Statistical Mechanics - Homework 9

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PROBLEM 1

By Likharev equation (5.24) we have

$$\delta N = \sqrt{\left\langle N^2 \right\rangle - \left\langle N \right\rangle^2} = \sqrt{T \frac{\partial \left\langle N \right\rangle}{\partial \mu}}$$

For fermions/bosons, we have

$$\langle N \rangle = \frac{1}{e^{(\varepsilon - \mu)/T} + 1}$$

and for classical particles,

$$\langle N \rangle = e^{-(\varepsilon - \mu)/T}$$

This gives

$$\delta N = \sqrt{\frac{e^{(\varepsilon - \mu)/T}}{\left(e^{(\varepsilon - \mu)/T} \pm 1\right)^2}} = \sqrt{\left\langle N \right\rangle^2 e^{(\varepsilon - \mu)/T}} = \sqrt{\left\langle N \right\rangle \mp \left\langle N \right\rangle^2}$$

for fermions/bosons and

$$\delta N = \sqrt{\frac{1}{e^{(\varepsilon - \mu)/T}}} = \sqrt{\langle N \rangle}$$

for classical particles. When $\langle N \rangle$ is small so that $\langle N \rangle^2$ is negligible, the three expressions become equal.

PROBLEM 2

The partition function is

$$Z = \sum_{k=0}^{N} {N \choose k} e^{\beta h(2k-N)} = e^{-\beta Nh} \left(1 + e^{2\beta h} \right)^{N}$$

The magnetization is

$$M = \left\langle \sum_{i} s_{i} \right\rangle = \frac{1}{Z} \sum_{k=0}^{N} (2k - N) {N \choose k} e^{\beta h(2k - N)} = \frac{T}{Z} \frac{\partial Z}{\partial h} = N \tanh(\beta h)$$

The RMS fluctuation of M is

$$\delta M = \sqrt{\langle M^2 \rangle - \langle M \rangle^2} = \sqrt{N} \operatorname{sech}(\beta h)$$

Mathematica calculations for these quantities are below:

In[132]:=
$$(*Partition Function*)$$

$$Z = e^{-h N \beta} (1 + e^{2 h \beta})^{N};$$

$$(*Magnetization*)$$

$$FullSimplify \left[\frac{1}{\beta Z} D[Z, h] \right]$$

$$Ou[133]:= N Tanh[h \beta]$$

$$In[134]:= (*M is also*)$$

$$FullSimplify \left[\frac{1}{Z} Sum \left[(2 k - N) Binomial[N, k] e^{\beta h (2 k - N)}, \{k, 0, N\} \right] \right]$$

$$Ou[134]:= N Tanh[h \beta]$$

$$(*M^{2}*)$$

$$FullSimplify \left[\frac{1}{Z} Sum \left[(2 k - N)^{2} Binomial[N, k] e^{\beta h (2 k - N)}, \{k, 0, N\} \right] \right]$$

$$Ou[129]:= N (N - (-1 + N) Sech[h \beta]^{2})$$

$$In[135]:= (*-^{2}*)$$

$$FullSimplify \left[(N - (N - 1) Sech[h \beta]^{2}) N - (N Tanh[h \beta])^{2} \right]$$

$$Ou[135]:= N Sech[h \beta]^{2}$$

PROBLEM 3

The partition function is

$$Z = \sum_{i} e^{-\beta E_i} = 2e^{3\beta J} + 6e^{-\beta J}$$

The energy and RMS fluctuation are

$$E = -\frac{1}{Z} \frac{\partial Z}{\partial \beta} \qquad (\delta E)^2 = -\frac{\partial E}{\partial \beta}$$

Calculating these in Mathematica:

In[136]:=

(*Partition Function*)

$$Z = 2 e^{3 \beta J} + 6 e^{-\beta J};$$

In[140]:=

(*Energy*)

 $e = \frac{-1}{Z} D[Z, \beta];$

FullSimplify[e]

Out[141]:=

 $3 \left(-1 + \frac{4}{3 + e^{4 J \beta}}\right) J$

In[142]:=

(* δE *)

FullSimplify[-D[e, β]]

Out[142]:=

 $\frac{48 e^{4 J \beta} J^2}{(3 + e^{4 J \beta})^2}$

PROBLEM 4

In class, we derived $S_F = \frac{T\eta}{\pi}$. We also know the relationship between S_q and S_F :

$$S_{q} = \frac{S_{F}}{m^{2}(\omega_{0}^{2} - \omega^{2})^{2} + \eta^{2}\omega^{2}}$$

Using these two results, we find

$$K_q = \int_{-\infty}^{\infty} d\omega e^{-i\omega t} S_q(\omega) = \frac{T\eta}{\pi} \int_{-\infty}^{\infty} d\omega \frac{e^{-i\omega t}}{m^2(\omega_0^2 - \omega^2)^2 + \eta^2 \omega^2} = \frac{T\eta}{2\pi^2 i} \sum_i \lim_{\omega \to \omega_i} (\omega - \omega_i) \frac{e^{-i\omega t}}{m^2(\omega_0^2 - \omega^2)^2 + \eta^2 \omega^2}$$

where the ω_i are the poles of S_q . Evaluating, we find

FullSimplify[Solve[
$$m^2(\omega \theta^2 - \omega^2)^2 + \eta^2 \omega^2 == 0, \omega$$
]]

Out[26]=

In[42]:=

$$\left\{ \left\{ \omega \to -\frac{\sqrt{-\frac{\eta^{2}-2\,\,\mathrm{m}^{2}\,\,\omega\theta^{2}+\sqrt{\eta^{4}-4\,\,\mathrm{m}^{2}\,\,\eta^{2}\,\,\omega\theta^{2}}}}{\sqrt{2}}} \right\}, \,\, \left\{ \omega \to \frac{\sqrt{-\frac{\eta^{2}-2\,\,\mathrm{m}^{2}\,\,\omega\theta^{2}+\sqrt{\eta^{4}-4\,\,\mathrm{m}^{2}\,\,\eta^{2}\,\,\omega\theta^{2}}}}{\sqrt{2}}} \right\}, \,\, \left\{ \omega \to -\sqrt{-\frac{\eta^{2}}{2\,\,\mathrm{m}^{2}} + \omega\theta^{2} + \frac{\sqrt{\eta^{4}-4\,\,\mathrm{m}^{2}\,\,\eta^{2}\,\,\omega\theta^{2}}}}{2\,\,\mathrm{m}^{2}}} \right\}, \,\, \left\{ \omega \to \sqrt{-\frac{\eta^{2}}{2\,\,\mathrm{m}^{2}} + \omega\theta^{2} + \frac{\sqrt{\eta^{4}-4\,\,\mathrm{m}^{2}\,\,\eta^{2}\,\,\omega\theta^{2}}}}{2\,\,\mathrm{m}^{2}}} \right\} \right\}$$

$$zeros = \left\{ -\frac{\sqrt{-\frac{\eta^{2}-2\,m^{2}\,\omega\theta^{2}+\sqrt{\eta^{4}-4\,m^{2}\,\eta^{2}\,\omega\theta^{2}}}{m^{2}}}}{\sqrt{2}}, \frac{\sqrt{-\frac{\eta^{2}-2\,m^{2}\,\omega\theta^{2}+\sqrt{\eta^{4}-4\,m^{2}\,\eta^{2}\,\omega\theta^{2}}}{m^{2}}}}{\sqrt{2}}, -\sqrt{-\frac{\eta^{2}}{2\,m^{2}}+\omega\theta^{2}+\frac{\sqrt{\eta^{4}-4\,m^{2}\,\eta^{2}\,\omega\theta^{2}}}{2\,m^{2}}}}, \sqrt{-\frac{\eta^{2}}{2\,m^{2}}+\omega\theta^{2}+\frac{\sqrt{\eta^{4}-4\,m^{2}\,\eta^{2}\,\omega\theta^{2}}}}{2\,m^{2}}}\right\};$$

$$Eull Simplify \left[\frac{T\,\eta}{m^{2}} \, Sum[Limit] \, e^{-i\,\omega\,\tau} \, (\omega - zeros[[i]]) \right] = 0.5 \, Zeros[[i]] \, distance (\omega - zeros[[i]])$$

$$\frac{ \operatorname{m} \operatorname{T} \sqrt{\eta^4 - 4\operatorname{m}^2 \eta^2 \omega \Theta^2} \left(\frac{ \sin \left[\frac{\tau \sqrt{-\eta_2^2 + \operatorname{m}^2 \omega \Theta^2 - \frac{1}{2} \sqrt{\eta^4 - 4\operatorname{m}^2 \eta^2 \omega \Theta^2}}{\operatorname{m}} \right]}{\sqrt{-\eta^2 + 2\operatorname{m}^2 \omega \Theta^2 - \sqrt{\eta^4 - 4\operatorname{m}^2 \eta^2 \omega \Theta^2}}} - \frac{ \sin \left[\frac{\tau \sqrt{-\eta^2 + 2\operatorname{m}^2 \omega \Theta^2 + \sqrt{\eta^4 - 4\operatorname{m}^2 \eta^2 \omega \Theta^2}}}{\sqrt{2\operatorname{m}}} \right]}{\sqrt{-\eta^2 + 2\operatorname{m}^2 \omega \Theta^2 + \sqrt{\eta^4 - 4\operatorname{m}^2 \eta^2 \omega \Theta^2}}} \right]}{\sqrt{2\operatorname{m}}} \right]}{\sqrt{2\operatorname{m}}}$$

PROBLEM 5

We can find the spectral density from

$$(\delta I)^2 = K_I(0) = \int_{-\infty}^{\infty} d\omega \ S_I(\omega)$$

The fluctuations in I are given by

$$(\delta I)^2 = \left(\delta[qN]\right)^2 = q^2(\delta N)^2 = q^2 \langle N \rangle = q\bar{I}$$

so we find

$$q\bar{I} = \int_{-\infty}^{\infty} d\omega S_I(\omega) \implies S_I(\omega) = q\bar{I}\delta(\omega)$$