

1) The strong interactions are blind to flavor. All quarks carry the same color charge. Thus, the Σ triplet of baryons, which all have identical wave functions, except for different numbers of u and d quarks, should have the same mass up to two effects: the electromagnetic energy, and the difference in mass between the u and d quark.

a) First, calculate the electromagnetic binding (or anti-binding) energy for each of these baryons, using the following simple model: Assume that the three quarks are, on average, each 1 fm from each other in the Σ . Express your answer in MeV.

The quark content of the Σ baryons is listed below in Table 1

Q	qqq
0	uds
+	uus
-	dds

Table 1: Quark content of Σ baryons by charge.

The electric potential energy between charges q_1 and q_2 is given as $U(r) = q_1 q_2 / (4\pi\epsilon_0 \langle r \rangle)$, so for an average separation distance of $\langle r \rangle = 1$ fm we have

$$\boxed{\begin{aligned} V_{\Sigma^0} &= \frac{q_d q_s + q_d q_u + q_s q_u}{4\pi\epsilon_0 \langle r \rangle} = -0.48 \text{ MeV} \\ V_{\Sigma^+} &= \frac{q_u (2q_s + q_u)}{4\pi\epsilon_0 \langle r \rangle} = 0 \text{ MeV} \\ V_{\Sigma^-} &= \frac{q_d (2q_s + q_d)}{4\pi\epsilon_0 \langle r \rangle} = 0.48 \text{ MeV} \end{aligned}}, \quad (1)$$

b) Just to get rid of gravity once and for all, calculate the gravitational binding energy for the three quarks in this model. For simplicity, assume that each quark has about $1/3$ the mass of the Σ . MeV again.

The expressions above may be translated into a gravitational potential energy by taking $q \rightarrow m$ and $(4\pi\epsilon_0)^{-1} \rightarrow G$. Hence, we have (for $\langle r \rangle = 1$ fm and $m_u = m_d = m_s = m_\Sigma/3$, where m_Σ is the average mass of the Σ baryons)

$$\boxed{V = -\frac{Gm_\Sigma^2}{3\langle r \rangle} = -6.28 \times 10^{-37} \text{ MeV}}, \quad (2)$$

which applies to all the Σ baryons since we assume each quark flavor has the same mass as the others. As is clearly seen here, this is a tiny binding energy. Additionally, since mass differences come from different binding energies, the gravitational binding energy cannot produce any effect in the approximation we made above. Even if we were more careful

about using exact masses, however, the gravitational binding energy differences would still be orders of magnitude smaller than the MeV scale that we care about.

c) Your answer to (a) may not be accurate, since we really do not know if 1 fm is the right distance. However, the relative binding energies may be more reliable. Use them and the three Σ masses to solve for both the electromagnetic binding and the $d - u$ mass difference.

The masses of the Σ baryons are given as follows:

$$m_{\Sigma^0} = m_u + m_d + m_s - V \quad (3)$$

$$m_{\Sigma^+} = 2m_u + m_s \quad (4)$$

$$m_{\Sigma^-} = 2m_d + m_s + V, \quad (5)$$

where $V = e^2/12\pi\epsilon_0\langle r \rangle$. Since we have 3 equations, we can solve for m_u , m_d , and V in terms of the Σ baryon masses and m_s :

$$m_u = -\frac{m_s - m_{\Sigma^+}}{2} \quad (6)$$

$$m_d = -\frac{m_s}{2} - \frac{m_{\Sigma^+}}{6} + \frac{m_{\Sigma^-}}{3} + \frac{m_{\sigma^0}}{3} \quad (7)$$

$$V = \frac{m_{\Sigma^+} + m_{\Sigma^-} - 2m_{\Sigma^0}}{3}. \quad (8)$$

If we subtract the expression for the mass of the d quark from that for the u quark and plug in the masses of the Σ baryons, we find

$$m_d - m_u = -\frac{2m_{\Sigma^+} - m_{\Sigma^-} - m_{\Sigma^0}}{3} = 3.78 \text{ MeV}$$

$$V = 0.51 \text{ MeV}$$

(9)

Note that this implies an average separation between the quarks of $\langle r \rangle = 0.94$ fm, which is not too far off of the 1 fm scale we assumed in part (a).

d) The $K^0 - K^\pm$ mass difference is 4.0 MeV/c² and the $B^0 - B^\pm$ mass difference is 0.33 MeV/c². Does this suggest to you that the K mesons are larger or smaller in size than the B mesons?

For the K mesons we have

$$m_{K^0} = m_d + m_s + V_K \quad (10)$$

$$m_{K^\pm} = m_u + m_s - 2V_K, \quad (11)$$

where $V = e^2/36\pi\epsilon_0\langle r \rangle_K$, so the Kaon mass difference is given as

$$m_{K^0} - m_{K^\pm} = (m_d - m_u) + 3V_K. \quad (12)$$

For the B mesons we have

$$m_{B^0} = m_d + m_b - V_B \quad (13)$$

$$m_{B^\pm} = m_u + m_b + 2V_B, \quad (14)$$

and as in the case with kaons

$$m_{B^0} - m_{B^\pm} = (m_d - m_u) - 3V_B. \quad (15)$$

Thus, the binding factors $V_K = 0.22 \text{ MeV}$ and $V_B = 3.45 \text{ MeV}$, which would lead to a prediction for the K system to be larger.

- 2)** Classically, light cannot scatter from light, but it is allowed by relativistic quantum mechanics. Called Delbrück scattering, after the physicist who predicted it in 1933, it was first observed in 1953 by R.R. Wilson. Draw the lowest-order Feynman diagram for Delbrück scattering.

We know that photons, which are chargeless, cannot interact directly via the electromagnetic interaction. However, they may produce virtual charged fermion-antifermion pairs (e.g. e^+e^-) that interact with each other, allowing the photons to indirectly interact with each other via the electromagnetic force. The lowest order diagram is shown below.

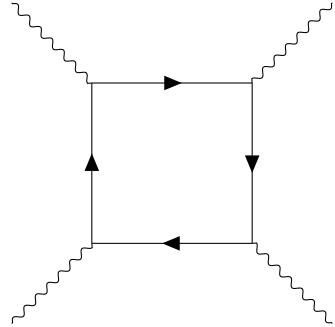


Figure 1: Lowest order feynman diagram for $\gamma\gamma$ interactions. Note that we could draw similar diagrams by interchanging adjacent photon vertices.

- 3)** In e^+e^- annihilation experiments, a second process occurs in which the electron and positron do not annihilate, but produce a fermion-antifermion pair by the exchange of two photons, as shown to the right.

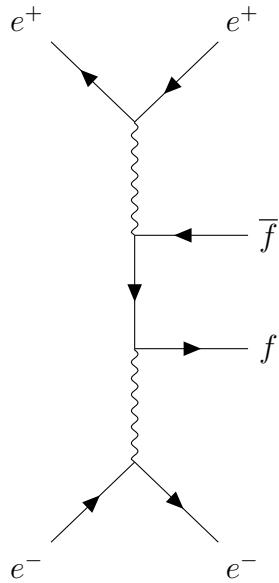


Figure 2: Feynman diagram for $e^+e^- \rightarrow e^+e^-\bar{f}f$ without annihilation of the initial state e^+e^- particles.

(a) For very high energy, and for a very high mass fermion-antifermion pair (these conditions allow you to ignore the mass of the fermions in your calculation), what do you expect for the ratio

$$R = \frac{\sigma(e^+e^- \rightarrow e^+e^- \text{hadrons})}{\sigma(e^+e^- \rightarrow e^+e^-\mu^+\mu^-)}. \quad (16)$$

Use u , d , s , c , and b quarks in your calculation (The world does not have an e^+e^- colliding beam facility that works at high enough energy to include t quarks).

For the reaction $e^+e^- \rightarrow e^+e^-\bar{f}f$, the difference in the squared amplitudes for different fermions f only arises from differences in the vertex factors where f and \bar{f} are produced, which is where the charge of the fermion, e_f , appears. For quarks, we have

$$\sigma(e^+e^- \rightarrow e^+e^- q\bar{q}) \propto 3 \sum_q e_q^4, \quad (17)$$

while for $\mu^+\mu^-$ production

$$\sigma(e^+e^- \rightarrow e^+e^-\mu^+\mu^-) \propto e^4. \quad (18)$$

Hence,

$$R = 3 \sum_q \left(\frac{e_q}{e} \right)^2 = 3 \left[3 \left(-\frac{1}{3} \right)^4 + 2 \left(\frac{2}{3} \right)^4 \right] = \frac{1}{27} (3 + 32) = \frac{35}{27}. \quad (19)$$

- (b) There are several other electromagnetic diagrams of the same order that lead to the identical final state, but would give a different answer than you have just calculated. Draw one of them. (Hint: let the e^+e^- annihilate.) (Note that since these are indistinguishable processes, the amplitudes have to be added before squaring. In most cases, the diagram for part (a) dominates, so your answer in part (a) is close to correct for the stated conditions.)

If we have the electron and positron annihilate then we may have a diagram of the same order as Fig. 2 like the one shown in Fig. 3.

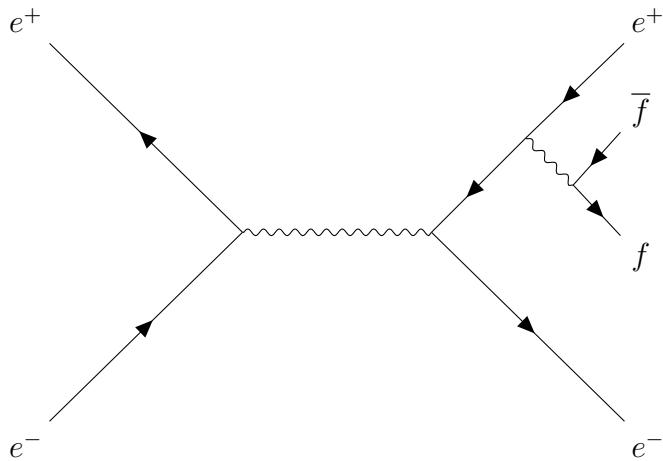


Figure 3: Diagram for $e^+e^- \rightarrow e^+e^- f\bar{f}$ process of the same order as Fig. 2. Note that the labels of f and e may be interchanged in the final state and the virtual photon line could have been drawn coming off of the lower leg of the final state fermion as well.

- 4) Consider the following possible decays and reactions. If the reaction is forbidden, state what conservation law is violated. (Don't forget to check energy conservation for the decays. For the reactions, you can assume that there is enough energy in the initial state to make the reaction go.) If the reaction or decay is allowed, draw a diagram showing the leading mechanism, (i.e. strong, if possible, otherwise electromagnetic, if possible, otherwise weak).

(a) $e^+e^- \rightarrow \tau^+\tau^-$

This is an allowed electromagnetic process.

(b) $e^+e^- \rightarrow \pi^+\pi^-$

This is an allowed process.

(c) $e^-p \rightarrow e^-n\pi^+$

This is an allowed strong process.

(d) $e^- p \rightarrow \nu_e n \pi^0$

This is an allowed weak process.

(e) $\nu_\mu e^- \rightarrow \nu_e \mu^-$

This is an allowed weak process.

(f) $\bar{\nu}_\mu e^- \rightarrow \bar{\nu}_e \mu^-$

This is forbidden by conservation of lepton number.

(g) $\bar{\nu}_e e^- \rightarrow \pi^- \pi^0$

This is an allowed weak process.

(h) $\bar{\nu}_e e^- \rightarrow \nu_e e^+$

This is an allowed weak process.

(i) $\pi^+ n \rightarrow p \gamma$

This is an allowed electromagnetic process.

(j) $e^+ e^- \rightarrow \nu_\mu \bar{\nu}_\mu$

This is an allowed weak process.

(k) $p \bar{p} \rightarrow \Sigma^+ \Sigma^-$

This is an allowed weak process.

(l) $\pi^+ n \rightarrow \Lambda K^+$

This is an allowed strong process.

(m) $\tau^- \rightarrow K^- \nu_\tau$

This is an allowed weak process.

(n) $\Sigma^0 \rightarrow \Lambda \pi^0$

This is forbidden by kinematic considerations.

(o) $\psi \rightarrow \rho^+ \pi^-$

This is an allowed strong process.

(p) $\pi^+ \rightarrow e^+ \nu_e$

This is an allowed weak process.

(q) $\Omega^- \rightarrow \Xi^0 K^-$

This is forbidden by kinematic considerations.

(r) $\Omega^- \rightarrow \Lambda K^-$

This is an allowed strong process.

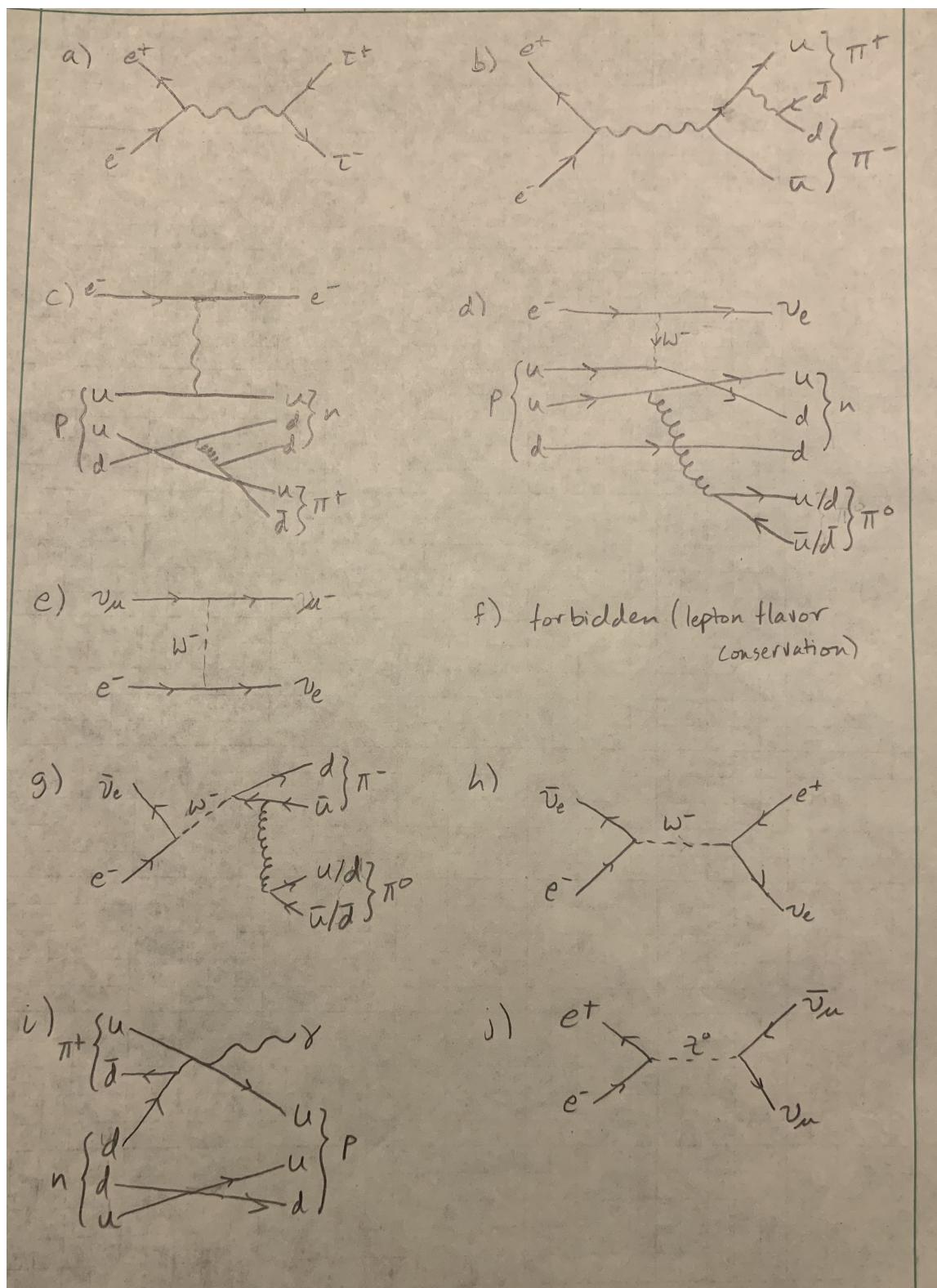
(s) $n \rightarrow e^- \bar{\nu}_e \pi^+$

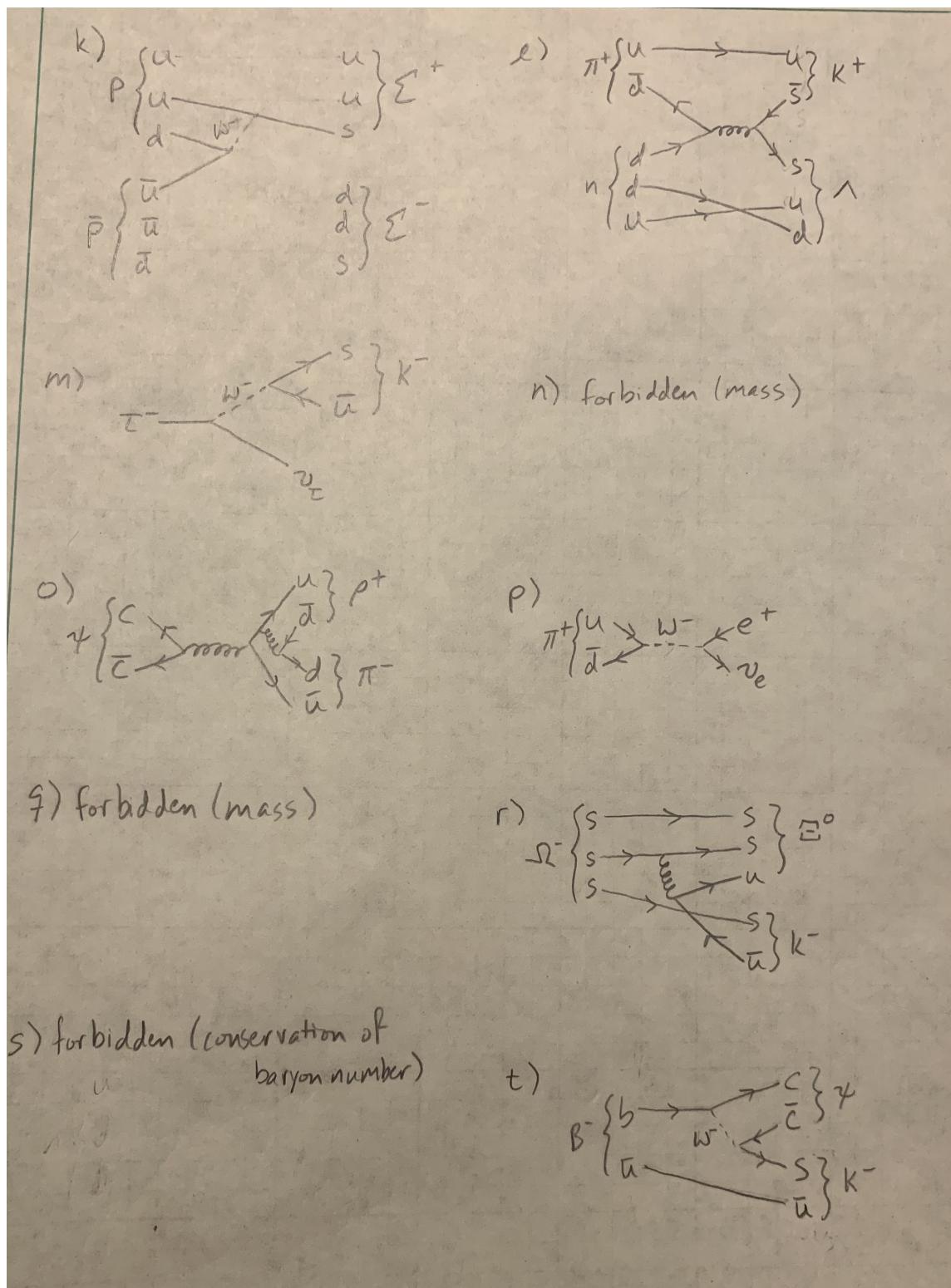
This is forbidden by conservation of baryon number.

(t) $B^- \rightarrow \psi K^-$

This is an allowed weak process.

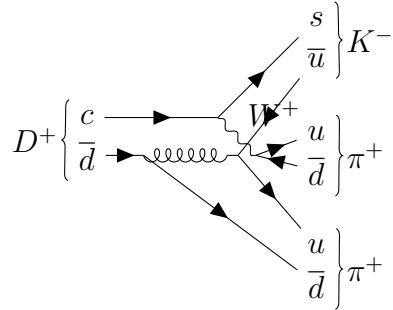
Possible contributing diagrams for each allowed process are shown below:





- 5) The second charmed particle to be discovered, the D^+ , was found first in the decay mode $D^+ \rightarrow K^- \pi^+ \pi^+$.

(a) Draw a diagram for this decay.



- (b) Ignoring its narrow width, the experimenters who discovered this mode knew immediately that it could not come from the strong decay of a meson. Why?

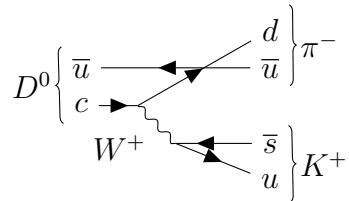
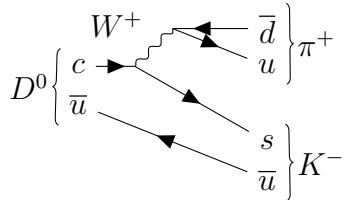
Writing the decay in terms of quark content we have

$$c\bar{d} \rightarrow s\bar{u} + u\bar{d} + u\bar{d}. \quad (20)$$

Obviously, flavor is not conserved in this decay, and since the strong force conserves quark flavor, this decay cannot proceed strongly.

- 6) The decays $D^0 \rightarrow K^- \pi^+$ and $D^0 \rightarrow K^+ \pi^-$ are both allowed. However, the former has a much higher branching ratio than the latter. Draw the diagrams for these two decays and estimate the ratio of branching fractions due to the Cabibo mixing angle factors. Compare your answer to the data summary in the PDG tables.

The first order diagrams are as follows:



The primary difference between the graphs is the vertex where c changes flavor through a weak process. For the generic process we have $c \rightarrow s' = s \cos \theta_C - d \sin \theta_C$, where $\sin^2 \theta_C \approx 0.05$. Thus, the branching ratio is

$$B = \frac{\Gamma(D^0 \rightarrow K^- \pi^+)}{\Gamma(D^0 \rightarrow K^+ \pi^-)} = \frac{\mathcal{P}(c \rightarrow s)}{\mathcal{P}(c \rightarrow d)} = \frac{\cos^2 \theta_c}{\sin^2 \theta_c} \approx 19. \quad (21)$$