

The quantum monad: towards quantum finite model theory



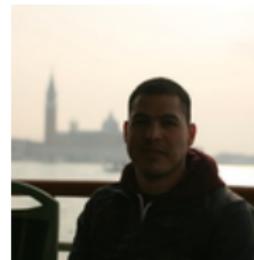
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- ▶ ‘The quantum monad on relational structures’
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- ▶ \rightsquigarrow quantum finite model theory?

1. Introduction

Motivation

With the advent of quantum computation and information:

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- ▶ delineate the scope of **quantum advantage**

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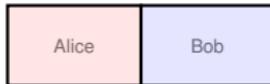
- ▶ use **quantum resources** for information-processing tasks
- ▶ delineate the scope of **quantum advantage**

- ▶ A setting in which this has been explored is **non-local games**

Non-local games

Alice and Bob cooperate in solving a task set by Verifier

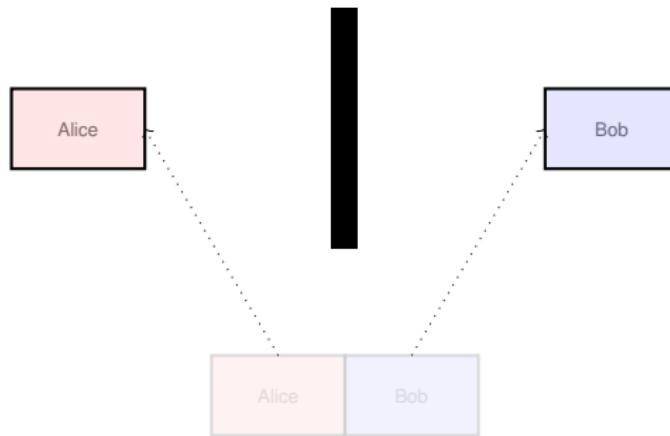
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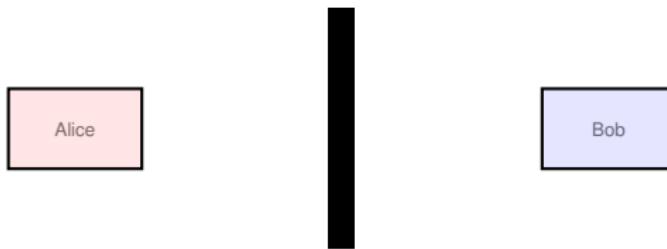
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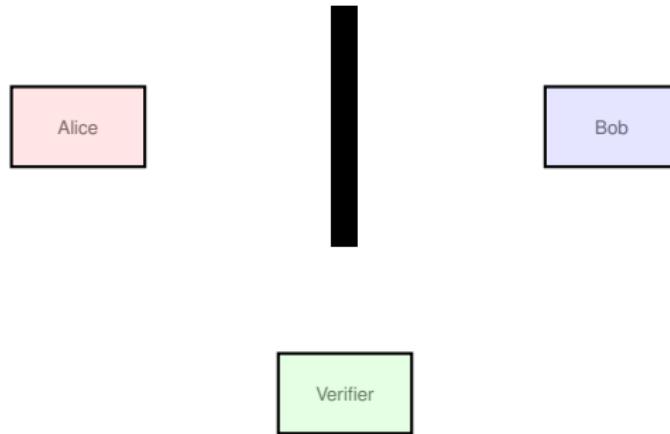
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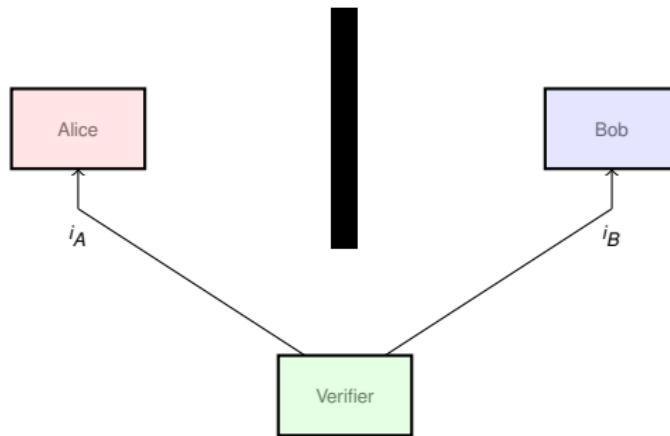
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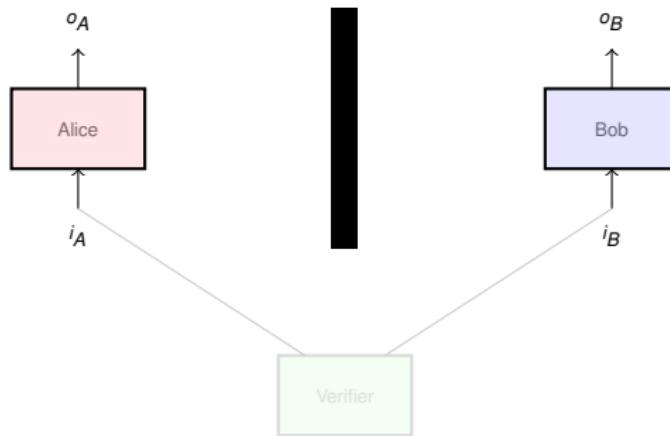
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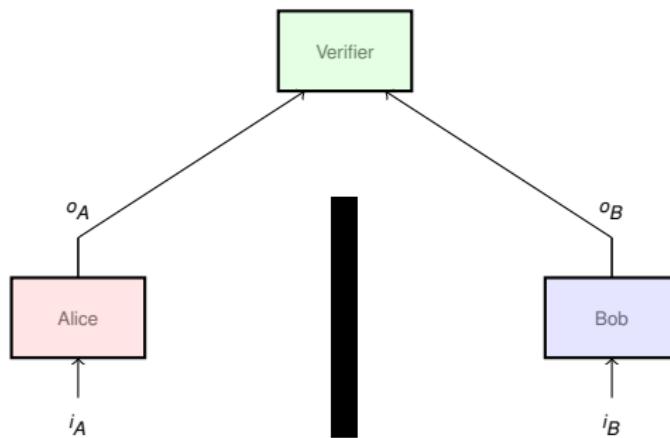
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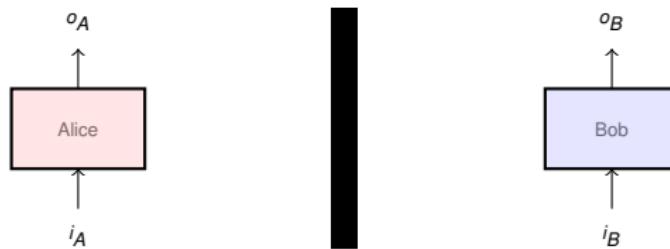
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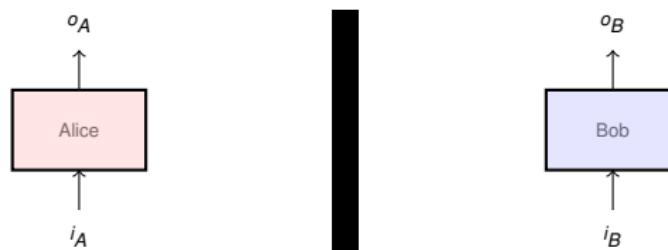
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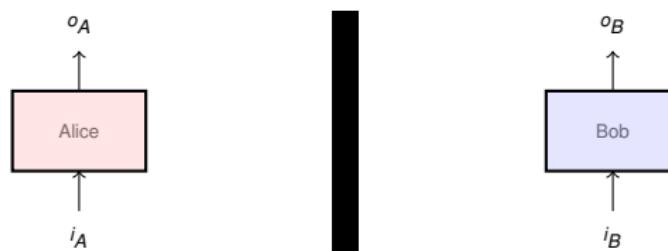


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A **perfect strategy** is one that wins with probability 1.

E.g.: Binary constraint systems

Magic square:

- ▶ Fill with 0s and 1s
- ▶ rows and first two columns: even parity
- ▶ last column: odd parity

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System of linear equations over \mathbb{Z}_2 :

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Clearly, this is not satisfiable in \mathbb{Z}_2 .

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The system has a **quantum solution** but no classical solution!

Examples of non-local games

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Many of these works have some aspects in common. We aim to flesh this out by subsuming them under a common framework.

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Many relevant questions in these areas can be phrased in terms of (existence, number of, ...) homomorphisms between finite relational structures.

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Motivation

What could it mean to quantise these fundamental structures?

- ▶ We formulate the task of constructing a homomorphism between two relational structures as a **non-local game**
- ▶ **uniformly** obtain quantum analogues *for free* for a whole range of classical notions from CS, logic, ...
- ▶ We then show that the use of quantum resources for information tasks is captured in a high-level way as **quantum homomorphisms**
- ▶ which can be integrated into a typed functional programming language through a **monadic** interface.

Outline

- ▶ Introduce homomorphism game for relational structures
- ▶ Arrive at the notion of quantum homomorphism,
which removes the two-player aspect of the game
(generalises Cleve & Mittal and Mančinska & Roberson)
- ▶ Quantum monad: capture quantum homomorphisms as classical
homomorphisms to a *quantised* version of a relational structure
(inspired on Mančinska & Roberson for graphs)
- ▶ Non-locality and state-independent strong contextuality
- ▶ Towards quantum finite model theory and descriptive complexity

2. Homomorphism game for relational structures

Relational structures and homomorphisms

A relational vocabulary σ consists of relational symbols R_1, \dots, R_p where R_l has an arity $k_l \in \mathbb{N}$ for each $l \in [p] := \{1, \dots, p\}$.

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A σ -**structure** is $\mathcal{A} = (A; R_1^{\mathcal{A}}, \dots, R_p^{\mathcal{A}})$ where:

- ▶ A is a non-empty set,
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where $f(\mathbf{x}) = \langle f(x_1), \dots, f(x_{k_l}) \rangle$ for $\mathbf{x} = \langle x_1, \dots, x_{k_l} \rangle$.

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(For simplicity, from now on consider a single relational symbol R of arity k)

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What about quantum resources?

Homomorphism game with quantum resources

Quantum resources:

- ▶ Finite-dimensional Hilbert spaces \mathcal{H} (Alice's) and \mathcal{K} (Bob's)
- ▶ A bipartite pure state ψ on $\mathcal{H} \otimes \mathcal{K}$

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These resources are used as follows:

- ▶ Given input $\mathbf{x} \in R^A$, Alice measures $\mathcal{E}_{\mathbf{x}}$ on her part of ψ
- ▶ Given input $x \in A$, Bob measures \mathcal{F}_x on his part of ψ
- ▶ Both output their respective measurement outcomes
- ▶ $P(\mathbf{y}, y \mid \mathbf{x}, x) = \psi^*(\mathcal{E}_{\mathbf{x}, \mathbf{y}} \otimes \mathcal{F}_{x, y})\psi$

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Perfect strategy conditions:

(QS1)

$$\psi^*(\mathcal{E}_{\mathbf{x}, \mathbf{y}} \otimes I_{\mathcal{K}})\psi = 0$$

if $\mathbf{y} \notin R^B$

(QS2)

$$\psi^*(\mathcal{E}_{\mathbf{x}, \mathbf{y}} \otimes \mathcal{F}_{x, y})\psi = 0$$

if $x = \mathbf{x}_i$ and $y \neq \mathbf{y}_i$

3. From quantum perfect strategies to quantum homomorphisms

Simplifying quantum strategies

Theorem¹ The existence of a quantum perfect strategy implies the existence of a strategy $(\psi, \{\mathcal{E}_x\}, \{\mathcal{F}_x\})$ with the following properties:

¹This generalises Cleve & Mittal and Mančinska & Roberson.

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$$\boxed{\begin{array}{ll} \psi^*(\mathcal{E}_{\mathbf{x},\mathbf{y}} \otimes I) \psi = 0 & \text{if } \mathbf{y} \notin R^B \\ \psi^*(\mathcal{E}_{\mathbf{x},\mathbf{y}} \otimes \mathcal{F}_{x,y}) \psi = 0 & \text{if } x = \mathbf{x}_i \text{ and } y \neq \mathbf{y}_i \end{array}}$$

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The reduction proceeds in three steps:

1. The state and strategies are projected down to the support of the Schmidt decomposition of the state. This reduces the dimension of the Hilbert space and preserves the probabilities of the strategy exactly.
2. It is shown that this strategy must already satisfy strong properties (PVMs and $\mathcal{E}_{\mathbf{x},y}^i = \mathcal{F}_{\mathbf{x}_i,y}^T$).
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N.B. In passing to the special form, the dimension is **reduced**; the process by which we obtain projective measurements is not at all akin to dilation.

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These $P_{x,y}$ are enough to determine the strategy!

Quantum homomorphisms

A quantum homomorphism from \mathcal{A} to \mathcal{B} is a family of projectors $\{P_{x,y}\}_{x \in A, y \in B}$ in some dimension $d \in \mathbb{N}$ satisfying:

- (QH1) For all $x \in A$, $\sum_{y \in B} P_{x,y} = I$.
- (QH2) For all $\mathbf{x} \in R^A$, $x = \mathbf{x}_i$, $x' = \mathbf{x}_j$,

$$[P_{x,y}, P_{x',y'}] = \mathbf{0} \quad \text{for any } y, y' \in B$$

so we can define a projective measurement $P_{\mathbf{x}} = \{P_{\mathbf{x},y}\}_y$,
where $P_{\mathbf{x},y} := P_{\mathbf{x}_1,y_1} \cdots P_{\mathbf{x}_k,y_k}$.

- (QH3) If $\mathbf{x} \in R^A$ and $\mathbf{y} \notin R^B$, then $P_{\mathbf{x},\mathbf{y}} = \mathbf{0}$.

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- (QH3) If $\mathbf{x} \in R^{\mathcal{A}}$ and $\mathbf{y} \notin R^{\mathcal{B}}$, then $P_{\mathbf{x},\mathbf{y}} = \mathbf{0}$.

Theorem For finite structures \mathcal{A} and \mathcal{B} , the following are equivalent:

1. The $(\mathcal{A}, \mathcal{B})$ -homomorphism game has a quantum perfect strategy.
2. There is a quantum homomorphism from \mathcal{A} to \mathcal{B} . ($\mathcal{A} \xrightarrow{q} \mathcal{B}$)

4. Quantum homomorphisms and the quantum monad

Quantum homomorphisms as Kleisli maps

For each $d \in \mathbb{N}$ and σ -structure \mathcal{A} , we can define a structure $\mathcal{Q}_d\mathcal{A}$ such that there is a one-to-one correspondence:²

$$\mathcal{A} \xrightarrow{q} \mathcal{B} \cong \mathcal{A} \longrightarrow \mathcal{Q}_d\mathcal{B}$$

- ▶ quantum homomorphisms from \mathcal{A} to \mathcal{B} of dimension d
- ▶ (classical) homomorphisms from \mathcal{A} to $\mathcal{Q}_d\mathcal{B}$

²Mančinska & Roberson: analogous construction for (their) graph homomorphisms.

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Universe of structure $\mathcal{Q}_d\mathcal{A}$: set of functions $p : A \longrightarrow \text{Proj}(d)$ such that $\sum_{x \in A} p(x) = I$. (Projector-valued distributions on A in dimension d .)

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For R of arity k , $R^{\mathcal{Q}_d\mathcal{A}}$ is the set of tuples $\langle p_1, \dots, p_k \rangle$ satisfying:

- (QR1) For all $1 \leq i, j \leq k$ and $x, x' \in A$, $[p_i(x), p_j(x')] = \mathbf{0}$.
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Quantum homomorphisms as Kleisli maps

\mathcal{Q}_d is a functor and moreover part of a **graded monad** on the category $R(\sigma)$ of relational structures and (classical) homomorphisms.

Monads play a major rôle in programming language theory, providing a uniform way of encapsulating various notions of computation:

- ▶ partiality
- ▶ exceptions
- ▶ non-determinism
- ▶ probabilistic
- ▶ state updates
- ▶ input/output
- ▶ ...

Monads

Functor $T : \mathcal{C} \rightarrow \mathcal{C}$ such that a T -program, a computation producing values of type B from values of type A with T -effects, is seen as a map $A \rightarrow TB$ in the category \mathcal{C} .

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$$\begin{array}{ccc} B & \xrightarrow{g} & TC \\ A & \xrightarrow{f} & TB \end{array}$$

Kleisli morphisms form a category.

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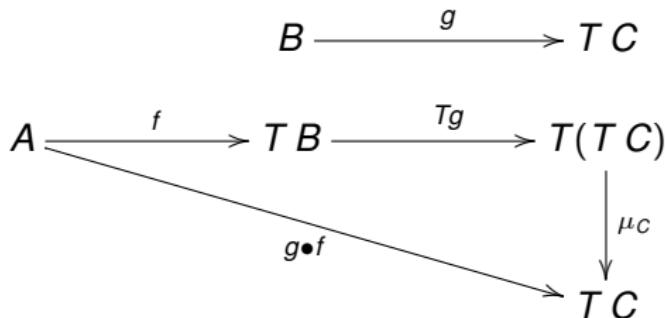
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- ~~~ composition of quantum homomorphisms
keeps track of the dimension of the resources

$$\mathcal{B} \xrightarrow{f} \mathcal{Q}_{d'} \mathcal{C}$$

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$$\begin{array}{ccc} \mathcal{B} & \xrightarrow{f} & \mathcal{Q}_{d'} \mathcal{C} \\ \mathcal{A} & \xrightarrow{f} & \mathcal{Q}_d \mathcal{B} \xrightarrow{\mathcal{Q}_{dg}} \mathcal{Q}_d(\mathcal{Q}_{d'} \mathcal{C}) \end{array}$$

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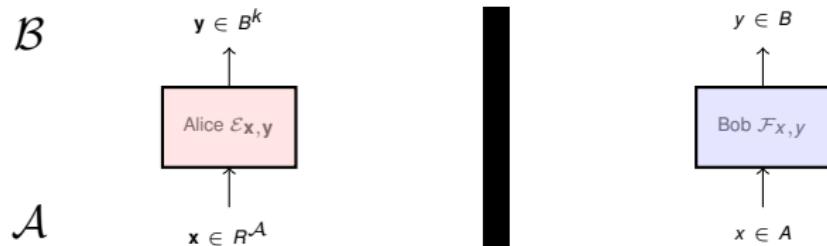
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Composition of perfect strategies



Composition of perfect strategies

 \mathcal{C}

$z \in C^k$

Alice $\mathcal{E}_{y,z}$

$y \in R^{\mathcal{B}}$

 \mathcal{B} 

$z \in C$

Bob $\mathcal{F}_{y,z}$

$y \in B$

o

 \mathcal{B}

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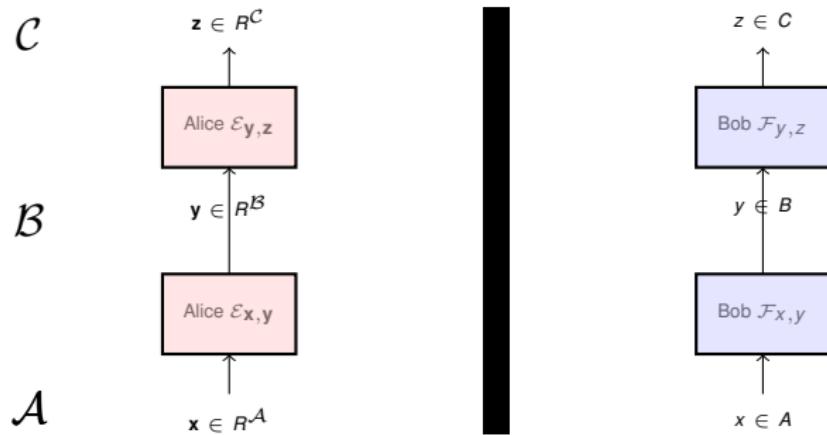
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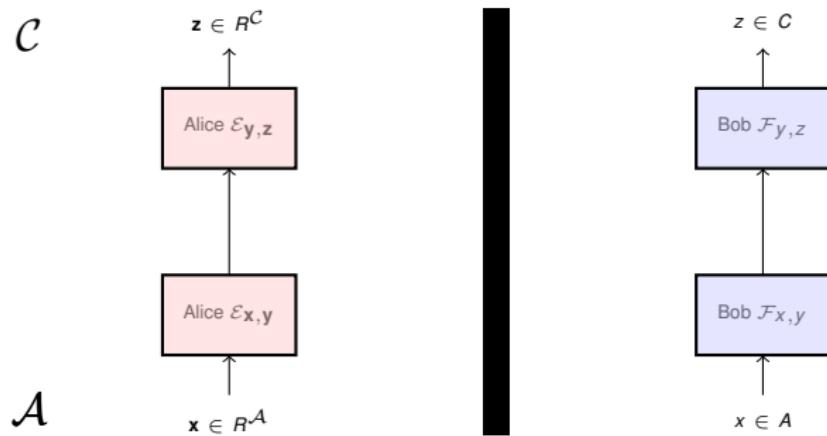
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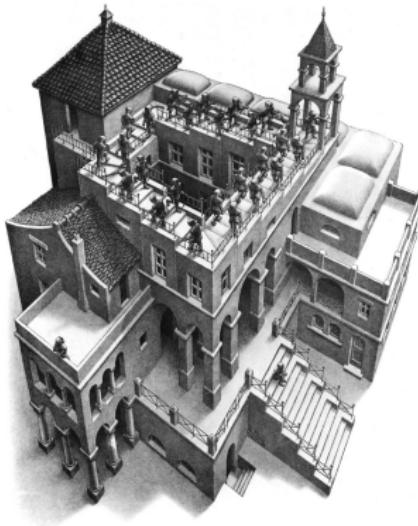


5. Contextuality and non-locality

Contextuality

Contextuality is a fundamental feature of quantum mechanics, which distinguishes it from classical physical theories.

It can be thought as saying that empirical predictions are inconsistent with all measurements having pre-determined outcomes.



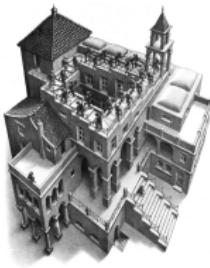
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Non-locality is a particular case of contextuality for Bell scenarios

...but here we show that certain contextuality proofs can always be underwritten by non-locality arguments.



Contextuality

Measurement scenario (X, \mathcal{M}, O) :

- ▶ X is a finite set of measurements
- ▶ O is a finite set of outcomes
- ▶ \mathcal{M} is a cover of X , where $C \in \mathcal{M}$ is a set of compatible measurements (context)

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Possibilistic information: for $C \in \mathcal{M}$ and $s \in O^C$, $e_C(s) \in \{0, 1\}$ indicates if joint outcome s for measurements C is possible or not.

Contextuality

Measurement scenario (X, \mathcal{M}, O) :

- ▶ X is a finite set of measurements
- ▶ O is a finite set of outcomes
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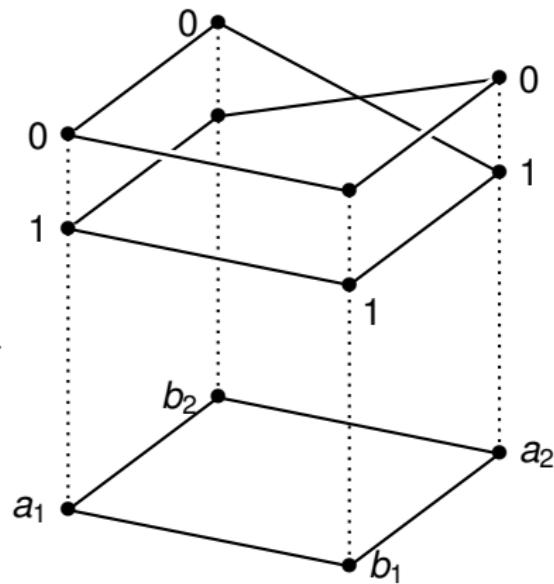
Strong contextuality: if there is no global assignment $g : X \rightarrow O$ such that for all $C \in \mathcal{M}$, $e_C(g|_C) = 1$. That is, no global assignment is consistent with the model in the sense of yielding **possible** outcomes in all contexts.

E.g.: GHZ, Kochen–Specker, (post-quantum) PR box

Strong contextuality

Strong Contextuality:
no consistent global
assignment.

A	B	(0, 0)	(0, 1)	(1, 0)	(1, 1)
a_1	b_1	✓	✗	✗	✓
a_1	b_2	✓	✗	✗	✓
a_2	b_1	✓	✗	✗	✓
a_2	b_2	✗	✓	✓	✗



Strong contextuality and constraint satisfaction

The support of e can be described as a CSP \mathcal{K}_e

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- ▶ consistent global assignments for e

Hence, e is strongly contextual iff \mathcal{K}_e has no (classical) solution.

Quantum correspondence

Quantum witness for e :

- ▶ state φ
- ▶ PVM $P_x = \{P_{x,o}\}_{o \in O}$ for each $x \in X$
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General way of turning state-independent contextuality proofs into Bell non-locality arguments (generalising Heywood & Redhead's construction).

6. Outlook

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- ▶ Quantising classical notions in the framework of relational structures

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Model theory:

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- ▶ \rightsquigarrow equivalences coarser than isomorphism:

$$\mathcal{A} \equiv^{\mathcal{L}} \mathcal{B} := \forall \phi \in \mathcal{L}. \mathcal{A} \models \phi \Leftrightarrow \mathcal{B} \models \phi$$

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- ▶ computational complexity \leftrightarrow expressive power of logics

E.g. PH \leftrightarrow SO (second-order logic)

NP \leftrightarrow \exists SO (existential second-order logic)

AC⁰ \leftrightarrow FO(+, \times) (first-order logic with + and \times)

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- ▶ Quantum descriptive complexity?
Do quantised versions of these logical equivalences correspond
to quantum computational complexity classes?

Thank you!

Questions...



The quantum monad on relational structures (MFCS'17, AQIS'17)
arXiv:1705.07310