

# PHYS2100 Semester 2 2021 Practice Exam

Siwan Li (s4583040)

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1.

$$\begin{aligned}
 F(r) &= -V'(r) = -\frac{GMm}{x^2} \\
 V'(r) &= \frac{GMm}{x^2} \\
 V(r) &= -\frac{GMm}{x} \\
 T &= \frac{1}{2}mv^2 = \frac{1}{2}m\dot{x}^2 \\
 E &= T + V = \frac{1}{2}mv^2 - \frac{GMm}{x}
 \end{aligned}$$

At infinite distance,

$$\begin{aligned}
 E &= 0 - 0 = 0 \\
 \therefore \frac{GMm}{x} &= \frac{1}{2}mv^2 \\
 \frac{2GM}{x} &= v^2 \\
 \frac{Rc^2}{x} &= v^2 \\
 v &= -c\sqrt{\frac{R}{x}}
 \end{aligned}$$

2. (i) Let  $y$  be the verticle length of string from pulley  $A$  to mass  $m_2$ . So then the remaining length of string on the incline would have length  $l - y$ .  $y$  will be the free variable in this solution. Let the vector space be centred at the tip of the inclined plane.

$$\begin{aligned}
 \mathbf{r}_1 &= h - (l - y) \cos \beta \hat{\mathbf{i}} + h \tan \beta - (l - y) \sin \beta \hat{\mathbf{j}} \\
 \mathbf{r}_2 &= h \hat{\mathbf{i}} + h \tan \beta - y \hat{\mathbf{j}} \\
 \mathbf{v}_1 &= \dot{y}(\cos \beta \hat{\mathbf{i}} + \sin \beta \hat{\mathbf{j}}) \\
 \mathbf{v}_2 &= -\dot{y} \hat{\mathbf{j}} \\
 T &= \frac{1}{2}m_1\dot{y}^2 + \frac{1}{2}m_2\dot{y}^2 \\
 V &= m_1g(h \tan \beta + (y - l) \sin \beta) + m_2g(h \tan \beta - y) \\
 L &= T - V = \frac{1}{2}m_1\dot{y}^2 + \frac{1}{2}m_2\dot{y}^2 - m_1g(h \tan \beta + (y - l) \sin \beta) - m_2g(h \tan \beta - y)
 \end{aligned}$$

(ii)

$$\begin{aligned} 0 &= \frac{d}{dt} \left( \frac{\partial L}{\partial \dot{y}} \right) - \frac{\partial L}{\partial y} \\ &= m_1 \ddot{y} + m_2 \ddot{y} + m_1 g \sin \beta - m_2 g \\ &= \ddot{y} + \frac{m_1 g \sin \beta - m_2 g}{m_1 + m_2} \\ \ddot{y} &= -\frac{m_1 g \sin \beta - m_2 g}{m_1 + m_2} \\ \Rightarrow y &= -\frac{m_1 g \sin \beta - m_2 g}{2(m_1 + m_2)} t^2 + c_1 t + c_2 \end{aligned}$$

(iii) The system is in equilibrium when  $\ddot{y} = \dot{y} = 0$ . I.e.

$$\begin{aligned} 0 &= \ddot{y} = -\frac{m_1 g \sin \beta - m_2 g}{m_1 + m_2} \\ 0 &= \dot{y} = -\frac{m_1 g \sin \beta - m_2 g}{m_1 + m_2} t + c_1 = c_1 \\ \Rightarrow y_0 &= -\frac{m_1 g \sin \beta - m_2 g}{2(m_1 + m_2)} t^2 + c_1 t + c_2 = c_2 \end{aligned}$$

3. (i)

$$\begin{aligned} \mathbf{F}(x, y) &= -k_1 x \hat{\mathbf{i}} - k_2 y \hat{\mathbf{j}} = -\nabla V \\ &= -\frac{\partial V}{\partial x} \hat{\mathbf{i}} - \frac{\partial V}{\partial y} \hat{\mathbf{j}} \\ \Rightarrow V &= \frac{k_1}{2} x^2 + g(y) = \frac{k_2}{2} y^2 + g(x) \\ \Rightarrow V &= \frac{k_1}{2} x^2 + \frac{k_2}{2} y^2 + c \text{ let } c = 0 \text{ so that potential is 0 at origin} \\ \dot{\mathbf{r}} &= -\frac{k_1}{m} x t \hat{\mathbf{i}} - \frac{k_2}{m} y t \hat{\mathbf{j}} \\ T &= \frac{(k_1^2 x^2 + k_2^2 y^2) t^2}{2m} \\ \therefore L &= T - V = \frac{(k_1^2 x^2 + k_2^2 y^2) t^2}{2m} - \frac{k_1 x^2 + k_2 y^2}{2} \\ &= \frac{m}{2} (\dot{x}^2 + \dot{y}^2) - \frac{k_1 x^2 + k_2 y^2}{2} \\ p_x &= \frac{\partial L}{\partial \dot{x}} = m \dot{x} \\ p_y &= \frac{\partial L}{\partial \dot{y}} = m \dot{y} \\ H &= p_x \dot{x} + p_y \dot{y} - L \\ &= m \dot{x}^2 + m \dot{y}^2 - \frac{m}{2} (\dot{x}^2 + \dot{y}^2) + \frac{k_1 x^2 + k_2 y^2}{2} \\ &= \frac{m}{2} (\dot{x}^2 + \dot{y}^2) + \frac{k_1 x^2 + k_2 y^2}{2} \\ &= \frac{p_x^2 + p_y^2}{2m} + \frac{k_1 x^2 + k_2 y^2}{2} \end{aligned}$$

(ii)

$$\begin{aligned}
\dot{x} &= \frac{\partial H}{\partial p_x} = \frac{p_x}{m} \\
\dot{y} &= \frac{\partial H}{\partial p_y} = \frac{p_y}{m} \\
\dot{p}_x &= -\frac{\partial H}{\partial x} = -k_1 x \\
\dot{p}_y &= -\frac{\partial H}{\partial y} = -k_2 y
\end{aligned}$$

(iii)

$$\begin{aligned}
\frac{d}{dt} \left( \frac{\partial L}{\partial \dot{x}} \right) - \frac{\partial L}{\partial x} &= m\ddot{x} + k_1 x = 0 \\
\frac{d}{dt} \left( \frac{\partial L}{\partial \dot{y}} \right) - \frac{\partial L}{\partial y} &= m\ddot{y} + k_2 y = 0 \\
\ddot{x} &= -\frac{k_1}{m} x, \quad \ddot{y} = -\frac{k_2}{m} y \\
\Rightarrow x &= A \sin \left( \sqrt{\frac{k_1}{m}} t \right) + B \cos \left( \sqrt{\frac{k_1}{m}} t \right) \\
1 &= A \sin(0) + B \cos(0) \\
\Rightarrow B &= 1 \\
\dot{x} &= -\sqrt{\frac{k_1}{m}} \sin \left( \sqrt{\frac{k_1}{m}} t \right) + \sqrt{\frac{k_1}{m}} A \cos \left( \sqrt{\frac{k_1}{m}} t \right) \\
-1 &= \sqrt{\frac{k_1}{m}} A \\
\Rightarrow A &= -\sqrt{\frac{m}{k_1}} \\
y &= C \sin \left( \sqrt{\frac{k_2}{m}} t \right) + D \cos \left( \sqrt{\frac{k_2}{m}} t \right) \\
0 &= C = D \\
\therefore x &= -\sqrt{\frac{m}{k_1}} \sin \left( \sqrt{\frac{k_1}{m}} t \right) + \cos \left( \sqrt{\frac{k_1}{m}} t \right), \quad y = 0
\end{aligned}$$

4.

$$\begin{aligned}
\dot{q} &= \frac{\partial H}{\partial p} = p \\
\dot{p} &= -\frac{\partial H}{\partial q} = -V'(q) = \begin{cases} -2q(|q| - a) - q^2 & \text{if } q \text{ is positive} \\ -2q(|q| - a) + q^2 & \text{if } q \text{ is negative} \\ \text{undefined} & \text{otherwise} \end{cases}
\end{aligned}$$

Consider the case when  $q > 0$

$$\begin{aligned}
(\dot{q}, \dot{p}) &= (0, 0) = (p, -2q(q - a) - q^2) \\
\Rightarrow p &= 0 \\
2q(q - a) &= -q^2 \\
q - a &= -\frac{q}{2} \\
q + \frac{q}{2} &= a \\
\frac{3q}{2} &= a \\
q &= \frac{2}{3}a
\end{aligned}$$

Otherwise if  $q$  is negative,

$$\begin{aligned}
\dot{p} &= 0 = -2q(-q - a) + q^2 \\
-q^2 &= 2q(q + a) \\
-q &= 2(q + a) \\
3q &= -2a \\
q &= -\frac{2}{3}a
\end{aligned}$$

Note, in either case,  $a > 0$

$$\begin{aligned}
H &= 0 + \frac{4}{9}a^2\left(\frac{2}{3}a - a\right) \\
&= -\frac{4}{27}a^3
\end{aligned}$$

5. (a) Note that  $\frac{df^n(x)}{dx} = \frac{df^{n-1}(x)}{dx} \frac{df(x)}{dx}$ , so

$$\begin{aligned}
&\lim_{n \rightarrow \infty} \frac{1}{n} \ln \left| \frac{df^n(x)}{dx} \right|_{x=x_0} \\
&\lim_{n \rightarrow \infty} \frac{1}{n} \ln \left| \prod_{i=0}^{n-1} \frac{df(x)}{dx} \right|_{x=x_i} \\
&\lim_{n \rightarrow \infty} \frac{1}{n} \sum_{i=0}^{n-1} \ln \left| \frac{df(x)}{dx} \right|_{x=x_i} \\
&\lim_{n \rightarrow \infty} \frac{1}{n} \sum_{i=0}^{n-1} \ln |f'(x_i)|
\end{aligned}$$

(b) In a fixed point,  $f'(x_i) = 0$ , so the exponent would be  $-\infty$ . Else for attractors,  $p$ -cycles and, in general, cases where  $x$  converges,  $\lambda < 0$ .

(c)

$$f'(x) = \begin{cases} r, & \text{if } 0 \leq x \leq \frac{1}{2} \\ -r, & \text{if } \frac{1}{2} \leq x \leq 1 \end{cases}$$

$$\begin{aligned} \lambda &= \lim_{n \rightarrow \infty} \frac{1}{n} \sum_{i=0}^{n-1} \ln |f'(x_i)| \\ &= \lim_{n \rightarrow \infty} \frac{1}{n} \sum_{i=0}^{n-1} \ln r \\ &= \lim_{n \rightarrow \infty} \ln r = \ln r \end{aligned}$$

So  $\lambda$  must diverge when  $r > 1$ .

6. (i)

$$\begin{aligned} d\tau^2 &= dt^2(1 - 0.8^2) = 0.36dt^2 \\ \tau^2 &= 0.36t^2 \\ \tau &= 0.6t \end{aligned}$$

So total time for Alice is  $\tau = 0.6 \cdot 4$  years, or 2.4 years.

- (ii) Due to length contraction, the journey that Alice actually takes is a lot shorter than what Bob perceives. And so Alice is covering less distance in the same amount of time, thus, isn't travelling faster than light.

7. (i)

$$\begin{aligned} d\tau &= dt_1 \sqrt{1 - \vec{v}_1^2} = dt_2 \sqrt{1 - \vec{v}_2^2} \\ \mathbf{u}_1 &= \left( \frac{dt_1}{d\tau}, \frac{d\vec{x}_1}{d\tau} \right) = \left( \frac{dt_1}{d\tau}, \frac{d\vec{x}_1}{dt_1} \frac{dt_1}{d\tau} \right) \\ &= \left( \frac{1}{\sqrt{1 - \vec{v}_1^2}}, \vec{v}_1 \frac{1}{\sqrt{1 - \vec{v}_1^2}} \right) = (\gamma_1, \gamma_1 \vec{v}_1) \end{aligned}$$

Similarly,  $\mathbf{u}_2 = (\gamma_2, \gamma_2 \vec{v}_2)$  where  $\gamma_2 = (1 - \vec{v}_2^2)^{-\frac{1}{2}}$

(ii)

$$\begin{aligned} V_{rel} &= \sqrt{1 - (\mathbf{u}_1 \cdot \mathbf{u}_2)^{-2}} \\ &= \sqrt{1 - (-\gamma_1 \gamma_2 + \gamma_1 \gamma_2 v_1 v_2)^{-2}} \\ &= \sqrt{1 - \left( \frac{-1 + v_1 v_2}{\sqrt{(1 - v_1)(1 - v_2)}} \right)^{-2}} \\ &= \sqrt{1 - \left( \frac{(1 - v_1)(1 - v_2)}{(v_1 v_2 - 1)^2} \right)} \\ &= \sqrt{\frac{(v_1 v_2 - 1)^2 - (1 - v_1)(1 - v_2)}{(v_1 v_2 - 1)^2}} \\ &= \sqrt{\frac{1 + v_1^2 v_2^2 - 2v_1 v_2 - 1 - v_1 v_2 + v_1 + v_2}{1 + v_1^2 v_2^2 - 2v_1 v_2}} \\ &= \sqrt{\frac{v_1^2 v_2^2 - 3v_1 v_2 + v_1 + v_2}{1 + v_1^2 v_2^2 - 2v_1 v_2}} \end{aligned}$$

Fuck this

8.

$$\begin{aligned} E &= \sqrt{m^2 + p^2} \\ &= \sqrt{m^2 \left( 1 + \left( \frac{p}{m} \right)^2 \right)} \\ &= m \sqrt{1 + (\gamma v)^2} \\ &\approx m \left( 1 + \frac{\gamma^2 v^2}{2} \right) \\ &\approx m + \frac{mv^2}{2} \end{aligned}$$

The  $\frac{mv^2}{2}$  term contributes kinetic energy, while  $m$  contributes mass energy