

Rigid Geometric Transformations

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November 6, 2017

This note is a quick refresher of the geometry of rigid transformations in three-dimensional space, expressed in Cartesian coordinates.

1 Cartesian Coordinates

Let us assume the notions of the *distance* between two points and the *angle* between lines to be known from geometry. The *law of cosines* is also stated without proof¹: if a, b, c are the sides of a triangle and the angle between a and b is θ , then

$$c^2 = a^2 + b^2 - 2ab \cos \theta .$$

The special case for $\theta = \pi/2$ radians is known as Pythagoras' theorem.

The definitions that follow focus on three-dimensional space. Two-dimensional geometry can be derived as a special case when the third coordinate of every point is set to zero.

A *Cartesian reference system* for three-dimensional space is a point in space called the *origin* and three mutually perpendicular, directed lines through the origin called the *axes*. The order in which the axes are listed is fixed, and is part of the definition of the reference system. The plane that contains the second and third axis is the *first reference plane*. The plane that contains the third and first axis is the *second reference plane*. The plane that contains the first and second axis is the *third reference plane*.

It is customary to mark the axis directions by specifying a point on each axis and at unit distance from the origin. These points are called the *unit points* of the system, and the *positive direction* of an axis is from the origin towards the axis' unit point. A Cartesian reference system is *right-handed* if the smallest rotation that brings the first unit point to the second is counterclockwise when viewed from the third unit point. The system is *left-handed* otherwise.

The *Cartesian coordinates* of a point in three-dimensional space are the signed distances of the point from the first, second, and third reference plane, in this order, and collected into a vector. The sign for coordinate i is positive if the point is in the half-space (delimited by the i -th reference plane) that contains the positive half of the i -th reference axis. It follows that the Cartesian coordinates of the origin are $\mathbf{t} = (0, 0, 0)^T$, those of the unit points are the vectors $\mathbf{e}_x = (1, 0, 0)^T$, $\mathbf{e}_y = (0, 1, 0)^T$, and $\mathbf{e}_z = (0, 0, 1)^T$, and the vector $\mathbf{p} = (x, y, z)^T$ of coordinates of an arbitrary point in space can also be written as follows:

$$\mathbf{p} = x\mathbf{e}_x + y\mathbf{e}_y + z\mathbf{e}_z .$$

The point \mathbf{p} can be reached from the origin \mathbf{t} by the following polygonal path:

$$\mathbf{t}, x\mathbf{e}_x, x\mathbf{e}_x + y\mathbf{e}_y, \mathbf{p} .$$

¹A proof based on trigonometry is straightforward but tedious, and a useful exercise.

Each segment of the path is followed by a right-angle turn, so Pythagoras' theorem can be applied twice to yield the distance of \mathbf{p} from the origin:

$$d(\mathbf{t}, \mathbf{p}) = \sqrt{x^2 + y^2 + z^2} .$$

From the definition of norm of a vector we see that

$$d(\mathbf{t}, \mathbf{p}) = \|\mathbf{p}\| .$$

So the norm of the vector of coordinates of a point is the distance of the point from the origin. A vector is often drawn as an arrow pointing from the origin to the point whose coordinates are the components of the vector. Then, the result above shows that the *length* of that arrow is the norm of the vector. Because of this, the words “length” and “norm” are often used interchangeably.

2 Orthogonality

The law of cosines yields a geometric interpretation of the inner product of two vectors \mathbf{a} and \mathbf{b} :

Theorem 2.1.

$$\mathbf{a}^T \mathbf{b} = \|\mathbf{a}\| \|\mathbf{b}\| \cos \theta$$

where θ is the acute angle between the two arrows that represent \mathbf{a} and \mathbf{b} geometrically.

So the inner product of two vectors is the product of the lengths of the two arrows that represent them and of the cosine of the angle between them. See the appendix for a proof.

Setting $\theta = \pi/2$ in the result above yields another important corollary:

Corollary 2.2. *The arrows that represent two vectors \mathbf{a} and \mathbf{b} are mutually perpendicular if and only if the two vectors are orthogonal:*

$$\mathbf{a}^T \mathbf{b} = 0 .$$

Because of this result, the words “perpendicular” and “orthogonal” are often used interchangeably.

3 Orthogonal Projection

Given two vectors \mathbf{a} and \mathbf{b} , the *orthogonal projection* of \mathbf{a} onto \mathbf{b} is the vector \mathbf{p} that represents the point p on the line through \mathbf{b} that is nearest to the endpoint of \mathbf{a} . See Figure 1.

Theorem 3.1. *The orthogonal projection of \mathbf{a} onto \mathbf{b} is the vector*

$$\mathbf{p} = P_{\mathbf{b}} \mathbf{a}$$

where $P_{\mathbf{b}}$ is the following square, symmetric, rank-1 matrix:

$$P_{\mathbf{b}} = \frac{\mathbf{b} \mathbf{b}^T}{\mathbf{b}^T \mathbf{b}} .$$

The signed magnitude of the orthogonal projection is

$$p = \frac{\mathbf{b}^T \mathbf{a}}{\|\mathbf{b}\|} = \|\mathbf{p}\| \operatorname{sign}(\mathbf{b}^T \mathbf{a}) .$$

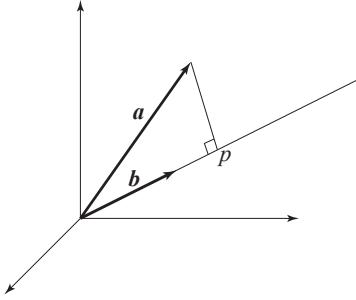


Figure 1: The vector from the origin to point p is the orthogonal projection of \mathbf{a} onto \mathbf{b} . The line from the endpoint of \mathbf{a} to p is orthogonal to \mathbf{b} .

From the definition of orthogonal projection we also see the following fact.

Corollary 3.2. *The coordinates of a point in space are the signed magnitudes of the orthogonal projections of the vector of coordinates of the point onto the three unit vectors that define the coordinate axes.*

This result is trivial in the basic Cartesian reference frame with unit points $\mathbf{e}_x = (1, 0, 0)^T$, $\mathbf{e}_y = (0, 1, 0)^T$, $\mathbf{e}_z = (0, 0, 1)^T$. If $\mathbf{p} = (x, y, z)^T$, then obviously

$$\mathbf{e}_x \mathbf{p} = x , \quad \mathbf{e}_y \mathbf{p} = y , \quad \mathbf{e}_z \mathbf{p} = z .$$

The result becomes less trivial in Cartesian reference systems where the axes have different orientations, as we will see soon.

4 The Cross Product

The *cross product* of two 3-dimensional vectors $\mathbf{a} = (a_x, a_y, a_z)^T$ and $\mathbf{b} = (b_x, b_y, b_z)^T$ is the 3-dimensional vector

$$\mathbf{c} = \mathbf{a} \times \mathbf{b} = (a_y b_z - a_z b_y, a_z b_x - a_x b_z, a_x b_y - a_y b_x)^T .$$

The following geometric interpretation is proven in the Appendix:

Theorem 4.1. *The cross product of two three-dimensional vectors \mathbf{a} and \mathbf{b} is a vector \mathbf{c} orthogonal to both \mathbf{a} and \mathbf{b} , oriented so that the triple $\mathbf{a}, \mathbf{b}, \mathbf{c}$ is right-handed, and with magnitude*

$$\|\mathbf{c}\| = \|\mathbf{a} \times \mathbf{b}\| = \|\mathbf{a}\| \|\mathbf{b}\| \sin \theta$$

where θ is the acute angle between \mathbf{a} and \mathbf{b} .

From its expression, we see that the magnitude of $\mathbf{a} \times \mathbf{b}$ is the area of a rectangle with sides \mathbf{a} and \mathbf{b} .

It is immediate to verify that the cross product of two vectors is a linear transformation of either vector:

$$\mathbf{a} \times (\mathbf{b}_1 + \mathbf{b}_2) = \mathbf{a} \times \mathbf{b}_1 + \mathbf{a} \times \mathbf{b}_2 \quad \text{and similarly} \quad (\mathbf{a}_1 + \mathbf{a}_2) \times \mathbf{b} = \mathbf{a}_1 \times \mathbf{b} + \mathbf{a}_2 \times \mathbf{b} .$$

So there must be a 3×3 matrix $[\mathbf{a}]_{\times}$ such that

$$\mathbf{a} \times \mathbf{b} = [\mathbf{a}]_{\times} \mathbf{b} .$$

This matrix is convenient for computing cross products of the form $\mathbf{a} \times \mathbf{p}$ where \mathbf{a} is a fixed vector but \mathbf{p} changes. Spelling out the definition of the cross product yields the following matrix:

$$[\mathbf{a}]_{\times} = \begin{bmatrix} 0 & -a_z & a_y \\ a_z & 0 & -a_x \\ -a_y & a_x & 0 \end{bmatrix} .$$

This matrix is skew-symmetric:

$$[\mathbf{a}]_{\times}^T = -[\mathbf{a}]_{\times} .$$

Of course, similar considerations hold for \mathbf{b} : Since

$$\mathbf{a} \times \mathbf{b} = -\mathbf{b} \times \mathbf{a} ,$$

we have

$$\mathbf{a} \times \mathbf{b} = -[\mathbf{b}]_{\times} \mathbf{a} = [\mathbf{b}]_{\times}^T \mathbf{a} .$$

5 The Triple Product

The *triple product* of three-dimensional vectors $\mathbf{a}, \mathbf{b}, \mathbf{c}$ is defined as follows:

$$\mathbf{a}^T (\mathbf{b} \times \mathbf{c}) = a_x(b_y c_z - b_z c_y) - a_y(b_x c_z - b_z c_x) + a_z(b_x c_y - b_y c_x) .$$

It is immediate to verify that

$$\mathbf{a}^T (\mathbf{b} \times \mathbf{c}) = \mathbf{b}^T (\mathbf{c} \times \mathbf{a}) = \mathbf{c}^T (\mathbf{a} \times \mathbf{b}) = -\mathbf{a}^T (\mathbf{c} \times \mathbf{b}) = -\mathbf{c}^T (\mathbf{b} \times \mathbf{a}) = -\mathbf{b}^T (\mathbf{a} \times \mathbf{c}) .$$

Again, from its expression, we see that the triple product of vectors $\mathbf{a}, \mathbf{b}, \mathbf{c}$ is, up to a sign, the volume of a parallelepiped with edges $\mathbf{a}, \mathbf{b}, \mathbf{c}$: The cross product $\mathbf{p} = \mathbf{b} \times \mathbf{c}$ is a vector orthogonal to the plane of \mathbf{b} and \mathbf{c} , and with magnitude equal to the base area of the parallelepiped. The inner product of \mathbf{p} and \mathbf{a} is the magnitude of \mathbf{p} times that of \mathbf{a} times the cosine of the angle between them, that is, the base area of the parallelepiped times its height (or the negative of its height). This gives the volume of the solid, up to a sign. The sign is positive if the three vectors form a right-handed triple. See Figure 2.

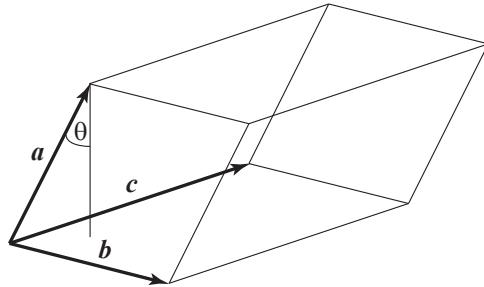


Figure 2: Up to a sign, the triple product of the vectors $\mathbf{a}, \mathbf{b}, \mathbf{c}$ is the volume of the parallelepiped with edges $\mathbf{a}, \mathbf{b}, \mathbf{c}$.

6 Multiple Reference Systems

When two or more reference systems are involved, notation must be introduced to avoid possible ambiguities as to which reference system coordinates are referred to. For instance, we may want to express the coordinates of the origin of one system with respect to the other system.

Reference systems will be identified with natural numbers, and the number zero is reserved for a privileged system called the *world reference system*. A left superscript is used to identify the reference system that a vector or a transformation is written in. A left superscript of zero can be optionally omitted.

Thus, ${}^2\mathbf{p}$ is the vector of the coordinates of point \mathbf{p} in reference frame 2. The same point in world coordinates can be written as either ${}^0\mathbf{p}$ or just \mathbf{p} . The origin \mathbf{t} of a reference system has always zero coordinates in that reference system, so

$${}^i\mathbf{t}_i = \begin{bmatrix} 0 \\ 0 \\ 0 \end{bmatrix}$$

for all natural numbers i , and therefore also

$$\mathbf{t}_0 = \begin{bmatrix} 0 \\ 0 \\ 0 \end{bmatrix}$$

since a zero left superscript is implied. Similarly, if $\mathbf{i}, \mathbf{j}, \mathbf{k}$ are the unit points of a reference system, we have

$$[{}^i\mathbf{i}_i \quad {}^i\mathbf{j}_i \quad {}^i\mathbf{k}_i] = [\mathbf{i}_0 \quad \mathbf{j}_0 \quad \mathbf{k}_0] = I,$$

the 3×3 identity matrix.

7 Rotation

A *rotation* is a transformation between two Cartesian references systems of equal origin and handedness. Let the two systems be S_0 (the world reference system) and S_1 . Then, the common origin is

$$\mathbf{t}_0 = \mathbf{t}_1 = \begin{bmatrix} 0 \\ 0 \\ 0 \end{bmatrix}$$

where the left superscripts were omitted because coordinates refer to the world reference system. The unit points of S_1 are $\mathbf{i}_1, \mathbf{j}_1, \mathbf{k}_1$ when their coordinates are expressed in S_0 . Then a point with coordinates $\mathbf{p} = (x, y, z)^T$ in S_0 can be reached from the origin \mathbf{t}_1 common to the two systems by a polygonal path with the following four vertices:

$$\mathbf{t}_1 \quad , \quad \mathbf{a} = {}^1x \mathbf{i}_1 \quad , \quad \mathbf{b} = {}^1x \mathbf{i}_1 + {}^1y \mathbf{j}_1 \quad , \quad \mathbf{p} = {}^1x \mathbf{i}_1 + {}^1y \mathbf{j}_1 + {}^1z \mathbf{k}_1 .$$

The steps of this path are along the axes of S_1 . The numbers ${}^1x, {}^1y, {}^1z$ are the magnitudes of the steps, and also the coordinates of the point in S_1 . These step sizes are the signed magnitudes of the orthogonal projections of the point onto $\mathbf{i}_1, \mathbf{j}_1, \mathbf{k}_1$, and from Theorem 3.1 we see that

$${}^1x = \mathbf{i}_1^T \mathbf{p} \quad , \quad {}^1y = \mathbf{j}_1^T \mathbf{p} \quad , \quad {}^1z = \mathbf{k}_1^T \mathbf{p}$$

because the vectors $\mathbf{i}_1, \mathbf{j}_1, \mathbf{k}_1$ have unit norm. These three equations can be packaged into a single matrix equation that expresses the vector ${}^1\mathbf{p} = ({}^1x, {}^1y, {}^1z)^T$ as a function of \mathbf{p} :

$${}^1\mathbf{p} = R_1 \mathbf{p} \quad \text{where} \quad R_1 = {}^0R_1 = \begin{bmatrix} \mathbf{i}_1^T \\ \mathbf{j}_1^T \\ \mathbf{k}_1^T \end{bmatrix}$$

where the 3×3 matrix 0R_1 is called a *rotation* matrix.

This result was obtained without using the privileged status of the world reference system S_0 (except to omit some of the left superscripts). Therefore, the result must be general: Given any two Cartesian reference systems S_a and S_b with a common origin,

$${}^b\mathbf{p} = {}^aR_b {}^a\mathbf{p} \quad \text{where} \quad {}^aR_b = \begin{bmatrix} {}^a\mathbf{i}_b^T \\ {}^a\mathbf{j}_b^T \\ {}^a\mathbf{k}_b^T \end{bmatrix} .$$

A rotation is a reversible transformation, and therefore the matrix aR_b must have an *inverse*, another matrix bR_a that transforms back from S_b to S_a :

$${}^a\mathbf{p} = {}^bR_a {}^b\mathbf{p} .$$

The proof of the following fact is given in the Appendix.

Theorem 7.1. *The inverse R^{-1} of a rotation matrix R is its transpose:*

$$R^T R = R R^T = I .$$

Equivalently, if aR_b is the rotation whose rows are the unit points of reference systems b expressed in reference system a , then

$${}^bR_a = {}^aR_b^T .$$

Note that R^T , being the inverse of R , is also a transformation between two Cartesian systems with the same origin and handedness, so R^T is a rotation matrix as well, and its rows must be mutually orthogonal unit vectors. Since the rows of R^T are the columns of R , we conclude that both the rows and columns of a rotation matrix are unit norm and orthogonal. This makes intuitive sense: Just as the rows of $R = {}^aR_b$ are the unit vectors of S_b expressed in S_a , so its columns (the rows of the inverse transformation $R^T = {}^bR_a$) are the unit vectors of S_a expressed in S_b .

The equations in Theorem 7.1 characterize combinations of rotations and possible inversions. An *inversion* (also known as a *mirror flip*) is a transformation that changes the direction of some of the axes. This is represented by a matrix of the form

$$S = \begin{bmatrix} s_x & 0 & 0 \\ 0 & s_y & 0 \\ 0 & 0 & s_z \end{bmatrix}$$

where s_x, s_y, s_z are equal to either 1 or -1 , and there is either one or three negative elements. It is easy to see that

$$S^T S = S S^T = I .$$

If there were zero or two negative elements, then S would be a rotation matrix, because the flip of two axes can be achieved by a rotation. For instance, the directions of both the x and the y axis can be flipped simultaneously by a 180-degree rotation around the z axis. No rotation can flip the directions of an odd number of axes.

The *determinant* of a 3×3 matrix is the triple product of its rows. Direct manipulation shows that this is the same as the triple product of its columns. It is immediate to see that the determinant of a rotation matrix is 1:

$$\det(R) = \mathbf{i}^T(\mathbf{j} \times \mathbf{k}) = \mathbf{i}^T\mathbf{i} = 1$$

because

$$\mathbf{i} \times \mathbf{j} = \mathbf{k}, \quad \mathbf{j} \times \mathbf{k} = \mathbf{i}, \quad \mathbf{k} \times \mathbf{i} = \mathbf{j}.$$

These equalities can be verified by the geometric interpretation of the cross product: each of the three vectors $\mathbf{i}, \mathbf{j}, \mathbf{k}$ is orthogonal to the other two, and its magnitude is equal to 1. The order of the vectors in the equalities above preserves handedness.

It is even easier to see that the determinant of an inversion matrix S is equal to -1 . Thus, the following conclusion can be drawn.

A matrix R is a rotation if and only if $R^T R = R R^T = I$ and $\det(R) = 1$.

A diagonal matrix S is an inversion if and only if $S^T S = S S^T = I$ and $\det(S) = -1$.

Note that in particular the identity matrix I is a rotation, and $-I$ is an inversion.

Geometric Interpretation of Orthogonality.

$$R^{-1} = R^T$$

is very simple, and yet was derived in the Appendix through a comparatively lengthy sequence of algebraic steps. This Section reviews orthogonality of rotation matrices from a geometric point of view, and derives the result above by conceptually simpler means. The rows of the rotation matrix

$${}^a R_b = \begin{bmatrix} r_{11} & r_{12} & r_{13} \\ r_{21} & r_{22} & r_{23} \\ r_{31} & r_{32} & r_{33} \end{bmatrix} = \begin{bmatrix} {}^a \mathbf{i}_b^T \\ {}^a \mathbf{j}_b^T \\ {}^a \mathbf{k}_b^T \end{bmatrix}$$

are the unit vectors of the reference system S_b , expressed in the reference system S_a . This means that its entry r_{mn} is the signed magnitude of the orthogonal projection of the m -th unit vector in S_b onto the n -th unit vector in S_a . For instance,

$$r_{12} = {}^a \mathbf{i}_b^T {}^a \mathbf{j}_a \quad \text{and} \quad r_{31} = {}^a \mathbf{k}_b^T {}^a \mathbf{i}_a.$$

However, the signed magnitude of the orthogonal projection of a unit vector onto another unit vector is simply the cosine of the angle between them:

$$r_{ij} = \cos \alpha_{ij}$$

where α_{ij} is the angle between the i -th axis in the new system and the j -th axis in the old.

Thus, the entries of a rotation matrix are *direction cosines*: they are all cosines of well-defined angles. This result also tells us that signed orthogonal projection magnitude is symmetric for unit vectors: For instance, the signed magnitude of the orthogonal projection of ${}^a\mathbf{i}_b$ onto ${}^a\mathbf{j}_a$ is the same as the signed magnitude of the orthogonal projection of ${}^a\mathbf{j}_a$ onto ${}^a\mathbf{i}_b$. Since angles do not depend on reference system, the projection is the same when expressed in S_b :

$$r_{12} = {}^a\mathbf{i}_b^T {}^a\mathbf{j}_a = {}^a\mathbf{j}_a^T {}^a\mathbf{i}_b = {}^b\mathbf{j}_a^T {}^b\mathbf{i}_b$$

where in the second equality we merely switched the two vectors with each other and in the third we changed the reference system (*i.e.*, both left superscripts) from S_a to S_b .

This symmetry is the deep reason for orthogonality: When we want to go from the “new” system S_b back to the “old” system S_a through the inverse matrix $R^{-1} = {}^bR_a$, we seek to express the unit vectors of S_a in the system S_b , that is, we seek the signed magnitudes of the orthogonal projections of each unit vector of the “old” system S_a onto each of the unit vectors of the “new” system S_b . Because of symmetry, these orthogonal projections are already available in the matrix R , just in a different arrangement: what we want in the rows of R^{-1} can be found in the columns of R . *Voilà*:

$$R^{-1} = R^T.$$

The Cross Product is Covariant with Rotations. We saw that the cross product of two vectors \mathbf{a} and \mathbf{b} is a third vector \mathbf{c} that is orthogonal to both \mathbf{a} and \mathbf{b} , and whose signed magnitude is the product $\|\mathbf{a}\| \|\mathbf{b}\|$ times the sine of the angle θ between \mathbf{a} and \mathbf{b} . If \mathbf{a} and \mathbf{b} are simultaneously rotated by the same rotation R to produce the new vectors \mathbf{a}' and \mathbf{b}' , then the line orthogonal to \mathbf{a} and \mathbf{b} rotates the same way, because rotating \mathbf{a} and \mathbf{b} is the same as rotating the reference system in the opposite direction. Thus, the direction of $\mathbf{c}' = \mathbf{a}' \times \mathbf{b}'$ is that of $R\mathbf{c}$. In addition, a rotation does not change the magnitudes of vectors it is applied to, nor does it change the angle between any pair of vectors. Therefore, $\mathbf{c}' = R\mathbf{c}$. To summarize, this argument shows that the cross product is covariant with rotations:

$$(R\mathbf{a}) \times (R\mathbf{b}) = R(\mathbf{a} \times \mathbf{b}).$$

In words, if you rotate the inputs to a cross product, the output rotates the same way.

8 Coordinate Transformation

A right-handed Cartesian system of reference S_1 can differ from the world reference system S_0 by a translation of the origin from $\mathbf{t}_0 = (0, 0, 0)^T$ to \mathbf{t}_1 and a rotation of the axes from unit points $\mathbf{i}_0 = \mathbf{e}_x$, $\mathbf{j}_0 = \mathbf{e}_y$, $\mathbf{k}_0 = \mathbf{e}_z$ to unit points \mathbf{i}_1 , \mathbf{j}_1 , \mathbf{k}_1 . Suppose that the origin of frame S_0 is first translated to point \mathbf{t}_1 and *then* the resulting frame is rotated by R_1 (see Figure 3). Given a point with coordinates $\mathbf{p} = (x, y, z)^T$ in S_0 , the coordinates ${}^1\mathbf{p} = ({}^1x, {}^1y, {}^1z)^T$ of the same point in S_1 are then

$${}^1\mathbf{p} = R_1(\mathbf{p} - \mathbf{t}_1). \quad (1)$$

The translation is applied first, to yield the new coordinates $\mathbf{p} - \mathbf{t}_1$ in an intermediate frame. This translation does not change the directions of the coordinate axes, so the rotation from the intermediate frame to S_1 is the same rotation R_1 as from S_0 to S_1 , which is applied thereafter.

The inverse transformation applies the inverse operations in reverse order:

$$\mathbf{p} = R_1^T {}^1\mathbf{p} + \mathbf{t}_1. \quad (2)$$

This can also be verified algebraically from equation (1): Multiplying both sides by R_1^T from the left yields

$$R_1^{T-1}\mathbf{p} = R_1^T R_1(\mathbf{p} - \mathbf{t}_1) = \mathbf{p} - \mathbf{t}_1$$

and adding \mathbf{t}_1 to both sides yields equation (2). Thus,

$${}^1R_0 = {}^0R_1^T \quad \text{and} \quad {}^1\mathbf{t}_0 = -{}^0R_1{}^0\mathbf{t}_1$$

since

$$\mathbf{p} = R_1^{T-1}\mathbf{p} + \mathbf{t}_1 = R_1^T({}^1\mathbf{p} - (-R_1\mathbf{t}_1)) .$$

More generally, if

$${}^b\mathbf{p} = {}^aR_b({}^a\mathbf{p} - {}^a\mathbf{t}_b) \quad (3)$$

then

$${}^a\mathbf{p} = {}^bR_a({}^b\mathbf{p} - {}^b\mathbf{t}_a) \quad \text{where} \quad {}^bR_a = {}^aR_b^T \quad \text{and} \quad {}^b\mathbf{t}_a = -{}^aR_b{}^a\mathbf{t}_b . \quad (4)$$

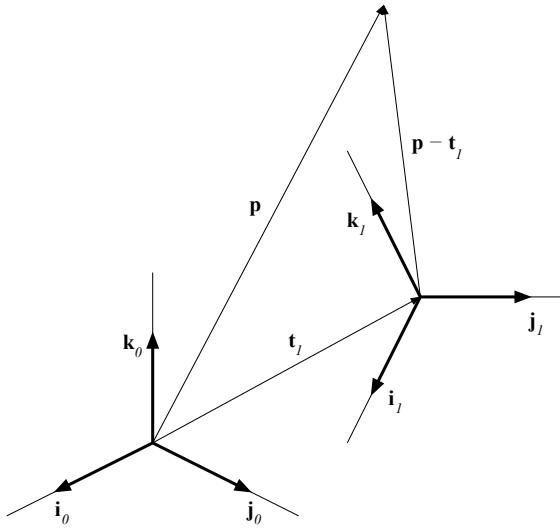


Figure 3: Transformation between two reference systems.

The transformations (3) and (4) are said to be *rigid*, in that they preserve distances. They are also sometimes referred to as *special Euclidean*, where the attribute “special” refers to the fact that mirror flips are not included—*i.e.*, the determinant of the rotation matrix is 1, rather than 1 just in magnitude.

Chaining Rigid Transformations If two right-handed Cartesian systems S_a and S_b are given, say, in the world reference system S_0 , then the transformation from S_a to S_b can be obtained in two steps, just as we did for rotations: First transform from S_a to S_0 , then from S_0 to S_b :

$$\mathbf{p} = {}^0\mathbf{p} = {}^aR_0({}^a\mathbf{p} - {}^a\mathbf{t}_0) = {}^0R_a^T({}^a\mathbf{p} + {}^0R_a{}^0\mathbf{t}_a) = R_a^T{}^a\mathbf{p} + \mathbf{t}_a$$

and

$${}^b\mathbf{p} = {}^0R_b({}^0\mathbf{p} - {}^0\mathbf{t}_b) = R_b(\mathbf{p} - \mathbf{t}_b)$$

so that

$${}^b\mathbf{p} = R_b(R_a^T {}^a\mathbf{p} + \mathbf{t}_a - \mathbf{t}_b) = R_bR_a^T[{}^a\mathbf{p} + R_a(\mathbf{t}_a - \mathbf{t}_b)] ,$$

that is,

$${}^b\mathbf{p} = {}^aR_b({}^a\mathbf{p} - {}^a\mathbf{t}_b) \quad \text{where} \quad {}^aR_b = R_bR_a^T \quad \text{and} \quad {}^a\mathbf{t}_b = R_a(\mathbf{t}_b - \mathbf{t}_a) .$$

The following box summarizes these results:

If

$$\mathbf{p}_a = R_b(\mathbf{p} - \mathbf{t}_b) \quad \text{and} \quad \mathbf{p}_b = R_a(\mathbf{p} - \mathbf{t}_a)$$

are the transformations between world coordinates and reference frames S_a and S_b , then the transformation from S_a to S_b is

$${}^b\mathbf{p} = {}^aR_b({}^a\mathbf{p} - {}^a\mathbf{t}_b) \quad \text{where} \quad {}^aR_b = R_bR_a^T \quad \text{and} \quad {}^a\mathbf{t}_b = R_a(\mathbf{t}_b - \mathbf{t}_a)$$

and the reverse transformation, from S_b to S_a , is

$${}^a\mathbf{p} = {}^bR_a({}^b\mathbf{p} - {}^b\mathbf{t}_a) \quad \text{where} \quad {}^bR_a = R_aR_b^T \quad \text{and} \quad {}^b\mathbf{t}_a = R_b(\mathbf{t}_a - \mathbf{t}_b) .$$

Consistently with these relations,

$${}^bR_a = {}^aR_b^T \quad \text{and} \quad {}^b\mathbf{t}_a = -{}^aR_b {}^a\mathbf{t}_b .$$

A Proofs

Theorem 2.1

$$\mathbf{a}^T \mathbf{b} = \|\mathbf{a}\| \|\mathbf{b}\| \cos \theta$$

where θ is the acute angle between the two arrows that represent \mathbf{a} and \mathbf{b} geometrically.

Proof. Consider a triangle with sides

$$a = \|\mathbf{a}\| , \quad b = \|\mathbf{b}\| , \quad c = \|\mathbf{b} - \mathbf{a}\|$$

and with an angle θ between a and b . Then the law of cosines yields

$$\|\mathbf{b} - \mathbf{a}\|^2 = \|\mathbf{a}\|^2 + \|\mathbf{b}\|^2 - 2\|\mathbf{a}\| \|\mathbf{b}\| \cos \theta .$$

From the definition of norm we then obtain

$$\|\mathbf{a}\|^2 + \|\mathbf{b}\|^2 - 2\mathbf{a}^T \mathbf{b} = \|\mathbf{a}\|^2 + \|\mathbf{b}\|^2 - 2\|\mathbf{a}\| \|\mathbf{b}\| \cos \theta .$$

Canceling equal terms and dividing by -2 yields the desired result.

Theorem 3.1

The orthogonal projection of \mathbf{a} onto \mathbf{b} is the vector

$$\mathbf{p} = P_{\mathbf{b}} \mathbf{a}$$

where $P_{\mathbf{b}}$ is the following square, symmetric, rank-1 matrix:

$$P_{\mathbf{b}} = \frac{\mathbf{b} \mathbf{b}^T}{\mathbf{b}^T \mathbf{b}} .$$

The signed magnitude of the orthogonal projection is

$$p = \frac{\mathbf{b}^T \mathbf{a}}{\|\mathbf{b}\|} = \|\mathbf{p}\| \operatorname{sign}(\mathbf{b}^T \mathbf{a}) .$$

Proof. To prove this, observe that since by definition point p is on the line through \mathbf{b} , the orthogonal projection vector \mathbf{p} has the form $\mathbf{p} = x\mathbf{b}$, where x is some real number. From elementary geometry, the line between p and the endpoint of \mathbf{a} is shortest when it is perpendicular to \mathbf{b} :

$$\mathbf{b}^T (\mathbf{a} - x\mathbf{b}) = 0$$

which yields

$$x = \frac{\mathbf{b}^T \mathbf{a}}{\mathbf{b}^T \mathbf{b}}$$

so that

$$\mathbf{p} = x\mathbf{b} = \mathbf{b} x = \frac{\mathbf{b} \mathbf{b}^T}{\mathbf{b}^T \mathbf{b}} \mathbf{a}$$

as advertised. The magnitude of \mathbf{p} can be computed as follows. First, observe that

$$P_{\mathbf{b}}^2 = \frac{\mathbf{b}\mathbf{b}^T}{\mathbf{b}^T\mathbf{b}} \frac{\mathbf{b}\mathbf{b}^T}{\mathbf{b}^T\mathbf{b}} = \frac{\mathbf{b}\mathbf{b}^T\mathbf{b}\mathbf{b}^T}{(\mathbf{b}^T\mathbf{b})^2} = \frac{\mathbf{b}\mathbf{b}^T}{\mathbf{b}^T\mathbf{b}} = P_{\mathbf{b}}$$

so that the orthogonal-projection matrix² $P_{\mathbf{b}}$ is *idempotent*:

$$P_{\mathbf{b}}^2 = P_{\mathbf{b}}.$$

This means that applying the matrix once or multiple times has the same effect. Then,

$$\|\mathbf{p}\|^2 = \mathbf{p}^T \mathbf{p} = \mathbf{a}^T P_{\mathbf{b}}^T P_{\mathbf{b}} \mathbf{a} = \mathbf{a}^T P_{\mathbf{b}} P_{\mathbf{b}} \mathbf{a} = \mathbf{a}^T P_{\mathbf{b}} \mathbf{a} = \mathbf{a}^T \frac{\mathbf{b}\mathbf{b}^T}{\mathbf{b}^T\mathbf{b}} \mathbf{a} = \frac{(\mathbf{b}^T \mathbf{a})^2}{\mathbf{b}^T\mathbf{b}}$$

which, once the sign of $\mathbf{b}^T \mathbf{a}$ is taken into account, yields the promised expression for the signed magnitude of \mathbf{p} .

Theorem 4.1

The cross product of two three-dimensional vectors \mathbf{a} and \mathbf{b} is a vector \mathbf{c} orthogonal to both \mathbf{a} and \mathbf{b} , oriented so that the triple $\mathbf{a}, \mathbf{b}, \mathbf{c}$ is right-handed, and with magnitude

$$\|\mathbf{c}\| = \|\mathbf{a} \times \mathbf{b}\| = \|\mathbf{a}\| \|\mathbf{b}\| \sin \theta$$

where θ is the acute angle between \mathbf{a} and \mathbf{b} .

Proof. That the cross product \mathbf{c} of \mathbf{a} and \mathbf{b} is orthogonal to both \mathbf{a} and \mathbf{b} can be checked directly:

$$\begin{aligned} \mathbf{c}^T \mathbf{a} &= (a_y b_z - a_z b_y) a_x + (a_z b_x - a_x b_z) a_y + (a_x b_y - a_y b_x) a_z = 0 \\ \mathbf{c}^T \mathbf{b} &= (a_y b_z - a_z b_y) b_x + (a_z b_x - a_x b_z) b_y + (a_x b_y - a_y b_x) b_z = 0 \end{aligned}$$

(verify that all terms do indeed cancel). We also have

$$(\mathbf{a}^T \mathbf{b})^2 + \|\mathbf{a} \times \mathbf{b}\|^2 = \|\mathbf{a}\|^2 \|\mathbf{b}\|^2$$

as can be shown by straightforward manipulation:

$$\begin{aligned} (\mathbf{a}^T \mathbf{b})^2 &= (a_x b_x + a_y b_y + a_z b_z)(a_x b_x + a_y b_y + a_z b_z) \\ &= a_x^2 b_x^2 + a_x b_x a_y b_y + a_x b_x a_z b_z \\ &\quad + a_y^2 b_y^2 + a_x b_x a_y b_y + a_y b_y a_z b_z \\ &\quad + a_z^2 b_z^2 + a_x b_x a_z b_z + a_y b_y a_z b_z \\ &= a_x^2 b_x^2 + a_y^2 b_y^2 + a_z^2 b_z^2 + 2a_x b_x a_y b_y + 2a_y b_y a_z b_z + 2a_x b_x a_z b_z \end{aligned}$$

and

$$\begin{aligned} \|\mathbf{a} \times \mathbf{b}\|^2 &= (a_y b_z - a_z b_y)^2 + (a_z b_x - a_x b_z)^2 + (a_x b_y - a_y b_x)^2 \\ &= a_y^2 b_z^2 + a_z^2 b_y^2 - 2a_y b_y a_z b_z \\ &\quad + a_x^2 b_z^2 + a_z^2 b_x^2 - 2a_x b_x a_z b_z \\ &\quad + a_x^2 b_y^2 + a_y^2 b_x^2 - 2a_x b_x a_y b_y \\ &= a_x^2 b_y^2 + a_y^2 b_x^2 + a_y^2 b_z^2 + a_z^2 b_y^2 + a_x^2 b_z^2 + a_z^2 b_x^2 \\ &\quad - 2a_x b_x a_y b_y - 2a_y b_z a_y b_y - 2a_x b_x a_z b_z \end{aligned}$$

²The matrix that describes orthogonal projection is not an orthogonal matrix. It could not possibly be, since it is rank-deficient.

so that

$$(\mathbf{a}^T \mathbf{b})^2 + \|\mathbf{a} \times \mathbf{b}\|^2 = a_x^2 b_x^2 + a_x^2 b_y^2 + a_x^2 b_z^2 + a_y^2 b_x^2 + a_y^2 b_y^2 + a_y^2 b_z^2 + a_z^2 b_x^2 + a_z^2 b_y^2 + a_z^2 b_z^2$$

but also

$$\|\mathbf{a}\|^2 \|\mathbf{b}\|^2 = a_x^2 b_x^2 + a_x^2 b_y^2 + a_x^2 b_z^2 + a_y^2 b_x^2 + a_y^2 b_y^2 + a_y^2 b_z^2 + a_z^2 b_x^2 + a_z^2 b_y^2 + a_z^2 b_z^2$$

so that

$$(\mathbf{a}^T \mathbf{b})^2 + \|\mathbf{a} \times \mathbf{b}\|^2 = \|\mathbf{a}\|^2 \|\mathbf{b}\|^2 \quad (5)$$

as desired. The result on the magnitude is a consequence of equation (5). From this equation we obtain

$$\|\mathbf{a} \times \mathbf{b}\|^2 = \|\mathbf{a}\|^2 \|\mathbf{b}\|^2 - (\mathbf{a}^T \mathbf{b})^2 = \|\mathbf{a}\|^2 \|\mathbf{b}\|^2 - \|\mathbf{a}\|^2 \|\mathbf{b}\|^2 \cos^2 \theta = \|\mathbf{a}\|^2 \|\mathbf{b}\|^2 \sin^2 \theta$$

or

$$\|\mathbf{a} \times \mathbf{b}\| = \pm \|\mathbf{a}\| \|\mathbf{b}\| \sin \theta.$$

Since the angle θ is acute (from the law of cosines), all quantities in the last equation are nonnegative, so that the $-$ sign yields an impossible equation. This yields the promised result.

Theorem 7.1

The inverse R^{-1} of a rotation matrix R is its transpose:

$$R^T R = R R^T = I.$$

Equivalently, if ${}^a R_b$ is the rotation whose rows are the unit points of reference systems b expressed in reference system a , then

$${}^b R_a = {}^a R_b^T.$$

Proof. Assume that $a = 0$ and $b = 1$, so that left superscripts equal to 0 can be omitted. This assumption can be made without loss of generality, because the following proof makes no use of the privileged nature of system S_0 , other than for simplifying away left superscripts. Also, for further brevity, let

$$R = {}^0 R_1.$$

When we rotate point \mathbf{p} through R we obtain a vector ${}^1 \mathbf{p}$ of coordinates in S_1 . We then look for a new matrix $R^{-1} = {}^1 R_0$ that applied to ${}^1 \mathbf{p}$ gives back the original vector \mathbf{p} :

$${}^1 \mathbf{p} = R \mathbf{p} \rightarrow \mathbf{p} = R^{-1} {}^1 \mathbf{p}$$

that is, by combining these two equations,

$$\mathbf{p} = R^{-1} R \mathbf{p}.$$

Since this equality is to hold for *any* vector \mathbf{p} , we need to find R^{-1} such that

$$R^{-1} R = I = \begin{bmatrix} 1 & 0 & 0 \\ 0 & 1 & 0 \\ 0 & 0 & 1 \end{bmatrix}.$$

The matrix R^{-1} is called the *left inverse* of R . However, even a *right inverse*, that is, a matrix Q such that

$$RQ = I$$

will do. This is because for *any* square matrix A , if the matrix B is the right inverse of A , that is, if $AB = I$, then B is also the left inverse:

$$BA = I.$$

The proof of this fact is a single line: suppose that B is the right inverse of A and the left inverse is a matrix C , so that $CA = I$. Then

$$C = CI = C(AB) = (CA)B = IB = B,$$

which forces us to conclude that B and C are the same matrix. So we can drop “left” or “right” and merely say *inverse*.

The inverse R^{-1} of the rotation matrix R is more easily found by looking for a right inverse. The three vectors $\mathbf{i}, \mathbf{j}, \mathbf{k}$ that make up the rows of R have unit norm,

$$\mathbf{i}^T \mathbf{i} = \mathbf{j}^T \mathbf{j} = \mathbf{k}^T \mathbf{k} = 1,$$

and are mutually orthogonal:

$$\mathbf{i}^T \mathbf{j} = \mathbf{j}^T \mathbf{k} = \mathbf{k}^T \mathbf{i} = 0.$$

Because of this,

$$RR^T = \begin{bmatrix} \mathbf{i}^T \\ \mathbf{j}^T \\ \mathbf{k}^T \end{bmatrix} \begin{bmatrix} \mathbf{i} & \mathbf{j} & \mathbf{k} \end{bmatrix} = \begin{bmatrix} \mathbf{i}^T \mathbf{i} & \mathbf{i}^T \mathbf{j} & \mathbf{i}^T \mathbf{k} \\ \mathbf{j}^T \mathbf{i} & \mathbf{j}^T \mathbf{j} & \mathbf{j}^T \mathbf{k} \\ \mathbf{k}^T \mathbf{i} & \mathbf{k}^T \mathbf{j} & \mathbf{k}^T \mathbf{k} \end{bmatrix} = \begin{bmatrix} 1 & 0 & 0 \\ 0 & 1 & 0 \\ 0 & 0 & 1 \end{bmatrix}$$

as promised.