

Unbiased sampling of HMC schemes for non separable Hamiltonian systems



RÉGIS SANTET, GABRIEL STOLTZ, TONY LELIÈVRE

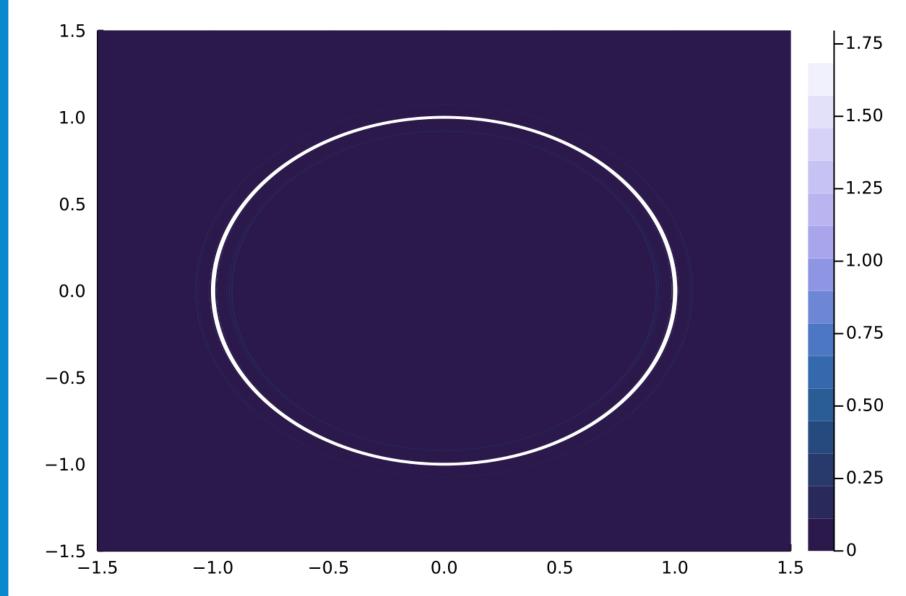
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DIFFUSION DEPENDENT OVERDAMPED LANGEVIN DYNAMICS

Using a position dependent diffusion coefficient¹ D: favoring exploration in **anisotropic** or **metastable** potential landscapes \Longrightarrow faster convergence to steady-state.

$$dq_t = \left(-D(q_t)\nabla V(q_t) + \frac{\operatorname{div} D(q_t)}{\beta}\right)dt + \sqrt{2\beta^{-1}D(q_t)}dW_t$$
(1)

• Aim: Unbiased estimation of $\mathbb{E}(\varphi) = \int_{\Omega} \varphi(x) \, \pi(\mathrm{d}x), \quad \pi = \mathrm{e}^{-\beta V}, \quad \beta^{-1} = k_B T.$



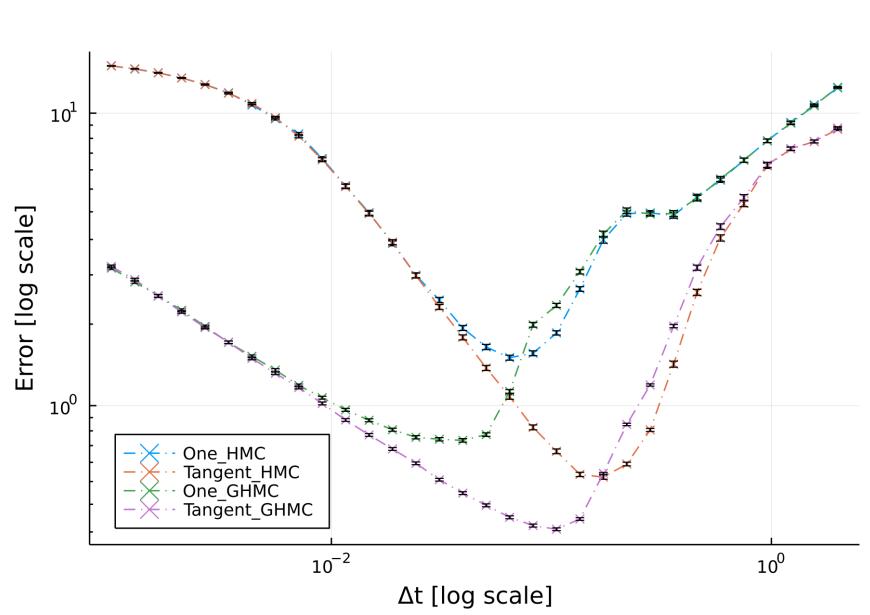
$$V(x,y) = 100(x^2 + y^2 - 1)^2$$

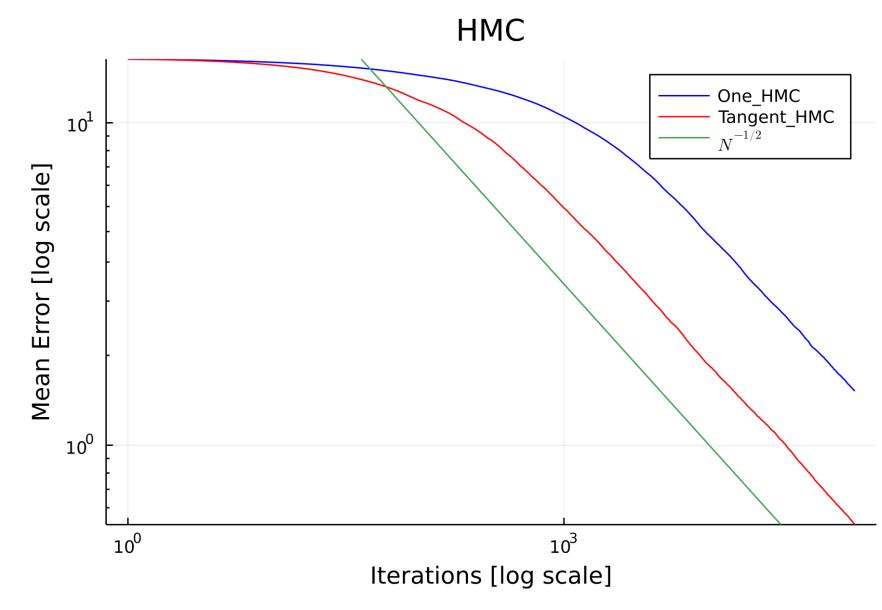
Anisotropic diffusion coefficient

$$D_{\text{Tangent}}(q) = \varepsilon \mathbf{I}_2 + \tilde{q}\tilde{q}^{\mathsf{T}}/\|q\|^2, \ \tilde{q} = (-y \ x)^{\mathsf{T}}$$

Isotropic diffusion coefficient

$$D_{\text{One}} \equiv (1 + \varepsilon)I_2, \ \varepsilon = 0.1$$





Error between empirical angle distribution and uniform distribution on $[0, 2\pi]$ after 10^5 iterations.

Mean error for optimal time steps in HMC case.

• Issue: Rejection rates scales as $\mathcal{O}\left(\Delta t^{1/2}\right)$ for Euler–Maruyama + Metropolis–Hastings.

BETTER UNBIASED NUMERICAL SAMPLING?

• Solution: Riemann Manifold (Generalized) Hamiltonian Monte Carlo^{2,3} scheme based on Langevin dynamics integration.

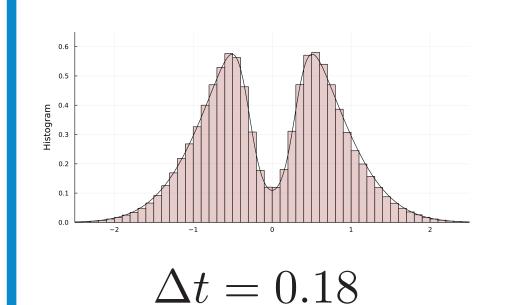
$$\begin{cases} dq_t = \nabla_p H(q_t, p_t) dt \\ dp_t = -\nabla_q H(q_t, p_t) dt - \gamma \nabla_p H(q_t, p_t) dt + \sqrt{\frac{2\gamma}{\beta}} dW_t \end{cases}$$

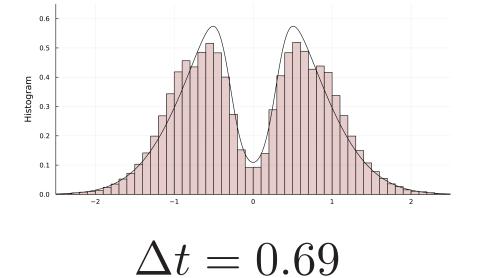
- Rejection rates scale as $\mathcal{O}(\Delta t^3)$. But need to have a time-reversible and volume-preserving numerical integrator⁵.
- Generalized Störmer–Verlet [GSV]: implicit symplectic time-reversible integrator⁴

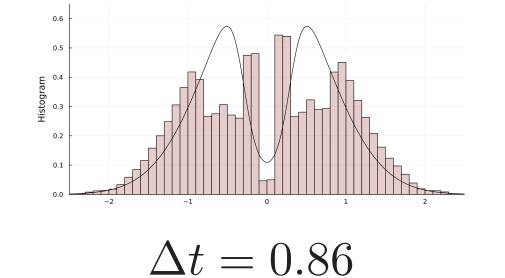
$$\begin{cases} p^{n+\frac{1}{2}} = p^{n} - \frac{\Delta t}{2} \nabla_{q} H\left(q^{n}, p^{n+\frac{1}{2}}\right) \\ q^{n+1} = q^{n} + \frac{\Delta t}{2} \left(\nabla_{p} H\left(q^{n}, p^{n+\frac{1}{2}}\right) + \nabla_{p} H\left(q^{n+1}, p^{n+\frac{1}{2}}\right)\right) \\ p^{n+1} = p^{n+\frac{1}{2}} - \frac{\Delta t}{2} \nabla_{q} H\left(q^{n+1}, p^{n+\frac{1}{2}}\right) \end{cases}$$

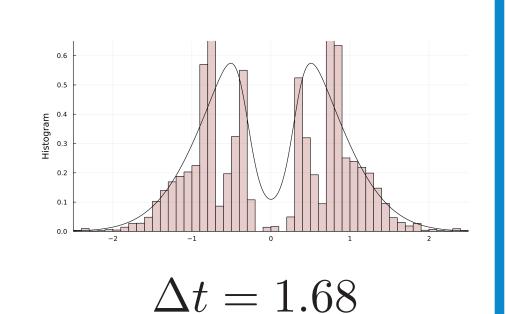
Proposition. GSV is first order weakly consistent with the O.L. dynamics with multiplicative noise (1) when using the Hamiltonian $H(q,p) = V(q) - \frac{1}{2} \ln\left(\det(D(q))\right) + \frac{1}{2} p^{\mathsf{T}} D(q) p$.

Sampling results for a double-well confining potential $V(q) = q^2 - 1 + Ke^{-q^2/(2\sigma)}$, oscillating diffusion coefficient $D(q) = \left(\frac{1+\cos(\pi q)}{2}\right)^2$.









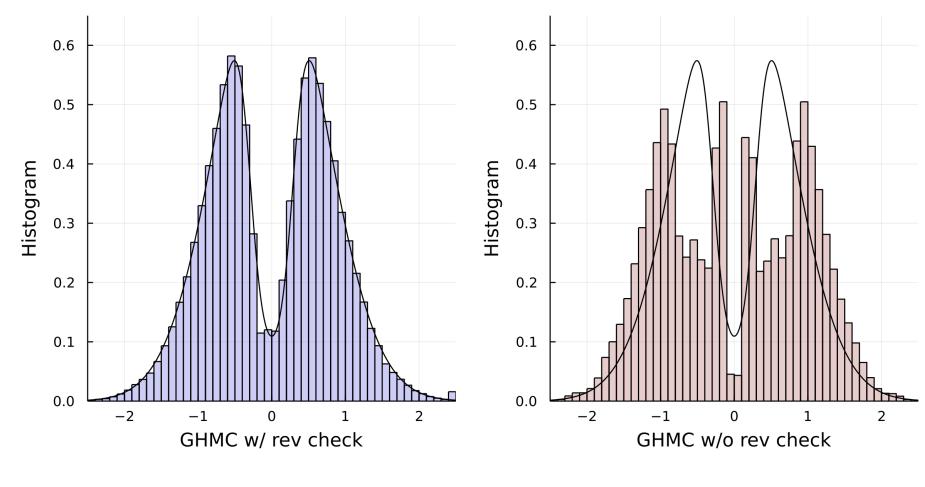
• Unbiased sampling with Metropolis–Hastings scheme: naive approach is biased!

IMPLICIT INTEGRATION

- Need to check for **forward/backward convergence** of implicit method and **reversibil-ity check** when using Newton's/fixed-point iteration methods^{6,7}
- Starting from q_0 , given a time step Δt ,
 - 1. Simulate $p_0 \sim \mathcal{N}(0, D(q_0)^{-1})$,
 - 2. Integrate the Hamiltonian dynamics using GSV during Δt
 - if this *forward* integration does not converge, stay in place: return $(q_1, p_1) = (q_0, p_0)$
 - if it converges to (q_{\star}, p_{\star}) , integrate the Hamiltonian dynamics starting from $(q_{\star}, -p_{\star})$:
 - if this *backward* integration does not converge, stay in place
 - else if the result differs from $(q_0, -p_0)$, stay in place
 - 3. Apply the M–H procedure between (q_0, p_0) and (q_{\star}, p_{\star}) and return the position component.

GHMC can be recovered by integrating an Ornstein-Ulhenbeck process for the first step.

- There are 4 ways to reject the proposal: no forward/backward convergence, no numerical reversibility, M–H ratio computation.
- Even for large time steps, unbiased sampling of the configuration space.



 $\Delta t = 0.86$

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