

Homework 5

Problem 19

Matrix elements of spin 1/2 operator $\hat{S}_{\mathbf{n}}$ in $|\pm \mathbf{x}\rangle$ basis form 2×2 matrix

$$\hat{S}_{\mathbf{n}} = \begin{pmatrix} \langle +\mathbf{x} | \hat{S}_{\mathbf{n}} | +\mathbf{x} \rangle & \langle +\mathbf{x} | \hat{S}_{\mathbf{n}} | -\mathbf{x} \rangle \\ \langle -\mathbf{x} | \hat{S}_{\mathbf{n}} | +\mathbf{x} \rangle & \langle -\mathbf{x} | \hat{S}_{\mathbf{n}} | -\mathbf{x} \rangle \end{pmatrix}.$$

Find this matrix and verify that it yields the correct commutator $[\hat{S}_{\mathbf{y}}, \hat{S}_{\mathbf{z}}] = i\hbar \hat{S}_{\mathbf{x}}$.

Solution

The diagonal matrix elements of $\hat{S}_{\mathbf{n}}$ are the expectation values

$$\langle \pm \mathbf{x} | \hat{S}_{\mathbf{n}} | \pm \mathbf{x} \rangle = \langle S_{\mathbf{n}} \rangle_{\pm \mathbf{x}} = \pm (\hbar/2) n_{\mathbf{x}}.$$

Taking into account that in $|\pm \mathbf{z}\rangle$ basis

$$|\pm \mathbf{x}\rangle = \frac{1}{\sqrt{2}} \begin{pmatrix} 1 \\ \pm 1 \end{pmatrix}, \quad \hat{S}_{\mathbf{n}} = \frac{\hbar}{2} \begin{pmatrix} n_{\mathbf{z}} & n_{-} \\ n_{+} & -n_{\mathbf{z}} \end{pmatrix}, \quad n_{\pm} = n_{\mathbf{x}} \pm i n_{\mathbf{y}},$$

we get the off-diagonal matrix elements:

$$\langle \pm \mathbf{x} | \hat{S}_{\mathbf{n}} | \mp \mathbf{x} \rangle = \frac{\hbar}{4} (1, \pm 1) \begin{pmatrix} n_{\mathbf{z}} & n_{-} \\ n_{+} & -n_{\mathbf{z}} \end{pmatrix} \begin{pmatrix} 1 \\ \mp 1 \end{pmatrix} = \frac{\hbar}{4} [2n_{\mathbf{z}} \pm (n_{+} - n_{-})] = \frac{\hbar}{2} (n_{\mathbf{z}} \pm i n_{\mathbf{y}}).$$

The matrix we seek is then given by

$$\hat{S}_{\mathbf{n}} = \frac{\hbar}{2} \begin{pmatrix} n_{\mathbf{x}} & n_{\mathbf{z}} + i n_{\mathbf{y}} \\ n_{\mathbf{z}} - i n_{\mathbf{y}} & -n_{\mathbf{x}} \end{pmatrix}, \quad \text{in } |\pm \mathbf{x}\rangle \text{ basis.}$$

Note that $\text{tr} \hat{S}_{\mathbf{n}} = 0$ and $\det \hat{S}_{\mathbf{n}} = -\hbar^2/4$, as it should be. As an ultimate test, we evaluate the commutator $[\hat{S}_{\mathbf{y}}, \hat{S}_{\mathbf{z}}]$:

$$[\hat{S}_{\mathbf{y}}, \hat{S}_{\mathbf{z}}] = \frac{\hbar^2}{4} \left[\begin{pmatrix} 0 & i \\ -i & 0 \end{pmatrix}, \begin{pmatrix} 0 & 1 \\ 1 & 0 \end{pmatrix} \right] = i \frac{\hbar^2}{2} \begin{pmatrix} 1 & 0 \\ 0 & -1 \end{pmatrix} = i\hbar \hat{S}_{\mathbf{x}}.$$

Problem 20

Because spin 1/2 operator $\hat{S}_{\mathbf{n}}$ has only two eigenvalues, any function of this operator can be written as the first-degree polynomial $\hat{f}(\hat{S}_{\mathbf{n}}) = c_0 \hat{\mathbb{1}} + c_1 \hat{S}_{\mathbf{n}}$. Find this polynomial for $\hat{f}(\hat{S}_{\mathbf{n}}) = e^{-i\theta \hat{S}_{\mathbf{n}}/\hbar}$, where θ is a real dimensionless number.

Solution

The projectors $\hat{\mathcal{P}}_{\pm} = |\pm \mathbf{n}\rangle \langle \pm \mathbf{n}|$ onto the two orthogonal eigenspaces of $\hat{S}_{\mathbf{n}}$ corresponding to eigenvalues $\pm \hbar/2$ satisfy

$$\hat{\mathcal{P}}_{+} + \hat{\mathcal{P}}_{-} = \hat{\mathbb{1}}, \quad \hat{\mathcal{P}}_{+} - \hat{\mathcal{P}}_{-} = 2\hat{S}_{\mathbf{n}}/\hbar.$$

The spectral decomposition of the operator $e^{-i\theta \hat{S}_{\mathbf{n}}/\hbar}$ then reads

$$e^{-i\theta \hat{S}_{\mathbf{n}}/\hbar} = e^{-i\theta/2} \hat{\mathcal{P}}_{+} + e^{i\theta/2} \hat{\mathcal{P}}_{-} = \cos(\theta/2) (\hat{\mathcal{P}}_{+} + \hat{\mathcal{P}}_{-}) - i \sin(\theta/2) (\hat{\mathcal{P}}_{+} - \hat{\mathcal{P}}_{-}) = \cos(\theta/2) \hat{\mathbb{1}} - i \sin(\theta/2) (2\hat{S}_{\mathbf{n}}/\hbar).$$

The right-hand side of this relation is the polynomial we seek.

Alternatively, one can get this result by expanding $e^{-i\theta \hat{S}_{\mathbf{n}}/\hbar}$ in Taylor series and using the relation $\hat{S}_{\mathbf{n}}^2 = (\hbar/2)^2 \hat{\mathbb{1}}$, which implies that $\hat{S}_{\mathbf{n}}^{2k} = (\hbar/2)^{2k} \hat{\mathbb{1}}$ for $k = 0, 1, 2, \dots$ and thus $\hat{S}_{\mathbf{n}}^{2k+1} = \hat{S}_{\mathbf{n}}^{2k} \hat{S}_{\mathbf{n}} = (\hbar/2)^{2k} \hat{S}_{\mathbf{n}} = (\hbar/2)^{2k+1} (2\hat{S}_{\mathbf{n}}/\hbar)$, so that

$$e^{-i\theta \hat{S}_{\mathbf{n}}/\hbar} = \sum_{k=0}^{\infty} \left[\frac{(-i\theta/2)^{2k}}{(2k)!} \hat{\mathbb{1}} + \frac{(-i\theta/2)^{2k+1}}{(2k+1)!} (2\hat{S}_{\mathbf{n}}/\hbar) \right] = \cos(\theta/2) \hat{\mathbb{1}} - i \sin(\theta/2) (2\hat{S}_{\mathbf{n}}/\hbar).$$

Problem 21

Show that for spin 1/2 operators $(\mathbf{a} \cdot \hat{\mathbf{S}})(\mathbf{b} \cdot \hat{\mathbf{S}}) = \alpha \hat{\mathbf{1}} + \boldsymbol{\beta} \cdot \hat{\mathbf{S}}$, where α and $\boldsymbol{\beta}$ are, respectively, a scalar and a vector that you need to find.

Solution

Because $(\mathbf{a} \cdot \hat{\mathbf{S}})(\mathbf{b} \cdot \hat{\mathbf{S}})$ is a geometric scalar bilinear in \mathbf{a} and \mathbf{b} and linear in $\hat{\mathbf{S}}$, it must have the form

$$(\mathbf{a} \cdot \hat{\mathbf{S}})(\mathbf{b} \cdot \hat{\mathbf{S}}) = \alpha \hat{\mathbf{1}} + \boldsymbol{\beta} \cdot \hat{\mathbf{S}} \quad \text{with } \alpha \propto \hbar^2(\mathbf{a} \cdot \mathbf{b}), \quad \boldsymbol{\beta} \propto \hbar(\mathbf{a} \times \mathbf{b}),$$

where \hbar 's take care of the units. The missing numerical coefficients can be found by comparing this expression with, e.g., $\hat{S}_z^2 = (\hbar/2)^2 \hat{\mathbf{1}}$ and $\hat{S}_x \hat{S}_y = i(\hbar/2) \hat{S}_z$. This gives

$$\alpha = \frac{\hbar^2}{4}(\mathbf{a} \cdot \mathbf{b}), \quad \boldsymbol{\beta} = \frac{i\hbar}{2}(\mathbf{a} \times \mathbf{b}).$$

Alternatively, this result can be obtained by a brute force calculation. Indeed, taking into account that

$$\hat{S}_x \hat{S}_y = -\hat{S}_y \hat{S}_x = i(\hbar/2) \hat{S}_z, \quad \hat{S}_y \hat{S}_z = -\hat{S}_z \hat{S}_y = i(\hbar/2) \hat{S}_x, \quad \hat{S}_z \hat{S}_x = -\hat{S}_x \hat{S}_z = i(\hbar/2) \hat{S}_y, \quad \hat{S}_x^2 = \hat{S}_y^2 = \hat{S}_z^2 = (\hbar/2)^2 \hat{\mathbf{1}},$$

we obtain

$$\begin{aligned} (\mathbf{a} \cdot \hat{\mathbf{S}})(\mathbf{b} \cdot \hat{\mathbf{S}}) &= (a_x \hat{S}_x + a_y \hat{S}_y + a_z \hat{S}_z)(b_x \hat{S}_x + b_y \hat{S}_y + b_z \hat{S}_z) \\ &= (\hbar/2)^2 (a_x b_x + a_y b_y + a_z b_z) \hat{\mathbf{1}} + i(\hbar/2) \{ (a_y b_z - a_z b_y) \hat{S}_x + (a_z b_x - a_x b_z) \hat{S}_y + (a_x b_y - a_y b_x) \hat{S}_z \} \\ &= (\hbar/2)^2 (\mathbf{a} \cdot \mathbf{b}) \hat{\mathbf{1}} + i(\hbar/2) (\mathbf{a} \times \mathbf{b}) \cdot \hat{\mathbf{S}}. \end{aligned}$$

This relation holds only for spin 1/2 operators. Remarkably, it remains valid even if vectors \mathbf{a} and \mathbf{b} are replaced with vector operators $\hat{\mathbf{A}}$ and $\hat{\mathbf{B}}$ that commute with $\hat{\mathbf{S}}$ but not necessarily with each other.

Problem 22

(a) Evaluate ΔS_{n_1} , ΔS_{n_2} , and $\langle [\hat{S}_{n_1}, \hat{S}_{n_2}] \rangle \equiv \langle \mathbf{n} | [\hat{S}_{n_1}, \hat{S}_{n_2}] | \mathbf{n} \rangle$ for spin 1/2 in the state $|\mathbf{n}\rangle$.

Feel free to use the relations

$$\langle \mathbf{a} \cdot \mathbf{S} \rangle_{\mathbf{n}} = \langle \mathbf{n} | \mathbf{a} \cdot \hat{\mathbf{S}} | \mathbf{n} \rangle = \frac{\hbar}{2} (\mathbf{a} \cdot \mathbf{n}), \quad \Delta_{\mathbf{n}}(\mathbf{a} \cdot \mathbf{S}) = \frac{\hbar}{2} |\mathbf{a} \times \mathbf{n}|, \quad [\mathbf{a} \cdot \hat{\mathbf{S}}, \mathbf{b} \cdot \hat{\mathbf{S}}] = i\hbar (\mathbf{a} \times \mathbf{b}) \cdot \hat{\mathbf{S}}.$$

(b) Verify that the expressions obtained in part (a) obey the uncertainty relation

$$\Delta S_{n_1} \Delta S_{n_2} \geq \frac{1}{2} \left| \langle [\hat{S}_{n_1}, \hat{S}_{n_2}] \rangle \right|.$$

Suggestion: As you will discover, part (b) reduces to demonstrating a certain relation between unit dimensionless vectors \mathbf{n} , \mathbf{n}_1 , and \mathbf{n}_2 . This can be done, for example, by switching to the spherical polar coordinates with $\mathbf{z} = \mathbf{n}$.

Solution

(a) We have

$$\begin{aligned} \Delta S_{n_1} &= \Delta(\mathbf{n}_1 \cdot \mathbf{S}) = \frac{\hbar}{2} |\mathbf{n}_1 \times \mathbf{n}|, \quad \Delta S_{n_2} = \Delta(\mathbf{n}_2 \cdot \mathbf{S}) = \frac{\hbar}{2} |\mathbf{n}_2 \times \mathbf{n}|, \\ \langle [\hat{S}_{n_1}, \hat{S}_{n_2}] \rangle &= i\hbar \langle \mathbf{n} | (\mathbf{n}_1 \times \mathbf{n}_2) \cdot \hat{\mathbf{S}} | \mathbf{n} \rangle = \frac{i\hbar^2}{2} \mathbf{n} \cdot (\mathbf{n}_1 \times \mathbf{n}_2). \end{aligned}$$

(b) Substituting ΔS_{n_1} , ΔS_{n_2} , and $\langle [\hat{S}_{n_1}, \hat{S}_{n_2}] \rangle$ found in part (a) into the uncertainty relation, we obtain

$$|\mathbf{n} \times \mathbf{n}_1| |\mathbf{n} \times \mathbf{n}_2| \geq |\mathbf{n} \cdot (\mathbf{n}_1 \times \mathbf{n}_2)|.$$

We wish to show that this inequality holds for all unit dimensionless vectors \mathbf{n} , \mathbf{n}_1 , and \mathbf{n}_2 , as it must. Following the suggestion, we switch to the spherical polar coordinates with $\mathbf{z} = \mathbf{n}$. In these coordinates, vectors $\mathbf{n}_{1,2}$ are specified by the pairs of angles $(\theta_{1,2}, \phi_{1,2})$. Taking into account that $0 \leq \theta_{1,2} \leq \pi$ and thus $\sin \theta_{1,2} \geq 0$, we find

$$|\mathbf{n} \times \mathbf{n}_{1,2}| = \sin \theta_{1,2},$$

$$|\mathbf{n} \cdot (\mathbf{n}_1 \times \mathbf{n}_2)| = |n_{1x}n_{2y} - n_{1y}n_{2x}| = \sin \theta_1 \sin \theta_2 |\cos \phi_1 \sin \phi_2 - \sin \phi_1 \cos \phi_2| = \sin \theta_1 \sin \theta_2 |\sin(\phi_1 - \phi_2)|.$$

The above inequality then assumes the form

$$\sin \theta_1 \sin \theta_2 \geq \sin \theta_1 \sin \theta_2 |\sin(\phi_1 - \phi_2)| \quad - \text{ obviously true! } \blacksquare$$

Alternatively, the above inequality can be derived by purely geometric means. We have

$$(\mathbf{n} \times \mathbf{n}_1) \times (\mathbf{n} \times \mathbf{n}_2) = \underbrace{((\mathbf{n} \times \mathbf{n}_1) \cdot \mathbf{n}_2)}_{\mathbf{n} \cdot (\mathbf{n}_1 \times \mathbf{n}_2)} \mathbf{n} - \underbrace{((\mathbf{n} \times \mathbf{n}_1) \cdot \mathbf{n})}_{0} \mathbf{n}_2 = (\mathbf{n} \cdot (\mathbf{n}_1 \times \mathbf{n}_2)) \mathbf{n},$$

where we used vector identities $\boldsymbol{\alpha} \times (\boldsymbol{\beta} \times \boldsymbol{\gamma}) = (\boldsymbol{\alpha} \cdot \boldsymbol{\gamma}) \boldsymbol{\beta} - (\boldsymbol{\alpha} \cdot \boldsymbol{\beta}) \boldsymbol{\gamma}$ and $\boldsymbol{\alpha} \cdot (\boldsymbol{\beta} \times \boldsymbol{\gamma}) = (\boldsymbol{\alpha} \times \boldsymbol{\beta}) \cdot \boldsymbol{\gamma}$. This yields the relation

$$|\mathbf{n} \cdot (\mathbf{n}_1 \times \mathbf{n}_2)| = |(\mathbf{n} \times \mathbf{n}_1) \times (\mathbf{n} \times \mathbf{n}_2)|.$$

The inequality $|\mathbf{n} \times \mathbf{n}_1| |\mathbf{n} \times \mathbf{n}_2| \geq |\mathbf{n} \cdot (\mathbf{n}_1 \times \mathbf{n}_2)|$ (see above) can now be written as

$$|\mathbf{N}_1| |\mathbf{N}_2| \geq |\mathbf{N}_1 \times \mathbf{N}_2|, \quad \mathbf{N}_{1,2} = \mathbf{n} \times \mathbf{n}_{1,2},$$

which is obviously valid for all $\mathbf{N}_{1,2}$.