

Coefficients for higher Landau levels

Fourier Transform of $\frac{1}{r}$:

Assuming $\left[r \in \text{Reals} \ \&\& \ r > 0 \ \&\& \ k \in \text{Reals} \ \&\& \ k > 0 \ \&\& \ d \in \text{Reals} \ \&\& \ d > 0, \frac{1}{2\pi} \int_0^{2\pi} \frac{1}{r} e^{-i(kr \cos[\theta])} r \, d\theta \right]$

$\text{BesselJ}[0, kr]$

Assuming $\left[r \in \text{Reals} \ \&\& \ r > 0 \ \&\& \ k \in \text{Reals} \ \&\& \ k > 0 \ \&\& \ d \in \text{Reals} \ \&\& \ d > 0, \int_0^\infty \text{BesselJ}[0, kr] \, dr \right]$

$\frac{1}{k}$

r has units of length, while k has units of 1/length.

Haldane Pseudopotential

Use Haldane pseudopotential. Jain's E3.28 does not have l^3 and omits e , the electron charge. We can view q as the unitless $q = kl$. We expect the energy to have units of $\frac{e^2}{l}$. See Jain problem 3.23--`Vm_Coulomb.nb` for an example that shows our k, l placements are correct.

$$\begin{aligned} V_m^{(n)} &= \frac{e^2}{(2\pi)} \int V(q) e^{-(q)^2} L_m((q)^2) \left(L_n\left(\frac{(q)^2}{2}\right) \right)^2 d^2 \vec{q} \\ &= \frac{e^2}{(2\pi)} \int V(k) e^{-(kl)^2} L_m((kl)^2) \left(L_n\left(\frac{(kl)^2}{2}\right) \right)^2 d^2 \vec{k} \end{aligned}$$

Try m=3, n=0:

Assuming $\left[k \in \text{Reals} \ \&\& \ k > 0 \ \&\& \ d \in \text{Reals} \ \&\& \ d \geq 0 \ \&\& \ l \in \text{Reals} \ \&\& \ l > 0, \right.$

$$\frac{1}{(2\pi)} \int_0^{2\pi} \int_0^\infty \frac{1}{k} \text{Exp}[-(kl)^2] \text{LaguerreL}[3, (kl)^2] \left(\text{LaguerreL}\left[0, \frac{(kl)^2}{2}\right] \right)^2 k \, dk \, d\theta \quad // \text{ Simplify}$$

$$\frac{5\sqrt{\pi}}{32l}$$

Here it is apparent that we have the correct units of 1/length--specifically, $1/l$.

General equation:

$\text{Vmn}[m_ , n_] := \text{Assuming}\left[k \in \text{Reals} \ \&\& \ k > 0 \ \&\& \ d \in \text{Reals} \ \&\& \ d \geq 0 \ \&\& \ l \in \text{Reals} \ \&\& \ l > 0, \right.$

$$\frac{1}{(2\pi)} \int_0^{2\pi} \int_0^\infty \frac{1}{k} \text{Exp}[-(kl)^2] \text{LaguerreL}[m, (kl)^2] \left(\text{LaguerreL}\left[n, \frac{(kl)^2}{2}\right] \right)^2 k \, dk \, d\theta \quad // \text{ Simplify}$$

Find c_i for n .

We find the constants c_i of the effective potential by equating the Haldane pseudopotential of $V_{\text{eff}}(k) = \frac{1}{k} + \frac{1}{l} \sum_{i=0}^{\text{imax}} c_i f_i(k)$ to the Haldane pseudopotential of $V_{\text{ZDS}}(k)$.

Reference: Frame potential--Haldane pseudopotential (9-19-13).nb”

The effective potential is:

$$V(r)^{\text{eff}} = \frac{1}{r} + \frac{1}{l} \sum_{i=0}^{\text{imax}} c_i \left(\frac{r}{l}\right)^i e^{-\left(\frac{r}{l}\right)^2}. \quad \text{The placement of the } l\text{'s ensures that the second part has units of 1/length.}$$

Take the 2-D Fourier transform using the formula (see 2-D Fourier transform integral form.nb,” Higher Mathematics for Physics and Engineering” by Shima, Nakayama and p. 690, 680; Arfken and Weber)

$$F(\vec{k}) = \int_0^\infty r f(r) J_0(k r) dr$$

so that:

$$\begin{aligned} V(k) &= \int_0^\infty r V(r) J_0(k r) dr \\ &= \int_0^\infty r \left(\frac{1}{r} + \sum_{i=0}^{\text{imax}} c_i r^i e^{-r^2} \right) J_0(k r) dr \\ &= \int_0^\infty J_0(k r) dr + \sum_{i=0}^{\text{imax}} c_i \int_0^\infty r^{i+1} e^{-r^2} J_0(k r) dr. \end{aligned}$$

The Fourier transform of $\frac{1}{r}$ is $\frac{1}{k}$:

$$\text{Assuming}\left[k \in \text{Reals} \ \&\& \ k > 0, \int_0^\infty \text{BesselJ}[0, k r] dr\right]$$

$$\frac{1}{k}$$

The Fourier transform of the other terms:

Method 1:

$$\text{Assuming}\left[r \in \text{Reals} \ \&\& \ r > 0 \ \&\& \ k \in \text{Reals} \ \&\& \ k > 0 \ \&\& \ l \in \text{Reals} \ \&\& \ l > 0, \frac{1}{2\pi} \int_0^{2\pi} \frac{1}{1} \left(\frac{r}{1}\right)^i \text{Exp}\left[-\left(\frac{r}{1}\right)^2\right] e^{-i(k r \cos[\theta])} r d\theta\right]$$

$$e^{-\frac{r^2}{1^2}} \left(\frac{r}{1}\right)^{1+i} \text{BesselJ}[0, k r]$$

$$\text{Assuming}\left[r \in \text{Reals} \ \&\& \ r > 0 \ \&\& \ k \in \text{Reals} \ \&\& \ k > 0 \ \&\& \ l \in \text{Reals} \ \&\& \ l > 0, \int_0^\infty e^{-\frac{r^2}{1^2}} \left(\frac{r}{1}\right)^{1+i} \text{BesselJ}[0, k r] dr\right]$$

$$\text{ConditionalExpression}\left[\frac{1}{2} \Gamma\left[1 + \frac{i}{2}\right] \text{LaguerreL}\left[-1 - \frac{i}{2}, -\frac{1}{4} k^2 l^2\right], \text{Re}[i] > -2\right]$$

Method 2:

`Assuming` $\left[k \in \text{Reals} \ \&\& \ k > 0 \ \&\& \ i \in \text{Integers} \ \&\& \ i \geq 0 \ \&\& \ l \in \text{Reals} \ \&\& \ l > 0, \int_0^\infty \frac{1}{1} \left(\frac{r}{1}\right)^i \text{Exp}\left[-\left(\frac{r}{1}\right)^2\right] \text{BesselJ}[0, k r] r \, dr\right]$

$\frac{1}{2} l \, \Gamma\left[1 + \frac{i}{2}\right] \text{LaguerreL}\left[-1 - \frac{i}{2}, -\frac{1}{4} k^2 l^2\right]$

`Vkiterm` $[i_]:= \frac{1}{2} l \, \Gamma\left[1 + \frac{i}{2}\right] \text{LaguerreL}\left[-1 - \frac{i}{2}, -\frac{1}{4} k^2 l^2\right];$

Example:

`Vkiterm` $[1] \text{ // FullSimplify}$

$\frac{1}{4} l \sqrt{\pi} \text{LaguerreL}\left[-\frac{3}{2}, -\frac{1}{4} k^2 l^2\right]$

The Psuedopotential energies

$\frac{1}{k}$ term:

`Vmnterml` $[m_, n_] := \text{Assuming}\left[k \in \text{Reals} \ \&\& \ k > 0 \ \&\& \ m \in \text{Integers} \ \&\& \ m \geq 0 \ \&\& \ n \in \text{Integers} \ \&\& \ n \geq 0 \ \&\& \ l \in \text{Reals} \ \&\& \ l > 0,$

$\frac{1}{(2 \pi)} \int_0^{2 \pi} \int_0^\infty \frac{1}{k} \text{Exp}\left[-(k l)^2\right] \text{LaguerreL}[m, (k l)^2] \left(\text{LaguerreL}\left[n, \frac{(k l)^2}{2}\right]\right)^2 k \, dk \, d\theta\right] \text{ // Simplify}$

`Vmnterml` $[2, 2]$

$\frac{3451 \sqrt{\pi}}{16384 l}$

$i=0$ to i_{max} terms:

```
Vmniterm[m_, n_, i_] := Assuming[k ∈ Reals && k > 0 && m ∈ Integers && m ≥ 0 && n ∈ Integers && n ≥ 0 && i ∈ Integers && i ≥ 0,  $\frac{1}{(2\pi)}$ 

$$\int_0^{2\pi} \int_0^\infty \frac{1}{2} \Gamma\left[1 + \frac{i}{2}\right] \text{LaguerreL}\left[-1 - \frac{i}{2}, -\frac{1}{4} k^2 l^2\right] \text{Exp}\left[-(k l)^2\right] \text{LaguerreL}\left[m, (k l)^2\right] \left(\text{LaguerreL}\left[n, \frac{(k l)^2}{2}\right]\right)^2 k \, dk \, d\theta \, // \text{Simplify}$$

```

Example:

```
Vmniterm[2, 1, 3]
```

$$\frac{9159 \sqrt{\frac{\pi}{5}}}{2000001}$$

```
(* Numerical integration: Vmniterm[m_,n_,i_] := Assuming[k ∈ Reals && k > 0 && m ∈ Integers && m ≥ 0 && n ∈ Integers && n ≥ 0 && i ∈ Integers && i ≥ 0,  $\frac{1}{(2\pi)}$ 
NIntegrate[ $\frac{1}{2} \Gamma\left[1 + \frac{i}{2}\right] \text{LaguerreL}\left[-1 - \frac{i}{2}, -\frac{k^2}{4}\right] \text{Exp}\left[-(k)^2\right] \text{LaguerreL}\left[m, (k)^2\right] \left(\text{LaguerreL}\left[n, \frac{(k)^2}{2}\right]\right)^2 k, \{k, 0, \infty\}, \{\theta, 0, 2\pi\}] \, // \text{Simplify*}$ 
```

Find the effective potential for the $n=0$ level that corresponds to a $V(k)$ at a higher n :

$$V_m^0\left(\frac{1}{k} + \sum_{i=0}^4 c_i f_i(k)\right) = V_m^n(V(k)).$$

We let m run from 0 to 4 so that there will be 5 equations with 5 unknowns (c_0 to c_4):

$$V_0^0\left(\frac{1}{k} + \sum_{i=0}^4 c_i f_i(k)\right) = V_0^n(V(k))$$

$$\vdots$$

$$V_4^0\left(\frac{1}{k} + \sum_{i=0}^4 c_i f_i(k)\right) = V_4^n(V(k)).$$

We can thus set up a matrix to solve for the c_i . The classical term, i.e. the first term in the effective potential, is moved to the right side so that the left side is composed entirely of the c_i terms.

$$c_0 V_0^0(f_0(k)) + c_1 V_0^0(f_1(k)) \dots c_4 V_0^0(f_4(k)) = V_0^n(V(k)) - V_0^0\left(\frac{1}{k}\right)$$

$$\vdots$$

$$c_0 V_4^0(f_0(k)) + c_1 V_4^0(f_1(k)) \dots c_4 V_4^0(f_4(k)) = V_4^n(V(k)) - V_4^0\left(\frac{1}{k}\right).$$

Now this can be written in the matrix form:

$$A \vec{c} = \vec{v}$$

where:

$$A = \begin{pmatrix} V_0^0(f_0(k)) & \cdots & \cdots & V_0^0(f_4(k)) \\ \vdots & \vdots & \vdots & \vdots \\ \vdots & \vdots & \vdots & \vdots \\ V_4^0(f_0(k)) & \cdots & \cdots & V_4^0(f_4(k)) \end{pmatrix}$$

$$\vec{c} = \begin{pmatrix} c_0 \\ \vdots \\ c_4 \end{pmatrix}$$

$$\vec{v} = \begin{pmatrix} V_0^n(V(k)) - V_0^0\left(\frac{1}{k}\right) \\ \vdots \\ V_4^n(V(k)) - V_4^0\left(\frac{1}{k}\right) \end{pmatrix}.$$

Solve for \vec{c} :

See Solving "Linear Systems" in the *Mathematica* tutorial.

```
A[imax_] := Parallelize[Table[Vmnterm[m, 0, i], {m, 0, imax}, {i, 0, imax}]]
```

```
A[4]
```

$$\left\{ \left\{ \frac{1}{5 \cdot 1}, \frac{\sqrt{\frac{\pi}{5}}}{5 \cdot 1}, \frac{4}{25 \cdot 1}, \frac{6 \sqrt{\frac{\pi}{5}}}{25 \cdot 1}, \frac{32}{125 \cdot 1} \right\}, \left\{ \frac{1}{25 \cdot 1}, \frac{3 \sqrt{\frac{\pi}{5}}}{50 \cdot 1}, \frac{8}{125 \cdot 1}, \frac{3 \sqrt{\frac{\pi}{5}}}{25 \cdot 1}, \frac{96}{625 \cdot 1} \right\}, \left\{ \frac{1}{125 \cdot 1}, \frac{3 \sqrt{\frac{\pi}{5}}}{200 \cdot 1}, \frac{12}{625 \cdot 1}, \frac{21 \sqrt{\frac{\pi}{5}}}{500 \cdot 1}, \frac{192}{3125 \cdot 1} \right\}, \right. \\ \left. \left\{ \frac{1}{625 \cdot 1}, \frac{7 \sqrt{\frac{\pi}{5}}}{2000 \cdot 1}, \frac{16}{3125 \cdot 1}, \frac{63 \sqrt{\frac{\pi}{5}}}{5000 \cdot 1}, \frac{64}{3125 \cdot 1} \right\}, \left\{ \frac{1}{3125 \cdot 1}, \frac{63 \sqrt{\frac{\pi}{5}}}{80000 \cdot 1}, \frac{4}{3125 \cdot 1}, \frac{693 \sqrt{\frac{\pi}{5}}}{200000 \cdot 1}, \frac{96}{15625 \cdot 1} \right\} \right\}$$

Calculation of v for various d

```
v[imax_, n_] := Parallelize[Table[Vmn[m, n] - Vmnterm1[m, 0], {m, 0, imax}]]
```

v[4, 1]

$$\left\{ -\frac{5\sqrt{\pi}}{32\,1}, -\frac{\sqrt{\pi}}{64\,1}, \frac{17\sqrt{\pi}}{256\,1}, \frac{11\sqrt{\pi}}{512\,1}, \frac{49\sqrt{\pi}}{4096\,1} \right\}$$

v[4, 2]

$$\left\{ -\frac{439\sqrt{\pi}}{2048\,1}, -\frac{191\sqrt{\pi}}{4096\,1}, \frac{379\sqrt{\pi}}{16384\,1}, \frac{709\sqrt{\pi}}{32768\,1}, \frac{14075\sqrt{\pi}}{262144\,1} \right\}$$

Calculation of coefficients

```
c[imax_, n_] := LinearSolve[A[imax] /. 1 -> 1, v[imax, n] /. 1 -> 1]
```

```
d[imax_, n_] := LinearSolve[A[imax], v[imax, n]]
```

the d function is the same because the 1's cancel out in the calculations

d[4, 0]

$$\{0, 0, 0, 0, 0\}$$

c[4, 0]

$$\{0, 0, 0, 0, 0\}$$

As expected.

imax=4, n=1:

d[4, 1]

$$\left\{ \frac{8125\sqrt{\pi}}{16}, -\frac{30525\sqrt{5}}{16}, \frac{108125\sqrt{\pi}}{32}, -\frac{91625\sqrt{5}}{64}, \frac{165625\sqrt{\pi}}{512} \right\}$$

c[4, 1]

$$\left\{ \frac{8125\sqrt{\pi}}{16}, -\frac{30525\sqrt{5}}{16}, \frac{108125\sqrt{\pi}}{32}, -\frac{91625\sqrt{5}}{64}, \frac{165625\sqrt{\pi}}{512} \right\}$$

$$\left\{ \frac{8125 \sqrt{\pi}}{16}, -\frac{30525 \sqrt{5}}{16}, \frac{108125 \sqrt{\pi}}{32}, -\frac{91625 \sqrt{5}}{64}, \frac{165625 \sqrt{\pi}}{512} \right\} // N$$

{900.074, -4266., 5988.96, -3201.25, 573.365}

cn1 = {900.0742211629573`, -4265.998438323818`, 5988.955394661216`, -3201.245756850285`, 573.3645880004416`};

c[8, 1]

Parallelize::nopar1:

LinearSolve[A[8] /. l → 1, v[8, 1] /. l → 1] cannot be parallelized; proceeding with sequential evaluation. >>

$$\left\{ \frac{9090375 \sqrt{\pi}}{32}, -\frac{18412775 \sqrt{5}}{8}, \frac{326664375 \sqrt{\pi}}{32}, -\frac{2651522875 \sqrt{5}}{192}, \right. \\ \left. \frac{8413115625 \sqrt{\pi}}{512}, -\frac{1367618125 \sqrt{5}}{192}, \frac{8689446875 \sqrt{\pi}}{3072}, -\frac{72153125 \sqrt{5}}{192}, \frac{6425609375 \sqrt{\pi}}{196608} \right\}$$

$$\left\{ \frac{9090375 \sqrt{\pi}}{32}, -\frac{18412775 \sqrt{5}}{8}, \frac{326664375 \sqrt{\pi}}{32}, -\frac{2651522875 \sqrt{5}}{192}, \right. \\ \left. \frac{8413115625 \sqrt{\pi}}{512}, -\frac{1367618125 \sqrt{5}}{192}, \frac{8689446875 \sqrt{\pi}}{3072}, -\frac{72153125 \sqrt{5}}{192}, \frac{6425609375 \sqrt{\pi}}{196608} \right\} // N$$

{503508., -5.14653 × 10⁶, 1.80937 × 10⁷, -3.08801 × 10⁷, 2.91247 × 10⁷, -1.59275 × 10⁷, 5.01356 × 10⁶, -840309., 57927.9}

imax=4, n=2:

c[4, 2]

$$\left\{ \frac{4438075 \sqrt{\pi}}{2048}, -\frac{1110375 \sqrt{5}}{128}, \frac{66728375 \sqrt{\pi}}{4096}, -\frac{29742325 \sqrt{5}}{4096}, \frac{111994375 \sqrt{\pi}}{65536} \right\}$$

$$\left\{ \frac{4438075 \sqrt{\pi}}{2048}, -\frac{1110375 \sqrt{5}}{128}, \frac{66728375 \sqrt{\pi}}{4096}, -\frac{29742325 \sqrt{5}}{4096}, \frac{111994375 \sqrt{\pi}}{65536} \right\} // N$$

{3840.96, -19397.5, 28875.2, -16236.8, 3028.94}

cn2 = {3840.9585568151842`, -19397.45297278382`, 28875.235652689782`, -16236.782350803573`, 3028.94380567179`};

c[4, 2.5] // N

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After n=2, the correlation energies stop becoming more positive and jump down below the n=0 level. The curves that are greater than n=3 will be below n=3.

c[4, 3] // N

{2461.9, -12412.9, 18439.9, -10345.8, 1925.12}

c[4, 4] // N

{1760.71, -8841.47, 13072.1, -7296.46, 1349.96}

c[4, 5] // N

{1282.66, -6399.65, 9391.39, -5199.41, 953.232}

cn5 = {1282.658703949136`, -6399.65137024662`, 9391.394237810113`, -5199.410444225987`, 953.2319876007417`};

c[4, 6] // N

{921.69, -4552.76, 6602.8, -3607.94, 651.638}

cn6 = {921.6895110934634`, -4552.763367933509`, 6602.796406745576`, -3607.9364288665274`, 651.6380762878204`};

c[4, 7]

$\left\{ \frac{6141729203475\sqrt{\pi}}{17179869184}, -\frac{94573751771375\sqrt{5}}{68719476736}, \frac{42382398924125\sqrt{\pi}}{17179869184}, -\frac{286886752192925\sqrt{5}}{274877906944}, \frac{127140398570625\sqrt{\pi}}{549755813888} \right\}$

$\left\{ \frac{6141729203475\sqrt{\pi}}{17179869184}, -\frac{94573751771375\sqrt{5}}{68719476736}, \frac{42382398924125\sqrt{\pi}}{17179869184}, -\frac{286886752192925\sqrt{5}}{274877906944}, \frac{127140398570625\sqrt{\pi}}{549755813888} \right\} // N$

{633.645, -3077.34, 4372.61, -2333.76, 409.91}

cn7 = {633.6446140146658`, -3077.3420854236706`, 4372.6087421923385`, -2333.7571465072588`, 409.9101516697673`};

c[4, 10]

$$\left\{ \frac{1\,523\,661\,083\,467\,558\,675\,\sqrt{\pi}}{144\,115\,188\,075\,855\,872}, \frac{38\,489\,450\,927\,660\,925\,\sqrt{5}}{1\,125\,899\,906\,842\,624}, \right. \\ \left. - \frac{65\,148\,254\,720\,069\,204\,625\,\sqrt{\pi}}{288\,230\,376\,151\,711\,744}, \frac{51\,144\,259\,670\,548\,626\,075\,\sqrt{5}}{288\,230\,376\,151\,711\,744}, - \frac{282\,946\,662\,333\,618\,450\,625\,\sqrt{\pi}}{4\,611\,686\,018\,427\,387\,904} \right\} \\ \left\{ \frac{1\,523\,661\,083\,467\,558\,675\,\sqrt{\pi}}{144\,115\,188\,075\,855\,872}, \frac{38\,489\,450\,927\,660\,925\,\sqrt{5}}{1\,125\,899\,906\,842\,624}, \right. \\ \left. - \frac{65\,148\,254\,720\,069\,204\,625\,\sqrt{\pi}}{288\,230\,376\,151\,711\,744}, \frac{51\,144\,259\,670\,548\,626\,075\,\sqrt{5}}{288\,230\,376\,151\,711\,744}, - \frac{282\,946\,662\,333\,618\,450\,625\,\sqrt{\pi}}{4\,611\,686\,018\,427\,387\,904} \right\} // N \\ \{18.7393, 76.4411, -400.625, 396.773, -108.748\}$$

cn10 = {18.739308402702555`, 76.44110117412255`, -400.62493238943216`, 396.7730355458878`, -108.74762489253389`};

c[4, 25]

$$\left\{ - \frac{16\,189\,982\,079\,673\,175\,202\,579\,413\,213\,454\,698\,429\,626\,515\,375\,\sqrt{\pi}}{22\,300\,745\,198\,530\,623\,141\,535\,718\,272\,648\,361\,505\,980\,416}, \right. \\ \frac{67\,681\,223\,489\,106\,848\,042\,628\,576\,973\,526\,535\,486\,264\,598\,875\,\sqrt{5}}{22\,300\,745\,198\,530\,623\,141\,535\,718\,272\,648\,361\,505\,980\,416}, - \frac{266\,143\,232\,952\,525\,608\,835\,049\,441\,216\,548\,388\,920\,580\,229\,375\,\sqrt{\pi}}{44\,601\,490\,397\,061\,246\,283\,071\,436\,545\,296\,723\,011\,960\,832}, \\ \left. \frac{248\,501\,861\,250\,419\,059\,726\,364\,692\,228\,031\,259\,753\,990\,056\,375\,\sqrt{5}}{89\,202\,980\,794\,122\,492\,566\,142\,873\,090\,593\,446\,023\,921\,664}, - \frac{490\,677\,959\,795\,921\,413\,375\,927\,980\,179\,111\,208\,770\,785\,971\,875\,\sqrt{\pi}}{713\,623\,846\,352\,979\,940\,529\,142\,984\,724\,747\,568\,191\,373\,312} \right\} \\ \left\{ - \frac{16\,189\,982\,079\,673\,175\,202\,579\,413\,213\,454\,698\,429\,626\,515\,375\,\sqrt{\pi}}{22\,300\,745\,198\,530\,623\,141\,535\,718\,272\,648\,361\,505\,980\,416}, \right. \\ \frac{67\,681\,223\,489\,106\,848\,042\,628\,576\,973\,526\,535\,486\,264\,598\,875\,\sqrt{5}}{22\,300\,745\,198\,530\,623\,141\,535\,718\,272\,648\,361\,505\,980\,416}, - \frac{266\,143\,232\,952\,525\,608\,835\,049\,441\,216\,548\,388\,920\,580\,229\,375\,\sqrt{\pi}}{44\,601\,490\,397\,061\,246\,283\,071\,436\,545\,296\,723\,011\,960\,832}, \\ \left. \frac{248\,501\,861\,250\,419\,059\,726\,364\,692\,228\,031\,259\,753\,990\,056\,375\,\sqrt{5}}{89\,202\,980\,794\,122\,492\,566\,142\,873\,090\,593\,446\,023\,921\,664}, - \frac{490\,677\,959\,795\,921\,413\,375\,927\,980\,179\,111\,208\,770\,785\,971\,875\,\sqrt{\pi}}{713\,623\,846\,352\,979\,940\,529\,142\,984\,724\,747\,568\,191\,373\,312} \right\} // N \\ \{-1286.77, 6786.31, -10576.5, 6229.24, -1218.71\}$$

cn25 = {-1286.7729677974512`, 6786.31207947123`, -10576.476120856927`, 6229.2431188327555`, -1218.7149348209432`};

```
c[4, 35]
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```
c[4, 50]
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50 is too high! Didn't calculate even after ~2hr 20min.

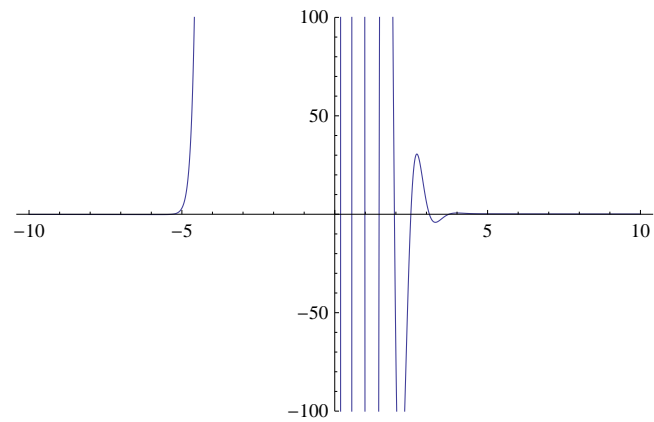
We see that as imax increases, so do the coefficients.

Plot of $\text{imax}=8$, $n=1$:

```
Plot[ $\frac{1}{r} + e^{-r^2} (503508.4429664134 - 5.1465270693010865 r +$   

 $1.80936727944498 r^2 - 3.0880132252060823 r^3 + 2.912472497586839 r^4 - 1.592753695187919 r^5 +$   

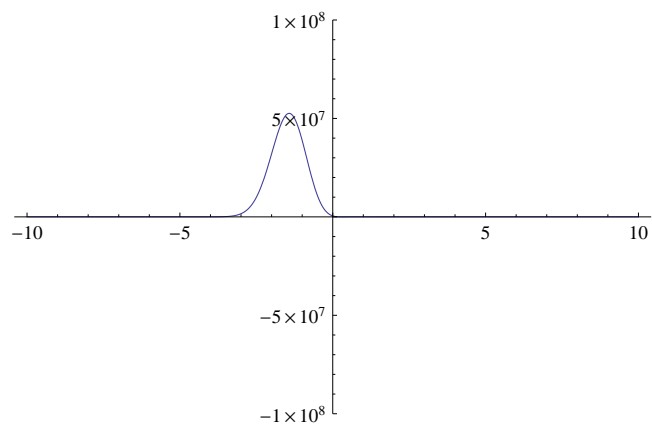
 $5.013555851508024 r^6 - 840308.8140054143 r^7 + 57927.93823818632 r^8)$ , {r, -10, 10}, PlotRange -> 100]
```



```
Plot[ $\frac{1}{r} + e^{-x^2} (503508.4429664134 - 5.1465270693010865 r +$   

 $1.80936727944498 r^2 - 3.0880132252060823 r^3 + 2.912472497586839 r^4 - 1.592753695187919 r^5 +$   

 $5.013555851508024 r^6 - 840308.8140054143 r^7 + 57927.93823818632 r^8)$ , {r, -10, 10}, PlotRange → 100 000 000]
```



Plot of imax=4, n=1:

```
Plot[ $\frac{1}{r} + e^{-x^2} (900.0742211629573 - 4265.998438323818 r + 5988.955394661216 r^2 - 3201.245756850285 r^3 + 573.3645880004416 r^4)$ ,  

{r, -10, 10}, PlotRange → 100]
```

