i}$$

$$\tag{9.6}$$

$$\mathbb{G}_{\alpha}^{\beta} = 1 + i \sum_{j} \epsilon_{j} (\mathbb{T}^{j})_{\alpha}^{\beta} \tag{9.7}$$

$$G^{b_i}_{c_i} = 1 + i \sum_j \epsilon_j (T^j)^{b_i}_{c_i}$$
 (9.8)

$$G_{a_i}^{d_i} = (G^*)_{a_i}^{d_i} = 1 - i \sum_j \epsilon_j (T^{j*})_{a_i}^{d_i} = 1 - i \sum_j \epsilon_j (T^j)_{a_i}^{d_i}$$
 (9.9)

When  $x'_{\alpha} = x_{\alpha}$ , to first order in  $\epsilon_i$ ,

$$0 = (\mathbb{T}^{j})_{\alpha}{}^{\beta} x_{\beta} = \left[ (T^{j})^{b_{i}}_{c_{i}} \frac{1}{\delta^{b_{i}}_{c_{i}}} - (T^{j})_{a_{i}}{}^{d_{i}} \frac{1}{\delta^{d_{i}}_{a_{i}}} \right] \delta^{b^{iq}}_{c^{iq}} \delta^{d^{ip}}_{a^{ip}} x_{d^{ip}}{}^{c^{iq}}$$
(9.10)

$$\left| (\mathbb{T}^j)_{\alpha}^{\ \beta} = \left[ (T^j)^{b_i}_{\ c_i} \frac{1}{\delta_{c_i}^{b_i}} - (T^j)_{a_i}^{\ d_i} \frac{1}{\delta_{a_i}^{d_i}} \right] \delta_{c:q}^{b:q} \delta_{a:p}^{d:p} \right|$$
(9.11)

Orthogonal Groups: COMING SOON

Quantum Shannon Information Theory: COMING SOON

Recoupling Equations: COMING SOON

### Reducibility

 $M \in \mathbb{C}^{d \times d}$ 

$$M|v\rangle = \lambda|v\rangle \tag{13.1}$$

If M is Hermitian  $(H^{\dagger} = H)$ , its eigenvalues are real.  $(\lambda = \langle \lambda | M \lambda \rangle \in \mathbb{R})$ 

$$cp(\lambda) \stackrel{\text{def}}{=} \det(M - \lambda) = 0$$
 (13.2)

If M is a Hermitian matrix, then there exists a unitary matric ( $CC^{\dagger}=C^{\dagger}C=1$ ) such that

$$CMC^{\dagger} = \begin{bmatrix} D_{\lambda_1} & 0 & 0 & 0 \\ 0 & D_{\lambda_2} & 0 & 0 \\ 0 & 0 & \ddots & 0 \\ 0 & 0 & 0 & D_{\lambda_r} \end{bmatrix}$$
 (13.3)

where

$$D_{\lambda_i} = \operatorname{diag}\left(\underbrace{\lambda_i, \lambda_i, \dots, \lambda_i}_{d_i \text{ times}}\right) \tag{13.4}$$

$$d = \sum_{i=1}^{r} d_i \tag{13.5}$$

$$CMC^{\dagger} = \begin{bmatrix} \lambda_1 & 0\\ 0 & \lambda_2 \end{bmatrix} \tag{13.6}$$

$$CP_1C^{\dagger} = \begin{bmatrix} 1 & 0 \\ 0 & 0 \end{bmatrix} = \frac{CMC^{\dagger} - \lambda_2}{\lambda_1 - \lambda_2}$$
 (13.7)

$$CP_2C^{\dagger} = \begin{bmatrix} 0 & 0 \\ 0 & 1 \end{bmatrix} = \frac{CMC^{\dagger} - \lambda_1}{\lambda_2 - \lambda_1}$$
 (13.8)

If  $I^{d_i \times d_i}$  is the  $d_i$  dimensional unit matrix,

$$P_i = C^{\dagger} diag(0, \dots, 0, I^{d_i \times d_i}, 0, \dots, 0)C$$
 (13.9)

$$= \prod_{j \neq i} \frac{M - \lambda_j}{\lambda_i - \lambda_j} \tag{13.10}$$

Note that  $P_i$  are Hermitian  $(P_i^{\dagger} = P_i)$  because M is Hermitian and its eigenvalues are real.)

Note that  $P_i$  and M commute

$$[P_i, M] = P_i M - M P_i = 0 (13.11)$$

orthogonal

$$P_i P_j = \delta(i, j) P_j \tag{13.12}$$

complete

$$\sum_{i} P_i = 1 \tag{13.13}$$

$$M = \sum_{i=1}^{r} P_i M P_i \tag{13.14}$$

$$d_i = \operatorname{tr} P_i \tag{13.15}$$

$$CMP_1C^{\dagger} = \begin{bmatrix} \lambda_1 & 0 \\ 0 & \lambda_2 \end{bmatrix} \begin{bmatrix} 1 & 0 \\ 0 & 0 \end{bmatrix}$$
 (13.16)

$$= \lambda_1 \begin{bmatrix} 1 & 0 \\ 0 & 0 \end{bmatrix} \tag{13.17}$$

$$MP_i = \lambda_i P_i \text{ (no } i \text{ sum)}$$
 (13.18)

$$f(M)P_i = f(\lambda_i)P_i \text{ (no } i \text{ sum)}$$
 (13.19)

 $M^{(1)}, M^{(2)}$ 

$$[M^{(1)}, M^{(2)}] = 0 (13.20)$$

Use  $M^{(1)}$  to decompose V into  $\bigoplus_i V_i$ . Use  $M^{(2)}$  to decompose  $V_i$  into  $\bigoplus_j V_{i,j}$ . If  $M^{(1)}$  and  $M^{(2)}$  don't commute, let  $P_i^{(1)}$  be the eigenvalue projection operators of  $M^{(1)}$ . The replace  $M^{(2)}$  by  $P_i^{(1)}M^{(2)}P_i^{(1)}$ 

$$[M^{(1)}, P_i^{(1)}M^{(2)}P_i^{(1)}] = 0 (13.21)$$

Spinors: COMING SOON

# Squashed Entanglement: COMING SOON

Symplectic Groups: COMING SOON

# Symmetrization and Antisymmetrization

(1,2) transposition, swaps 1 and 2,  $1 \to 2 \to 1$ . (3,2,1) means  $3 \to 2 \to 1 \to 3$ . A reordering of  $(1,2,3,\ldots,p)$  is a permutation on p letters. A permutation can be expressed as a product of transpositions (3,2,1)=(3,2)(2,1) is an even permutation because it can be expressed as a product of an even number of transpositions. An odd permutation can be expressed as a product of an odd number of permutations.

#### 17.1 Symmetrization

$$\mathbb{1}_{a_1,a_2}^{b_2,b_1} = \delta_{a_1}^{b_1} \delta_{a_2}^{b_2} = \begin{cases} a_1 \leftarrow b_1 \\ a_2 \leftarrow b_2 \end{cases}$$
 (17.1)

$$(\sigma_{(1,2)})_{a_1,a_2}^{b_2,b_1} = \delta_{a_1}^{b_2} \delta_{a_2}^{b_1} = \begin{cases} a_1 \leftarrow \bullet \leftarrow b_1 \\ \downarrow \\ a_2 \leftarrow \bullet \leftarrow b_2 \end{cases}$$
 (17.2)

$$1 = \underbrace{\hspace{1cm}} (17.3)$$

$$\sigma_{(1,2)} = \begin{array}{c} & \longleftarrow & \longleftarrow & \longleftarrow \\ & \uparrow \\ & \downarrow \\ & \longleftarrow & \sigma_{(2,3)} = \\ & \longleftarrow & \downarrow \\ & \longleftarrow & \longleftarrow \end{array}$$

$$(17.4)$$

#### Claim 1

**proof:** We only prove it for p = 3.

QED

$$= \frac{n+p-1}{p} \begin{pmatrix} \leftarrow \mathcal{S}_{p-1} \leftarrow \\ \leftarrow \\ \vdots \\ \vdots \\ \vdots \end{pmatrix}$$

$$(17.17)$$

$$\operatorname{tr}_{\underline{a}_{1}} \mathcal{S}_{p} = \frac{n+p-1}{p} \mathcal{S}_{p-1} \tag{17.18}$$

$$\operatorname{tr}_{\underline{a}_1,\underline{a}_2,\dots,\underline{a}_k} \mathcal{S}_p = \frac{(n+p-1)(n+p-2)\dots(n=p-k)}{p(p-1)\dots(p-k+1)} \mathcal{S}_{p-k}$$
 (17.19)

$$d_{S_p} = \operatorname{tr}_{\underline{a}^p} S_p = \frac{(n+p-1)!}{p!(n-1)!} = \binom{n+p-1}{p}$$
 (17.20)

For p=2,

$$d_{\mathcal{S}_2} = \frac{(n+1)n}{2} \tag{17.21}$$

#### 17.2 Antisymmetrization

$$\begin{array}{c|cccc}
\leftarrow \mathcal{A}_p \leftarrow & & \leftarrow \bullet \leftarrow & \\
\leftarrow & & \downarrow & \\
\leftarrow & & \leftarrow & \leftarrow & \\
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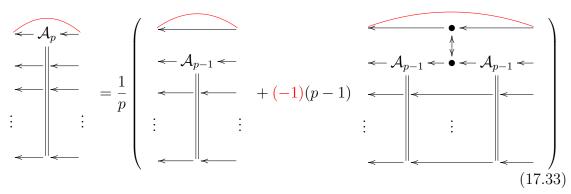
$$\begin{aligned}
S_{p}A_{[1,q]} &= A_{p}S_{[1,q]} = 0 \\
&\leftarrow S_{p} \leftarrow &\leftarrow A_{[1,q]} \leftarrow &\leftarrow &\leftarrow S_{[1,q]} \leftarrow \\
&\leftarrow &\parallel \leftarrow &\leftarrow &\parallel \leftarrow &\leftarrow &\parallel \leftarrow \\
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&\leftarrow &\parallel \leftarrow &\leftarrow &\leftarrow &\parallel \leftarrow &\leftarrow &\parallel \leftarrow &\leftarrow \\
&\vdots &\vdots &\vdots &\vdots &\vdots &\vdots &\vdots
\end{aligned}$$

$$(17.27)$$

#### Claim 2

**proof:** We only prove it for p = 3.

#### QED



$$= \frac{n + (-1)(p-1)}{p} \begin{pmatrix} \leftarrow \mathcal{A}_{p-1} \leftarrow \\ \leftarrow \\ \vdots \\ \vdots \\ \cdots \\ \vdots \end{pmatrix}$$

$$(17.34)$$

$$\operatorname{tr}_{\underline{a}_1} \mathcal{A}_p = \frac{n-p+1}{p} \mathcal{A}_{p-1} \tag{17.35}$$

$$\operatorname{tr}_{\underline{a}_{1},\underline{a}_{2},\dots,\underline{a}_{k}} \mathcal{A}_{p} = \frac{(n-p+1)(n-p+2)\dots(n-p+k)}{p(p-1)\dots(p-k+1)} \mathcal{A}_{p-k}$$
 (17.36)

$$d_{\mathcal{A}_p} = \operatorname{tr}_{\underline{a}^p} \mathcal{A}_p = \frac{\prod_{i=n-p+1}^n i}{p!}$$
(17.37)

$$= \frac{\prod_{i=n}^{n-p+1} i}{p!} \tag{17.38}$$

$$= \begin{cases} \frac{n!}{p!(n-p)!} = \binom{n}{p} & \text{if } p \le n\\ 0 & \text{otherwise} \end{cases}$$
 (17.39)

For  $p = 2 \le n$ ,

$$d_{\mathcal{A}_2} = \binom{n}{2} \tag{17.40}$$

$$\mathcal{A}_p = 0 \text{ if } n$$

For example, for n=2 and p=3

$$\mathcal{A}_{3}|a,a,b\rangle = \frac{1}{6} \begin{pmatrix} |a,a,b\rangle + |a,b,a\rangle + |b,a,a\rangle \\ -|a,b,a\rangle - |a,a,b\rangle - |b,a,a\rangle \end{pmatrix}$$

$$= 0$$

$$(17.43)$$

## Unitary Groups: COMING SOON

#### SU(n) 18.1

$$m(p,q) = \delta_b^a \sum_{a=1}^n (p_a)^* q_a$$
 (18.1)

$$d \leftarrow c$$

$$\mathbb{1}_{d,b}^{a,c} = \delta_b^a \delta_d^c =$$

$$a \rightarrow b$$
(18.2)

$$\uparrow\downarrow_{d,b}^{a,c} = \delta_d^a \delta_b^c = \begin{pmatrix} d & c \\ \uparrow & \downarrow \\ a & b \end{pmatrix}$$
 (18.3)

$$P_i = \sum_{j \neq i} \frac{M - \lambda_j}{\lambda_i - \lambda_j} \tag{18.5}$$

 $\lambda_1 = n$ 

$$P_{1} = \frac{\uparrow \downarrow -n}{0-n} = 1 - \frac{1}{n} \uparrow \downarrow$$

$$c \qquad d \qquad e \qquad b \qquad a \qquad b \qquad b \qquad d$$

$$c \leftarrow -\frac{1}{n} \uparrow \qquad d$$

$$\lambda_{2} = 0 \qquad (18.6)$$

$$P_{2} = \frac{\uparrow \downarrow -0}{n-0} = \frac{1}{n} \uparrow \downarrow$$

$$c$$

$$d$$

$$b$$

$$d$$

$$d$$

$$d$$

$$d$$

$$d$$

$$d$$

$$(18.7)$$

$$tr P_1 = \frac{1}{n}$$

$$(18.8)$$

$$= n^2 - 1 (18.9)$$

$$trP_2 = \frac{1}{n}$$
 (18.10)

$$= 1 \tag{18.11}$$

$$(T_i)_a^b = i \sim T_i$$

$$\downarrow$$

$$a$$
(18.12)

$$T_i^{\dagger} = T_i \tag{18.13}$$

Claim 3

$$C_F \delta_a^b = (T_i T_i)_a^b = \frac{n^2 - 1}{n} \delta_a^b$$
 (18.14)

proof:

$$(T_{i}T_{i})_{a}^{b} = \sum_{i} i \sim T_{i} \qquad T_{i} \sim i$$

$$= \sum_{i} i \sim T_{i} \qquad T_{i} \sim i$$

$$(18.15)$$

$$= \sum_{i} i \sim T_{i} T_{i} \sim i$$
 (18.16)

**QED** 

Wigner Coefficients: COMING SOON

Wigner-Ekart Theorem: COMING SOON

Young Tableau: COMING SOON

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