

## PHY 623 - Homework 4

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### PROBLEM 1

The contraction of  $X$  with itself is

$$-\frac{l^2}{4} [\ln(z-w) - \ln(\bar{z}-\bar{w}) - \ln(z-\bar{w}) - \ln(\bar{z}-w)]$$

The OPE is

$$\begin{aligned} \partial_z X(z, \bar{z}) \partial_z X(z, \bar{z}) e^{ikX(w, \bar{w})} &= \left( \partial_z [\ln(z-w) - \ln(z-\bar{w})] \right)^2 \left[ \frac{(ik)^2}{2} \left( \frac{-l^2}{4} \right) + \frac{(ik)^3}{3!} 3 \left( \frac{-l^2}{4} \right) X(w, \bar{w}) + \dots \right] \\ &\quad + 2ik \partial_z [\ln(z-w) - \ln(z-\bar{w})] \left( 1 + ikX(w, \bar{w}) + \frac{(ik)^2}{2} X(w, \bar{w}) X(w, \bar{w}) + \dots \right) \\ &= \frac{l^2 k^2}{8} \left( \frac{1}{z-w} - \frac{1}{z-\bar{w}} \right)^2 e^{ikX(w, \bar{w})} + 2ik \left( \frac{1}{z-w} - \frac{1}{z-\bar{w}} \right) e^{ikX(w, \bar{w})} \end{aligned}$$

We see that the conformal dimension of  $e^{ikX}$  is  $\frac{l^2 k^2}{8}$ .

### PROBLEMS 2,3,4

The stress tensor for a general  $bc$  system is

$$T = (1-\lambda)(\partial b)c - \lambda b \partial c$$

Now we consider the OPE of  $T(z)T(w)$ . The terms with two contractions are

$$\begin{aligned} (1-\lambda)^2 \partial_z \left( \frac{\epsilon}{z-w} \right) \partial_w \left( \frac{1}{z-w} \right) - \lambda(1-\lambda) \left[ \partial_z \partial_w \left( \frac{\epsilon}{z-w} \right) \frac{1}{z-w} + \frac{\epsilon}{z-w} \partial_z \partial_w \left( \frac{1}{z-w} \right) \right] + \lambda^2 \partial_w \left( \frac{\epsilon}{z-w} \right) \partial_z \left( \frac{1}{z-w} \right) \\ = \frac{-\epsilon(1-\lambda)^2}{(z-w)^2} \frac{1}{(z-w)^2} + \frac{\epsilon\lambda(1-\lambda)}{(z-w)^4} + \frac{-\epsilon\lambda^2}{(z-w)^4} = \frac{-\epsilon(12\lambda^2 - 12\lambda + 2)}{2(z-w)^4} \end{aligned}$$

The terms with one contraction are

$$\begin{aligned} (1-\lambda)^2 \partial_z \left( \frac{\epsilon}{z-w} \right) c(z) \partial_w b + (1-\lambda)^2 \partial_w \left( \frac{1}{z-w} \right) (\partial_z b) c(w) \\ - \lambda(1-\lambda) \partial_z \partial_w \left( \frac{\epsilon}{z-w} \right) c(z) b(w) - \lambda(1-\lambda) \frac{1}{z-w} (\partial_z b) (\partial_w c) \\ - \lambda(1-\lambda) \frac{\epsilon}{z-w} (\partial_z c) \partial_w b - \lambda(1-\lambda) b(z) c(w) \partial_z \partial_w \left( \frac{1}{z-w} \right) \\ + \lambda^2 \partial_w \left( \frac{\epsilon}{z-w} \right) (\partial_z c) b(w) + \lambda^2 b(z) (\partial_w c) \partial_z \left( \frac{1}{z-w} \right) \end{aligned}$$

$$= -(1-\lambda)\epsilon \frac{1}{(z-w)^2} c(z) \partial_w b + (1-\lambda) \frac{1}{(z-w)^2} (\partial_z b) c(w) + \lambda \epsilon \frac{1}{(z-w)^2} (\partial_z c) b(w) - \lambda \frac{1}{(z-w)^2} b(z) \partial_w c$$

Taylor expanding gives

$$\begin{aligned} & \frac{1}{(z-w)} \left\{ (1-\lambda)^2 \left[ \partial_w b + (z-w) \partial_w^2 b + \dots \right] c(w) - (1-\lambda) \epsilon \left[ c(w) + (z-w) \partial_w c + \dots \right] \partial_w b \right\} \\ & + \frac{1}{(z-w)^2} \left\{ \lambda \epsilon \left[ \partial_w c + (z-w) \partial_w^2 c + \dots \right] b(w) - \lambda \left[ b(w) + (z-w) \partial_w b \right] \partial_w c \right\} \\ = & \frac{1}{(z-w)^2} \left[ (1-\lambda)(\partial b)c - \epsilon(1-\lambda)c\partial b + \lambda\epsilon(\partial c)b - \lambda b\partial c \right] + \frac{1}{z-w} \left[ (1-\lambda)(\partial^2 b)c - \epsilon(1-\lambda)(\partial c)\partial b + \lambda\epsilon(\partial^2 c)b - (\partial b)\partial c \right] \\ = & \frac{2}{(z-w)^2} \left[ (1-\lambda)(\partial b)c - \lambda b\partial c \right] + \frac{1}{z-w} \left[ (1-\lambda)\partial((\partial b)c) - \lambda\partial(b\partial c) \right] \\ = & \frac{2T(w)}{(z-w)^2} + \frac{\partial T(w)}{z-w} \end{aligned}$$

All together, the  $TT$  OPE is

$$T(z)T(w) = \frac{-\epsilon(12\lambda^2 - 12\lambda + 2)}{2(z-w)^4} + \frac{2T(w)}{(z-w)^2} + \frac{\partial T(w)}{z-w}$$

We can now look at specific  $b$  and  $c$ . If we take  $\psi \equiv \frac{1}{\sqrt{2}}(b+c)$ ,  $\epsilon = 1$ , and  $\lambda = 1/2$ , then we obtain the result for the  $TT$  OPE with  $T^{(\psi)}(z) = -\frac{1}{2i}\psi(z)\partial_z\psi(z)$ . We find a central charge of  $1/2$  in this case, thanks to the extra factor of  $1/2$  contributed by our definition of  $\psi$ .

For the susy ghosts, we take  $b = \beta$ ,  $c = \gamma$ ,  $\lambda = 3/2$  and  $\epsilon = -1$ , which gives a central charge of  $11$ . For the coordinate ghosts, we choose  $\lambda = 2$  and  $\epsilon = 1$ , which gives a central charge of  $-26$ .

We already know that the  $T^{(x)}T^{(x)}$  OPE gives a central charge of  $1$ . For a system with  $D$  coordinate scalars and the corresponding coordinate ghosts, we find central charge cancellation only in  $D = 26$ . If we include fermions and susy ghosts as well, then we find that  $D$  must satisfy  $D + D/2 - 26 + 11 = 0$ . The only solution is  $D = 10$ .