

Quantum Mechanics II - Solutions

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1 Time Independent Perturbation Theory

1. The Hamiltonian of the system is

$$H = \frac{p^2}{2m} + \varepsilon_0 q x$$

with eigenvalues

$$E_n = \frac{\hbar^2 \pi^2 n^2}{2ma^2}$$

and eigenvectors

$$\psi_n = \begin{cases} \sqrt{\frac{2}{a}} \sin\left(\frac{n\pi x}{a}\right) & x \in [0, a] \\ 0 & x \notin [0, a] \end{cases}.$$

- (a) The perturbation on the first energy level, is first order, is

$$\Delta E = \langle \phi_1 | W | \phi_1 \rangle = \frac{2\varepsilon_0 q}{a} \int_0^a x \sin^2\left(\frac{\pi x}{a}\right) dx = \frac{\varepsilon_0 q}{a} \int_0^a x \left(1 - \cos\left(\frac{2\pi x}{a}\right)\right) dx = \frac{\varepsilon_0 q a}{2}$$

- (b) The perturbation on the wavefunction gives

$$\Delta \phi_1 = \sum_{n \neq 1}^{\infty} \frac{\langle \phi_n | W | \phi_1 \rangle}{E_1 - E_n} = \sum_{n=2}^{\infty} \frac{4\varepsilon_0 q m a}{\hbar^2 \pi^2 (1 - n^2)} \int_0^a x \sin\left(\frac{\pi x}{a}\right) \sin\left(\frac{n\pi x}{a}\right) dx$$

as every even term is null, it simplifies to

$$= \sum_{n=1}^{\infty} \frac{4\varepsilon_0 q m a}{\hbar^2 \pi^2 n(n+1)} \int_0^a x \sin\left(\frac{\pi x}{a}\right) \sin\left(\frac{(2n+1)\pi x}{a}\right) dx$$

2. The Hamiltonian of the unperturbed system is

$$H_0 = \hbar\omega \left(N + \frac{1}{2} \mathbb{I} \right)$$

with eigenvalues

$$E_n = \hbar\omega \left(n + \frac{1}{2} \right).$$

Let us introduce a perturbation $V = \varepsilon qx$. Recalling ladder operators defined as

$$\begin{cases} a = \sqrt{\frac{m\omega}{2\hbar}} \left(x + \frac{i}{m\omega} p \right) \\ a^\dagger = \sqrt{\frac{m\omega}{2\hbar}} \left(x - \frac{i}{m\omega} p \right) \end{cases} \Leftrightarrow \begin{cases} x = \sqrt{\frac{\hbar}{2m\omega}} (a^\dagger + a) \\ p = i\sqrt{\frac{\hbar m\omega}{2}} (a^\dagger - a) \end{cases}$$

we may write

$$V = \varepsilon q \sqrt{\frac{\hbar}{2m\omega}} (a^\dagger + a).$$

(a) The energy perturbation on first order is

$$\Delta E = \langle n | V | n \rangle = \varepsilon q \sqrt{\frac{\hbar}{2m\omega}} \langle n | a^\dagger + a | n \rangle = 0$$

and on second order

$$\begin{aligned} \Delta E_n &= \sum_{m \neq n} \frac{|\langle m | V | n \rangle|^2}{E_n - E_m} = \frac{q^2 \varepsilon^2 \hbar}{2m\omega} \sum_{m \neq n} \frac{|\langle m | a^\dagger + a | n \rangle|^2}{\hbar\omega (n - m)} \\ &= \frac{q^2 \varepsilon^2 \hbar}{2m\omega^2} \sum_{m \neq n} \frac{n \langle m | n - 1 \rangle^2 + (n + 1) \langle m | n + 1 \rangle^2}{(n - m)} = -\frac{q^2 \varepsilon^2 \hbar}{2m\omega^2} \end{aligned}$$

(b) The problem admits an analytical solution by noticing

$$H = H_0 + V = \frac{p^2}{2m} + \frac{1}{2} m\omega^2 x^2 + \varepsilon qx = \frac{p^2}{2m} + \frac{1}{2} m\omega^2 \left(x + \frac{\varepsilon q}{m\omega^2} \right)^2 - \frac{\varepsilon^2 q^2}{2m\omega^2}$$

by setting $X = x + \frac{\varepsilon q}{m\omega^2}$ we get

$$\frac{p^2}{2m} + \frac{1}{2} m\omega^2 X^2 - \frac{\varepsilon^2 q^2}{2m\omega^2} = \hbar\omega (N + 1) - \frac{\varepsilon^2 q^2}{2m\omega^2}$$

with eigenvalues

$$E_n' = E_n - \frac{\varepsilon^2 q^2}{2m\omega^2}$$

so the perturbed solution energy on second order is exact.

(c) The first order correction to the wave function is

$$\begin{aligned} |\psi_n\rangle &= |n\rangle + \sum_{m \neq n} \frac{\langle m | V | n \rangle}{E_n - E_m} |m\rangle = \frac{q^2 \varepsilon^2 \hbar}{2m\omega} \sum_{m \neq n} \frac{\langle m | a^\dagger + a | n \rangle}{\hbar\omega (n - m)} |m\rangle = \\ &= |n\rangle + \frac{q\varepsilon}{\omega} \sqrt{\frac{\hbar}{2m\omega}} (\sqrt{n} |n-1\rangle - \sqrt{n+1} |n+1\rangle). \end{aligned}$$

3. The Hamiltonian of the system is

$$H = \frac{p^2}{2m} + \frac{1}{2}m\omega^2 x^2 + \alpha x^4 = \hbar\omega \left(N + \frac{1}{2} \mathbb{I} \right) + \alpha \frac{\hbar}{2m\omega} (a^{\dagger 2} + a^2 + 2N + \mathbb{I})$$

(a) The perturbation to the first energy level is given by

$$\Delta E_n = \frac{\hbar}{2m\omega} \langle n | a^{\dagger 2} + a^2 + 2N + \mathbb{I} | n \rangle = \frac{\hbar}{2m\omega} (2\hbar n + 1).$$

(b) The perturbed wave function is

$$\begin{aligned} |\psi_n\rangle &= |n\rangle + \frac{\hbar}{2m\omega} \sum_{m \neq n} \frac{\langle m | a^{\dagger 2} + a^2 + 2N + \mathbb{I} | n \rangle}{E_n - E_m} |m\rangle \\ &= |n\rangle + \frac{\hbar^2}{4m\omega^2} \left[\sqrt{(n+1)(n+2)} |n+2\rangle - \sqrt{n(n-1)} |n-2\rangle \right]. \end{aligned}$$