

## PHY 622 - Homework 6

M. Ross Tagaras  
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### PART (A)

*Note:* I can't figure out how to typeset the double-plus symbol, so I'm just going to write ++ everywhere instead.

The spinning string action is

$$S = \frac{1}{\pi l^2} \int d\sigma^{++} d\sigma^= \left( 2\partial_{++} X \cdot \partial_= X + \frac{i}{2} (\psi^+ \cdot \partial_= \psi^+ + \psi^- \cdot \partial_{++} \psi^-) + \frac{1}{2} F \cdot F \right)$$

The individual terms transform as

$$(\partial_{++} X \partial_= X)' = \Lambda [\partial_{++} X] \Lambda^{-1} [\partial_= X] = \partial_{++} X \partial_= X$$

$$(\psi^+ \partial_= \psi^+)' = \Lambda^{1/2} \psi^+ \Lambda^{-1} \partial_= \Lambda^{1/2} \psi^+ = \psi^+ \partial_= \psi^+$$

$$(\psi^- \partial_{++} \psi^-)' = \Lambda^{-1/2} \psi^- \Lambda \partial_{++} \Lambda^{-1/2} \psi^- = \psi^- \partial_{++} \psi^-$$

$$(F^\mu F^\nu \eta_{\mu\nu})' = \Lambda^\mu{}_\rho F^\rho \Lambda^\nu{}_\sigma F^\sigma \eta_{\mu\nu} = F^\rho F^\sigma \eta_{\rho\sigma}$$

so the entire action is invariant.

### PART (B)

Since  $X^\mu$  is a scalar field, it transforms as  $(X^\mu(\sigma'))' = X^\mu(\sigma)$ . If we Taylor expand the left side, we find (to first order)

$$X^{\mu'}(\sigma) + \partial_\nu X^\mu(\sigma)(\sigma' - \sigma)^\nu = X^\mu(\sigma) \implies \delta X^\mu = \lambda (\sigma^= \partial_= - \sigma^{++} \partial_{++}) X^\mu$$

From the rule  $\psi^\pm(\sigma') = \Lambda^{\pm 1/2} \psi^\pm(\sigma)$  for spinors, Taylor expanding gives

$$\psi^{\pm'}(\sigma) + \partial_\mu \psi^\pm(\sigma)(\sigma' - \sigma)^\mu = \left(1 \pm \frac{1}{2}\right) \psi^\pm(\sigma) \implies \delta \psi^\pm = \lambda \left(\pm \frac{1}{2} + \sigma^= \partial_= - \sigma^{++} \partial_{++}\right) \psi^\pm$$

### PART (C)

The variation of the action is

$$\begin{aligned} \delta S = \frac{1}{\pi l^2} \int & \left[ 2\partial_{++} (\lambda \sigma^= \partial_= X - \lambda \sigma^{++} \partial_{++} X) \cdot \partial_= X + \partial_{++} X \cdot \partial_= (\lambda \sigma^= \partial_= X - \lambda \sigma^{++} \partial_{++} X) \right. \\ & + \frac{i}{2} \left( \left[ \frac{\lambda}{2} + \lambda \sigma^= \partial_= - \lambda \sigma^{++} \partial_{++} \right] \psi^+ \partial_= \psi^+ + \psi^+ \partial_= \left( \left[ \frac{\lambda}{2} + \lambda \sigma^= \partial_= - \lambda \sigma^{++} \partial_{++} \right] \psi^+ \right) \right) \\ & \left. + \frac{i}{2} \left( \left[ -\frac{\lambda}{2} + \lambda \sigma^= \partial_= - \lambda \sigma^{++} \partial_{++} \right] \psi^- \partial_{++} \psi^- + \psi^- \partial_{++} \left( \left[ -\frac{\lambda}{2} + \lambda \sigma^= \partial_= - \lambda \sigma^{++} \partial_{++} \right] \psi^- \right) \right) \right] \end{aligned}$$

If we look at the bosonic terms, we see that we can make them all proportional to  $\partial_{++}\partial_{=}X$  if we integrate the first and third terms by parts. If we do this and drop the total derivatives, we get

$$\delta S_{bosonic} = \frac{2\lambda}{\pi l^2} \int [\sigma^{++}\partial_{++}X + \sigma^=\partial_{=}X - \sigma^=\partial_{=}X - \sigma^{++}\partial_{++}X] \partial_{++}\partial_{=}X = 0$$

In the fermion terms, we immediately see that the terms with  $(\partial_{++}\psi^+)^2$  and  $(\partial_{=}\psi^-)^2$  are identically zero. The first line of fermionic terms is then

$$\begin{aligned} & \psi^+\partial_{=}\psi^+ - \sigma^{++}(\partial_{=}\psi^+)\partial_{++}\psi^+ + \psi^+\partial_{=}\sigma^=\partial_{=}\psi^+ - \psi^+\partial_{=}\sigma^{++}\partial_{++}\psi^+ \\ &= \psi^+\partial_{=}\psi^+ - \sigma^{++}(\partial_{++}\psi^+)\partial_{=}\psi^+ + \partial_{=}\psi^+\sigma^=\partial_{=}\psi^+ - (\partial_{=}\psi^+)^2\sigma^= - \psi^+\partial_{=}\psi^+ - \sigma^{++}\psi^+\partial_{=}\partial_{++}\psi^+ \\ &= -\sigma^{++}(\partial_{++}\psi^+)\partial_{=}\psi^+ - \partial_{=}\sigma^{++}\psi^+\partial_{++}\psi^+ + \sigma^{++}(\partial_{=}\psi^+)\partial_{++}\psi^+ = 0 \end{aligned}$$

The second line vanishes by similar manipulations.

#### **PART (D)**

$SO(1,1)$  is abelian, so all commutators are trivially zero.

#### **PART (E)**

I'm not sure that the idea of semi-local Lorentz symmetry makes any physical sense.