

PROOF OF BOUND OF VARIANCE

Proof: $\text{tr}(\hat{\rho}) = 1 \rightarrow \text{Var}(\hat{O})$ only depends on traceless

part: $O_0 = O - \frac{\text{tr}(O)}{2^n} \mathbb{I}$

$$\hat{O} - E(\hat{O}) = \text{tr}(O\hat{\rho}) - \text{tr}(O\rho) = \text{tr}\left(O_0 + \frac{\text{tr}(O)}{2^n} \mathbb{I}\right)\hat{\rho} - \text{tr}\left(\rho\left(O_0 + \frac{\text{tr}(O)}{2^n} \mathbb{I}\right)\right) = \text{tr}(O_0\hat{\rho}) - \text{tr}(O_0\rho)$$

M^{-1} is self adjoint

M is invertible, if M is self adjoint, M^{-1} will be too.

$$\text{tr}(\underline{X}M(Y)) = \sum_i \langle i | \sum_b E \sum_b \underbrace{\langle b | U Y U^\dagger | b \rangle}_{1} U^\dagger | b \rangle X | b \rangle U | i \rangle$$

$$= \sum_i E \sum_b \langle i | X U^\dagger | b \rangle X | b \rangle U | i \rangle \underbrace{\langle b | U Y U^\dagger | b \rangle}_{1}$$

$$= \sum_i E \sum_b \langle b | U | i \rangle X | i \rangle X U^\dagger | b \rangle \langle b | U Y U^\dagger | b \rangle$$

$$= E \sum_b \langle b | U X U^\dagger | b \rangle \langle b | U Y U^\dagger | b \rangle$$

$$\text{tr}(M(X)Y) = \sum_i \langle i | E \sum_b \langle b | U X U^\dagger | b \rangle U^\dagger | b \rangle X U | i \rangle Y | i \rangle$$

$$= \sum_i E \sum_b \langle b | U Y | i \rangle X | i \rangle U^\dagger | b \rangle \langle b | U X U^\dagger | b \rangle$$

$$= E \sum_b \langle b | U Y U^\dagger | b \rangle \langle b | U X U^\dagger | b \rangle$$

$$M M^{-1} = I$$

$$(M^{-1})^\dagger M^\dagger = I$$

$$M^\dagger = M$$

$$(M^{-1})^\dagger = (M^\dagger)^{-1} = M^{-1}$$

$$\text{Var}(\hat{o}) = E(\hat{o} - E(\hat{o}))^2 = E(\hat{o}^2) - E(\hat{o})^2$$

$$= E(\text{tr } O_0 \hat{\rho})^2 - \text{tr}(O_0 E(\hat{\rho}))^2 = -$$

$$\begin{aligned} & \text{tr}(O_0 \tilde{M}^{-1}(U^\dagger |\hat{b} \times \hat{b}| U)) = \text{tr}(\tilde{M}^{-1}(O_0) U^\dagger |\hat{b} \times \hat{b}| U) \\ & = \left(\sum_i \langle i | \tilde{M}^{-1}(O_0) U^\dagger |\hat{b} \times \hat{b}| U | i \rangle \right)^2 = \left(\sum_i \langle \hat{b} | U \tilde{M}^{-1}(O_0) U^\dagger | \hat{b} \rangle \right)^2 \end{aligned}$$

$$= E(\langle \hat{b} | U \tilde{M}^{-1}(O_0) U^\dagger | \hat{b} \rangle^2) - [\text{tr}(O_0 \rho)]^2$$

Randomness comes from $\left\{ \begin{array}{l} U \sim \mathcal{U} \text{ classically} \\ \text{which } |b \times b| \text{ we measure (Quant.)} \end{array} \right.$

$$E(\langle \hat{b} | U \tilde{M}^{-1}(O_0) U^\dagger | \hat{b} \rangle^2) =$$

$$E \sum_{U \sim \mathcal{U} \text{ before}} \langle \hat{b} | U \rho U^\dagger | \hat{b} \rangle \langle \hat{b} | U \tilde{M}^{-1}(O_0) U^\dagger | \hat{b} \rangle^2$$

setting prob of $|b \times b|$

dependence on ρ !

We just take the max

By removing $(\text{tr}(\rho_p))^2$

$$\rightarrow \text{Var}(\hat{o}) \leq \left\| O - \frac{\text{tr} O}{2^n} \mathbb{I} \right\|_{\text{shadow}}^2$$

PROOF USING MEDIAN OF MEANS

→ Median of means.

Given $\hat{\mu}_{\text{nom}}$ from a total of $n_T = \underbrace{K}_{\text{Part}} \cdot \underbrace{N}_{\text{elements in each part}}$
and assuming $\text{Var}(X_1) < \infty$

$$P(|\hat{\mu}_{\text{nom}} - \mu_0| > \varepsilon) \leq e^{-2K \left(\frac{1}{2} - \frac{K}{n_T} \frac{\sigma^2}{\varepsilon^2} \right)^2}$$

In the paper, they look for $P \leq 2e^{-\frac{K}{2}}$ which is supposed to happen in the worst situation.

$$\leq e^{-2K \left(\frac{1}{2} - \underbrace{\frac{K}{n_T} \frac{\sigma^2}{\varepsilon^2}}_{\ll 1} \right)^2} \leq 2e^{-\frac{K}{2}}$$

n_T Worst case

$$\frac{K}{n_T} \frac{\sigma^2}{\varepsilon^2} \ll 1 \rightarrow \frac{K}{K \cdot N} \frac{\sigma^2}{\varepsilon^2} \ll 1 \quad N \sim \text{Const} \cdot \frac{\sigma^2}{\varepsilon^2}$$

In the paper, they use $N = 34 \frac{\sigma^2}{\varepsilon^2}$ as a safe number.

$$\text{Using } \Pr\left(\bigvee_{i=1}^M |\hat{\sigma}_i - \text{tr}(\sigma_i \rho)| \geq \varepsilon\right) \leq \sum_{i=1}^M \Pr(|\hat{\sigma}_i - \text{tr}(\sigma_i \rho)| \geq \varepsilon)$$

$$= M \underbrace{\Pr(|\hat{\sigma}_i - \text{tr}(\sigma_i \rho)| \geq \varepsilon)}_{\substack{\text{We want this} \\ \frac{\delta}{M}}}$$

$$\underbrace{\Pr(|\hat{\mu} - \mu_0| \geq \varepsilon)}_{\frac{\delta}{M}} \leq 2e^{-\frac{K}{2}} = \frac{\delta}{M}$$

$$\Rightarrow K = 2 \log \frac{2M}{\delta}$$

We conclude that we need $n_T = N \cdot K$

$$\text{with } N = \frac{34}{\varepsilon^2} \sigma^2 = \frac{34}{\varepsilon^2} \max_{1 \leq i \leq M} \left\| D_i - \frac{\text{tr}(D_i)}{2^n} I \right\|_{\text{shadow}}^2$$

and

$$K = 2 \log\left(\frac{2M}{\delta}\right)$$

to learn

$$|\hat{O}_i(N, K) - \text{tr}(D_i p)| \leq \varepsilon \quad \forall 1 \leq i \leq M$$

with prob at least $1 - \delta$

In general:

$$N_{\text{Tot}} = O\left(\frac{\log M}{\varepsilon^2} \max_{1 \leq i \leq M} \left\| D_i - \frac{\text{tr}(D_i)}{2^n} I \right\|_{\text{shadow}}^2\right)$$

LET'S TALK ABOUT UNITARY T-DESIGNS

Unitary t -designs

→ Classical spherical t -designs

$$p_t: S(\mathbb{R}^d) \rightarrow \mathbb{R}$$



$p_t(x)$
max degree

$$x \in S(\mathbb{R}^d)$$

We want to avg p_t over all points of the sphere.
Can we do that with a set of finite points?

YES!

$X = \{x: x \in S(\mathbb{R}^d)\}$ is a spherical t -design if

$$\frac{1}{|X|} \sum_{x \in X} p_t(x) = \int_{S(\mathbb{R}^d)} p_t(u) d\gamma(u) \quad \text{f.a. } p_t(u)$$

Higher t imply bigger set X

E.g. in \mathbb{R}^3 , a cube is a 3-design
a dodecahedron is a 5-design.

→ Complex projective designs

Extend the idea to $S(\mathbb{C}^d)$. The polynomials have degree at most t for the entries and complex conj.

X , a complex projective t -design, is a subset of X of $S(\mathbb{C}^d)$ s.t.

$$\frac{1}{|X|} \sum_{x \in X} p_{t,t}(x) = \int_{S(\mathbb{C}^d)} p_{t,t}(u) d\gamma(u) \quad \text{f.a. } p_{t,t}(u)$$

→ Unitary t -design.

Our Sphere is now the unitary group $U(d)$

Polynomials act on elements $\in U(d)$.

A unitary t -design is the set of points $\in U(d)$ that suffice to compute the average of $P_{t,t}$ (max degree of t in the entries of U or its complex conj) over all the unitary transformations.

$\{U_k\}_1^K$ is unitary t -design if

$$\frac{1}{K} \sum_{k=1}^K P_{t,t}(U_k) = \int_{U(d)} \underbrace{d\gamma(U)}_{\text{Haar measure of the group}} P_{t,t}(U)$$

→ Where does it appear?

Avg fidelity:

A noisy channel $\sim V$ (perfect)

Let's see how it affect $|0\rangle$ state: $\left\{ \begin{array}{l} V|0\rangle\langle 0|V^\dagger \\ \Lambda(|0\rangle\langle 0|) \end{array} \right.$

Fidelity for the 0 state: $\text{Tr}(V|0\rangle\langle 0|V^\dagger \Lambda(|0\rangle\langle 0|))$

$$= \langle 0|V^\dagger \Lambda(|0\rangle\langle 0|)V|0\rangle$$

Only $|0\rangle$ state ? → Avg over all the states

$$| \text{Rand} \rangle = \underbrace{V|0\rangle}_{\text{random} \sim U(2)}$$

$$\bar{F}(\Lambda, V) = \int_{\mathcal{U}} d\mu(U) \langle \underbrace{0|} V^\dagger V^\dagger \Lambda(\overbrace{V|0\rangle\langle 0|V^\dagger}) V \underbrace{|0\rangle} \rangle$$

This is a polynomial of degree 2 in the entries of U & U^\dagger \rightarrow We can use a U. 2 design.

Clifford group is a 3-design (also a 2, 1-design) (NOT 4)

Clifford group: Normalizer of the Pauli group. Sends Paulis to Pauli (up to a phase) via conjugation

$$C P C^\dagger = \pm P' \quad C \in \text{Cl}(2^n) \quad P, P' \in \mathcal{P}(2^n)$$

$\underbrace{\text{Cl}(2^n)}_{\text{Clifford g. of } n \text{ qubit}}$ $\underbrace{\mathcal{P}(2^n)}_{\text{Pauli group on } n \text{ qubits}}$

$$P = \sigma_1 \otimes \dots \otimes \sigma_n \in \mathcal{P}(1)$$

Generated by H, S, CNOT

\rightarrow Gottesman-Knill theorem. Circuits with preparation and measurements in the comp. basis + Clifford gates can be efficiently simulated on a classical comp.

The key idea is to keep track of how operators change instead of the state.

We only need to evaluate the fidelity on the states that $V \in \text{Cl}(2^n)$ generates by $V|0\rangle$!

There are 24 elements in $\text{Cl}(2)$

" 11520 in $\text{Cl}(4)$

For our interests

$$E_{U \sim C(2^n)} (U A U)^{\otimes K} = \int_{U(d)} (U A U^\dagger)^{\otimes K} d\mu(U) \quad \forall 2^n \times 2^n \text{ mat } A$$

$K=1,2,3$

There are explicit formulas for:

$$E_{U \sim C(2^n)} U^\dagger |x\rangle\langle x| U \langle x| U A U^\dagger |x\rangle = \frac{A + \text{tr}(A)I}{(2^n+1)2^n} = \frac{1}{2^n} \mathcal{D}_{\frac{1}{2^n+1}}(A)$$

for $A \in H_{2^n}$

$$E_{U \sim C(2^n)} U^\dagger |x\rangle\langle x| U \langle x| U B_0 U^\dagger |x\rangle \langle x| U C_0 U^\dagger |x\rangle = \frac{\text{tr}(B_0 C_0)I + B_0 C_0 + C_0 B_0}{(2^n+2)(2^n+1)2^n}$$

for $B_0, C_0 \in H_{2^n}$

* this is eq to a depolarizing ch. of n qubits with loss probp.

$$\mathcal{D}_p(A) = pA + (1-p) \frac{\text{tr}(A)}{2^n} I$$

THIS CAN BE INVERTED!

$$\mathcal{D}_{\frac{1}{2^n+1}}^{-1}(A) = (2^n+1)A - \text{tr}(A)I$$

SHADOWS NORM FOR GLOBAL CLIFFORD UNITARIES

Proof:

We already know that

$$\hat{\rho} = M^{-1}(U^\dagger |\hat{b} \times b\rangle U) = \underline{(2^n + 1) U^\dagger |\hat{b} \times b\rangle U - \mathbb{1}}$$

and $M^{-1}(O_0) = (2^n + 1) O_0$ $\underbrace{O_0}_{\text{traceless}} \in H_{2^n}$

→ Let's compute the shadow norm

$$\|O_0\|_{\text{sh.}}^2 = \max_{\sigma} \left(\frac{1}{E} \sum_b \langle b | U \sigma U^\dagger | b \rangle \langle b | U M^{-1}(O) U^\dagger | b \rangle^2 \right)^{1/2}$$

$\sum_i |i\rangle \langle i| = \mathbb{1}$

$$= \max_{\sigma} \left[\text{tr} \left(\sigma \sum_b E \underline{U^\dagger |b\rangle \langle b| U} \underline{\langle b | U M^{-1}(O) U^\dagger | b \rangle^2} \right) \right]^{1/2}$$

We can use the fact that CE is a unitary 3-design

$$E_{U \in \mathcal{U}(2^n)} \underline{U^\dagger |x\rangle \langle x| U} \underline{\langle x | U B_0 U^\dagger | x \rangle \langle x | U C_0 U^\dagger | x \rangle} =$$

$$= \frac{\text{tr}(B_0 C_0) \mathbb{I} + B_0 C_0 + C_0 B_0}{(2^n + 2)(2^n + 1) 2^n} \quad \text{for } B_0, C_0 \text{ traceless } \in H_{2^n}$$

If $B_0 = C_0 = O_0$

from the sum

$$= \max_{\sigma} \text{tr} \left(\sigma \frac{2^n (2^n + 1)^2 (\text{tr}(O_0^2) \mathbb{I} + 2 O_0^2)}{(2^n + 2) \cdot (2^n + 1) 2^n} \right)$$

$$= \left(\frac{2^n + 1}{2^n + 2} \right) \max_{\sigma} \left(\underbrace{\text{tr}(\sigma)}_1 \text{tr}(O_0^2) + 2 \text{tr}(\sigma O_0^2) \right)$$

Now we use $\max_{\text{state } \sigma} (\sigma O_0^2) = \|O_0^2\|_{\infty} \leq \text{tr}(O_0^2)$
state spectr. norm

$$\rightarrow \text{tr}(O_0^2) \leq \|O_0\|_{\text{shadow}}^2 \leq 3 \text{tr}(O_0^2)$$

SCALING FULL PROCESS TOMOGRAPHY

for a E acting on n -qubits, we need to learn a ρ on a $2n$ -qubits system.

A reduced k -qubit ρ_k can be expressed as

$$\rho_k = \frac{1}{2^{2k}} \sum_i \alpha_i O_i^{(2k)}$$

w. $\alpha_i = \text{tr}(\rho_k O_i^{(2k)}) = \text{tr}(\rho O_i^{(2k)} \otimes \mathbb{1}_{2n-2k})$

We estimate this w. classical shadows

$$\hat{\rho}_k = \frac{1}{4^k} \sum_i \hat{\alpha}_i O_i^{(2k)}$$

$$\langle \hat{\rho}_k - \rho_k, \hat{\rho}_k - \rho_k \rangle_F = \langle \frac{1}{4^k} \sum_i (\hat{\alpha}_i - \alpha_i) O_i^{(2k)}, \frac{1}{4^k} \sum_j (\hat{\alpha}_j - \alpha_j) O_j^{(2k)} \rangle_F$$

$$= \frac{1}{4^{2k}} \sum_{i,j} (\hat{\alpha}_i - \alpha_i)(\hat{\alpha}_j - \alpha_j) \langle O_i^{(2k)}, O_j^{(2k)} \rangle_F$$

$$= \frac{1}{4^{2k}} \cdot 4^k \sum_i (\hat{\alpha}_i - \alpha_i)^2 = \frac{1}{4^k} \sum_i (\hat{\alpha}_i - \alpha_i)^2$$

Including the normalization in the previous theorem;

$$|\hat{\alpha}_i - \alpha_i| = |\hat{\alpha}_i - \text{tr}(\rho O_i^{(2k)} \otimes \mathbb{1}_{2n-2k})| \leq \frac{\epsilon}{2^k} \quad \forall i$$

By substitution:

$$\langle \hat{\rho}_k - \rho_k, \hat{\rho}_k - \rho_k \rangle_F \leq \frac{16^k}{4^k} \frac{\epsilon^2}{2^{2k}} = \epsilon^2$$

$$\Rightarrow \|\hat{\rho}_k - \rho_k\|_F \leq \epsilon$$

For Random global Clifford:

$$\left\| O - \frac{\text{tr}(O)}{2} \mathbb{1} \right\|_{\text{shadow}}^2 \leq 3 \text{tr}(O^2) = 3 \cdot 4^k \text{tr}(O^2) = 3 \cdot 4^k \sum_i O_i^2$$

$$O_i = \begin{pmatrix} \text{---} & & & \\ & 2^{2n} & & \\ & & 1 & \\ & & & 1 & \\ & & & & 1 & \\ & & & & & 1 & \\ & & & & & & 1 & \\ & & & & & & & 1 & \\ & & & & & & & & 1 & \\ & & & & & & & & & 1 \end{pmatrix} \rightarrow O_i^2 = \begin{pmatrix} 1 & & & & & & & & & \\ & 1 & & & & & & & & \\ & & 1 & & & & & & & \\ & & & 1 & & & & & & \\ & & & & 1 & & & & & \\ & & & & & 1 & & & & \\ & & & & & & 1 & & & \\ & & & & & & & 1 & & \\ & & & & & & & & 1 & \\ & & & & & & & & & 1 \end{pmatrix}$$

2^{2n} 2^{2n}

$$\text{tr}(O_i^2) = 2^{2n} = 4^n$$

—/

$$N = \frac{68}{\epsilon^2} \log(2M/\delta) \max_i \left\| O_i - \frac{\text{tr}(O_i)}{d} I \right\|_{\text{Fro}}^2$$

→ global Clifford

$$N = \frac{68}{\epsilon^2} 3 \cdot 4^n \log(2M/\delta) \sim O(4^n) \text{ with } n = \text{no. of qubits}$$