

PHY 622 - Homework 2

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PROBLEM 1

The induced metric is $h_{\alpha\beta} = \partial_\alpha X^\mu \partial_\beta X^\nu \delta_{\mu\nu}$. The components are

$$h_{\theta\theta} = \partial_\theta X^\mu \partial_\theta X_\mu = [\partial_\theta (R \cos \phi \sin \theta)]^2 + [\partial_\theta (R \sin \phi \sin \theta)]^2 + [\partial_\theta (R \cos \theta)]^2 = R^2$$

$$h_{\theta\phi} = h_{\phi\theta} = \partial_\phi X^\mu \partial_\theta X_\mu = \partial_\theta (R \cos \phi \sin \theta) \partial_\phi (R \cos \phi \sin \theta) + \partial_\theta (R \sin \phi \sin \theta) \partial_\phi (R \sin \phi \sin \theta) = 0$$

$$h_{\phi\phi} = \partial_\phi X^\mu \partial_\phi X_\mu = R^2 \sin^2 \theta$$

So as a matrix, h is

$$h_{\alpha\beta} = R^2 \begin{pmatrix} 1 & 0 \\ 0 & \sin^2 \theta \end{pmatrix}$$

with $\sqrt{\det h_{\alpha\beta}} = R^2 \sin \theta$. The action is

$$\int d\theta d\phi \sqrt{h} = 2\pi R^2 \int_0^\pi d\theta \sin \theta = 4\pi R^2$$

In stereographic coordinates, the points $X^\mu(x, y, z)$ on S^2 are the points that lie on the line connecting the points $(0, 0, R)$ and $(\xi_1, \xi_2, -R)$ such that $x^2 + y^2 + z^2 = R^2$. The equation of this line is

$$\frac{x}{\xi_1} = \frac{y}{\xi_2} = \frac{R - z}{2R}$$

From these two equations, we can find expressions for x, y , and z . The first equality gives

$$x^2 = \frac{R^2 - z^2}{1 + \frac{\xi_2^2}{\xi_1^2}}$$

and the second gives

$$z = R - \frac{2Rx}{\xi_1}$$

Combining these:

$$x^2 = -\frac{4R^2 x^2}{\xi^2} + \frac{4R^2 \xi_1 x}{\xi^2}$$

which simplifies to

$$x = \frac{4R^2 \xi_1}{4R^2 + \xi^2}$$

In a similar manner, we find that

$$y = \frac{4R^2\xi_2}{4R^2 + \xi^2}$$

To find $z(\xi_1, \xi_2, R)$, we can simply combine the previous two expressions using $x^2 + y^2 + z^2 = R^2$:

$$z^2 = R^2 - \frac{16R^4\xi^2}{(4R^2 + \xi^2)^2} = \frac{R^2(16R^4 + \xi^4 - 8R^2\xi^2)}{(4R^2 + \xi^2)^2}$$

so we see that

$$z = \frac{R(\xi^2 - 4R^2)}{4R^2 + \xi^2}$$

With these new coordinates, we can now calculate the induced metric. A few useful formulas are:

$$\frac{\partial x^i}{\partial \xi^j} = \frac{4R^2\delta_j^i(\xi^2 + 4R^2) - 8R^2\xi^i\xi_j}{(\xi^2 + 4R^2)^2} \quad \frac{\partial z}{\partial \xi^i} = \frac{16R^3\xi_i}{(\xi^2 + 4R^2)^2}$$

The determinant of the induced metric is

$$h = \left[\left(\frac{\partial x^1}{\partial \xi^1} \right)^2 + \left(\frac{\partial x^2}{\partial \xi^1} \right)^2 + \left(\frac{\partial x^3}{\partial \xi^1} \right)^2 \right] \left[\left(\frac{\partial x^1}{\partial \xi^2} \right)^2 + \left(\frac{\partial x^2}{\partial \xi^2} \right)^2 + \left(\frac{\partial x^3}{\partial \xi^2} \right)^2 \right] - \left[\frac{\partial x^1}{\partial \xi^1} \frac{\partial x^1}{\partial \xi^2} + \frac{\partial x^2}{\partial \xi^1} \frac{\partial x^2}{\partial \xi^2} + \frac{\partial x^3}{\partial \xi^1} \frac{\partial x^3}{\partial \xi^2} \right]^2$$

which after some tedious algebra can be shown to be

$$h = \frac{256R^8}{(4R^2 + \xi^2)^4}$$

The action is

$$S = \int_{-\infty}^{\infty} d\xi_1 \int_{-\infty}^{\infty} d\xi_2 \frac{16R^4}{(4R^2 + \xi^2)^2} = \int_{-\infty}^{\infty} d\xi_1 \int_{-\infty}^{\infty} d\xi_2 \frac{16R^4}{(4R^2 + \xi_1^2 + \xi_2^2)^2} = \int_{-\infty}^{\infty} d\xi_1 \frac{8\pi R^4}{(4R^2 + \xi_1^2)^{3/2}} = 4\pi R^2$$

PROBLEM 2

We define the induced metric as

$$h_{\alpha\beta} = \partial_\alpha X^\mu \partial_\beta X^\nu \eta_{\mu\nu}$$

Its determinant is then

$$\det h_{\alpha\beta} = h_{00}h_{11} - h_{01}h_{10}$$

$$= \ell^4 \left\{ \left[-(\partial_0 t)^2 + (\partial_0 \sigma)^2 \right] \left[-(\partial_1 t)^2 + (\partial_1 \sigma)^2 \right] - \left[-(\partial_0 t)(\partial_1 t) + (\partial_0 \sigma)(\partial_1 \sigma) \right] \left[-(\partial_1 t)(\partial_0 t) + (\partial_1 \sigma)(\partial_0 \sigma) \right] \right\} = -\ell^4$$

Therefore, we see that the minus sign under the square root is needed to ensure that $\sqrt{-h}$ is real.

PROBLEM 3

The action is

$$\begin{aligned} S &= -T \int dt \int_0^\pi d\sigma \sqrt{-h} = -T \int dt \int_0^\pi d\sigma \sqrt{[\partial_t(\ell t)]^2 [\partial_\sigma(f(\sigma))]^2} = -T\ell \int dt \int_0^\pi d\sigma \partial_\sigma f(\sigma) \\ &= -T\ell \int dt [f(\pi) - f(0)] = \int dt [-T\ell a] \end{aligned}$$

Since our string has no dynamics by choice of gauge, there should be no kinetic energy, and therefore the remaining term in the Lagrangian is $-V$.

To obtain the equations of motion for the Nambu-Goto action, it is easier to use the form

$$\mathcal{L} = -T \sqrt{\left(\frac{\partial X^\mu}{\partial t} \frac{\partial X_\mu}{\partial \sigma}\right)^2 - \frac{\partial X^\mu}{\partial t} \frac{\partial X_\mu}{\partial t} \frac{\partial X^\nu}{\partial \sigma} \frac{\partial X_\nu}{\partial \sigma}}$$

The equations of motion are then

$$0 = \partial_\mu \frac{\partial \mathcal{L}}{\partial(\partial_\mu X)} = \partial_t \left(\frac{\left(\frac{\partial X^\nu}{\partial t} \frac{\partial X_\nu}{\partial \sigma}\right) \frac{\partial X^\mu}{\partial \sigma} - \frac{\partial X^\nu}{\partial \sigma} \frac{\partial X_\nu}{\partial \sigma} \frac{\partial X^\mu}{\partial t}}{\sqrt{\left(\frac{\partial X^\mu}{\partial t} \frac{\partial X_\mu}{\partial \sigma}\right)^2 - \frac{\partial X^\mu}{\partial t} \frac{\partial X_\mu}{\partial t} \frac{\partial X^\nu}{\partial \sigma} \frac{\partial X_\nu}{\partial \sigma}}} \right) + \partial_\sigma \left(\frac{\left(\frac{\partial X^\nu}{\partial t} \frac{\partial X_\nu}{\partial \sigma}\right) \frac{\partial X^\mu}{\partial t} - \frac{\partial X^\nu}{\partial t} \frac{\partial X_\nu}{\partial t} \frac{\partial X^\mu}{\partial \sigma}}{\sqrt{\left(\frac{\partial X^\mu}{\partial t} \frac{\partial X_\mu}{\partial \sigma}\right)^2 - \frac{\partial X^\mu}{\partial t} \frac{\partial X_\mu}{\partial t} \frac{\partial X^\nu}{\partial \sigma} \frac{\partial X_\nu}{\partial \sigma}}} \right)$$

In the static gauge, we find that

$$\frac{\partial X^\nu}{\partial t} \frac{\partial X_\nu}{\partial \sigma} = 0 \quad \frac{\partial X^\nu}{\partial \sigma} \frac{\partial X_\nu}{\partial \sigma} = (\partial_\sigma f)^2 \quad \frac{\partial X^\nu}{\partial t} \frac{\partial X_\nu}{\partial t} = -\ell^2$$

so the equations of motion become

$$0 = \partial_t \left(\frac{-(\partial_\sigma f)^2 \frac{\partial X^\mu}{\partial t}}{\ell} \right) + \partial_\sigma \left(\frac{\ell \frac{\partial X^\mu}{\partial \sigma}}{\partial_\sigma f} \right)$$

These are trivially satisfied for $\mu = 2, \dots, d-1$. For $\mu = 0$, we find

$$0 = -\partial_t (\partial_\sigma f)^2 + \partial_\sigma (0) = 0$$

For $\mu = 1$, we find

$$0 = \partial_t (0) + \partial_\sigma \left(\frac{\ell \partial_\sigma f}{\partial_\sigma f} \right) = 0$$