

PSWF¹-Radon approach to reconstruction from band-limited Hankel transform

Rodion Zaytsev
Supervisor: Roman G. Novikov

December 6, 2023

¹Prolate Spheroidal Wave Functions

Abstract

- ▶ New formulas for reconstructions from band-limited Hankel transform of integer and half-integer order
- ▶ PSWF-Radon approach to super-resolution in multidimensional Fourier analysis
- ▶ Numerical examples to illustrate super-resolution

Presentation Plan

1 Introduction

- Preliminaries
- Applications of Hankel Transform

2 Band-Limited Hankel Transform

- Problem statement
- Statement of our method
- Outline of the proof
- Numerical examples

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Bessel Functions (of the first kind)

Integral representation

Bessel function of *order* $\nu \in \mathbb{R}$

$$J_\nu(x) = \frac{1}{\pi} \int_0^\pi \cos(\nu\tau - x \sin \tau) d\tau - \frac{\sin \nu\pi}{\pi} \int_0^\infty e^{-x \sinh t - \nu t} dt \quad (1.1)$$

Bessel Functions of Integer Order

For $\nu \in \mathbb{Z}$

$$\blacktriangleright J_\nu(x) = \frac{1}{\pi} \int_0^\pi \cos(\nu\tau - x \sin \tau) d\tau = \frac{1}{2\pi} \int_{-\pi}^\pi e^{i(\nu\tau - x \sin \tau)} d\tau$$

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$$\blacktriangleright J_{-\nu}(x) = (-1)^\nu J_\nu(x)$$

Bessel Functions of Half-integer Order

For $\nu \in \frac{1}{2}\mathbb{N}$ the ordinary Bessel function is related to the spherical Bessel function by

$$j_{\nu-\frac{1}{2}}(x) = \sqrt{\frac{\pi}{2x}} J_{\nu}(x) \quad (1.2)$$

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Rayleigh's formula for the spherical Bessel functions

$$j_n(x) = (-x)^n \left(\frac{1}{x} \frac{d}{dx} \right)^n \frac{\sin x}{x} \quad n \in \mathbb{Z}_{\geq 0} \quad (1.3)$$

Hankel Transform

Definition

Hankel transform of order $\nu \geq -\frac{1}{2}$

$$\begin{aligned}\mathcal{H}_\nu : \mathcal{L}^2(\mathbb{R}_+) &\rightarrow \mathcal{L}^2(\mathbb{R}_+) \\ \mathcal{H}_\nu[f](x) &:= \int_0^\infty f(y) J_\nu(xy) \sqrt{xy} \, dy, \quad x \geq 0\end{aligned}\tag{1.4}$$

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Remark

It is an involution, i.e. $\mathcal{H}_\nu = \mathcal{H}_\nu^{-1}$

Hankel Transform and the 2D Laplacian

Theorem

If $\lim_{r \rightarrow \infty} rf'(r) = \lim_{r \rightarrow \infty} rf(r) = 0$ then

$$\mathcal{H}_\nu \left[\left(\frac{1}{r} \frac{d}{dr} \left(r \frac{d}{dr} \right) - \frac{\nu^2}{r^2} \right) f(r) \right] (p) = -p^2 \mathcal{H}_\nu[f](p) \quad (1.5)$$

Remark

Consider the 2-dimensional Laplacian

$$\nabla^2 = \frac{\partial^2}{\partial x_1^2} + \frac{\partial^2}{\partial x_2^2}$$

Its radial part is

$$\frac{1}{r} \frac{\partial}{\partial r} \left(r \frac{\partial}{\partial r} \right)$$

Hankel transform in PDEs

We will consider the following example

Axisymmetric Steady-State Heat Equation with a Symmetric Source

$$u_{rr} + \frac{1}{r}u_r + u_{zz} = -q(r), \quad 0 < r < \infty, \quad 0 < z < \infty \quad (1.6)$$

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Solution

Apply the zeroth Hankel transform over r and use the theorem from the previous slide. Denoting $\tilde{u}(p, z) := \mathcal{H}_0[u]$ and $\tilde{q}(p) := \mathcal{H}_0[q](p)$, we have

$$\tilde{u}_{zz} - p^2\tilde{u} = -\tilde{q} \quad (1.8)$$

$$\tilde{u}(p, 0) = 0 \quad (1.9)$$

Hankel transform in PDEs

Solution (continued)

The bounded solution of this boundary value problem is

$$\tilde{u}(p, z) = \frac{\tilde{q}(p)}{p^2} (1 - e^{-pz}) \quad (1.10)$$

Taking the inverse Hankel transform (which is the Hankel transform itself) we get the answer to the original problem

$$u(r, z) = \mathcal{H}_0 \left[\frac{\mathcal{H}_0[q](p)}{p} (1 - e^{-pz}) \right] (r, z) \quad (1.11)$$

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Let $\sigma, r > 0$ be given. Find $f \in \mathcal{L}^2(\mathbb{R}_+)$ from $w = \mathcal{H}_\nu[f]$ given on $[0, r]$ (possibly with some noise), under a priori assumption that $\text{supp } f \subset [0, \sigma]$.

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Remark (noiseless case)

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We will only consider the problem in case ν is integer or half-integer.

Naive Approach

$$f \approx f_{\text{naive}} := \mathcal{H}_{\nu}^{-1} [w^{\text{ext}}] \quad \text{on } [0, \sigma], \quad (2.1)$$

where

$$w^{\text{ext}}(x) := \begin{cases} w(x), & \text{for } x \in [0, r], \\ 0, & \text{otherwise.} \end{cases}$$

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- ▶ Stable and accurate reconstruction for sufficiently large r
- ▶ *diffraction limit*: small details (especially less than π/r) are blurred

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Definition

Super-resolution: techniques that allow reconstruction beyond this diffraction limit. This is the main purpose of our work.

Operator \mathcal{F}_c

Definition

$$\begin{aligned}\mathcal{F}_c : \mathcal{L}^2([-1, 1]) &\rightarrow \mathcal{L}^2([-1, 1]) \\ \mathcal{F}_c[f](x) &:= \mathcal{F}_c[f](x) := \int_{-1}^1 e^{icxy} f(y) dy\end{aligned}\tag{2.2}$$

$c > 0$ is the *bandwidth parameter*

Prolate Spheroidal Wave Functions

SVD decomposition for \mathcal{F}_c

$$\mathcal{F}_c[f](x) = \sum_{j=0}^{\infty} \mu_{j,c} \psi_{j,c}(x) \int_{-1}^1 \psi_{j,c}(y) f(y) dy$$

$$\mathcal{F}_c^{-1}[g](y) = \sum_{j \in \mathbb{N}} \frac{1}{\mu_{j,c}} \psi_{j,c}(y) \int_{-1}^1 \psi_{j,c}(x) g(x) dx$$

Definition

The eigenfunctions $\{\psi_{j,c}, j = 0, 1, 2, \dots\}$ of \mathcal{F}_c are *prolate spheroidal wave functions* (PSWFs)

Remark

Their eigenvalues obey $0 < |\mu_{j+1,c}| < |\mu_{j,c}|$

Exact formula for integer ν

Let T_n , $n = 0, 1, 2 \dots$ denote the Chebyshev polynomial of the first kind

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Theorem (Cormack-type formula for integer ν)

Let $r, \sigma > 0$, $c = r\sigma$, and $f \in \mathcal{L}^2(\mathbb{R}_+)$ be supported in $[0, \sigma]$. Then, for $\nu \in \mathbb{Z}$, the Hankel transform $\mathcal{H}_\nu[f]$ on $[0, r]$ uniquely determines f by the formula

$$f(y) = -\frac{2i^\nu}{\sigma} \sqrt{y} \frac{d}{dy} \int_y^\sigma \frac{y T_{|\nu|}\left(\frac{x}{y}\right)}{x(x^2 - y^2)^{\frac{1}{2}}} \mathcal{F}_c^{-1}[g_{r,\nu}](x/\sigma) dx, \quad y \in [0, \sigma]$$

$$g_{r,\nu}(x) := \begin{cases} \frac{1}{\sqrt{r|x|}} \mathcal{H}_\nu[f](r|x|), & \text{if } x \geq 0, \\ (-1)^\nu \frac{1}{\sqrt{r|x|}} \mathcal{H}_\nu[f](r|x|), & \text{otherwise.} \end{cases}$$

Exact formula for half-integer ν

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$$f(y) = \frac{\sqrt{2\pi}i^{2\nu-1}}{\sigma} \frac{d^2}{dy^2} \int_y^\sigma \frac{y^2}{x^2} P_{2\nu-1}\left(\frac{x}{y}\right) \mathcal{F}_c^{-1}[g_{r,\nu}](x/\sigma) dx, \quad y \in [0, \sigma]$$

$$g_{r,\nu}(x) := \begin{cases} \frac{1}{r|x|} \mathcal{H}_\nu[f](r|x|), & \text{if } x \geq 0, \\ (-1)^{2\nu-1} \frac{1}{r|x|} \mathcal{H}_\nu[f](r|x|), & \text{otherwise.} \end{cases}$$

Plane-wave expansion

$$e^{ipq} = \sum_{l \in \mathbb{Z}} i^l J_l(|p||q|) e^{il(\phi_p - \phi_q)} \quad (d = 2)$$

$$e^{ipq} = 4\pi \sum_{l=0}^{\infty} \sum_{m=-l}^l i^l j_l(|p||q|) Y_{lm}^*(\theta_q, \phi_q) Y_{lm}(\theta_p, \phi_p) \quad (d = 3)$$

Connection between the Fourier and Hankel Transforms

► $d = 2$

$$\mathcal{F}[f(|\mathbf{q}|)e^{il\phi_{\mathbf{q}}}] (|\mathbf{p}|, \phi_{\mathbf{p}}) = \frac{1}{2\pi} i^l e^{il\phi_{\mathbf{p}}} H_l[f](|\mathbf{p}|) \quad (2.3)$$

► $d = 3$

$$\begin{aligned} \mathcal{F}[f(|\mathbf{q}|)Y_{lm}(\theta_{\mathbf{q}}, \phi_{\mathbf{q}})] (|\mathbf{p}|, \theta_{\mathbf{p}}, \phi_{\mathbf{p}}) = \\ \frac{i^l}{2\pi^2|\mathbf{p}|} \sqrt{\frac{\pi}{2}} Y_{lm}(\theta_{\mathbf{p}}, \phi_{\mathbf{p}}) H_{l+\frac{1}{2}}[|\mathbf{q}|f(|\mathbf{q}|)] \end{aligned} \quad (2.4)$$

Projection Theorem

► Classical

$$\mathcal{F}[v](p) = \hat{v}(p) := \frac{1}{(2\pi)^d} \int_{\mathbb{R}^d} e^{ipq} v(q) dq, \quad p \in \mathbb{R}^d \quad (2.5)$$

$$\mathcal{R}_\theta[u](y) := \int_{q \in \mathbb{R}^d : q\theta = y} u(q) dq, \quad y \in \mathbb{R} \quad (2.6)$$

$$\mathcal{F}[u](s\theta) = \frac{1}{(2\pi)^d} \int_{-\infty}^{\infty} e^{ist} \mathcal{R}[u](t, \theta) dt, \quad s \in \mathbb{R}, \theta \in \mathbb{S}^{d-1} \quad (2.7)$$

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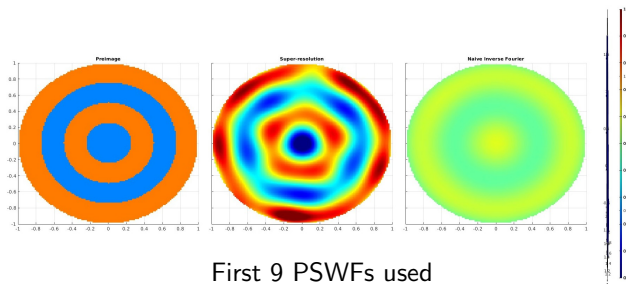
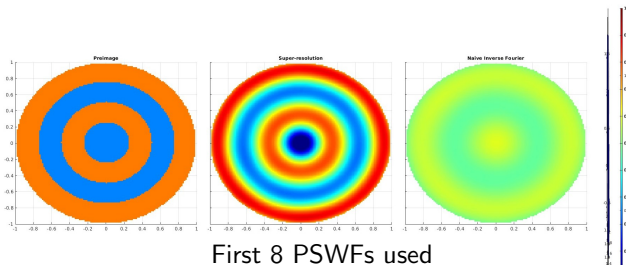
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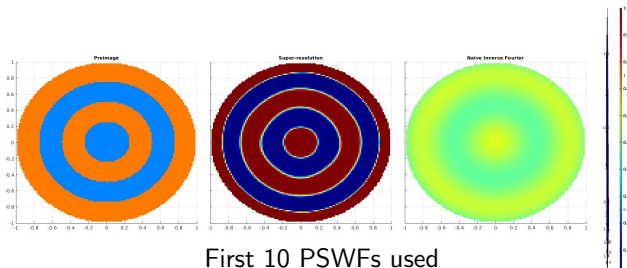
► Band-limited analogue

$$\hat{v}(r x \theta) = \left(\frac{\sigma}{2\pi} \right)^d \mathcal{F}_c [\mathcal{R}_\theta[v_\sigma]](x), \quad \text{for } d \geq 2, \quad (2.8)$$

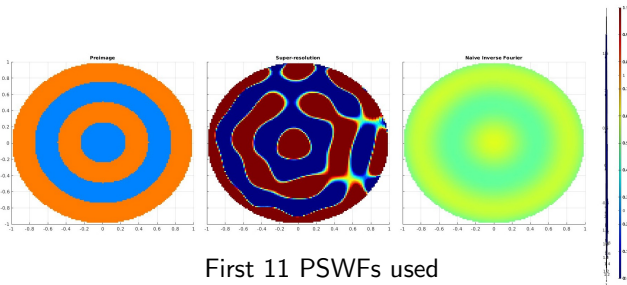
20% Noise Reconstruction



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First 10 PSWFs used



First 11 PSWFs used

Q&A

Thank you for attention!

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