

Homework 2

Problem 5

(a) Derive the triangle inequality $\|\varphi + \psi\| \leq \|\varphi\| + \|\psi\|$. (Here $\|\Psi\| = \sqrt{\langle \Psi | \Psi \rangle}$ is the norm of vector $|\Psi\rangle$.)

Suggestion: write $\|\varphi + \psi\|^2 = \langle \varphi + \psi | \varphi + \psi \rangle$ and use the relation $\text{Re} \langle \varphi | \psi \rangle \leq |\langle \varphi | \psi \rangle|$ and the Schwartz inequality.

(b) When the triangle inequality becomes an equality?

(c) Show that $\|\varphi - \psi\| \geq \left| \|\varphi\| - \|\psi\| \right|$.

Solution

(a) We have

$$\|\varphi + \psi\|^2 = \langle \varphi + \psi | \varphi + \psi \rangle = \langle \varphi | \varphi \rangle + \langle \psi | \psi \rangle + \langle \varphi | \psi \rangle + \langle \psi | \varphi \rangle = \|\varphi\|^2 + \|\psi\|^2 + 2\text{Re} \langle \varphi | \psi \rangle.$$

The last term here satisfies

$$\text{Re} \langle \varphi | \psi \rangle \leq |\langle \varphi | \psi \rangle| \leq \sqrt{\langle \varphi | \varphi \rangle \langle \psi | \psi \rangle} = \|\varphi\| \|\psi\|,$$

where we used the Schwartz inequality. Therefore,

$$\|\varphi + \psi\|^2 \leq \|\varphi\|^2 + \|\psi\|^2 + 2\|\varphi\| \|\psi\| = (\|\varphi\| + \|\psi\|)^2.$$

Taking a square root, we end up with the triangle inequality $\|\varphi + \psi\| \leq \|\varphi\| + \|\psi\|$. ■

(b) It is obvious from the derivation in part (a) that $\|\varphi + \psi\| = \|\varphi\| + \|\psi\|$ if and only if

$$\text{Re} \langle \varphi | \psi \rangle = |\langle \varphi | \psi \rangle| = \|\varphi\| \|\psi\|,$$

which happens when $|\varphi\rangle = c|\psi\rangle$ with a **real and non-negative** coefficient c .

(c) Proceeding as in part (a), we write

$$\|\varphi - \psi\|^2 = \langle \varphi - \psi | \varphi - \psi \rangle = \|\varphi\|^2 + \|\psi\|^2 - 2\text{Re} \langle \varphi | \psi \rangle \geq \|\varphi\|^2 + \|\psi\|^2 - 2\|\varphi\| \|\psi\| = (\|\varphi\| - \|\psi\|)^2,$$

which gives

$$\|\varphi - \psi\| = \sqrt{\|\varphi - \psi\|^2} \geq \sqrt{(\|\varphi\| - \|\psi\|)^2} = \left| \|\varphi\| - \|\psi\| \right|. \quad \blacksquare$$

Problem 6

What is the number of independent *real* parameters $N(\mathcal{N})$ needed to specify (up to a phase factor) a state vector in \mathcal{N} -dimensional Hilbert space?

Solution

Let $\{|\phi_n\rangle\}$ be an arbitrary orthonormal basis. Any state vector $|\psi\rangle$ can be expanded in this basis as $|\psi\rangle = \sum_n \psi_n |\phi_n\rangle$. Substituting here $\psi_n = |\psi_n| e^{i\theta_n}$ and taking into account that $\|\psi\| = \sum_n |\psi_n|^2 = 1$, we obtain

$$|\psi\rangle = \sum_{n=1}^{n=\mathcal{N}} |\psi_n| e^{i\theta_n} |\phi_n\rangle = e^{i\theta_1} \left\{ \left(1 - \sum_{n=2}^{n=\mathcal{N}} |\psi_n|^2 \right)^{1/2} |\phi_1\rangle + \sum_{n=2}^{n=\mathcal{N}} |\psi_n| e^{i(\theta_n - \theta_1)} |\phi_n\rangle \right\}.$$

If we ignore the overall phase factor $e^{i\theta_1}$ here, the remaining parameters are $\mathcal{N}-1$ numbers $|\psi_2|, |\psi_3|, \dots, |\psi_{\mathcal{N}}|$ and $\mathcal{N}-1$ relative phases $\theta_2 - \theta_1, \theta_3 - \theta_1, \dots, \theta_{\mathcal{N}} - \theta_1$. That is, we need

$$N(\mathcal{N}) = 2(\mathcal{N} - 1)$$

real parameters altogether. For example, for $\mathcal{N} = 2$ (as it is the case for spin 1/2), we need only 2 parameters, but for $\mathcal{N} = 3$ (e.g., spin 1) the corresponding figure is $N(3) = 4$.

Problem 7

For any pure state ψ of spin $1/2$ there exists a unique Bloch vector \mathbf{n} such that $\text{Prob}_\psi(S_{\mathbf{n}} = \hbar/2) = 1$. Find the angles θ and ϕ specifying the Bloch vector in the spherical coordinates for the state represented by the state vector

$$|\psi\rangle = \frac{1}{\sqrt{2}}|+\mathbf{x}\rangle + \frac{e^{2\pi i/3}}{\sqrt{2}}|-\mathbf{x}\rangle.$$

Suggestion: write $|\psi\rangle$ in $|\pm\mathbf{z}\rangle$ basis and compare with $|\psi\rangle = \cos(\theta/2)|+\mathbf{z}\rangle + e^{i\phi}\sin(\theta/2)|-\mathbf{z}\rangle$ (up to a phase factor).

Solution

Using $\langle+\mathbf{z}|\pm\mathbf{x}\rangle = 1/\sqrt{2}$ and $\langle-\mathbf{z}|\pm\mathbf{x}\rangle = \pm 1/\sqrt{2}$, we find

$$\langle\pm\mathbf{z}|\psi\rangle = \frac{1}{2}(1 \pm e^{2\pi i/3}) = e^{i\pi/3} \times \begin{cases} \cos(\pi/3), \\ e^{3\pi i/2} \sin(\pi/3). \end{cases}$$

Substitution into $|\psi\rangle = |+\mathbf{z}\rangle\langle+\mathbf{z}|\psi\rangle + |-\mathbf{z}\rangle\langle-\mathbf{z}|\psi\rangle$ yields the state vector in $|\pm\mathbf{z}\rangle$ basis,

$$|\psi\rangle = e^{i\pi/3} \left[\cos(\pi/3)|+\mathbf{z}\rangle + e^{3\pi i/2} \sin(\pi/3)|-\mathbf{z}\rangle \right].$$

Comparison with the "canonical" form quoted above shows that the two angles specifying the Bloch vector are

$$\theta = 2\pi/3, \quad \phi = 3\pi/2.$$

Problem 8

Verify that the inner product of two spin $1/2$ state vectors

$$|\mathbf{n}_1\rangle = \cos(\theta_1/2)|+\mathbf{z}\rangle + e^{i\phi_1}\sin(\theta_1/2)|-\mathbf{z}\rangle, \quad |\mathbf{n}_2\rangle = \cos(\theta_2/2)|+\mathbf{z}\rangle + e^{i\phi_2}\sin(\theta_2/2)|-\mathbf{z}\rangle$$

with arbitrary angles $\theta_{1,2}$ and $\phi_{1,2}$ satisfies

$$|\langle\mathbf{n}_1|\mathbf{n}_2\rangle|^2 = \frac{1}{2}(1 + \mathbf{n}_1 \cdot \mathbf{n}_2).$$

For reference: Cartesian components of a unit vector \mathbf{n} specified by angles θ and ϕ is spherical polar coordinates read

$$n_{\mathbf{x}} = \sin\theta \cos\phi, \quad n_{\mathbf{y}} = \sin\theta \sin\phi, \quad n_{\mathbf{z}} = \cos\theta.$$

You will also need the trigonometric identities

$$\cos(\alpha + \beta) = \cos\alpha \cos\beta - \sin\alpha \sin\beta, \quad 1 + \cos(2\alpha) = 2\cos^2\alpha, \quad 1 - \cos(2\alpha) = 2\sin^2\alpha, \quad \sin(2\alpha) = 2\sin\alpha \cos\alpha.$$

Solution

The inner product of state vectors $|\mathbf{n}_1\rangle$ and $|\mathbf{n}_2\rangle$ is

$$\langle\mathbf{n}_1|\mathbf{n}_2\rangle = \cos(\theta_1/2)\cos(\theta_2/2) + \sin(\theta_1/2)\sin(\theta_2/2)e^{-i\phi}, \quad \phi = \phi_1 - \phi_2.$$

This gives

$$\begin{aligned} |\langle\mathbf{n}_1|\mathbf{n}_2\rangle|^2 &= \langle\mathbf{n}_1|\mathbf{n}_2\rangle\langle\mathbf{n}_1|\mathbf{n}_2\rangle^* \\ &= [\cos(\theta_1/2)\cos(\theta_2/2) + \sin(\theta_1/2)\sin(\theta_2/2)e^{-i\phi}] [\cos(\theta_1/2)\cos(\theta_2/2) + \sin(\theta_1/2)\sin(\theta_2/2)e^{+i\phi}] \\ &= \cos^2(\theta_1/2)\cos^2(\theta_2/2) + \sin^2(\theta_1/2)\sin^2(\theta_2/2) + 2\sin(\theta_1/2)\cos(\theta_1/2)\sin(\theta_2/2)\cos(\theta_2/2)\cos\phi \\ &= \frac{1}{4}(1 + \cos\theta_1)(1 + \cos\theta_2) + \frac{1}{4}(1 - \cos\theta_1)(1 - \cos\theta_2) + \frac{1}{2}\sin\theta_1\sin\theta_2\cos\phi \\ &= \frac{1}{2}(1 + \cos\theta_1\cos\theta_2 + \sin\theta_1\sin\theta_2\cos\phi). \end{aligned}$$

Comparing this expression with

$$\begin{aligned}
 \mathbf{n}_1 \cdot \mathbf{n}_2 &= n_{1\mathbf{x}}n_{2\mathbf{x}} + n_{1\mathbf{y}}n_{2\mathbf{y}} + n_{1\mathbf{z}}n_{2\mathbf{z}} \\
 &= \sin \theta_1 \sin \theta_2 (\cos \phi_1 \cos \phi_2 + \sin \phi_1 \sin \phi_2) + \cos \theta_1 \cos \theta_2 \\
 &= \sin \theta_1 \sin \theta_2 \cos \phi + \cos \theta_1 \cos \theta_2,
 \end{aligned}$$

we arrive at

$$|\langle \mathbf{n}_1 | \mathbf{n}_2 \rangle|^2 = \frac{1}{2}(1 + \mathbf{n}_1 \cdot \mathbf{n}_2). \quad \blacksquare$$