

PH223: Physics for Chemists and Mechanical Engineers

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Preface

Preface

This document contains my lecture notes for a new, experimental course. The goal of the course is to teach the introductory physics of waves, optics, and electricity and magnetism for mechanical engineering students.

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Acknowledgments

BYU-I

R. Todd Lines.

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1 Where We Start

Fundamental Concepts

What is this class?

This class is designed to teach the physics of wave motion, electricity and magnetism, and optics. We have three major goals. One is to teach the physics that is not covered by Statics, Dynamics, and the Engineering Electronics Course. This physics can affect the mechanical systems you will design, build, or test, so knowing this physics is a very good thing. The second objective is to teach a different method of thinking about how things work. The third goal is to describe electrical and wave motion enough that the quantum nature of atoms and molecules make sense as our chemists take physical chemistry.

In engineering, the design parameters are often the goal. In physics, the physical relationship is the goal. For design engineers, both views are useful and important. The design is no good if the underlying principles preclude it from working!

As an example, I once worked on an optics project with a strong mechanical component. The system had scanning mechanisms that were fantastic mechanical devices. It was part of an aircraft and integrated into the aircraft system. But the optical system required two lasers that were separated in wavelength by only a few nanometers. The chief engineer knew how to build all the systems, but did not understand the physics that required the close wavelength spacing. He judged that the difficulty in building the device at that wavelength spacing outweighed any benefit, and he changed the specs to give two wavelengths that were fifty nanometers apart. Fifty nanometers is a pretty small tolerance. Surely it would be good enough! The resulting product did not work. For two years he tried to fine tune the scanners, and servos to make it work. After ten million dollars and two years, he finally moved the wavelengths closer. The cost of the change was an extra \$100,000 dollars, about 1/100 of the cost of the mistake. The

2 Chapter 1 Where We Start

system worked, but since this was a race to market, the time lost and the reputation lost on the faulty product destroyed the viability of the business. It is a bad day when you and your friends lose your jobs because you made a fundamental physics mistake!

Physics courses stress how we know what we know. They support the discipline called *system engineering*, which deals with the design of new and innovative products. As a more positive example, the National Weather service often releases requests for proposed weather sensing equipment. Their request might say something like the following:

Measure the moisture of the soil globally from an altitude of 800 km with an accuracy of 5%. The suggested instrument is a passive microwave radiometer.

The job of a system engineer is to determine what type of instrument to build. What is the underlying principle that it will use to do its job? What signal processing will it need? What mechanical and electrical systems will support this? This must all be determined before the bearings and slip-rings, and structures can be designed and built.

The radiometer design that came out of this project is flying today (or one very like it based on the original design) and is a major part of the predictive models that tell us what the weather will be in a few days.

Because this type of reasoning is our goal, we will not only do typical homework problems, but we will also work on our conceptual understanding.

I will also emphasize a problem solving method that I used with my engineering team in industry. It is a structured approach to finding a solution that emphasizes understanding as well as providing a numeric answer for a particular design. When you are part of an innovative design team, you will have to repeat a calculation over and over again each time some other part of the sign changes. If you have produced a symbolic solution, a numerical model, or at least a curve, you are ready for any changes in specifications. But if you have just “found the answer” you will have to find that answer again every time the overall design specs change. This approach is too slow, and, at least in my team, would have you finding a new job because our design efforts were always done against exacting schedules and budgets. By thinking in a structured method, with an eye toward symbolic answers or relationships rather than end numbers, you will learn to be a more valuable engineer. The process we will use is the same approach I used to teach my new engineers in the defense industry. It has been proven useful over and over for decades.

This same problem solving process is useful in chemistry, particularly as you study physical chemistry.

So let's get started. To understand waves, we need to get the waves moving. You studied Oscillation in Dynamics or PH121. Oscillating systems are often the disturbance that starts a wave. We will begin with a review of oscillation.

Simple Harmonic Motion

You are, no doubt, an expert in simple harmonic motion (SHM) after your Dynamics class. But this will get us warmed up for the semester. In class we will use our clickers and go through a few questions. We will usually use the clicker system to answer a few questions to test your understanding of the reading material. This allows me to not waste time on things you already know, and to help me find the ones you don't. Most lectures will consist of me asking you if you have questions, and then if you don't, I will ask you "clicker questions." Where there is reason to believe you don't understand (with a normal cutoff of 80% of the class answering correctly being our definition of "understanding"), I will use the material from these written lectures to teach the concepts. So we won't always go through all the ideas and skills demonstrated in these written lectures. If you feel you would have liked more explanation on something but we did not cover that concept in class because most people were "getting it," you can come and see me in my office.

SHM

Question 223.1.1

Question 223.1.2

QQuestion 223.1.3

Question 223.1.4

Question 223.1.5

Question 223.1.6

Question 15.3.8

Question 15.3.9.2

Question 15.3.9.3

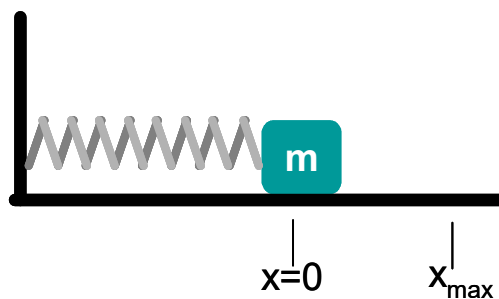
Question 15.3.9.5

Question 15.4

Let's consider a mass attached to a spring resting on a frictionless surface¹. This mass-spring system can oscillate.

In the position shown the spring is neither pushing nor pulling on the mass. We will call this position the *equilibrium position* for the mass.

¹ Yes, I know there are no actual frictionless surfaces, but we are starting out at freshman level physics, so we will make the math simple enough that a freshman could do it by making simplifying assumptions. In this case, that the surface is frictionless.



Definition 1.1 *Equilibrium Position:* The position of the mass when the spring is neither stretched nor compressed.

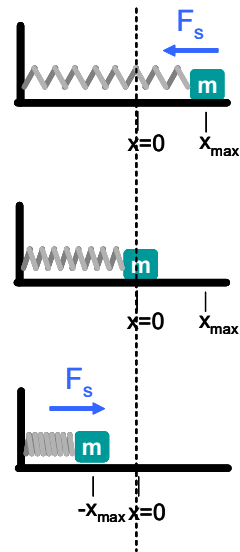
Hooke's Law

A law in physics is a mathematical expression of a mental model of how the universe works. Long ago it was noticed that the pull of a spring grew in strength as the spring was pulled out of equilibrium. The mathematical expression of this is

$$F_s = -kx \quad (1.1)$$

The force, F_s is directly proportional to the displacement from equilibrium, x . Since a man named Hooke wrote this down, it is called Hooke's law.

Hooke's Law is, strictly speaking, not a law that is always obeyed. It is a good model for most springs as long as we don't stretch them too far. We will often use the word "law" to mean *an equation that gives a basic relationship*. In that sense, Hook's law is a law.



Lets write Hooke's Law using Newton's second Law

$$\Sigma F_x = ma_x$$

If we assume no friction, we have just

$$-kx = ma_x$$

We can write this as

$$a_x = -\frac{k}{m}x \quad (1.2)$$

This expression says the acceleration is directly proportional to the position, and opposite the direction of the displacement from equilibrium. We can see that the spring force tries to oppose the change in displacement. We call such a force a *restoring force*.

Definition 1.2 *Restoring force: A force that is always directed toward the equilibrium position*

This is a good definition of *simple harmonic motion*.

Mathematical Representation of Simple Harmonic Motion

Recall from your Dynamics or PH121 classes that acceleration is the second derivative of position

$$a = \frac{dv}{dt} = \frac{d^2x}{dt^2}$$

Hook's Law tells us

$$\begin{aligned} F &= ma = -kx \\ m \frac{d^2x}{dt^2} &= -kx \end{aligned}$$

We have a new kind of equation. If you are taking this freshman physics class as a... well... freshman, you may not have seen this kind of equation before. It is called a differential equation. But really the chances are that you are a sophomore or junior (or even a senior) and have lot of experience with differential equations. The solution of this equation is a function or functions that will describe the motion of our mass-spring system as a function of time. We will need to know this function, so let's see how we can find it.

Start by defining a quantity ω as

$$\omega^2 = \frac{k}{m} \quad (1.3)$$

why define ω^2 ? Because experience has shown that it is useful to define ω this way! But you probably remember ω as having to do with rotational speed, and from trigonometry (trig) you may remember using ω to mean angular frequency

$$\omega = 2\pi f$$

so our definition of ω may hint that k/m will have something to do with the frequency of oscillation of the mass-spring system.

We can write our differential equation as

$$\frac{d^2x}{dt^2} = -\omega^2 x \quad (1.4)$$

To solve this differential equation we need a function who's second derivative is the negative of itself. We know a few of these

$$\begin{aligned} x(t) &= A \cos(\omega t + \phi_o) \\ x(t) &= A \sin(\omega t + \phi_o) \end{aligned} \quad (1.5)$$

where A , ω , and ϕ_o are constants that we must find. Let's choose the cosine function

and explicitly take its derivatives.

$$\begin{aligned}x(t) &= A \cos(\omega t + \phi_o) \\ \frac{dx(t)}{dt} &= -\omega A \sin(\omega t + \phi_o) \\ \frac{d^2x(t)}{dt^2} &= -\omega^2 A \cos(\omega t + \phi_o)\end{aligned}$$

Let's substitute these expressions into our differential equation for the motion

$$\begin{aligned}\frac{d^2x}{dt^2} &= -\omega^2 x \\ -\omega^2 A \cos(\omega t + \phi_o) &= -\omega^2 A \cos(\omega t + \phi_o)\end{aligned}$$

As long as the constant ω^2 is our $\omega^2 = k/m$ we have a solution (now you know why we defined it as ω^2 !). Since from trig we remember ω as the angular frequency.

$$\omega = 2\pi f$$

Thus

$$\omega = \sqrt{\frac{k}{m}} = 2\pi f \quad (1.6)$$

The frequency of oscillation depends on the mass and the stiffness of the spring.

$$f = \frac{1}{2\pi} \sqrt{\frac{k}{m}} \quad (1.7)$$

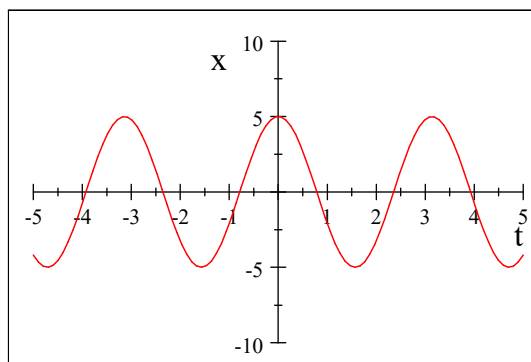
Let's see if this is reasonable. Imagine driving along in your student car (say, a 1972 Gremlin). You go over a bump, and the car oscillates. Your car is a mass, and your shock absorbers are springs. You have an oscillation. But suppose you load your car with everyone in your apartment². Now as you hit the bump the car oscillates at a different frequency, a lower frequency. That is what our frequency equation tells us. Note also that if we changed to a different set of shocks, the k would change, and we would get a different frequency.

We still don't have a complete solution to our differential equation, because we don't know A and ϕ_o . From trigonometry, we recognize ϕ_o as the initial phase angle. We will call it the *phase constant* in this class. We will have to find this by knowing the initial conditions of the motion. We will do this in a minute.

A is the amplitude. We can find its value when the motion has reached its maximum displacement. Let's look at a specific case

$$\begin{aligned}A &= 5 \\ \phi_o &= 0 \\ \omega &= 2\end{aligned}$$

² If you are married, imagin taking two other couples with you in your car.



We can easily see that the amplitude A corresponds to the maximum displacement x_{\max} .

Other useful quantities we can identify

We know from trigonometry that a cosine function has a period T .

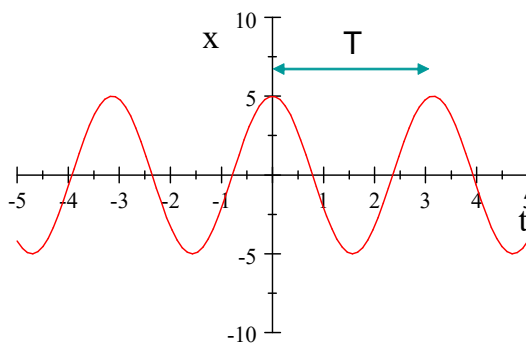


Figure 1.1.

The period is related to the frequency

$$T = \frac{1}{f} = \frac{2\pi}{\omega} \quad (1.8)$$

We can write the period and frequency in terms of our mass and spring constant

$$T = 2\pi\sqrt{\frac{m}{k}} \quad (1.9)$$

$$f = \frac{1}{2\pi}\sqrt{\frac{k}{m}} \quad (1.10)$$

Velocity and Acceleration

Since we know the derivatives of

$$x(t) = A \cos(\omega t + \phi_o) \quad (1.11)$$

we can identify the velocity of the mass and its acceleration

$$v(t) = \frac{dx(t)}{dt} = -\omega A \sin(\omega t + \phi_o)$$

Recall that $A = x_{\max}$

$$v(t) = \frac{dx(t)}{dt} = -\omega x_{\max} \sin(\omega t + \phi_o) \quad (1.12)$$

We identify

$$v_{\max} = \omega x_{\max} = x_{\max} \sqrt{\frac{k}{m}} \quad (1.13)$$

Likewise for the acceleration

$$a(t) = \frac{dv(t)}{dt} \quad (1.14)$$

$$= \frac{d}{dt} (-\omega x_{\max} \sin(\omega t + \phi_o)) \quad (1.15)$$

$$= -\omega^2 x_{\max} \cos(\omega t + \phi_o)$$

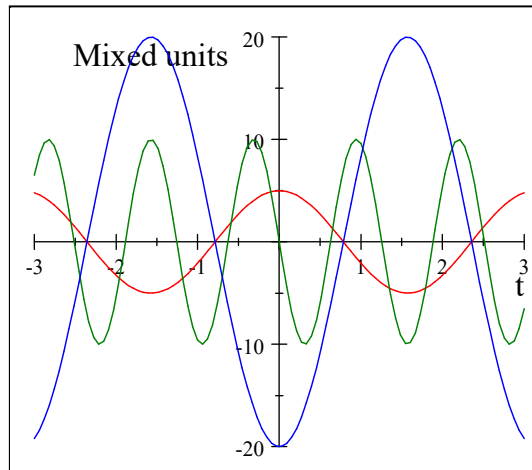
where we can identify

$$a_{\max} = \omega^2 x_{\max} = \frac{k}{m} x_{\max} \quad (1.16)$$

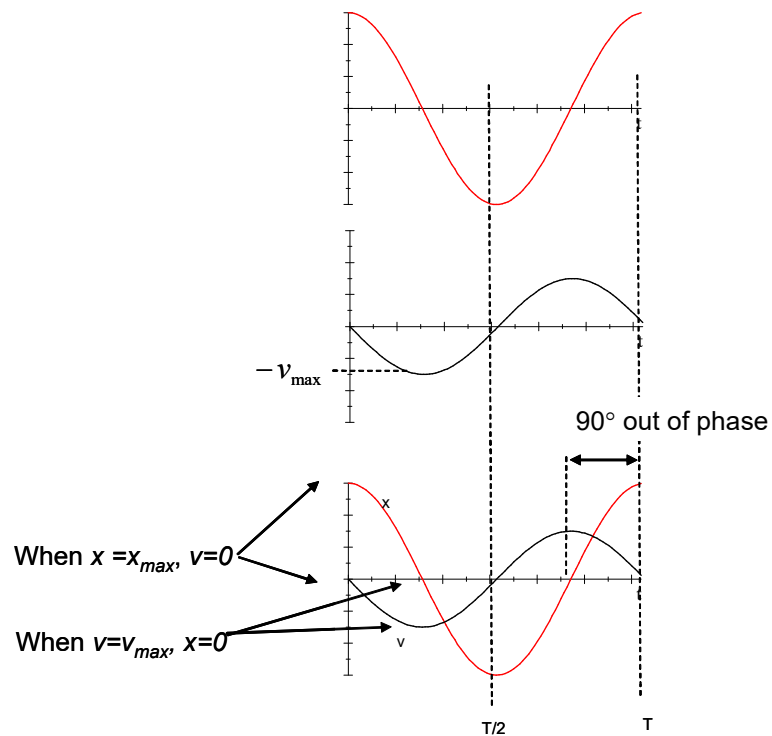
Comparison of position, velocity, acceleration

Don't do in class

Let's plot $x(t)$, $v(t)$, and $a(t)$ for a specific case



Red is the displacement, green is the velocity, and blue is the acceleration. Note that each has a different maximum amplitude. That is a bit confusing until we recognize that they each have different units. We have just plotted them on the same graph to make it easy to compare their phases. Note that they are not in phase!



The acceleration is 90° out of phase from the velocity.

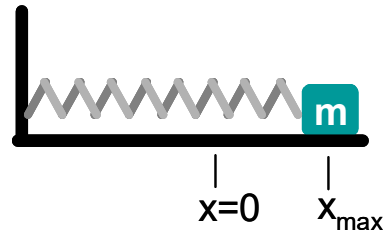
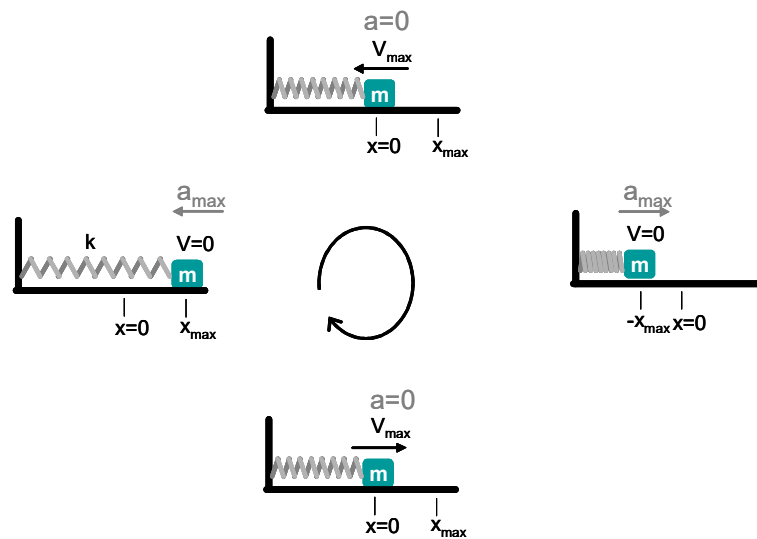


Figure 1.2.



An example of oscillation

We want to see how to find A , ω , and especially ϕ_o . These quantities will be important in our study of waves. So let's do a problem.

Let's take as our system a horizontal mass-spring system where the mass is on a frictionless surface.

Initial Conditions

Now let's find A and ϕ_o . To do this we need to know how we started the mass-spring motion. We call the information on how the system starts its motion the *initial*

conditions.

Suppose we start the motion by pulling the mass to $x = x_{\max}$ and releasing it at $t = 0$. These are our initial conditions. Let's see if we can find the phase. Our initial conditions require

$$\begin{aligned}x(0) &= x_{\max} \\v(0) &= 0\end{aligned}\tag{1.17}$$

Using our formula for $x(t)$ and $v(t)$ we have

$$\begin{aligned}x(0) &= x_{\max} = x_{\max} \cos(0 + \phi_o) \\v(0) &= 0 = -v_{\max} \sin(0 + \phi_o)\end{aligned}\tag{1.18}$$

From the first equation we get

$$1 = \cos(\phi_o)$$

which is true if

$$\phi_o = 0, 2\pi, 4\pi, \dots$$

from the second equation we have

$$0 = \sin \phi_o$$

which is true for

$$\phi_o = 0, \pi, 2\pi, \dots$$

If we choose $\phi_o = 0$, these conditions are met. Of course we could choose $\phi_o = 2\pi$, or $\phi_o = 4\pi$, but we will follow the rule to take the smallest value for ϕ_o that meets the initial conditions.

A second example

Using the same equipment, let's start with

$$\begin{aligned}x(0) &= 0 \\v(0) &= +v_i\end{aligned}\tag{1.19}$$

that is, the mass is moving, and we start watching just as it passes the equilibrium point.

$$\begin{aligned}x(0) &= 0 = x_{\max} \cos(0 + \phi_o) \\v(0) &= v_i = -v_{\max} \sin(0 + \phi_o)\end{aligned}\tag{1.20}$$

from

$$0 = x_{\max} \cos(\phi_o)$$

(first equation above) we see that³

$$\phi_o = \pm \frac{\pi}{2}$$

but we don't know the sign. Using our initial velocity condition

$$v_i = -v_{\max} \sin\left(\pm \frac{\pi}{2}\right)$$

$$v_i = -\omega x_{\max} \sin\left(\pm \frac{\pi}{2}\right)$$

We defined the initial velocity as positive, and we insist on having positive amplitudes, so x_{\max} is positive. Thus we need a minus sign from $\sin(\phi_o)$ to make v_i positive. This tells us to choose

$$\phi_o = -\frac{\pi}{2}$$

with a minus sign.

Our solutions are

$$x(t) = \frac{v_i}{\omega} \cos\left(\omega t - \frac{\pi}{2}\right)$$

$$v(t) = v_i \sin\left(\omega t - \frac{\pi}{2}\right)$$

Remark 1.1 *Generally to have a complete solution to a differential equation, you must find all the constants (like A and ϕ_o) based on the initial conditions.*

A third example

So far we have concentrated on finding ϕ_o . Let's do a more complete example where we find ϕ_o , A , and ω .

A particle moving along the x axis in simple harmonic motion starts from its equilibrium position, the origin, at $t = 0$ and moves to the right. The amplitude of its motion is 4.00 cm, and the frequency is 1.50 Hz.

a) show that the position of the particle is given by

$$x = (4.00 \text{ cm}) \sin(3.00\pi t)$$

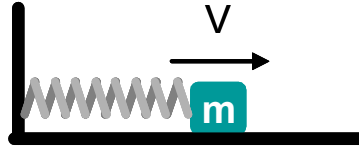
determine

b) the maximum speed and the earliest time ($t > 0$) at which the particle has this speed,

³ Really there are more possibilities, but we are taking the smallest value for ϕ_o as we discussed above.

c) the maximum acceleration and the earliest time ($t > 0$) at which the particle has this acceleration, and

d) the total distance traveled between $t = 0$ and $t = 1.00$ s



Type of problem

We can recognize this as an oscillation problem. This leads us to a set of basic equations

Basic Equations

$$x(t) = A \cos(\omega t + \phi_o)$$

$$v(t) = -\omega x_{\max} \sin(\omega t + \phi_o)$$

$$a(t) = -\omega^2 A \cos(\omega t + \phi_o)$$

$$\omega = 2\pi f$$

$$v_m = \omega x_m$$

$$a_m = \omega^2 x_m$$

$$T = \frac{1}{f}$$

We should write down what we know and give a set of variables

Variables

t	time, initial time = 0	$t_o = 0$
x	Position, Initial position = 0	$x(0) = 0$
v		
a		
x_m	x amplitude	$x_m = 4.00 \text{ cm}$
v_m	v amplitude	
a_m	a amplitude	
ω	angular frequency	
ϕ_o	phase	
f	frequency	$f = 1.50 \text{ Hz}$

Now we are ready to start solving the problem. We do this with algebraic symbols first

Symbolic Solution

Part (a)

We can start by recognizing that we can find ω because we know the frequency. We just use the basic equation.

$$\omega = 2\pi f$$

We also know the amplitude $A = x_{\max}$ which is given. Knowing that

$$x(0) = 0 = A \cos(0 + \phi_o)$$

we can guess that

$$\phi_o = \pm \frac{\pi}{2}$$

Using

$$v(0) = -\omega x_{\max} \sin\left(0 \pm \frac{\pi}{2}\right)$$

again and demanding that amplitudes be positive values, and noting that at $t = 0$ the velocity is positive from the initial conditions:

$$\phi_o = -\frac{\pi}{2}$$

We also note from trigonometry that

$$x(t) = x_{\max} \cos\left(2\pi ft - \frac{\pi}{2}\right)$$

which is a perfectly good answer. However, if we remember our trig, we could write this using

$$\cos\left(\theta - \frac{\pi}{2}\right) = \sin(\theta)$$

Then we have

$$\begin{aligned} x(t) &= x_{\max} \cos\left(2\pi ft - \frac{\pi}{2}\right) \\ &= x_{\max} \sin(2\pi ft) \end{aligned}$$

Part (b)

We have a basic equation for v_{\max}

$$\begin{aligned} v_m &= \omega x_{\max} \\ &= 2\pi f x_{\max} \end{aligned}$$

to find when this happens, take

$$v(t) = v_{\max} = -\omega x_{\max} \sin\left(2\pi ft - \frac{\pi}{2}\right)$$

and recognize that $\sin(\theta) = 1$ is at a maximum when $\theta = \pi/2$ so the entire argument of the sine function must be $\pi/2$ when we are at the maximum displacement, so

$$\frac{\pi}{2} = \left(2\pi ft - \frac{\pi}{2}\right)$$

or

$$\pi = 2\pi ft$$

then the time is

$$\frac{1}{2f} = t$$

Part (c)

Like with the velocity we must use a basic formula, this time

$$a(t) = -\omega^2 A \cos(\omega t + \phi_o)$$

but recognize that the maximum is achieved when $\cos(\omega t + \phi_o) = 1$ or when $\omega t + \phi_o = 0$

$$\begin{aligned} t &= \frac{\phi_o}{\omega} \\ &= \frac{-\frac{\pi}{2}}{2\pi f} \\ &= \frac{-1}{4f} \end{aligned}$$

The formula for a_{\max} is

$$\begin{aligned} a_{\max} &= -\omega^2 x_{\max} \\ &= -(2\pi f)^2 x_m \end{aligned}$$

Part (d)

We know the period is

$$T = \frac{1}{f}$$