



**MANE 6961:**

**Adjoint for Scientists  
and Engineers**

**Lecture 14**

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JEC 2036**

# Motivation

Last class I introduced the adjoint-weighted residual (AWR) method in the context of linear BVPs and linear functionals.

In this lecture, I will first generalize the AWR method to nonlinear problems. Subsequently, I will describe various approximations that are used to make the method useful in practice.

# AWR Method for Nonlinear Problems

# Problem Statement

Let  $R_H : \mathbb{R}^S \rightarrow \mathbb{R}^S$  be the residual corresponding to the discretization of a nonlinear BVP on a mesh of nominal size  $H$ . Furthermore, let  $u_H \in \mathbb{R}^S$  be the solution to

$$R_H(u_H) = 0.$$

Our goal is to estimate the functional error

$$\delta J_H = J_H(u_H) - J(u),$$

where  $J(u)$  is the exact (nonlinear) functional value based on the exact solution of the nonlinear BVP, and  $J_H(u_H)$  is the discretized functional value based on the discrete solution.

# Averaged derivatives

Proceeding as in the linear case, we have

$$\delta J_H \approx J_H(u_H) - J_h(u_h),$$

where  $u_h \in \mathbb{R}^s$  is the solution to  $R_h(u_h) = 0$ , the discretized BVP on a fine mesh with nominal element size  $h$ .

For the subsequent analysis, we define

$$\delta u_h \equiv u_h^H - u_h,$$

where, as before,  $u_h^H \equiv I_h^H u_H$  is the coarse solution represented on the fine mesh/space.

## Averaged derivatives (cont.)

Thus, we have

$$\begin{aligned} J_H(u_H) - J_h(u_h) &= J_h(u_h^H) - J_h(u_h) \\ &= J_h(u_h + \delta u_h) - J_h(u_h). \end{aligned}$$

Unlike the linear case, we cannot express this functional difference as a functional of the difference. However, notice that

## Averaged derivatives (cont.)

Consequently, the functional difference is equal to an averaged derivative times  $\delta u_h$ .

$$(g_h, \delta u_h)_h = g_h^T \delta u_h \quad \rightarrow \quad \left( \int_0^1 J'_h[u_h + s\delta u_h] ds \right)^T \delta u_h.$$

This suggests that we define the adjoint in a slightly different way from the linear case. . .

**Definition: Mean-value Adjoint [FD11]**

The (discrete) mean-value adjoint is denoted by  $\bar{\psi}_h$  and satisfies the linear equation

$$(\bar{R}_h[u_h, u_h^H])^T \bar{\psi}_h = (\bar{J}_h[u_h, u_h^H])^T,$$

where the averaged Jacobian and gradient are, respectively,

$$\begin{aligned}\bar{R}_h[u_h, u_h^H] &\equiv \int_0^1 R'_h[u_h + s\delta u_h] ds \\ \bar{J}_h[u_h, u_h^H] &\equiv \int_0^1 J'_h[u_h + s\delta u_h] ds.\end{aligned}$$



# Mean-value Adjoint

Notice that we have the following identities:

$$\overline{R}_h[u_h, u_h^H] \delta u_h =$$

Similarly

$$\overline{J}_h[u_h, u_h^H] \delta u_h =$$

## Mean-value Adjoint (cont.)

Using the above identities, as well as the definition of the mean-value adjoint, we find that

$$J_h(u_h + \delta u_h) - J_h(u_h) =$$

# Mean-value Adjoint (cont.)

## Theorem: AWR for nonlinear problems

Let  $u_H$  denote the solution to  $R_H(u_H) = 0$ , which corresponds to a discretized BVP on a coarse space. Analogously, let  $u_h$  denote the solution of  $R_h(u_h) = 0$ , the fine-space discretization of the same BVP. Then the difference

$$J_H(u_H) - J_h(u_h) = \bar{\psi}_h^T R_h(u_h^H),$$

where  $\bar{\psi}_h$  is the solution to the mean-value adjoint equation on the fine-space.

# AWR Approximations

# Problems with the AWR Method

There is a significant computational cost when applying the AWR method to nonlinear problems: **the mean-value adjoint**.

- The mean-value adjoint requires the integral of potentially expensive derivatives.

The AWR method, whether it is used on linear or nonlinear problems, has another drawback.

- The adjoint,  $\psi_h$  or  $\overline{\psi}_h$ , is needed on the fine mesh/space.

Let's consider the approximations used to mitigate these costs.

# Avoiding the Mean-value Adjoint

We can avoid the mean-value adjoint by replacing the integrated Jacobian and gradient with point derivatives as follows:

$$J_h(u_h + \delta u_h) - J_h(u_h) =$$

# Avoiding the Mean-value Adjoint (cont.)

Based on the above approximation, we can use the adjoint corresponding to the linearization about  $u_h$ :

$$(R'_h[u_h])^T \psi_h = J'_h[u_h].$$

Thus,

$$J_h(u_h + \delta u_h) - J_h(u_h) =$$

## Avoiding the Mean-value Adjoint (cont.)

The above approximation for the mean-value adjoint is frequently used in the practice.

- The approximation is justified provided  $\delta u_h$  is sufficiently small.
- If the problem is highly nonlinear (e.g. shocks), then  $\delta u_h$  may become large.



# Avoiding Fine-space Solutions

The second problem with the exact AWR is the need for  $\psi_h$ , which implies the solution of

$$(R'_h[u_h])^T \psi_h = J'_h[u_h].$$

This is problematic for two reasons:

- The linear system for  $\psi_h$  may be significantly larger than the baseline problem for  $u_H$ .
- The Jacobian and gradient are linearized about  $u_h$ , which we do not have.

## Avoiding Fine-space Solutions (cont.)

One common approach to approximate  $\psi_h$  (and  $u_h$ ), is to use a patch-based reconstruction, in which  $\psi_H$  is interpolated onto the fine space using a larger stencil.

Examples:

## Avoiding Fine-space Solutions (cont.)

A disadvantage of the reconstruction approach, especially in the case of  $u_h$ , is that it does not account for the physics of the problem.

- A high-order interpolation is not justified near a discontinuity, such as a shock.

## Avoiding Fine-space Solutions (cont.)

Consequently, a second common approach is to approximately solve for  $\psi_h$  using a few iterations of an iterative method.

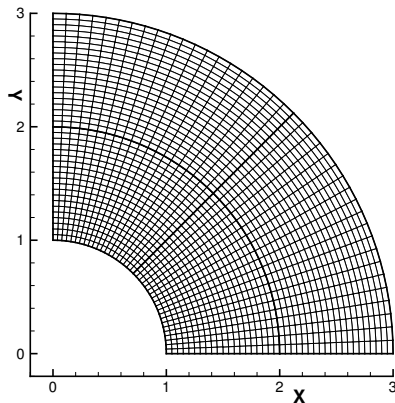
- This is justified, since the difference  $\psi_h^H - \psi_h$  is likely to contain mostly high-frequency terms that can be eliminated rapidly using a suitable stationary iterative method (i.e. a smoother).

# AWR Examples: Euler Equations

# Vortex Flow [Hic12]

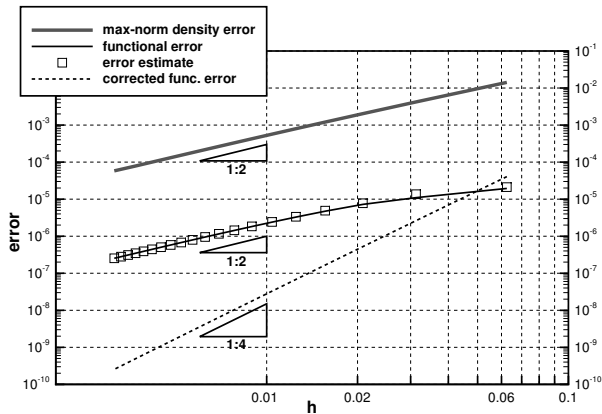
Isentropic Vortex flow:

- streamlines are concentric circles
- smooth solution (ideal case)
- functional is force in  $x$ -direction over inner radius.



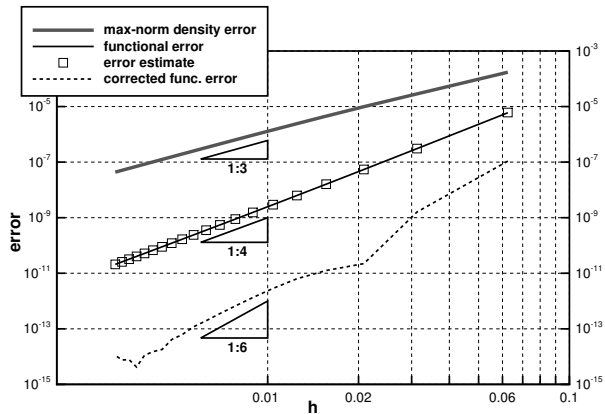
# Vortex Flow [Hic12] (cont.)

2nd-order coarse space; 3rd-order fine space



# Vortex Flow [Hic12] (cont.)

3rd-order coarse space; 4th-order fine space



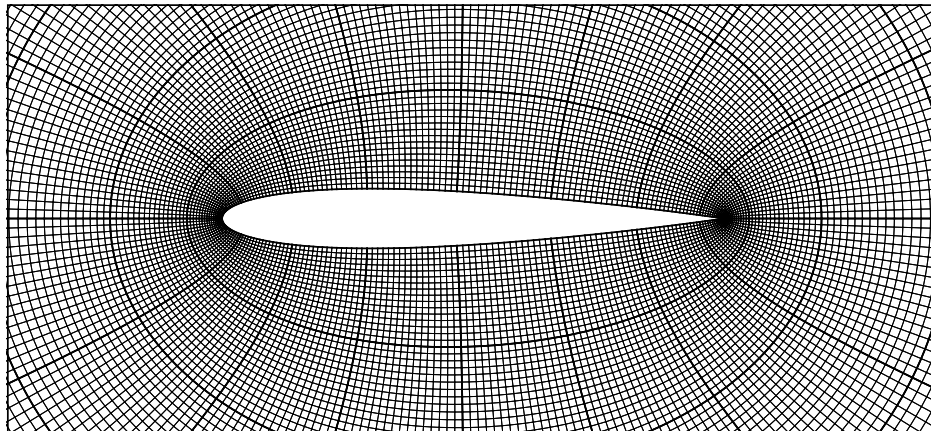


# Subsonic Airfoil [Hic12]

Mach 0.5 flow over NACA 0012

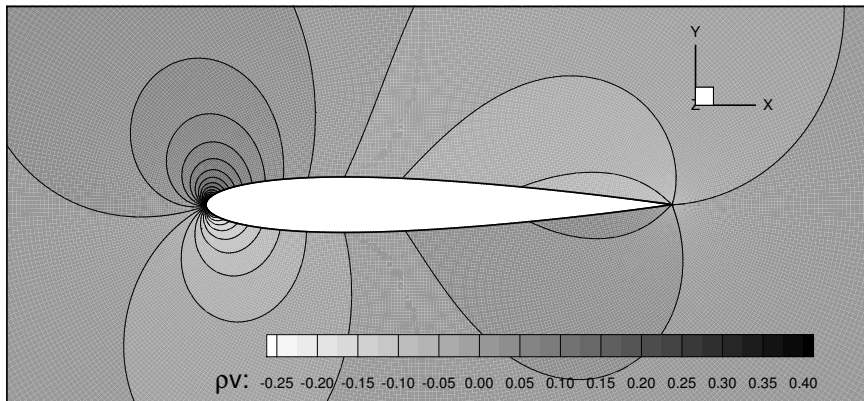
- functional is coefficient of drag
- flow and adjoint fields have singularities at trailing edge

# Subsonic Airfoil [Hic12] (cont.)



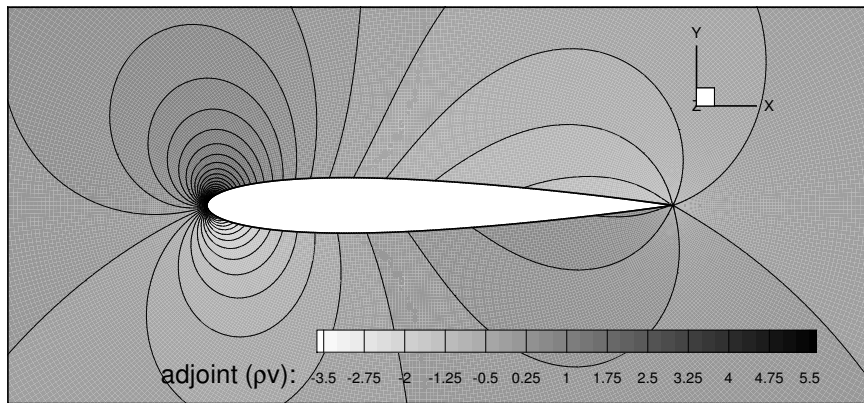
# Subsonic Airfoil [Hic12] (cont.)

$\rho v$  contours



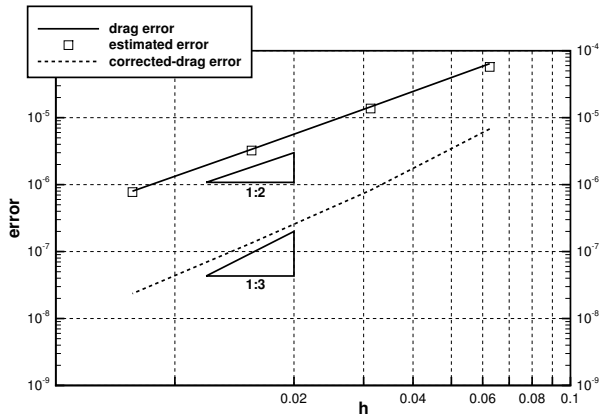
# Subsonic Airfoil [Hic12] (cont.)

$\psi_{pv}$  contours



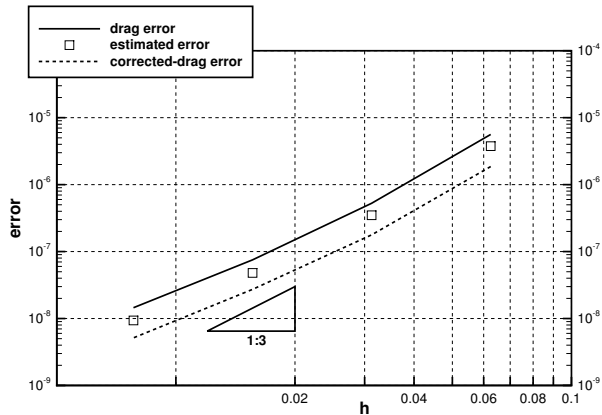
# Subsonic Airfoil [Hic12] (cont.)

2nd-order coarse space; 3rd-order fine space



# Subsonic Airfoil [Hic12] (cont.)

3rd-order coarse space; 4th-order fine space



# References

- [FD11] Krzysztof J. Fidkowski and David L. Darmofal, *Review of output-based error estimation and mesh adaptation in computational fluid dynamics*, AIAA Journal **49** (2011), no. 4, 673–694.
- [Hic12] Jason E. Hicken, *Output error estimation for summation-by-parts finite-difference schemes*, Journal of Computational Physics **231** (2012), no. 9, 3828–3848.