Device-Independent Quantum Key Distribution (DI-QKD) Secure Against Collective Attacks





Non-Locality and Contextuality (2nd Semester - 2023/2024) Técnico Lisboa, ULisboa

by

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Introduction

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Cryptography in the post-quantum era

Future quantum threats on cryptography include:

- Simon's Algorithm
 - Brute-force attacks on cryptographic key spaces
 - ► Impact on Advanced Standard Encryption (AES)!
- Grover's Algorithm
 - Brute-force attacks on cryptographic key spaces
 - ► AES-128 and AES-192 are no longer secure!
- Brassard-Høyer-Tapp (BHT) Algorithm
 - Combination of Grover's algorithm and Birthday Paradox
 - ► SHA-3-224 and SHA-3-256 are no longer secure!
- Shor's Algorithm
 - Solves factorization and discrete logarithm problems
 - Completely breaks Rivest-Shamir-Adleman (RSA) and Elliptic Curve Cryptography (ECC)!
- New public-key cryptography and key exchanges are needed!

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Cryptography in the post-quantum era

Two new main approaches arise...

	(Classical) Post-Quantum Cryptography	Quantum Cryptography
Foundation	(Still Believed) Hard Mathematical Problems	Quantum Mechanics and Physics
Type of Information	Classical	(Mainly) Quantum
Encoding	N/A	Discrete-Variables (DV) for qubits or Continuous-Variables (CV) for qumodes
Strategies	N/A	Prepare-and-Measure or Entanglement
Families	Lattice-based, Code-based, Hash-based, Isogeny-based, Multivariate, and Zero-Knowledge Proofs (ZKPs)	Quantum Key Distribution (QKD), Semi-Quantum Key Distribution (QKD), Quantum Conference Key Agreement (QCKA), Quantum Digital Signature Scheme (QDSS), Quantum Bit Commitment (QBC), Quantum Oblivious Transfer (QOT), and Quantum Multi-Party Computation (QMPC)
Popular Primitives	CRYSTALS-Kyber, CRYSTALS-Dilithium, FALCON, SPHINCS+, McEliece, HQC, and BIKE	BB84, B92, SSP, SARG04, E91, BBM92, KMB09, T12, Decoy State, Squeezed State, DPS, MSZ96, GG02

Table 1: Overview of the two main approaches for cryptography in the post-quantum era

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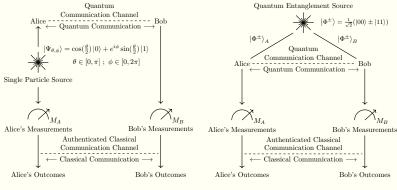


Figure 1a: High-level procedure schematic of a Prepare-and-Measure QKD protocol

Figure 1b: High-level procedure schematic of an Entanglement-based QKD protocol

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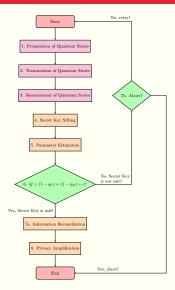


Figure 2: Flowchart of a QKD protocol

1. Preparation of Quantum States

Single Particles, Entangled Particles, Coherent States, Fock States, etc.

2. Transmission of Quantum States

- Uses a quantum communication channel with a certain efficiency η_C
- "Flying" quantum states can be eavesdropped, introducing a noise ε

3. Measurement of Quantum States

- Uses quantum measurement devices with a certain efficiency η_D
- Composes a raw key

4. Secret Key Sifting

Identifies which protocol rounds can be used to compose a sifted key

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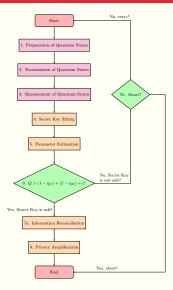


Figure 2: Flowchart of a QKD protocol

5. Parameter Estimation

- Samples and evaluates the Quantum Bit Error Rate (QBER) Q
- Estimates the key rate
- Estimates the secure mutual information between Alice and Bob
- Estimates the Holevo bound

6. (Eavesdropping detected?)

$$\Rightarrow$$
 Q > $(1 - \eta_{C}) + (1 - \eta_{D}) + \epsilon$?

7a. Yes! \Rightarrow Information Reconciliation

- Applies an Error Correction Code (ECC) to correct the sifted into an error-free key
- These ECC algorithms include:
 - Cascade protocol, Winnow protocol, and Low-Density Parity-Check (LDPC) codes
- Can be accelerated by classical software and hardware (e.g., OpenMP, CUDA, etc.)

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7b. No! \Rightarrow Abort? (or Retry?)

Introduction

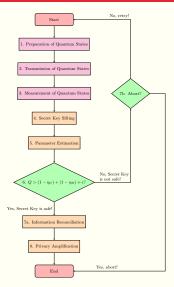


Figure 2: Flowchart of a OKD protocol

8. Privacy Amplification Estimation

- Applies a Universal Hash Function on the error-free key, composing the final (and amplified) secret key
- This Universal Hash Function can be a Toeplitz Hashing procedure, usually applied using a random seed as well
- Can be accelerated by classical software and hardware (e.g., OpenMP, CUDA, etc.)
- The Secret Key Sifting, Parameter Estimation, Information Reconciliation, and Privacy Amplification steps require an authenticated and interactive exchange of classical information/messages
 - We can achieve it with a Carter-Wegman Message Authentication Code (CW-MAC), requiring an initial small secret defined a priori

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Attacks on QKD protocols

Independent and identically distributed (i.i.d.) rounds:

- The devices behave independently and in the same way
- The quantum states distributed are always the same

There are three main attacks on QKD protocols:

- Individual Attacks:
 - The eavesdropper has no quantum memory
 - The eavesdropper can only attack individually each round

Collective Attacks:

- The eavesdropper has no quantum memory
- The eavesdropper can perform arbitrary global operations

Coherent Attacks:

- The eavesdropper has quantum memory (trace of rounds)
- The eavesdropper can perform arbitrary global operations
- The parties' quantum states can be arbitrarily correlated

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Problem

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Can a QKD protocol with untrusted quantum devices be secure?

Considering Entanglement-based QKD protocols:

- ▶ The entangled particles are emitted from a common source
- The parties measure each particle on a randomly chosen basis
- ▶ Here, we assume that:
 - ► The locations of the parties (Alice and Bob) are secure
 - Alice and Bob trust their measuring devices
- The source of the entangled particles:
 - Does not need to be trusted by Alice and Bob
 - Might be under the control of an eavesdropper (Eve)

And about untrusted quantum measurement devices?

- No guarantees on the expected measurement bases
- No assumptions on the dimension of the Hilbert Space

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Motivation

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- Device-Independent Quantum Key Distribution (DI-QKD) is a concept of security for QKD protocols that:
 - Seeks to ensure the security of QKD protocols:
 - Without considering any details about the internal working of the quantum devices being used:
 - The quantum devices can be imperfect, untrusted, or manipulated by a malicious party
 - Based on the violation of Bell Inequalities, ensuring:
 - Quantum correlations between the quantum devices
 - The security is inferred directly from those quantum correlations observed on the outcomes
 - Do not exist any local hidden variables
 - It is a "holy-grail" on Quantum Cryptography!

Motivation 12/38

- Device-Independent Quantum Key Distribution (DI-QKD) requires the following basic assumptions:
 - The physical locations of the parties are secure
 - No unwanted information can leak out to the outside
 - The parties have a Trusted Random Number Generator (TRNG), producing a classical random output
 - Possibly, one derived from thermal noise or based on a Quantum Random Number Generator (QRNG)
 - The parties have trusted classical devices
 - Capable of storing and processing the classical data generated by their quantum devices
 - The parties share a public authenticated classical communication channel
 - The parties can start with a small shared secret
 - Quantum Mechanics is correct (and well-defined)

Motivation 13/38

- Some reasons lead the (usual) QKD protocols to be insecure in Device-Independent (DI) scenarios:
 - Sometimes they produce classical correlations
 - ▶ We can reproduce them without quantum mechanics
 - We can generate them from a set of classical random data shared by the parties' systems
 - Those classical correlations can be written as:

$$-P(ab|XY) = \sum_{\lambda} P(\lambda) \times D(a|X,\lambda) \times D(b|Y,\lambda)$$

- Where:
- λ is a classical variable with probability distribution $P(\lambda)$, shared by the parties' quantum devices
- $D(a|X, \lambda)$ is a function that completely specifies Alice's outputs once the input X and the variable λ are given
- $D(b|Y, \lambda)$ is a function that completely specifies Bob's outputs once the input Y and the variable λ are given

Motivation 14/38

- Some reasons lead the (usual) QKD protocols to be insecure in Device-Independent scenarios:
 - A copy of λ will give the full information about the outputs a and b to Eve, once the inputs X and Y are announced
 - However... The strategy for these correlations is not available to the eavesdropper if the outputs a and b:
 - Are correlated in a non-local way
 - Violate a Bell Inequality
 - Therefore, the violation of a Bell Inequality is a key requirement for the security of DI-QKD protocols!

Motivation 15/38

Results

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Let's consider the following QKD protocol:

- ➤ Alice and Bob share an entangled Werner quantum state
 - $\begin{array}{l} \blacktriangleright \ \, \rho_{\mathit{AB}} = p |\Phi^{+}\rangle \langle \Phi^{+}| + (1-p)\frac{\mathbb{I}}{4} \\ \text{Where: } |\Phi^{+}\rangle = \frac{1}{\sqrt{2}}(|00\rangle + |11\rangle) \text{ and} \\ \text{the term } \frac{\mathbb{I}}{4} \text{ represents white noise} \end{array}$
- They choose a measurement to apply to their particles for each round, resulting on binary outcomes, where:
 - ► Alice has three measurements choices: $X \in \{A_0, A_1, A_2\}$ • $A_0 = \sigma_z$ • $A_1 = \frac{(\sigma_z + \sigma_x)}{\sqrt{2}}$ • $A_2 = \frac{(\sigma_z - \sigma_x)}{\sqrt{2}}$
 - ▶ Bob has two measurements choices: $Y \in \{B_0, B_1\}$
 - $B_1 = \sigma_z$ $B_2 = \sigma_x$
 - ▶ The binary outcomes are denoted as $\{+1, -1\}$

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Recall that:

$$\rho_{AB} = \begin{bmatrix} \frac{(1+\rho)}{4} & 0 & 0 & \frac{\rho}{2} \\ 0 & \frac{(1-\rho)}{4} & 0 & 0 \\ 0 & 0 & \frac{(1-\rho)}{4} & 0 \\ \frac{\rho}{2} & 0 & 0 & \frac{(1+\rho)}{4} \end{bmatrix}$$

$$A_0 = B_1 = \sigma_z = \begin{bmatrix} 1 & 0 \\ 0 & -1 \end{bmatrix}$$

▶
$$A_0 = B_1 = \sigma_z = \begin{bmatrix} 1 & 0 \\ 0 & -1 \end{bmatrix}$$
 ▶ $A_1 = \frac{(\sigma_z + \sigma_x)}{\sqrt{2}} = \begin{bmatrix} \frac{1}{\sqrt{2}} & \frac{1}{\sqrt{2}} \\ \frac{1}{\sqrt{2}} & -\frac{1}{\sqrt{2}} \end{bmatrix}$

$$B_2 = \sigma_{\mathsf{X}} = \begin{bmatrix} 0 & 1 \\ 1 & 0 \end{bmatrix}$$

$$\textbf{3} \quad B_2 = \sigma_{\textbf{X}} = \begin{bmatrix} 0 & 1 \\ 1 & 0 \end{bmatrix}$$

$$\textbf{3} \quad A_2 = \frac{(\sigma_{\textbf{Z}} - \sigma_{\textbf{X}})}{\sqrt{2}} = \begin{bmatrix} \frac{1}{\sqrt{2}} & \frac{1}{\sqrt{2}} \\ \frac{1}{\sqrt{2}} & -\frac{1}{\sqrt{2}} \end{bmatrix}$$

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Regarding this QKD protocol:

- The (initial) raw key is extracted from the resulting outcomes of the pair of measurements $\{A_0, B_1\}$:
 - For which the QBER Q is defined as follows:

•
$$Q = P(a \neq b|01) = P(a \neq b|A_0, B_1) =$$

= $P(a = 0, b = 1|A_0, B_1) + P(a = 1, b = 0|A_0, B_1)$

- In this context, the QBER Q is used for:
 - Estimating the amount of quantum correlations between the parties, using the same measurement
 - Quantifying the amount of classical communication required for the Error Correction protocol/code during the Information Reconciliation step

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Regarding this QKD protocol:

- ► The measurements A₁, A₂, B₁, and B₂ are used on a subset of the particles to estimate the Clauser-Horne-Shimony-Holt (CHSH) polynomial:
 - - Where the correlators are defined as $\langle a_i b_i \rangle = P(a = b|i,j) P(a \neq b|i,j)$
 - The CHSH polynomial is used by the parties to:
 - Bound Eve's potential partial information about the key
 - Define how much secret information leaked to Eve needs to be reduced during the Privacy Amplification step
- The parameters Q and S are used to estimate the information available to a potential eavesdropper

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Regarding this QKD protocol:

- The CHSH polynomial's correlations satisfy:

 - $ightharpoonup S = 2\sqrt{2}p = 2\sqrt{2}(1-2Q)$
- Regarding the CHSH polynomial, we have:
 - ► Classically correlated data, for $p \le \frac{1}{\sqrt{2}}$, and thus, $S \le 2$
 - In this case, secure DI-QKD protocol is not possible
 - Maximal quantum violation, for p=1, and thus, $S=2\sqrt{2}$
 - In this case, the potential eavesdropper has has no available information about the secret key
 - Now, we can interpolate for the range $rac{1}{\sqrt{2}} !$
- To bound the eavesdropper's information:
 - No assumptions about:
 - Behaviour of quantum measurements choices X and Y
 - Dimension of the quantum systems ρ_{AB}

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Let's consider some eavesdropping strategies:

- For the most general attacks:
 - The only data available to the parties to bound the eavesdropper's knowledge is:
 - The observed relation between the inputs and outputs
 - No assumptions on the type of quantum measurements and quantum physical systems used are made
 - ▶ Generally, we can model these attacks as a tripartite entangled quantum state $|\Psi\rangle_{ABF} \in \mathcal{H}_A^{\otimes n} \otimes \mathcal{H}_B^{\otimes n} \otimes \mathcal{H}_E$
 - Where: n is the number of bits of the raw key
 - ► The size of the Hilbert Space of the parties' systems is:
 - Unknown to the parties
 - Fixed (and known) to the eavesdropper

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Let's consider some eavesdropping strategies:

- > Focusing on collective attacks:
 - ► The eavesdropper applies the same attack to each quantum physical system of the parties
 - The quantum states are i.i.d., and thus, $|\Psi
 angle_{\mathit{ABE}}=|\psi
 angle_{\mathit{ABE}}^{\otimes n}$
 - The quantum measurement devices
 - Have no memory register
 - Behave i.i.d. in every round of the QKD protocol
 - The (asymptotic) secret key rate r has a lower bound given by the Devetak-Winter key rate r_{DW} formula:

•
$$r \geq r_{DW} = \underbrace{I(A_0:B_1)}_{\mbox{Mutual information between Alice and Bob}} - \underbrace{\chi(B_1:E)}_{\mbox{Holevo quantity between Eve and Bob}}$$

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Let's consider some eavesdropping strategies:

- Focusing on collective attacks:
 - ▶ The mutual information between Alice and Bob is given as:

•
$$I(A_0:B_1) = \underbrace{H(A_0)}_{\mbox{Individual (binary)}} + \underbrace{H(B_1)}_{\mbox{Shannon entropy}} - \underbrace{H(A_0,B_1)}_{\mbox{Shannon entropy}}$$

Shannon entropy

for Alice and Bob

- Since we assume uniform marginals, we also have:
 - $I(A_0:B_1)=1-H(Q)$ $\Big\}$ Individual (binary) Shannon entropy on QBER
- The Holevo quantity between Eve and Bob is given as:

•
$$\chi(B_1:E) = S(\rho_E) - \frac{1}{2} \sum_{b_1=\pm 1} S(\rho_{E|b_1})$$

Where:

- ρ_E denotes the Eve's quantum state after (partially) tracing out Alice and Bob's particles, i.e., $\rho_E = Tr_{AB} \left(|\psi\rangle_{ABF} \langle \psi|_{ABF} \right)$
- $\rho_{E|b_1}$ denotes the Eve's quantum state when Bob has obtained the outcome result b_1 for the measurement setting $B_1 = \sigma_Z$

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Let's consider some eavesdropping strategies:

- Security against collective attacks:
 - The optimal collective attack occurs when:
 - The tripartite entangled quantum state $|\psi\rangle_{ABE}$ is the purification of the (original) bipartite entangled quantum state ρ_{AB}
 - The Holevo quantity $\chi(B_1:E)$ achieves its possible largest value (compatible with the parameters Q and S)

When the parties symmetrize their uniform marginals:

$$\blacktriangleright \chi(B_1:E) \leq h\left(rac{1+\sqrt{(rac{5}{2})^2-1}}{2}
ight)$$
 Theorem for DI-QKD

Considering the optimal collective attack:

- We have to consider $\chi(B_1:E)=h\left(\frac{1+\sqrt{(\frac{5}{2})^2-1}}{2}\right)$
 - Without violating the Bell Inequality (for $S \leq 2$), the Holevo bound will be $\chi(B_1 : E) \leq h\left(\frac{1}{2}\right) \Leftrightarrow \chi(B_1 : E) \leq 1$ (full information for Eve)
- ► The key rate is given by $r \ge 1 h(Q) h\left(\frac{1 + \sqrt{\left(\frac{5}{2}\right)^2 1}}{2}\right)$

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- Simplifying the calculations for the DI-QKD protocol...
 - Recall that for the CHSH polynomial, we have:

►
$$S = 2\sqrt{2}(1-2Q)$$

• We can simplify the (maximum) Holevo bound $\chi(B_1:E)$ and Devetak-Winter key rate r_{DW} for the DI-QKD protocol:

$$\chi(B_1:E) = h\left(\frac{1+\sqrt{(\frac{5}{2})^2 - 1}}{2}\right) = h\left(\frac{1+\sqrt{\left(\frac{2\sqrt{2}(1-2Q)}{2}\right)^2 - 1}}{2}\right)$$

$$r \ge r_{DW} = 1 - h(Q) - h\left(\frac{1+\sqrt{(\frac{5}{2})^2 - 1}}{2}\right) =$$

$$r \ge r_{DW} = 1 - h(Q) - h\left(\frac{1 + \sqrt{\left(\frac{2\sqrt{2}(1 - 2Q)}{2}\right)^2 - 1}}{2}\right) = 1 - h(Q) - h\left(\frac{1 + \sqrt{\left(\frac{2\sqrt{2}(1 - 2Q)}{2}\right)^2 - 1}}{2}\right)$$

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- Simplifying the calculations for the (usual) Entanglement-based (E91) QKD protocol...
 - Recall that for the CHSH polynomial, we have:

$$S = 2\sqrt{2}(1-2Q)$$

The respective Holevo bound is given as follows:

$$\chi(B_1:E) \leq h\left(Q + \frac{s}{2\sqrt{2}}\right)$$

We can simplify the (maximum) Holevo bound $\chi(B_1:E)$ and the Devetak-Winter key rate r_{DW} for the (usual) Entanglement-based QKD protocol:

$$\chi(B_1:E) = h\left(Q + \frac{s}{2\sqrt{2}}\right) = h\left(Q + \frac{2\sqrt{2}(1-2Q)}{2\sqrt{2}}\right) =$$

$$= h\left(Q + (1-2Q)\right) = h\left(1-Q\right)$$

$$r \ge r_{DW} = 1 - h(Q) - h\left(Q + \frac{s}{2\sqrt{2}}\right) =$$

$$= 1 - h(Q) - h\left(Q + \frac{2\sqrt{2}(1-2Q)}{2\sqrt{2}}\right)$$

$$= 1 - h(Q) - h\left(Q + (1-2Q)\right) = 1 - h(Q) - h(1-Q)$$

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How do QBERs Q impact the (usual) Entanglement-based QKD and DI-QKD protocols?

· For Holevo bounds:

- Greater Holevo bounds for the DI-QKD protocol
- We can easily detect the presence of an eavesdropper for the DI-QKD protocol, allowing to tolerate better the QBER - For a QBER Q around 14%, the eavesdropper has all the information about the raw key in the DI-QKD protocol

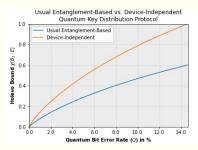


Figure 3: Holevo bounds with respect to QBER Q

For Devetak-Winter key rates:

- Lower Devetak-Winter key rates for the DI-QKD protocol
 The noise introduced by the eavesdropper will have greater impact on the DI-QKD protocol, reducing more the key rate
 - For a QBER Q around 7%, no extractable secure raw key
 will be possible in the DI-OKD protocol

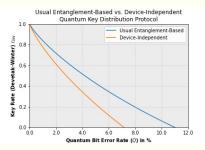


Figure 4: Devetak-Winter key rates with respect to QBER Q

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What is a loophole in Bell experiments?

- Flaw/gap in the experiment that may allow the results to:
 - Be explained by local-hidden variables theory
 - ▶ Be falsely interpreted as violating Bell's Inequality due to imperfections or limitation in the design of the experiment

There are two main loopholes in Bell experiments:

- Locality loophole
 - The quantum measurement devices can:
 - Be close to each other and/or delay their actions
 - Communicate between themselves and influence the outcomes

Detection loophole

- The quantum measurement devices can:
 - Fail to detect and measure all the incoming particles
 - Not receive the incoming particles due to transmission losses

Local models explain and determine these loopholes!

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What is a loophole-free Bell experiment?

- A Bell experiment requiring:
 - A party cannot know any information about the other party's input before producing its own output
 - 2. Measurement devices with high detection efficiencies Ex.: A detection efficiency $\eta_D > 82.8\%$, for the CHSH Inequality

Loopholes in the perspective of a DI-QKD protocol

- How to circumvent the locality loophole?
 - Considerable spatial separation between the parties
 - No sub-luminal signals travel between their quantum devices
 - Proper and secure isolation of the parties' locations
 - No unwanted information can leak to the outside

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Loopholes in the perspective of a DI-QKD protocol

- How to circumvent the detection loophole?
 - (Possible) Post-selection on the measurement data
 - Discard no-detection events and keep only the events on which both quantum measurement devices produce an outcome
 - Fair sampling assumption
 - The sample of detected particles is a fair, random and unbiased sample of the set of all particles
 - No correlations between the quantum state of the particles and their detection probability
 - Clearly unjustified for DI-QKD protocols, where we assume that the quantum devices can be provided by an untrusted party
 - This loophole still needs to be closed for DI-QKD protocols!

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- Handling detection loophole for a DI-QKD protocol
 - Addressing no-detection events for the security proof:
 - We need to consider all the measurement outcomes
 - Detection " ± 1 " and no-detection outcomes " \pm "
 - A possible strategy is to consider only two of them - Replace all no-detection outcomes "⊥" by outcome "-1"
 - How to overcome the detection loophole?
 - For quantum transmission losses:
 - Use heralded quantum memories (with quantum repeaters)
 - For detector losses:
 - Perform tomography tests only on the quantum detectors
 - Use trusted quantum detectors or trusted calibration devices
 - For a detector efficiency η_D :
 - ► The CHSH polynomial's correlations satisfy now:

•
$$Q = \eta_D (1 - \eta_D)$$
 • $S = 2\sqrt{2}\eta_D^2 + 2(1 - \eta_D)^2$

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- Reformulating both the Holevo bound $\chi(B_1:E)$ and key rate r_{DW} , taking into account a detector efficiency η_D :
 - ▶ For DI-QKD protocol:

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- Reformulating both the Holevo bound $\chi(B_1:E)$ and key rate r_{DW} , taking into account a detector efficiency η_D :
 - For (usual) Entanglement-based QKD protocol:

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How detector efficiencies η_D impact the (usual) Entanglement-based QKD and DI-QKD protocols?

For Holevo bounds:

- Greater Holevo bounds for the DI-QKD protocol - We can detect the presence of an eavesdropper easier for the DI-QKD protocol, but requiring higher detector efficiencies - For a detector efficiency $\eta_D=83\%$, the eavesdropper has all the information about the raw key in the DI-QKD protocol

Figure 5: Holevo bounds for a detector efficiency η_D

· For Devetak-Winter key rates:

- Lower Devetak-Winter key rates for the DI-QKD protocol
 The detector inefficiency will have greater impact on the DI-OKD protocol, reducing more the key rate
- For a detector efficiency $\eta_D=92.4\%$, no extractable secure raw key will be possible in the DI-OKD protocol

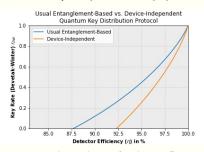


Figure 6: Devetak-Winter key rates for a detector efficiency η_{D}

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Conclusion

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Some possible directions and open questions

Possible directions:

- Consider other quantum cryptographic protocols:
 - Based on different Bell inequalities
 - Even under the assumption of collective attacks
- Consider situations in which the eavesdropper may:
 - Have partial information about measurement settings

Open questions:

- How is the security of the DI-QKD protocol modified for two-way Information Reconciliation techniques?
 - Is a Bell inequality violation sufficient for security?
- Is de Finetti theorem extendable to the DI scenario?
 - Does the security against collective attacks implies security against the most general type of attacks?

Conclusion 37/38

Thanks for your attention!