

Problem 1)

Two concentric spheres have radii a, b ($b > a$) and each is divided into two hemispheres by the same horizontal plane. The upper hemisphere of the inner sphere and the lower hemisphere of the outer sphere are maintained at potential V . The other hemispheres are at zero potential.

Determine the potential in the region $a \leq r \leq b$ as a series in Legendre polynomials. Include terms at least up to $l = 4$. Check your solution against known results in the limiting cases $b \rightarrow \infty$, and $a \rightarrow 0$.

Since the problem is rotationally symmetric about the axis perpendicular to the plane separating the hemispheres of the cylinder, passing through the center of the spheres, we can write the potential as

$$\Phi(r, \phi) = \sum_{l=0}^{\infty} (A_l r^l + B_l r^{-l-1}) P_l(\cos \theta). \quad (1)$$

Note that we must keep both terms with r -dependence since there is no singular behavior over the region of interest. We have two boundary conditions:

$$\Phi(a, \theta) = \begin{cases} V & -\pi/2 < \theta < \pi/2 \\ 0 & \pi/2 < \theta < 3\pi/2, \end{cases} \quad \Phi(b, \theta) = \begin{cases} V & \pi/2 < \theta < 3\pi/2 \\ 0 & -\pi/2 < \theta < \pi/2. \end{cases} \quad (2)$$

We can exploit the orthogonality of legendre polynomials

$$\int_{-1}^1 dx P_l(x) P_{l'}(x) = \int_0^\pi d\theta \sin \theta P_l(\cos \theta) P_{l'}(\cos \theta) = \frac{2}{2l+1} \delta_{ll'}. \quad (3)$$

Hence, at $r = a$ and $r = b$, respectively, we have

$$A_l a^l + B_l a^{-l-1} = \frac{2l+1}{2} V \int_0^{\pi/2} \sin \theta P_l(\cos \theta) d\theta = \frac{2l+1}{2} V \int_0^1 P_l(x) dx \quad (4)$$

$$A_l b^l + B_l b^{-l-1} = \frac{2l+1}{2} V \int_{\pi/2}^\pi \sin \theta P_l(\cos \theta) d\theta = \frac{2l+1}{2} (-1)^l V \int_0^1 P_l(x) dx. \quad (5)$$

Solving this system of equations gives

$$A_l = \frac{2l+1}{2} \frac{a^{l+1} - (-1)^l b^{l+1}}{a^{2l+1} - b^{2l+1}} V \int_0^1 P_l(x) dx \quad (6)$$

$$B_l = \frac{2l+1}{2} \frac{a^{l+1} b^{l+1} [(-1)^l a^l - b^l]}{a^{2l+1} - b^{2l+1}} V \int_0^1 P_l(x) dx. \quad (7)$$

Thus, the potential for $a \leq r \leq b$ becomes

$$\begin{aligned} \Phi(r, \theta) = \frac{V}{2} \sum_{l=0}^{\infty} \frac{2l+1}{a^{2l+1} - b^{2l+1}} & \left[(a^{l+1} - (-1)^l b^{l+1}) r^l + a^{l+1} b^{l+1} ((-1)^l a^l - b^l) r^{-l-1} \right] \\ & \times \left[\int_0^1 dy P_l(y) \right] P_l(\cos \theta). \end{aligned} \quad (8)$$

We can go further and evaluate the integral of the Legendre polynomial as follows:

$$\int_0^1 P_l(x) dx = \frac{1}{2l+1} \int_0^1 [P_{l+1}(x) - P_{l-1}(x)] dx = \frac{P_{l-1}(0) - P_{l+1}(0)}{2l+1}. \quad (9)$$

Note that for even l , and therefore odd $l \pm 1$, the integral is zero. Furthermore note that we must be careful in the case $l = 0$ since we have not defined P_l for $l < 0$:

$$\int_0^1 P_0(x) dx = 1. \quad (10)$$

Putting this into our expression for the potential:

$$\Phi(r, \theta) = \frac{V}{2} \left(1 + \sum_{l=0}^{\infty} [P_{2l}(0) - P_{2(l+1)}(0)] \left[\frac{a^{2(l+1)} + b^{2(l+1)}}{a^{4l+3} - b^{4l+3}} r^{2l+1} - \frac{a^{2(l+1)} b^{2(l+1)} (a^{2l+1} + b^{2l+1})}{a^{4l+3} - b^{4l+3}} r^{-2(l+1)} \right] P_{2l+1}(\cos \theta) \right), \quad (11)$$

where we have redefined $l \rightarrow 2l+1$ to only capture the non-zero terms. Through the first 4 terms (of the original sum indices)

$$\Phi(r, \theta) = \frac{V}{2} \left\{ 1 + \frac{3}{2} \frac{[a^2 - b^2]r - a^2 b^2 [a + b] r^{-2}}{a^3 - b^3} P_1(\cos \theta) - \frac{7}{8} \frac{[a^4 - b^4]r^3 - a^4 b^4 [a^3 + b^3] r^{-4}}{a^7 - b^7} P_3(\cos \theta) + \dots \right\}. \quad (12)$$

If we take $b \rightarrow \infty$, we have

$$A_l \rightarrow 0 \text{ and } B_l \rightarrow \frac{V}{2} [P_{l-1}(0) - P_{l+1}(0)] a^{l+1} V, \quad (13)$$

giving

$$\Phi(r, \theta) \Big|_{b \rightarrow \infty} = \frac{V}{2} + \frac{V}{2} \sum_{l=0}^{\infty} [P_{2l}(0) - P_{2(l+1)}(0)] \left(\frac{a}{r} \right)^{2(l+1)} P_{2l+1}(\cos \theta). \quad (14)$$

If we instead take $a \rightarrow 0$, we obtain

$$A_l \rightarrow \frac{V}{2} [P_{l-1}(0) - P_{l+1}(0)] (-1)^l b^{-l} \text{ and } B_l \rightarrow 0, \quad (15)$$

giving

$$\Phi(r, \theta) \Big|_{a \rightarrow 0} = \frac{V}{2} + \frac{V}{2} \sum_{l=0}^{\infty} [P_{2l}(0) - P_{2(l+1)}(0)] \left(\frac{r}{b} \right)^{2l+1} P_{2l+1}(\cos \theta). \quad (16)$$

Problem 2)

A spherical surface of radius R has charge uniformly distributed over its surface with a density $Q/4\pi R^2$, except for a spherical cap at the north pole, defined by the cone $\theta = \alpha$.

(a) Show that the potential inside the spherical surface can be expressed as

$$\Phi = \frac{Q}{8\pi\epsilon_0} \sum_{l=0}^{\infty} \frac{1}{2l+1} [P_{l+1}(\cos \alpha) - P_{l-1}(\cos \alpha)] \frac{r^l}{R^{l+1}} P_l(\cos \theta) \quad (17)$$

where, for $l = 0$, $P_{l-1}(\cos \alpha) = -1$. What is the potential outside?

(b) Find the magnitude and the direction of the electric field at the origin.

(c) Discuss the limiting forms of the potential (part (a)) and electric field (part (b)) as the spherical cap becomes (1) very small, and (2) so large that the area with charge on it becomes a very small cap at the south pole.

(a) This problem again has azimuthal symmetry because of the rotational invariance about the axis of the cone, meaning

$$\Phi_{\text{in}}(r, \theta) = \sum_{l=0}^{\infty} A_l r^l P_l(\cos \theta), \quad (18)$$

inside the sphere, and

$$\Phi_{\text{out}}(r, \theta) = \sum_{l=0}^{\infty} B_l r^{-l-1} P_l(\cos \theta). \quad (19)$$

Note that technically, we could include a constant term in the potential outside the sphere, but this must be zero since the potential at infinity is zero (the charge distribution far from the sphere is just that of a point charge).

Specifying the surface charge density amounts to specifying the derivative of the potential over the surface:

$$-\frac{\partial \Phi_{\text{out}}}{\partial r} + \frac{\partial \Phi_{\text{in}}}{\partial r} = \frac{\sigma}{\epsilon_0} = \frac{Q}{4\pi\epsilon_0 R^2} \Theta(\theta - \alpha), \quad (20)$$

where $\Theta(x)$ is the Heaviside step function. Taking derivatives, we find

$$\sum_{l=0}^{\infty} [(l+1)B_l r^{-l-2} + lA_l r^{l-1}] P_l(\cos \theta) = \frac{Q}{4\pi\epsilon_0 R^2} \Theta(\theta - \alpha). \quad (21)$$

Using the orthogonality relation for Legendre polynomials, we find

$$\begin{aligned} (l+1)B_l R^{-l} + lA_l R^{l+1} &= (2l+1) \frac{Q}{8\pi\epsilon_0} \int_{\alpha}^{\pi} \sin \theta P_l(\cos \theta) d\theta \\ &= \frac{Q}{8\pi\epsilon_0} [P_{l+1}(\cos \alpha) - P_{l-1}(\cos \alpha)], \end{aligned} \quad (22)$$

where we have defined $P_{-1}(0) = -1$ in order to include the $l = 0$ term implicitly.

Observe that we also have a BC at $r = R$, where we impose continuity of the potential:

$$\sum_{l=0}^{\infty} A_l R^l P_l(\cos \theta) = \sum_{l=0}^{\infty} B_l R^{-l-1} P_l(\cos \theta) \Rightarrow B_l = A_l R^{2l+1}. \quad (23)$$

Putting this into the first BC, we have

$$A_l = \frac{Q}{8\pi\epsilon_0} \frac{1}{(2l+1)R^{l+1}} [P_{l+1}(\cos \alpha) - P_{l-1}(\cos \alpha)], \quad (24)$$

and therefore

$$B_l = \frac{Q}{8\pi\epsilon_0} \frac{R^l}{2l+1} [P_{l+1}(\cos \alpha) - P_{l-1}(\cos \alpha)]. \quad (25)$$

Thus, the potential inside the sphere is

$$\Phi(r, \theta) = \frac{Q}{8\pi\epsilon_0} \sum_{l=0}^{\infty} \frac{P_{l+1}(\cos \alpha) - P_{l-1}(\cos \alpha)}{2l+1} \frac{r^l}{R^{l+1}} P_l(\cos \theta), \quad (26)$$

and outside the sphere

$$\Phi(r, \theta) = \frac{Q}{8\pi\epsilon_0} \sum_{l=0}^{\infty} \frac{P_{l+1}(\cos \alpha) - P_{l-1}(\cos \alpha)}{2l+1} \frac{R^l}{r^{l+1}} P_l(\cos \theta). \quad (27)$$

(b) The electric field inside the sphere is given by

$$\vec{E} = -\vec{\nabla}\Phi = -\hat{r} \frac{\partial \Phi}{\partial r} + \hat{\theta} \frac{1}{r} \frac{\partial \Phi}{\partial \theta}. \quad (28)$$

Hence,

$$\begin{aligned} \vec{E} = & \hat{r} \frac{Q}{8\pi\epsilon_0} \sum_{l=1}^{\infty} \frac{l[P_{l-1}(\cos \alpha) - P_{l+1}(\cos \alpha)]}{2l+1} \frac{r^{l-1}}{R^{l+1}} P_l(\cos \theta) \\ & + \hat{\theta} \frac{Q}{4\pi\epsilon_0} \sum_{l=0}^{\infty} \frac{P_{l-1}(\cos \alpha) - P_{l+1}(\cos \alpha)}{2l+1} \frac{r^{l-1}}{R^{l+1}} \frac{\partial P_l(\cos \theta)}{\partial \theta}. \end{aligned} \quad (29)$$

At the origin, the angular dependence only comes from the $l = 1$ term since the derivative of P_0 is zero and each higher-order term is proportional to a positive power of r . This leaves us with

$$\vec{E} = \frac{Q}{8\pi\epsilon_0} \frac{P_0(\cos \alpha) - P_2(\cos \alpha)}{3R^2} [\cos \theta \hat{r} - \sin \theta \hat{\theta}]. \quad (30)$$

Since,

$$\hat{\mathbf{r}} = \sin \theta \cos \phi \hat{\mathbf{x}} + \sin \theta \sin \phi \hat{\mathbf{y}} + \cos \theta \hat{\mathbf{z}} \quad (31)$$

$$\hat{\boldsymbol{\theta}} = \cos \theta \cos \phi \hat{\mathbf{x}} + \cos \theta \sin \phi \hat{\mathbf{y}} - \sin \theta \hat{\mathbf{z}}, \quad (32)$$

we have

$$\vec{\mathbf{E}} = \frac{Q}{8\pi\epsilon_0} \frac{2 - 3\cos^2 \alpha + 1}{6R^2} \hat{\mathbf{z}} = \frac{Q}{16\pi\epsilon_0 R^2} \sin^2 \alpha \hat{\mathbf{z}}. \quad (33)$$

From symmetry considerations, we can see that this at least points in the right direction.

(c) $\alpha \rightarrow 0$: In this case, the electric field at the center is

$$\vec{\mathbf{E}} \approx \frac{Q}{16\pi\epsilon_0 R^2} \alpha^2 \hat{\mathbf{z}}. \quad (34)$$

For the potential on the other hand, we have

$$P_l(\cos \alpha) \approx P_l(1 - \alpha^2/2) \approx P_l(1) - \frac{\alpha^2}{2} P'_l(1) = 1 - \frac{\alpha^2}{2} \frac{l(l+1)}{2}, \quad (35)$$

so

$$P_{l+1}(\cos \alpha) - P_{l+1}(\cos \alpha) \approx \left(1 - \frac{\alpha^2}{2} \frac{(l+1)(l+2)}{2}\right) - \left(1 - \frac{\alpha^2}{2} \frac{l(l-1)}{2}\right) = -\frac{2l+1}{2} \alpha^2. \quad (36)$$

The potential inside the sphere is then

$$\begin{aligned} \Phi &= \frac{Q}{8\pi\epsilon_0} \left[(\cos \alpha + 1) \frac{1}{R} - \frac{\alpha^2}{2} \sum_{l=1}^{\infty} \frac{r^l}{R^{l+1}} P_l(\cos \theta) \right] \\ &= \frac{Q}{8\pi\epsilon_0} \left[(2 - \alpha^2/2) \frac{1}{R} - \frac{\alpha^2}{2} \left(\frac{1}{|\vec{\mathbf{r}} - \vec{\mathbf{R}}|} - \frac{1}{R} \right) \right] \\ &= \boxed{\frac{Q}{4\pi\epsilon_0 R} - \frac{Q}{16\pi\epsilon_0} \frac{\alpha^2}{|\vec{\mathbf{r}} - \vec{\mathbf{R}}|}}. \end{aligned} \quad (37)$$

Note that this is actually almost the potential outside the sphere since this limit describes essentially a uniformly charged sphere with a point charge of opposite sign (relative to the surface charge) at the north pole. We just have to change the first term to reflect this:

$$\boxed{\Phi = \frac{Q}{4\pi\epsilon_0 r} - \frac{Q}{16\pi\epsilon_0} \frac{\alpha^2}{|\vec{\mathbf{r}} - \vec{\mathbf{R}}|}}. \quad (38)$$

$\alpha \rightarrow \pi$: In this case, we can define $\beta = \pi - \alpha$, which is very small in this limit. The electric field is then just

$$\vec{\mathbf{E}} \approx \frac{Q}{16\pi\epsilon R^2} \beta^2 \hat{\mathbf{z}}. \quad (39)$$

We notice that this limit describes essentially a point charge (with the same magnitude as in the previous part but opposite sign) at the south pole ($\vec{r} = -\vec{R}$) of the sphere, which just has potential

$$\Phi = \frac{Q}{16\pi\epsilon_0} \frac{\alpha^2}{|\vec{r} + \vec{R}|}, \quad (40)$$

which is valid both inside and outside the sphere.

Problem 3)

A thin, flat, conducting, circular disc of radius R is located in the xy plane with its center at the origin, and is maintained at a fixed potential V . With the information that the charge density on a disc at fixed potential is proportional to $(R^2 - \rho^2)^{-1/2}$, where ρ is the distance out from the center of the disc,

(a) show that for $r > R$, the potential is

$$\Phi(r, \theta, \phi) = \frac{2V}{\pi} \frac{R}{r} \sum_{l=0}^{\infty} \frac{(-1)^l}{2l+1} \left(\frac{R}{r}\right)^{2l} P_{2l}(\cos \theta). \quad (41)$$

(c) find the potential for $r < R$.

(c) What is the capacitance of the disc?

(a) The surface charge density has the form

$$\sigma(\rho) = \frac{\sigma_0}{\sqrt{R^2 - \rho^2}}. \quad (42)$$

We can relate σ_0 to the potential V by

$$V = \frac{1}{4\pi\epsilon_0} \int dS \frac{\sigma(\rho)}{\rho} = \frac{\sigma_0}{2\epsilon_0} \int_0^R \frac{d\rho}{\sqrt{R^2 - \rho^2}} = \frac{\pi}{4\epsilon_0} \sigma_0, \quad (43)$$

so

$$\sigma_0 = \frac{4\epsilon_0 V}{\pi}. \quad (44)$$

We can now calculate the potential anywhere in space as

$$\begin{aligned} \Phi(\vec{r}) &= \frac{1}{4\pi\epsilon_0} \int d^3\vec{r}' \frac{\rho(\vec{r}')}{|\vec{r} - \vec{r}'|} \\ &= \frac{V}{\pi^2} \int_0^R \int_0^{2\pi} \int_{-1}^1 dr' d\phi' d(\cos \theta') \frac{r'}{\sqrt{R^2 - r'^2}} \delta(\cos \theta') \sum_{l=0}^{\infty} \frac{r_{<}^l}{r_{>}^{l+1}} P_l(\cos \theta') \\ &= \frac{2V}{\pi} \int_0^R dr' \frac{r'}{\sqrt{R^2 - r'^2}} \sum_{l=0}^{\infty} \frac{r_{<}^l}{r_{>}^{l+1}} P_l(0). \end{aligned} \quad (45)$$

In general, the integration of $P_l(\cos \gamma)$ is not possible, but we can fix \vec{r} to lie along the z -axis such that $\gamma = \theta'$. Observe that for this problem, $r_{<} = \min\{r', r\} = r'$ and $r_{>} = \max\{r', r\} = r$ since we are interested in $r > R \geq r'$. Hence, the potential along the z -axis is

$$\begin{aligned}
 \Phi(\vec{r}) &= \frac{2V}{\pi} \sum_{l=0}^{\infty} \frac{P_l(0)}{r^{l+1}} \int_0^R dr' \frac{r'^{l+1}}{\sqrt{R^2 - r'^2}} \\
 &= \frac{2V}{\pi} \sum_{l=0}^{\infty} \frac{1}{r^{2l+1}} \frac{(-1)^l (2l-1)!!}{2^l l!} \int_0^R dr' \frac{r'^{2l+1}}{\sqrt{R^2 - r'^2}} \\
 &= \frac{2V}{\pi} \sum_{l=0}^{\infty} \frac{R^{2l+1}}{r^{2l+1}} \frac{(-1)^l (2l-1)!!}{2^l l!} \frac{\sqrt{\pi} \Gamma(l+1)}{2\Gamma(l+3/2)} \\
 &= \frac{2V}{\pi} \frac{R}{r} \sum_{l=0}^{\infty} \frac{(-1)^l (2l-1)!!}{2^l l!} \frac{\sqrt{\pi} l!}{2(l+1/2)\Gamma(l+1/2)} \frac{R^l}{r^l} \\
 &= \frac{2V}{\pi} \frac{R}{r} \sum_{l=0}^{\infty} \frac{(-1)^l}{2l+1} \left(\frac{R}{r}\right)^l.
 \end{aligned} \tag{46}$$

In order to obtain the angular dependence, we just append Legendre polynomials to this:

$$\Phi(\vec{r}) = \frac{2V}{\pi} \frac{R}{r} \sum_{l=0}^{\infty} \frac{(-1)^l}{2l+1} \left(\frac{R}{r}\right)^l P_{2l}(\cos \theta). \tag{47}$$

(b) We can determine the potential in the region $r < R$ fairly easily now that we have it in the region $r > R$. Since the setup is rotationally symmetric about the z -axis, we can write (enforcing that the potential be non-singular)

$$\Phi(\vec{r}) = \sum_{l=0}^{\infty} A_l r^l P_l(\cos \theta). \tag{48}$$

Note that the potential must be continuous at $r = R$, giving

$$\sum_{l=0}^{\infty} A_l R^l P_l(\cos \theta) = \sum_{l=0}^{\infty} \frac{2V}{\pi} \frac{(-1)^l}{2l+1} P_{2l}(\cos \theta). \tag{49}$$

Thus,

$$A_{2l} = \frac{2V}{\pi} \frac{(-1)^l}{2l+1} R^{-l}, \tag{50}$$

and $A_{2l+1} = 0$, which gives

$$\Phi(\vec{r}) = \frac{2V}{\pi} \sum_{l=0}^{\infty} \frac{(-1)^l}{2l+1} \left(\frac{r}{R}\right)^l P_{2l}(\cos \theta). \tag{51}$$

(c) Lastly, we just need to find the total charge Q on the disk in terms of V :

$$C = \frac{Q}{V} = 8\epsilon_0 R. \quad (52)$$

Problem 4)

Three point charges $(q, -2q, q)$ are located in a straight line with separation a and with the middle charge $(-2q)$ at the origin of a grounded conducting spherical shell of radius b .

(a) Write down the potential of the three charges in the absence of the grounded sphere. Find the limiting form of the potential as $a \rightarrow 0$, but the product $qa^2 = Q$ remains finite. Write this latter answer in spherical coordinates.

(b) The presence of the grounded sphere of radius b alters the potential for $r < b$. The added potential can be viewed as caused by the surface-charge density induced on the inner surface at $r = b$ or by image charges located at $r > b$. Use linear superposition to satisfy the boundary conditions and find the potential everywhere inside the sphere for $r < a$ and $r > a$. Show that in the limit $a \rightarrow 0$,

$$\Phi(r, \theta, \phi) \rightarrow \frac{Q}{2\pi\epsilon_0 r^3} \left(1 - \frac{r^5}{b^5}\right) P_2(\cos \theta). \quad (53)$$

(a) Without the sphere, we simply have the potential

$$\Phi(\vec{x}) = \frac{q}{4\pi\epsilon_0} \left(\frac{1}{|\vec{x} - \vec{a}|} - \frac{2}{|\vec{x}|} + \frac{1}{|\vec{x} + \vec{a}|} \right), \quad (54)$$

where $\vec{a} = a\hat{z}$. If we assume $a \rightarrow 0$ but also take $q \rightarrow \infty$ such that $qa^2 = Q$ is constant in the limit, the potential becomes

$$\begin{aligned} \Phi(\vec{x}) &= \frac{q}{4\pi\epsilon_0} \left(\frac{1}{x} \left[2 + \left(\frac{a}{x}\right)^2 [3\cos^2\theta - 1] + \frac{1}{32} \left(\frac{a}{x}\right)^4 [20\cos 4\theta + 35\cos 4\theta + 9] \right] - \frac{2}{x} \right) \\ &\approx \frac{Q}{2\pi\epsilon_0} \frac{1}{x^3} P_2(\cos \theta) \end{aligned} \quad (55)$$

(b) We can write the potential with the grounded sphere as

$$\Phi(\vec{x}) = \sum_{l=0}^{\infty} A_l r^l P_l(\cos \theta) + \frac{q}{4\pi\epsilon_0} \left(\frac{1}{|\vec{x} - \vec{a}|} - \frac{2}{|\vec{x}|} + \frac{1}{|\vec{x} + \vec{a}|} \right) \quad (56)$$

since we can expand any contribution from an image charge or the induced surface charge density in legendre polynomials. The terms r^{-l-1} are not included though since there are

no extra point charges inside the sphere which are not included in the second term. We can find the coefficients A_l, B_l such that $\Phi \equiv 0$ on the shell:

$$\begin{aligned}\Phi(b, \theta, \phi) &= \sum_{l=0}^{\infty} A_l b^l P_l(\cos \theta) + \frac{q}{4\pi\epsilon_0} \sum_{l=0}^{\infty} \frac{a^l}{b^{l+1}} [1 + (-1)^l] P_l(\cos \theta) - \frac{q}{2\pi\epsilon_0 b} = 0 \\ \sum_{l=0}^{\infty} A_l b^l P_l(\cos \theta) + \sum_{l=1}^{\infty} \frac{q}{2\pi\epsilon_0 b} \left(\frac{a}{b}\right)^{2l} P_{2l}(\cos \theta) &= 0.\end{aligned}\quad (57)$$

We then have

$$A_{2l} = -\frac{q}{2\pi\epsilon_0 b^{2l+1}} \left(\frac{a^{2l}}{b^{2l}}\right)^{2l} \quad (58)$$

for $l \geq 1$. We also have $A_0 = A_{2l+1} = 0$.

Thus,

$$\Phi(\vec{x}) = \frac{q}{2\pi\epsilon_0} \sum_{l=1}^{\infty} \left(\frac{a}{b}\right)^{2l} \left[-\frac{r^{2l}}{b^{2l+1}} + \left(\frac{b}{a}\right)^{2l} \frac{r_{<}^{2l}}{r_{>}^{2l+1}} \right] P_{2l}(\cos \theta). \quad (59)$$

Thus, for $r > a$

$$\boxed{\Phi(\vec{x}) = \frac{q}{2\pi\epsilon_0} \sum_{l=1}^{\infty} \left(\frac{a}{b}\right)^{2l} \left[-\frac{r^{2l}}{b^{2l+1}} + \frac{b^{2l}}{r^{2l+1}} \right] P_{2l}(\cos \theta)} \quad (60)$$

while for $r < a$

$$\begin{aligned}\boxed{\Phi(\vec{x}) &= \frac{q}{2\pi\epsilon_0} \sum_{l=1}^{\infty} \left(\frac{a}{b}\right)^{2l} \left[-\frac{1}{b^{2l+1}} + \frac{b^{2l}}{a^{4l+1}} \right] r^l P_{2l}(\cos \theta) \\ &= \frac{q}{2\pi\epsilon_0} \sum_{l=1}^{\infty} a^{2l} \left[-\frac{1}{b^{4l+1}} + \frac{1}{a^{4l+1}} \right] r^l P_{2l}(\cos \theta)}.\end{aligned} \quad (61)$$

In the limit $a \rightarrow 0$, the only region that matters is $r > a$, which gives

$$\boxed{\Phi \approx \frac{Q}{2\pi\epsilon_0} \frac{1}{b^2} \left[\frac{b^2}{r^3} - \frac{r^2}{b^3} \right] = \frac{Q}{2\pi\epsilon_0 r^3} \left[1 - \frac{r^5}{b^5} \right]} \quad (62)$$