

# Mechanics

May 20, 2025

## 1 Basic Elementary Principles

### Result 1 (Newton's law)

1. In an inertial reference frame, an object remains at rest or constant velocity unless acted upon by external force.

2.

$$\dot{\mathbf{p}} = \frac{d(m\mathbf{v})}{dt} = \mathbf{F}^{(e)}$$

3. Action and reaction:  $\mathbf{F}_{ij} = -\mathbf{F}_{ji}$  (Additional condition for strong version:  $\mathbf{r}_{ij} \times \mathbf{F}_{ji} = 0$ )

### 1.1 Single particle

#### Result 2 (Conservation of Linear Momentum)

$$\text{If } \mathbf{F} = 0, \text{ then } \dot{\mathbf{p}} = 0$$

Angular momentum of the particle around point  $O$  is  $\mathbf{L} = \mathbf{r} \times \mathbf{p}$ . Torque  $\mathbf{N} = \mathbf{r} \times \mathbf{F}$ .

$$\mathbf{N} = \frac{d}{dt}(\mathbf{r} \times m\mathbf{v}) = \frac{d\mathbf{L}}{dt} \equiv \dot{\mathbf{L}}$$

#### Result 3 (Conservation of Angular Momentum)

$$\text{If } \mathbf{N} = 0, \text{ then } \dot{\mathbf{L}} = 0$$

If force field  $\mathbf{F}$  is conservative ( $\oint \mathbf{F} \cdot d\mathbf{s} = 0$  or  $\int_a^b \mathbf{F} \cdot d\mathbf{s} = T_b - T_a$  is the same for any path between  $a$  and  $b$ ). There exists a potential scalar field  $V$  such that

$$\mathbf{F} = -\nabla V(\mathbf{r})$$

#### Result 4 (Conservation of Energy)

$E = T + V$  is conserved where  $T$  and  $V$  are kinetic energy and potential energy respectively

## 1.2 System of particles

Center of mass:

$$\mathbf{R} = \frac{\sum_i m_i \mathbf{r}_i}{\sum m_i} = \frac{\sum_i m_i \mathbf{r}_i}{M}$$

Let  $\mathbf{F}_i^{(e)}$  external force acting on particle  $i$ -th, and  $\mathbf{F}_{ji}$  is the force exerted by  $j$ -th particle on  $i$ -th particle in the system,

$$m_i \frac{d^2(\mathbf{r}_i)}{dt^2} = \dot{\mathbf{p}}_i = \sum_j \mathbf{F}_{ji} + \mathbf{F}_i^{(e)}$$

Summing over all particles, we get

$$\dot{\mathbf{p}} = M \frac{d^2 \mathbf{R}}{dt^2} = \sum_i \mathbf{F}_i^{(e)} \equiv \mathbf{F}^{(e)}$$

### Result 5 (Conservation of Linear Momentum for a system of particles)

if the total external force  $\mathbf{F}^{(e)} = 0$ , the total linear momentum  $\mathbf{p}$  is conserved.

The change of total angular momentum:

$$\begin{aligned} \frac{d}{dt} \mathbf{L} &= \sum_i \mathbf{r}_i \times \dot{\mathbf{p}}_i = \sum_i \mathbf{r}_i \times (\mathbf{F}_i^{(e)} + \sum_j \mathbf{F}_{ji}) \\ &= \mathbf{N}^{(e)} + \sum_{i,j} \mathbf{r}_i \times \mathbf{F}_{ji} \\ &= \mathbf{N}^{(e)} + \sum_{i < j} \mathbf{r}_i \times \mathbf{F}_{ji} + \sum_{j < i} \mathbf{r}_i \times \mathbf{F}_{ji} \\ &= \mathbf{N}^{(e)} + \sum_{i < j} \mathbf{r}_i \times \mathbf{F}_{ji} + \sum_{i < j} \mathbf{r}_j \times \mathbf{F}_{ij} \\ &= \mathbf{N}^{(e)} + \sum_{i < j} \mathbf{r}_i \times \mathbf{F}_{ji} - \sum_{i < j} \mathbf{r}_j \times \mathbf{F}_{ji} \\ &= \mathbf{N}^{(e)} + \sum_{i < j} (\mathbf{r}_i - \mathbf{r}_j) \times \mathbf{F}_{ji} \\ &= \mathbf{N}^{(e)} + \sum_{i < j} \mathbf{r}_{ij} \times \mathbf{F}_{ji} \end{aligned}$$

When the internal force abides strong version of Newton's 3rd law, internal force pair has the same direction as the two particles  $\mathbf{r}_{ij}$ . Then  $\mathbf{r}_{ij} \times \mathbf{F}_{ji} = 0$ . Forces satisfy strong Newton's 3rd law are **central**.

### Result 6 (Conservation of Angular Momentum for a system of particles)

$\mathbf{L}$  is constant in time if applied external torque  $\mathbf{N}^{(e)} = 0$  and internal forces are central

Let  $\mathbf{r}_i = \mathbf{R} + \mathbf{r}'_i$  where  $\mathbf{r}'_i$  is the position of  $i$ th particle wrt to  $\mathbf{R}$  instead of  $O$ . Similarly,  $\mathbf{v}_i = \mathbf{v} + \mathbf{v}'_i$  where  $\mathbf{v} = \dot{\mathbf{R}}$ . We can find  $\mathbf{R} = \frac{\sum_i m_i \mathbf{r}_i}{M} = \frac{\sum_i m_i (\mathbf{R} + \mathbf{r}'_i)}{M} = \mathbf{R} + \sum_i \mathbf{r}'_i$  which means

$$\sum_i \mathbf{r}'_i = 0, \quad \sum_i \dot{\mathbf{r}}'_i = \sum_i \mathbf{v}'_i = 0$$

**Result 7** The total angular momentum about  $O$  is the angular momentum of center of mass + angular momentum of each particle around the center of mass.

$$\begin{aligned} \mathbf{L} &= \sum_i \mathbf{r}_i \times \mathbf{p}_i \\ &= \mathbf{R} \times M\mathbf{v} + \sum_i \mathbf{r}'_i \times \mathbf{p}'_i + \sum_i \mathbf{R} \times m_i \mathbf{v}'_i + \sum_i m_i \mathbf{r}'_i \times \mathbf{v} \\ &= \mathbf{R} \times M\mathbf{v} + \sum_i \mathbf{r}'_i \times \mathbf{p}'_i \end{aligned}$$

In general  $\mathbf{L}$  depends on choice of  $O$  through  $\mathbf{R}$  except when the center of mass is at rest  $\mathbf{R} = 0$ .

**Result 8** The total kinetic energy of the system is

$$T = \sum_i \frac{1}{2} m_i v_i^2 = \sum_i \frac{1}{2} m_i (\mathbf{v} + \mathbf{v}'_i) \cdot (\mathbf{v} + \mathbf{v}'_i) = \frac{1}{2} M v^2 + \sum_i \frac{1}{2} m_i (v'_i)^2$$

The  $i$ th particle has trajectory  $\mathbf{r}_i(t)$  from time  $a$  to  $b$ . Then the work is the change in kinetic energy.

$$W_i = \int_a^b \mathbf{F}_i \cdot d\mathbf{r}_i = \int_a^b \frac{d(m_i v_i)}{dt} \cdot \frac{d\mathbf{r}_i}{dt} dt = \int_a^b m_i \dot{\mathbf{v}}_i \cdot \mathbf{v}_i dt = \int_a^b d\left(\frac{1}{2} m_i v_i^2\right) = T_i(b) - T_i(a)$$

Also we can write

$$W_i = \int_a^b \mathbf{F}_i \cdot d\mathbf{r}_i = \int_a^b \mathbf{F}_i^{(e)} \cdot d\mathbf{r}_i + \sum_j \int_a^b \mathbf{F}_{ji} \cdot d\mathbf{r}_i$$

If the external force  $\mathbf{F}_i^{(e)}$  is conservative, then  $\mathbf{F}_i^{(e)} = -\nabla V_i$

$$\int_a^b \mathbf{F}_i^{(e)} \cdot d\mathbf{r}_i = - \int_a^b \nabla V_i \cdot d\mathbf{r}_i = - \int_a^b \left( \frac{\partial V_i}{\partial x_i} dx_i + \frac{\partial V_i}{\partial y_i} dy_i + \frac{\partial V_i}{\partial z_i} dz_i \right) = - \int_a^b dV_i = -V_i|_a^b$$

If the action/reaction forces are conservative and central,  $V_{ji}(|\mathbf{r}_i - \mathbf{r}_j|)$  only depends on distance between  $i$ th and  $j$ th particle, hence it is symmetric  $V_{ji} = V_{ij}$ . Also

$$\mathbf{F}_{ji} = -\nabla_i V_{ij} = -\mathbf{F}_{ij} = \nabla_j V_{ji} = \nabla_j V_{ij}$$

We can find the gradients of  $V_{ij}$  with respect to  $\mathbf{r}_{ij} = \mathbf{r}_i - \mathbf{r}_j$ ,

$$\nabla_i V_{ij} = \frac{\partial V_{ij}}{\partial \mathbf{r}_i} = \frac{\partial V_{ij}}{\partial \mathbf{r}_{ij}} \frac{\partial \mathbf{r}_{ij}}{\partial \mathbf{r}_i} = \frac{\partial V_{ij}}{\partial \mathbf{r}_{ij}} = \nabla_{ij} V_{ij} = -\nabla_j V_{ij}$$

The work done by internal forces can be written in sums of pairs of potentials between  $i$ th and  $j$ th particle.

$$\begin{aligned}
\sum_{i,j} \int_a^b \mathbf{F}_{ji} \cdot d\mathbf{r}_i &= - \sum_{i,j} \int_a^b \nabla_j V_{ji} \cdot d\mathbf{r}_i \\
&= - \sum_{i < j} \int_a^b \nabla_j V_{ji} \cdot d\mathbf{r}_i + \nabla_i V_{ij} \cdot d\mathbf{r}_j \\
&= - \sum_{i < j} \int_a^b \nabla_j V_{ij} \cdot d\mathbf{r}_i - \nabla_j V_{ij} \cdot d\mathbf{r}_j \\
&= - \sum_{i < j} \int_a^b \nabla_j V_{ij} d(\mathbf{r}_i - \mathbf{r}_j) \\
&= - \sum_{i < j} \int_a^b \nabla_j V_{ij} d(\mathbf{r}_{ij}) \\
&= \sum_{i < j} \int_a^b \nabla_{ij} V_{ij} d(\mathbf{r}_{ij}) \quad (\text{note that } \nabla_{ij} V_{ij} = -\nabla_j V_{ij}) \\
&= \frac{1}{2} \sum_{i,j} \int_a^b \nabla_{ij} V_{ij} d(\mathbf{r}_{ij}) \\
&= \frac{1}{2} \sum_{i,j} V_{ij} \Big|_a^b
\end{aligned}$$

**Result 9** Total potential energy  $V$  when external and internal forces are conservative and internal forces also central is:

$$V = \sum_i V_i + \frac{1}{2} \sum_{i,j} V_{ij}$$

Such that total energy  $T + V$  is conserved

On a rigid body,  $|\mathbf{r}_{ij}|$  is constant between any two points. Then  $0 = d(r_{ij}^2) = d(\mathbf{r}_{ij} \cdot \mathbf{r}_{ij}) = 2\mathbf{r}_{ij} \cdot d(\mathbf{r}_{ij})$ . Therefore  $d(\mathbf{r}_{ij})$  is perpendicular to  $\mathbf{r}_{ij}$  which means the work done by internal force  $\mathbf{F}_{ji} \cdot d(\mathbf{r}_{ij}) = 0$