J. Phys. Soc. Japan 20 (1965) 1093

## Electron Effective Mass of SnO<sub>2</sub>

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(Received April 6, 1965)

The electron effective mass of  $SnO_2$  has been measured with thin films<sup>1)~3)</sup> and natural crystals<sup>4)</sup>. However, the reported values have been scattered between 0.1 and 1.0 m, and the reliable value is not yet obtained. We wish to report here the value of the electron effective mass estimated from the measurements of Hall effect and thermoelectric power with n-type  $SnO_2$  single crystals prepared by the vapor reaction method<sup>5)</sup>.

The temperature dependence of the Hall mobility is presented in Fig. 1.  $T^{-2}$  dependence is observed above 160°K. Since the power -2 is rather close to -3/2, it is likely that acoustical phonons are most important in the electron scattering process for this temperature region. The deviation from the  $T^{-3/2}$  dependence suggests that optical phonons also contribute to the scattering process to some extent.

The carrier concentration data shown by circles in Fig. 1 were calculated from the Hall coefficients, assuming that  $\mu_H/\mu=3\pi/8$ . The sample was found

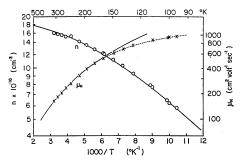


Fig. 1. Temperature dependence of the Hall mobility  $\mu_H$  and of the carrier concentration n for a  $\mathrm{SnO_2}$  single crystal. The experimental data are indicated by crosses for  $\mu_H$  and by circles for n. The solid line for  $\mu_H$  represents  $\mu_H = \mathrm{constant} \times T^{-2}$ . The solid line for n represents the calculated curve of eq. (1) with  $N_D = 1.95 \times 10^{17} \mathrm{~cm^{-3}}$ ,  $N_A = 1 \times 10^{16} \mathrm{~cm^{-3}}$ ,  $m^* = 0.40m$  and  $E_D = 0.024 \mathrm{~eV}$ .

to be nearly in the carrier saturation region over the whole temperature region investigated. Extraporating the measured points to  $T=\infty$ ,  $(N_D-N_A)$  was estimated to be  $(1.85\pm0.05)\times10^{17}\,\mathrm{cm}^{-3}$ , where  $N_A$  and  $N_D$  are the acceptor and donor concentrations, respectively. In analyzing the carrier concentration data, the following well known formula for the electron concentration n derived from nondegenerate statistics was used.

$$n(n+N_A)/(N_D-N_A-n) = (N_C/g) \exp(-E_D/kT)$$
. (1)

Here  $N_C=2(2\pi m^*kT/h^2)^{3/2}$ ,  $m^*$  being the density-of-states mass, and  $E_D$  is the ionization energy of the donor centers. The donor degeneracy factor g was assumed to be 2, corresponding to the simple spin degeneracy. Taking the experimental error in to account and assuming that  $N_A \leq N_D/10$ , it is possible to satisfactorily fit the data with eq. (1) using the values of  $N_A=(1\pm1)\times 10^{16}~{\rm cm}^{-3}$ ,  $N_D=(1.75\pm0.15)\times 10^{17}~{\rm cm}^{-3}$ ,  $m^*=(0.41\pm0.10)m$  and  $E_D=0.024\pm0.004~{\rm eV}$ , as is seen in Fig. 1.

Measurements of thermoelectric power Q at room temperature were also performed. The data were analyzed with the familiar formula,

$$Q = -(k/e)[(5/2 - s) + \ln(N_U/n)], \qquad (2)$$

where s is the scattering factor. The contribution of the phonon-drag effect to the thermoelectric power was neglected in eq. (2), since the temperatures are fairly high. Using the values of n obtained from the Hall effect measurements and assuming that s=-1/2 which corresponds to the acoustical phonon scattering, the density-of-states mass was found to be 0.33m. This agrees with that obtained above from the analysis of the carrier concentration within the experimental error.

It is concluded from our experiments that the electron effective mass of  $SnO_2$  lies near 0.35m.

More detailed study on the electron effective mass of SnO<sub>2</sub>, including free carrier absorption and reflection measurements, is now in progress.

## References

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