

# Quantum Walks, Algorithms and Implementations

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## Random Walks

- M. Dyer, A. Frieze, and R. Kannan (1991):  
A random polynomial-time algorithm for approximating the volume of convex bodies.
- R. Motwani and P. Raghavan (1995):  
Randomized Algorithms.
- U. Schöning (1999):  
A probabilistic algorithm for k-sat and constraint satisfaction problems.
- M. Jerrum, A. Sinclair, and E. Vigoda (2004):  
A polynomial-time approximation algorithm for permanent of a matrix with nonnegative entries.

## Quantum Walks

- Y. Aharonov, L. Davidovich, and N. Zagury (1993):  
Quantum random walks.
- E. Farhi and S. Gutmann (1998):  
Quantum computation and decision trees.
- M. Szegedy (2004):  
Quantum speed-up of markov chain based algorithms
- A. Patel, K. Raghunathan, and P. Rungta (2005). Quantum random walks do not need a coin toss.
- A. Childs (2009):  
Universal computation by quantum walk.
- R. Portugal (2016):  
The staggered quantum walk model.

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The space of our quantum walk is composed by coin  $\mathcal{H}_C$  and walker spaces  $\mathcal{H}_W$ , and we have  $\mathcal{H} = \mathcal{H}_C \otimes \mathcal{H}_W$ . The evolution operator consists of tossing a coin and performing a shift, and we say

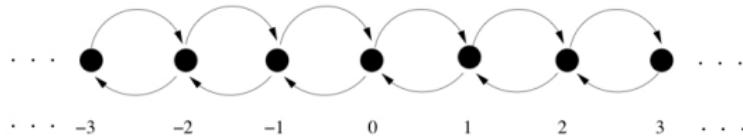
$$U = S(C \otimes I_W) \quad \rightarrow \quad |\psi(t)\rangle = U^t |\psi(0)\rangle \quad (1)$$

We can describe the shift operator as

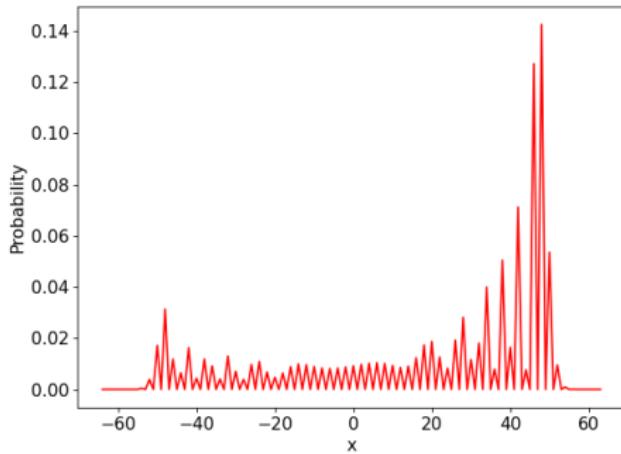
$$\begin{aligned} S |0\rangle |x\rangle &= |0\rangle |x+1\rangle \\ S |1\rangle |x\rangle &= |1\rangle |x-1\rangle \end{aligned} \quad (2)$$

and  $S$  in the computational basis has the format

$$S = |0\rangle \langle 0| \otimes \sum_x |x+1\rangle \langle x| + |1\rangle \langle 1| \otimes \sum_x |x-1\rangle \langle x| \quad (3)$$



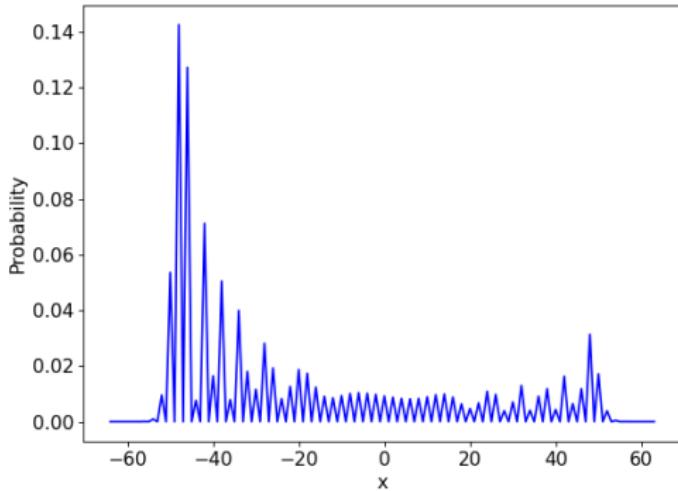
# Simulation & Remarks



We consider  $t = 70$ , initial condition, and the Hadamard coin, as

$$|\psi(0)\rangle = |0\rangle|0\rangle, C = \frac{1}{\sqrt{2}} \begin{bmatrix} 1 & 1 \\ 1 & -1 \end{bmatrix} \quad (4)$$

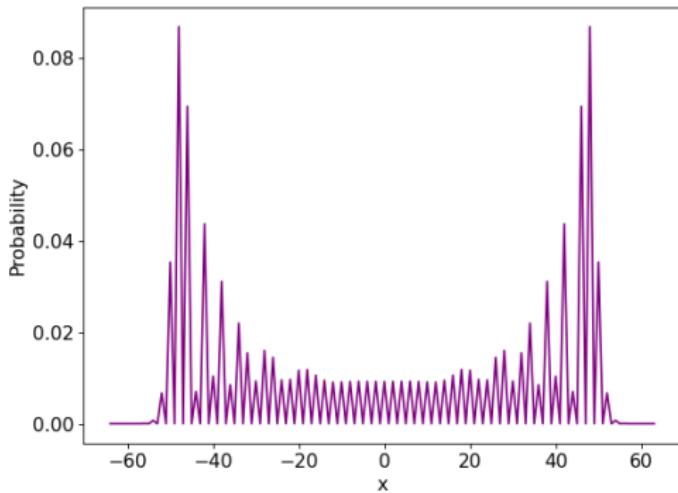
# Simulation & Remarks



We consider  $t = 70$ , initial condition, and the Hadamard coin, as

$$|\psi(0)\rangle = |1\rangle|0\rangle, C = \frac{1}{\sqrt{2}} \begin{bmatrix} 1 & 1 \\ 1 & -1 \end{bmatrix} \quad (5)$$

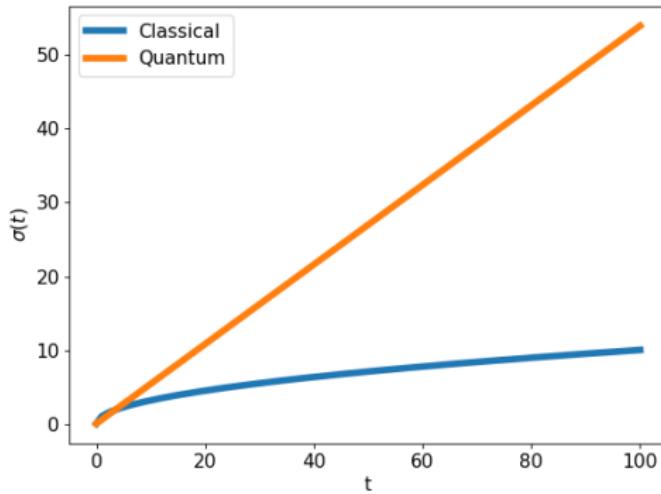
# Simulation & Remarks



We consider  $t = 70$ , initial condition, and the Hadamard coin, as

$$|\psi(0)\rangle = \frac{|0\rangle - i|1\rangle}{\sqrt{2}} |0\rangle, C = \frac{1}{\sqrt{2}} \begin{bmatrix} 1 & 1 \\ 1 & -1 \end{bmatrix} \quad (6)$$

# Simulation & Remarks



Classical:  $\sigma(t) \sim \sqrt{t}$

Quantum:  $\sigma(t) \sim t$

Given a graph  $G$  defined by its Laplacian matrix  $\mathcal{L}$ , we can define a relation between a random walk and a quantum walk by

$$\frac{\partial p(x, t)}{\partial t} = \gamma \mathcal{L} p(x, t) \quad \longrightarrow \quad i \frac{\partial |\psi(x, t)\rangle}{\partial t} = H |\psi(x, t)\rangle \quad (7)$$

Considering  $\gamma$  a jumping-rate (amplitude per time) and  $H = -\gamma \mathcal{L}$ , we have the solution

$$U = e^{i\gamma \mathcal{L} t} \quad \longrightarrow \quad |\psi(t)\rangle = U |\psi(0)\rangle \quad (8)$$

We may define the adjacency matrix for an infinite line

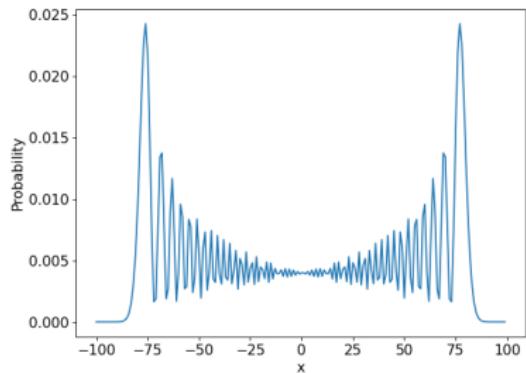
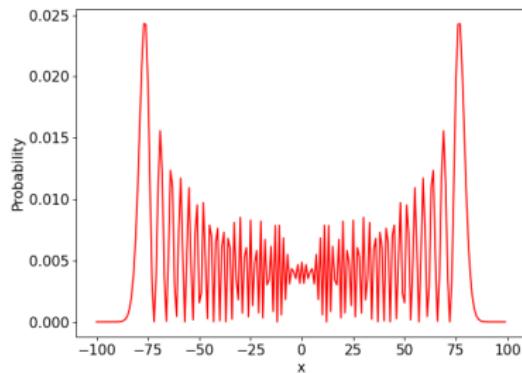
$$A = \sum_x |x+1\rangle\langle x| + |x\rangle\langle x+1| \quad (9)$$

and its Laplacian matrix as

$$\mathcal{L} = A - D = \sum_x |x+1\rangle\langle x| + |x\rangle\langle x+1| - 2I \quad (10)$$

where  $D$  denotes the degree matrix of a graph.

# Simulation & Remarks

(a)  $|\psi_a\rangle$ (b)  $|\psi_b\rangle$ 

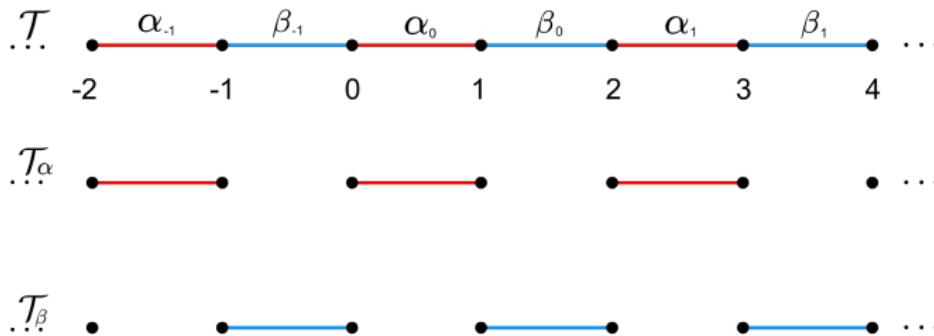
We consider here  $t = 40$ ,  $\gamma = 1$  and initial conditions

$$|\psi_a\rangle = |0\rangle \quad \text{and} \quad |\psi_b\rangle = |+\rangle \quad (11)$$

We construct this walk based on a graph **tesselation**, that is a set of disjoint cliques over all vertices; and a graph **covering**, that is a family of tesselations, where every edge belongs to, at least, one tesselation.

Considering an infinite line, we can tessellate this graph as

$$\begin{aligned}\mathcal{T}_\alpha &= \{\{2x, 2x+1\} : x \in \mathbb{Z}\}, \\ \mathcal{T}_\beta &= \{\{2x+1, 2x+2\} : x \in \mathbb{Z}\}.\end{aligned}$$



We may define states associated to each tesselation such as

$$\begin{aligned} |\alpha_x\rangle &= \frac{|2x\rangle + |2x+1\rangle}{\sqrt{2}}, \\ |\beta_x\rangle &= \frac{|2x+1\rangle + |2x+2\rangle}{\sqrt{2}}. \end{aligned}$$

and operators associated to each tesselation as

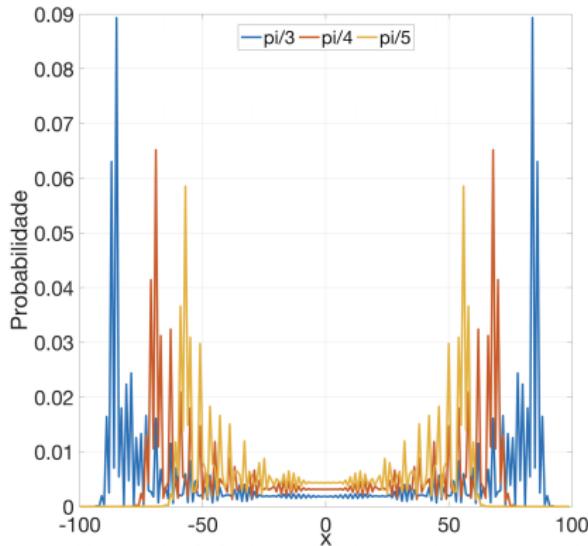
$$\begin{aligned} H_\alpha &= 2 \sum_{x=-\infty}^{\infty} |\alpha_x\rangle \langle \alpha_x| - I, \\ H_\beta &= 2 \sum_{x=-\infty}^{\infty} |\beta_x\rangle \langle \beta_x| - I. \end{aligned}$$

Finally, we can describe the evolution operator as

$$U = e^{i\theta_\beta H_\beta} e^{i\theta_\alpha H_\alpha}.$$

where  $\theta_\alpha, \theta_\beta \in [0, \pi]$

# Simulation & Remarks

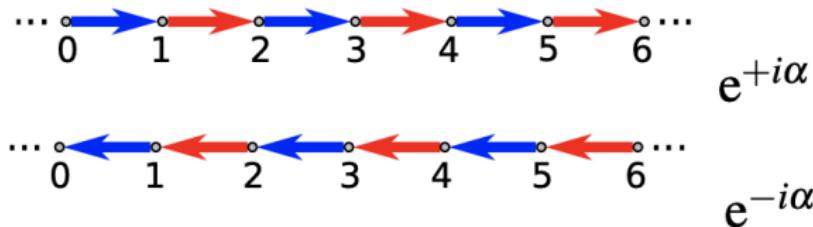


We consider  $\theta_\alpha = \theta_\beta = \theta$ , after  $t = 50$ , and initial condition

$$|\psi(0)\rangle = \frac{|0\rangle + |1\rangle}{\sqrt{2}} \quad (12)$$

# Simulation & Remarks

B. Chagas and R. Portugal (2020). Discrete-time quantum walks on oriented graphs.



We can modify the operators in order to have a sense of direction as

$$\begin{aligned} H_0 &= \sum_x e^{-i\alpha} |2x-1\rangle \langle 2x| + e^{+i\alpha} |2x\rangle \langle 2x-1|, \\ H_1 &= \sum_x e^{-i\alpha} |2x\rangle \langle 2x+1| + e^{+i\alpha} |2x+1\rangle \langle 2x|. \end{aligned}$$

and the evolution will be

$$U = e^{i\theta_\beta} H_1 e^{i\theta_\alpha} H_0.$$

# Simulation & Remarks

Defining the standard deviation as  $\sigma(t) = \sqrt{\langle x^2 \rangle - \langle x \rangle^2}$ , we have the moments

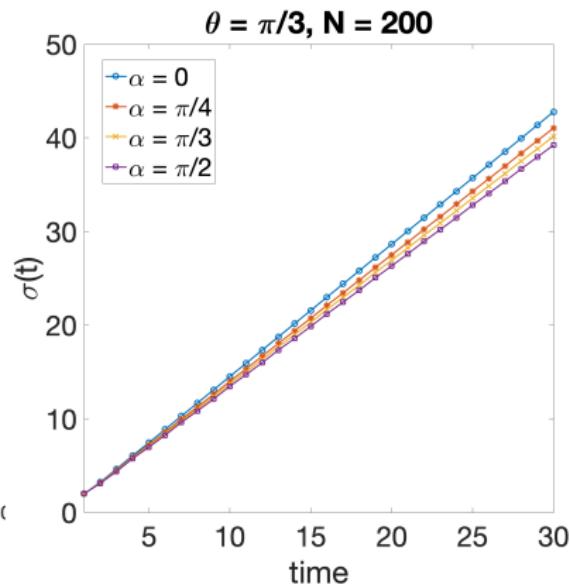
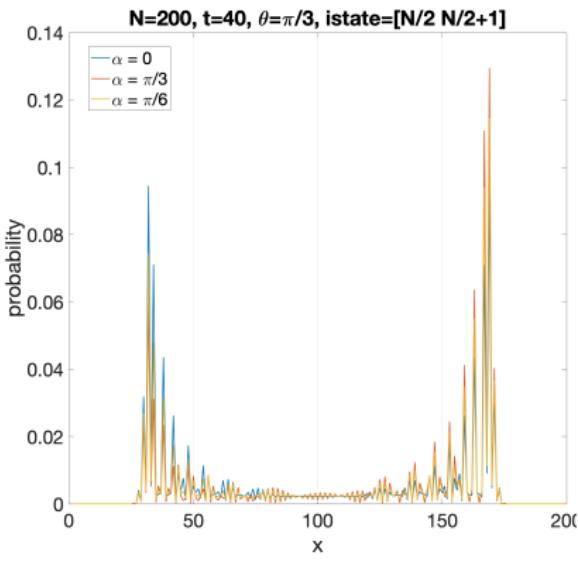
$$\begin{aligned}\frac{\langle x \rangle}{t} &= 2(|a|^2 - |b|^2)(1 - \cos \theta) + \frac{i \sin 2\theta (\bar{a}be^{i\alpha} - a\bar{b}e^{-i\alpha})}{1 + |\cos \theta|} + \mathcal{O}\left(\frac{1}{t}\right) \\ \frac{\langle x^2 \rangle}{t^2} &= 4(1 - |\cos \theta|) + \mathcal{O}\left(\frac{1}{t}\right)\end{aligned}$$

considering the initial condition

$$|\psi(0)\rangle = a|0\rangle + b|b\rangle \tag{13}$$

where  $|a|^2 + |b|^2 = 1$ .

# Simulation & Remarks



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Quantum Algorithms based on Quantum Walks

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## 4 Final Remarks

# Quantum Walk + Algorithms

- Triangle Finding
  - F. Magniez, M. Santha, and M. Szegedy (2003). Quantum algorithms for the triangle problem.
  - Problem: Given a graph  $G$  on  $n$  nodes, find a triangle, if there is any.
- Element Distinctness
  - A. Ambainis (2014). Quantum walk algorithm for element distinctness
  - Problem: determine whether the elements of a list are distinct.
- Matrix Product Verification
  - H. Buhrman, and R. Špalek (2005). Quantum verification of matrix products.
  - Problem: Let  $A, B, C$  be  $n \times n$  matrices over any integral domain. A verification of a matrix product is deciding whether  $AB = C$ .
- Group Commutativity
  - F. Magniez, and A. Nayak (2005). Quantum complexity of testing group commutativity
  - Problem: Given a black-box group  $G$  with generators  $g_1, \dots, g_k$ , decide if  $G$  is abelian

Searching Algorithms: Given a list of elements, find a marked element, if there is any.

- L. Grover (1996):  
A fast quantum mechanical algorithm for database search.
- M. Boyer, et al (1996):  
Tight bounds on quantum searching.
- C. Zalka (1999):  
Grover's quantum searching algorithm is optimal.

## Searching Algorithms Based on Quantum Walks

- A. Childs, and J. Goldstone (2004). Spatial search by quantum walk.
- R. Portugal (2018). Quantum walks and search algorithms.
- J. Janmark, D. Meyer, and T. Wong (2014). Global symmetry is unnecessary for fast quantum search.
- S. Chakraborty, L. Novo, A. Ambainis, and Y. Omar. Spatial search by quantum walk is optimal for almost all graphs.

Given a graph  $G$  defined by its Laplacian matrix  $\mathcal{L}$ , we can modify the Hamiltonian

$$H = -\gamma \mathcal{L} \quad (14)$$

by introducing the oracle Hamiltonian

$$H_w = -|w\rangle\langle w| \quad (15)$$

and now on we consider the time-independent Hamiltonian,

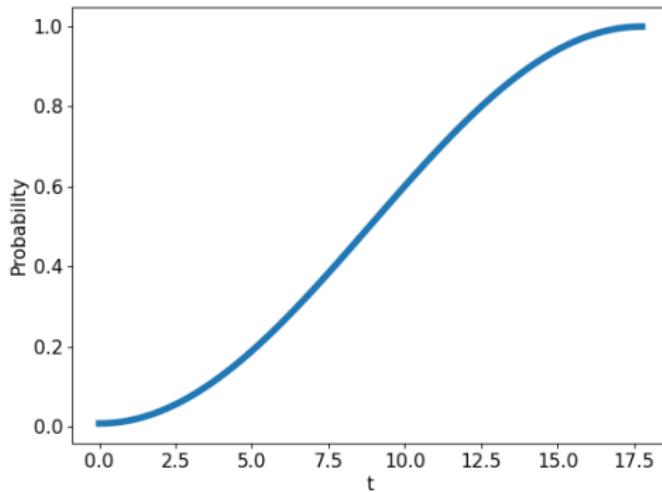
$$H = -\gamma L + H_w = -\gamma L - |w\rangle\langle w|. \quad (16)$$

Moreover, the evolution operator will be

$$|\psi(t)\rangle = e^{itH} |s\rangle \quad (17)$$

where  $|s\rangle$  denotes a uniform superposition.

# Implementation



The optimal case, for one marked element over a clique graph, occurs when

$$t = \frac{\pi}{2} \sqrt{N} \quad \text{and} \quad \gamma = \frac{1}{N} \quad (18)$$

Given the evolution operator for a general staggered quantum walk as

$$U = \prod_k e^{i\theta_k H_k} \quad (19)$$

we can add an oracle unitary operator, considering the marked element  $w$ , as

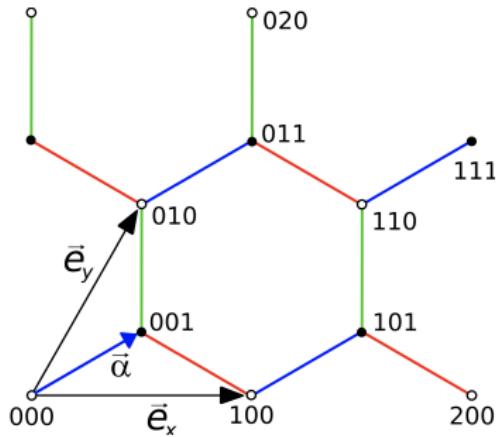
$$U_w = I - 2|w\rangle\langle w| \quad (20)$$

and the searching evolution consists of

$$|\psi(t)\rangle = (UU_w)^t |s\rangle, \quad (21)$$

where  $|s\rangle$  is the uniform superposition state.

# Hexagonal Lattice



$$|\alpha_{x,y}\rangle = \frac{1}{\sqrt{2}}(|x,y,1\rangle + |x+1,y,0\rangle)$$

$$|\beta_{x,y}\rangle = \frac{1}{\sqrt{2}}(|x,y,1\rangle + |x,y+1,0\rangle)$$

$$|\gamma_{x,y}\rangle = \frac{1}{\sqrt{2}}(|x,y,0\rangle + |x,y,1\rangle)$$

# Evolution Operator

The evolution operator for this graph will be

$$U = e^{i\theta H_\gamma} e^{i\theta H_\beta} e^{i\theta H_\alpha} \quad (22)$$

considering the operators

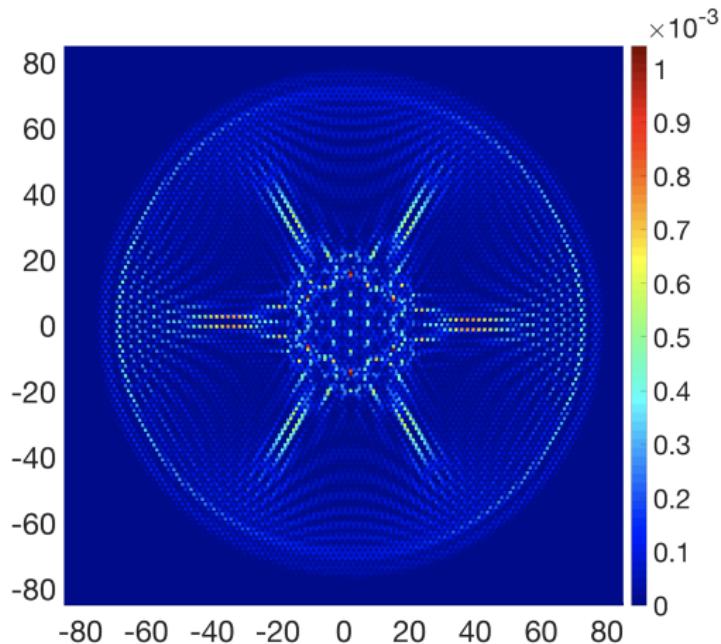
$$H_\alpha = 2 \sum_{x,y=0}^{n-1} |\alpha_{x,y}\rangle \langle \alpha_{x,y}| - I,$$

$$H_\beta = 2 \sum_{x,y=0}^{n-1} |\beta_{x,y}\rangle \langle \beta_{x,y}| - I,$$

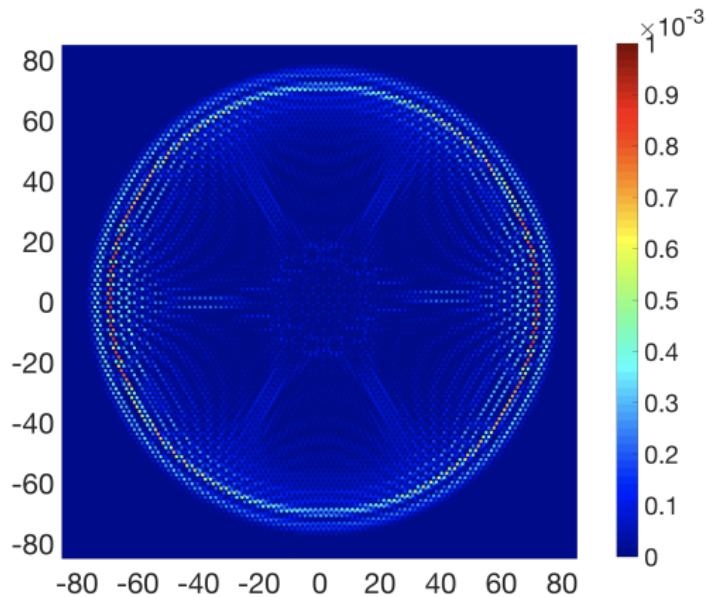
$$H_\gamma = 2 \sum_{x,y=0}^{n-1} |\gamma_{x,y}\rangle \langle \gamma_{x,y}| - I,$$

where

$$|\psi_t\rangle = U^t |\psi_0\rangle$$

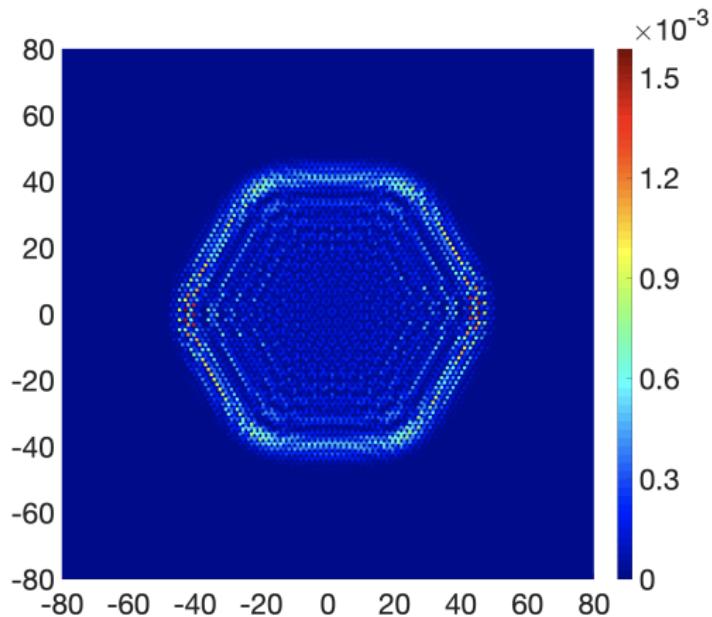


$$|\psi_0^b\rangle = \frac{1}{\sqrt{2}}(|1,1,0\rangle + |1,0,1\rangle)$$



$$|\psi_0^c\rangle = \frac{1}{\sqrt{6}}(|1,1,0\rangle + |1,0,1\rangle + |1,0,0\rangle + |0,0,1\rangle + |0,1,0\rangle + |0,1,1\rangle)$$

$$n=121, t=58, \theta = \pi/6, |\psi_0^c\rangle$$



$$|\psi_0^c\rangle = \frac{1}{\sqrt{6}}(|1,1,0\rangle + |1,0,1\rangle + |1,0,0\rangle + |0,0,1\rangle + |0,1,0\rangle + |0,1,1\rangle)$$

# Searching Problem

B. Chagas, R. Portugal, S. Boettcher, and E. Segawa (2018). Staggered Quantum Walk on Hexagonal Lattices.

We considered the evolution operator

$$|\psi(t)\rangle = (UU_w)^t |s\rangle, \quad (23)$$

where  $U$  is the staggered quantum walk operator for the hexagonal graph, and  $U_w$  the oracle operator. We've got the following time execution

$$t = \Theta(\sqrt{N\ell n N})$$

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# Some questions

What about planar graphs?

- B. Chagas, R. Portugal, S. Boettcher, and E. Segawa (2018). Staggered Quantum Walk on Hexagonal Lattices.
- R. Portugal, and T. Fernandes (2017). Quantum search on the two-dimensional lattice using the staggered model with Hamiltonians

# Some questions

What about physical realizations of continuous-time quantum walks?

- R. Balu, D. Castillo, and G. Siopsis (2018). Physical realization of topological quantum walks on IBM-Q and beyond Staggered Quantum Walk on Hexagonal Lattices.
- F. Acasiete, F. Agostini, J. Moqadam, and R. Portugal (2020). Experimental Implementation of Quantum Walks on IBM Quantum Computers

# Some questions

What's the relation between angles and tessellation in staggered quantum walks?

- A. Abreu, L. Cunha, T. Fernandes, C. de Figueiredo, L. Kowada, F. Marquezino, D. Posner, R. Portugal (2017). The tessellation problem of quantum walks
- R. Santos (2018). The role of tessellation intersection in staggered quantum walks.

# Questions?

# Obrigado!