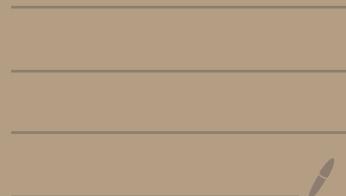


Lecture 09

- Asymptotic analysis of
Jeffrey-Klaumel ODE
- N-S BL



Office hours today → Zoom
 ↳ look canvas and Box

Recall Jeffrey - Hammel ODE

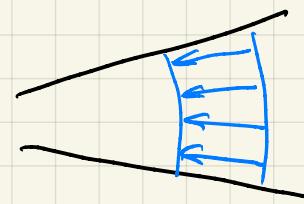
$$f''' + 2\alpha \text{Re} ff' + 4\alpha^2 f' = 0$$

$$f(\theta) = \frac{U(\theta)}{U_{\max}}$$

$$\gamma = \frac{\theta}{x}$$

$$f(\gamma)$$

recall α -small
 $\alpha |\text{Re}|$ - large "inviscid"



- great but no BL

Asymptotic → zoom in on

region of interest

make $\xi = O(1)$
 ξ in region of interest

$$\xi = [(1-\gamma) \sqrt{-\alpha \text{Re}}]^{-1} \quad \xi = 0 \text{ on wall} \quad \xi = A$$

we "think" this is the needed zoom factor

Change variable

$$f(\eta) = f(\xi)$$

$$f'(\eta) = \frac{df}{d\eta} = \frac{df}{d\xi} \frac{d\xi}{d\eta} = -f' A$$

$$f''(\eta) = \frac{d^2f}{d\eta^2} = \frac{d^2f}{d\xi^2} \left(\frac{d\xi}{d\eta} \right)^2 + \frac{df}{d\xi} \frac{d^2\xi}{d\eta^2} = A^2 f''$$

$$f''' = -A^3 f'''$$

$$f''' + 2\alpha Re ff' + 4\alpha^2 f' = 0$$

$$-(\sqrt{-\alpha Re})^2$$

$$-A^2$$

$$-A^3 f''' + 2A^2 A f' f - 4\alpha^2 A f' = 0$$

A large

$$f''' - 2f' f = 0$$

A, Re gone

$A \rightarrow \infty$ "mode 1"

$$f(0) = 0$$

$$f(\xi \rightarrow \infty) = 1$$

$$\left. \begin{array}{l} f'(\xi \rightarrow \infty) = 0 \\ \text{matching} \\ \text{"uniform flow"} \end{array} \right\}$$

time + just trying it

$$f(\xi) = 3 \tanh^{-1} \left[\frac{\xi}{\sqrt{2}} - \tanh^{-1} \sqrt{\frac{2}{3}} \right] - 2$$

$$f(\gamma) = 3 \tanh^{-1} \left[(-\frac{1}{2} \alpha Re)^{1/2} (1-\gamma) - \tanh^{-1} \sqrt{\frac{2}{3}} \right] - 2$$

Same, by a procedure
on gov. eq.

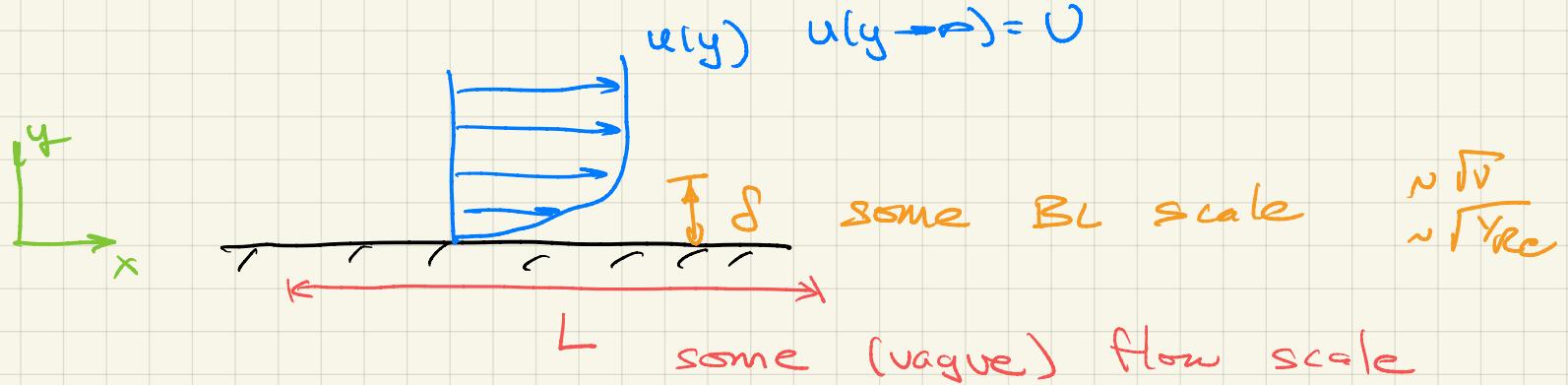
Navier - Stoke BL Hypothesis

Backer

Boundary Layer - a thin layer in which viscosity
is important no matter how high Re

- self-consistent
- seems to work in cases
- no proof
- not always applicable
- one guidelines, no specific class of applicability





expect :

$$\left| \frac{\partial u}{\partial x} \right| \ll \left| \frac{\partial u}{\partial y} \right| \quad \text{and} \quad \left| \frac{\partial^2 u}{\partial x^2} \right| \ll \left| \frac{\partial^2 u}{\partial y^2} \right|$$

→ can reason in advance

x-mom

$$\frac{\partial u}{\partial t} + u \frac{\partial u}{\partial x} + v \frac{\partial u}{\partial y} = -\frac{1}{\rho} \frac{\partial p}{\partial x} + \nu \left(\frac{\partial^2 u}{\partial x^2} + \frac{\partial^2 u}{\partial y^2} \right)$$

small?

large

0

mass

$$\frac{\partial u}{\partial x} + v \frac{\partial v}{\partial y} = 0$$

small?

large

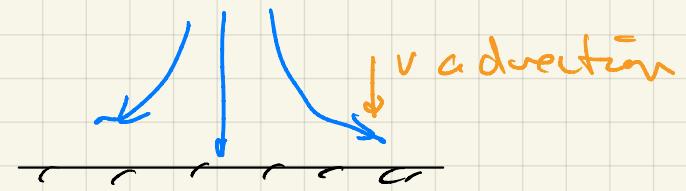
choose to preserve exactly

- "leaking" some small mass might be ok ... but maybe not for $t \rightarrow \infty$

- pocketwise condition - also "strong"
- tells us v small but needs to be retained

$$\text{so } \frac{\partial}{\partial y} \sim \frac{1}{S} \Rightarrow v \approx S$$

- v important on



- keep $u \frac{\partial u}{\partial x} + v \frac{\partial u}{\partial y}$ in x -momentum

• NOTE:

$$\boxed{\frac{\partial u}{\partial t} + u \frac{\partial u}{\partial x} + v \frac{\partial u}{\partial y} = - \frac{1}{f} \frac{\partial p}{\partial x} + v \frac{\partial^2 u}{\partial y^2}}$$

have we

- dropped $v \frac{\partial^2 u}{\partial x^2}$ from N-S ?

- added $v \frac{\partial^2 u}{\partial y^2}$ to Euler ?

- examine matrix advection versus wall \perp diffusion

$$u \frac{\partial u}{\partial x}$$

v.

$$v \frac{\partial^2 u}{\partial y^2}$$

\rightarrow must balance

$$\bigcirc \left[\frac{u \frac{\partial u}{\partial x}}{v \frac{\partial^2 u}{\partial y^2}} \right] = 1$$

$$u \sim v$$

$$x \sim L$$

$$\delta \sim \delta$$

$$\downarrow \quad \frac{U_L^2}{V U / \delta^2} = \frac{\delta^2 U}{VL} = \frac{\delta^2}{L^2} \quad \underbrace{\frac{UL}{V}}_{Re \#} = O(1)$$

$$\frac{\delta}{L} \sim \frac{1}{\sqrt{Re}} \iff \delta \sim \sqrt{V}$$

• Mass

$$\boxed{\frac{\partial u}{\partial x} + \frac{\partial v}{\partial y} = 0}$$

$$\sim \frac{u}{\delta} \quad \sim \frac{v}{\delta}$$

$$\rightarrow V \sim \frac{su}{\delta} \approx \frac{J}{R_C} \delta^2$$

• y-momentum

$$y \sim \delta \quad v \sim \delta \quad v \sim \delta^2$$

$$\frac{\partial v}{\partial t} + u \frac{\partial v}{\partial x} + v \frac{\partial u}{\partial y} = - \frac{1}{\rho} \frac{\partial p}{\partial y} + v \left(\frac{\partial^2 u}{\partial x^2} + \frac{\partial^2 u}{\partial y^2} \right)$$

$$\sim \delta$$

$$\sim \delta$$

$$\sim \delta^2 / \delta$$

$$\sim \frac{1}{\delta}$$

$$\delta^2 \left(\sim \delta \quad \sim \frac{\delta}{\delta^2} \right)$$

leading order

$$-\frac{1}{\rho} \frac{\partial p}{\partial y} = O(\delta)$$

$$\boxed{\frac{\partial p}{\partial y} = 0}$$

y momentum
to $O(\delta)$

BC:

$$u = v = 0 \quad \text{at } y = 0$$

$(P \rightarrow P_0)$

$$u \rightarrow U(x, t) \quad \text{for } y/\delta \rightarrow \infty$$

NOTES

- only have 1 ∇ BC
- whole system can be one 3rd order equation (∇ [Ockendon + Ockendon])
- $\nabla \cdot \mathbf{v}$ is "asymptotically" small
 - some presumption of $\nabla = 0$ where it's really just tiny
- what is S ? → exact value doesn't matter
 $S \sim \sqrt{D}$ and smooth transition to OUTER flow
- what is the "OUTER" flow
 - assume effectively miscid

x-mom

$$\frac{\partial U}{\partial t} + U \frac{\partial U}{\partial x} = - \gamma g \frac{\partial P}{\partial x}$$

$U=0$ — no U advection
↑
no viscous

→ streamline on edge
of B_2

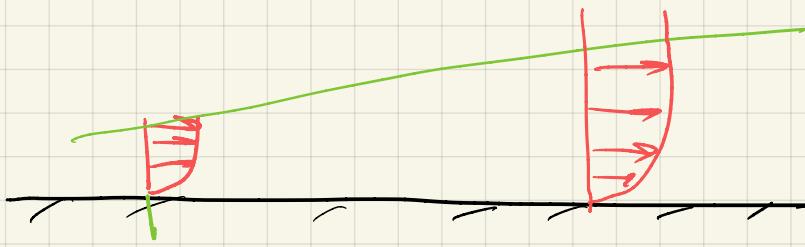
Steady

$$\frac{\partial}{\partial x} \left(\frac{1}{2} U^2 + \gamma g P \right) = 0$$

$$P + \frac{1}{2} \gamma g U^2 = \text{const}$$

↑ persist to wall $\frac{\partial P}{\partial y} = 0$
in BL

• also need x BC : $u(x_0, y)$



↑ need info at some x_0

- unsteady \rightarrow also need an IC

SYSTEMATIC

single ansatz \rightarrow result

start

$$\frac{D\bar{u}}{Dt} = -\frac{1}{\rho} \nabla p + \frac{1}{Re} \nabla^2 \bar{u} \quad \nabla \cdot \bar{u} = 0$$

$$y' = Re^{1/2} y \quad v' = Re^{1/2} v$$

so

(drop primes)

$$\frac{\partial \bar{u}}{\partial t} + u \frac{\partial \bar{u}}{\partial x} + v \frac{\partial \bar{u}}{\partial y} = - \frac{\partial p}{\partial x} + \frac{\partial^2 \bar{u}}{\partial y^2}$$

$$0 = - \frac{\partial p}{\partial y}$$

$$\frac{\partial \bar{u}}{\partial x} + \frac{\partial \bar{v}}{\partial y} = 0$$

A $\xrightarrow{\text{Re model}} \infty$
for high Re