

Interpretation of X-ray Spectra of Supernova Remnants

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INTRODUCTION

Spectral model fits to observed spectra of hot (X-ray emitting) plasmas provide us crucial information on the state and amount of hot plasma. The spectral fits inform us of the elemental composition of the plasma and its temperature and, in the case of non-equilibrium ionization plasmas, tell us the age of the plasma by the parameter $\tau = ne t$, with t the age since the plasma was shock heated.

The most common modeling software for X-ray spectra is the XSPEC software package (Arnaud 1996). The purpose of the current work is to explain the assumptions and provide corrections for cases where the assumptions are not accurate.

Definition of emission measure (EM) and definition of abundance factors (A_Z) given below, with A_{sol} a solar abundance set:

$$EM = \int n_e(\rho) dV = n_e^2 V \quad (\text{if uniform})$$

$$EM_Z = \int n_Z(\rho) dV = n_Z n_H V \quad (\text{if uniform})$$

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$$n_H = A_1 n_{H0} \quad n_Z = A_Z A_{H0} n_{H0}$$

For XSPEC, the value of the electron density, n_{eX} , is tied to the hydrogen density, n_{HX} , by:

$$n_{eX} = 1.2 n_{HX}$$

ANALYSIS

The variation with solar abundance set is small (~3.5%). However the variation of the electron density between the different abundance sets (A_Z) is large, with a number of examples given in Table 1. The ionization state and composition gives the ratio ne/t . The variation is nearly a factor of 1000, instead of the single value assumed by XSPEC.

Table 1. Electron-to-(fiducial)hydrogen density ratio n_e/n_{H0} for different cases of composition and different temperatures.

KT(kV)	0.10	0.25	0.50	1.0	2.0	4.0	10.0	40.0
COMPOSITION								
SOLAR CASES ^a								
avg ^b	1.20532	1.20652	1.20760	1.20836	1.20866	1.20880	1.20885	1.20886
sepl	1.17672	1.17756	1.17829	1.17871	1.17903	1.17907	1.17907	1.17908
fsd	1.20536	1.20647	1.20765	1.20829	1.20856	1.20867	1.20872	1.20873
asob	1.16981	1.17092	1.17198	1.17258	1.17297	1.17302	1.17303	1.17303
gna	1.17861	1.17962	1.18058	1.18133	1.18149	1.18154	1.18154	1.18155
wdm	1.20152	1.20224	1.20295	1.20333	1.20354	1.20362	1.20365	1.20366
lchd	1.18470	1.18552	1.18625	1.18666	1.18696	1.18709	1.18710	1.18710
lqsp	1.17596	1.17689	1.17768	1.17771	1.17776	1.17776	1.17775	1.17775
lqsp	1.20159	1.20254	1.20342	1.20394	1.20420	1.20431	1.20436	1.20437
NON-SOLAR CASES								
H 0 ^b	0.20532	0.20652	0.20760	0.20836	0.20866	0.20880	0.20885	0.20886
(H/He) 0.1 ^b	0.12046	0.12066	0.12183	0.12250	0.12280	0.12294	0.12299	0.12300
(H/He) 0 ^b	0.00992	0.01112	0.01229	0.01296	0.01326	0.01340	0.01345	0.01346
(H to N) 0 ^b	0.07143	0.08083	0.09005	0.09800	0.10100	0.10143	0.10149	0.10150
(H to Al) 0 ^b	0.00093	0.00142	0.00155	0.00168	0.00183	0.00192	0.00195	0.00196
N100Fe ^c	0.00108	0.00157	0.00172	0.00186	0.00202	0.00211	0.00215	0.00215
N100Fe+HHe ^d	1.1965	1.1970	1.1971	1.1971	1.1971	1.1971	1.1971	1.1971
N100Fe+10 ^e	0.01080	0.01374	0.01720	0.01862	0.02021	0.02112	0.02150	0.02152
N100Fe+10HHe ^f	1.2892	1.2111	1.2126	1.2140	1.2156	1.2165	1.2169	1.2169

^aAbundance references given in the XSPEC manual- see the abund command.

^bAll other elements taken to have solar abundance factors, $A_Z = 1$.

^cN100 model (Seitzmuhl et al. 2013) with $A_{Fe}=1$.

^dN100 model with $A_{Fe}=1$, H He added at solar abundance.

^eN100 model with $A_{Fe}=10$.

^fN100 model with $A_{Fe}=10$, H He added at solar abundance.

To simplify notation, hereafter we define the electron density ratio as r_e , given by:

$$r_e = \frac{n_e}{1.2 n_{H0}} = \text{ratio of electron density to XSPEC value}$$

the fiducial density is given in terms of $norm_{HX}$ by:

$$n_{H0} = r_e^{-1/2} \sqrt{10^{14} norm_{HX} \frac{4 \pi P_X}{1.2 V}}$$

The XSPEC norm for hot plasma emission models is proportional to EM_{HX} :

$$norm_{HX} = 10^{-14} EM_{HX} / (4 \pi D^2 (1+z)^2) \quad \text{with} \quad EM_{HX} = \int n_e(\rho) n_H(\rho) dV = n_e n_H V \quad (\text{if uniform})$$

From the detailed calculation (Leahy, Foster, Seitzmuhl, 2024) we show the correct density of the plasma is given by:

$$n_H = A_1 r_e^{-1/2} n_{HX}$$

$$n_Z = A_Z A_{H0} r_e^{-1/2} n_{HX}$$

Element masses depend on the element densities and emitting volume,

$$M_Z = n_Z \int n_H(\rho) dV = n_Z n_H V \quad (\text{if uniform}) \\ = n_Z EM_{HX} / n_e = n_Z A_{H0} A_Z EM_{HX} / n_e \\ = n_Z A_{H0} A_Z \sqrt{10^{14} norm_{HX} \frac{4 \pi P_X}{1.2 V}}$$

RESULTS

First, we review the contributions of various elements to the observed X-ray spectrum. In general, elements that strongly contribute (like Fe and Ni) are tightly constrained, but weakly contributing elements like H, He, Li, Be, B, C are poorly constrained, as shown in Fig.1 below.

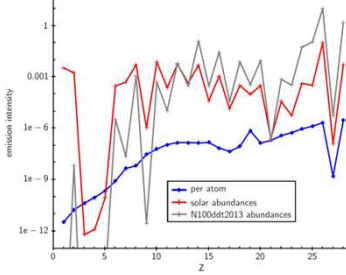


Figure 1. Illustration of the relative contributions of different elements to the 0.3 to 8 keV energy range for a spectrum, using the response matrices for the AstroSat SXT instrument (Singh et al. 2017). The maximum intensity in photons $\text{cm}^{-2} \text{s}^{-1} \text{keV}^{-1}$ is shown for elements $Z = 1$ to 28 for three cases: per atom, for solar abundances and for N100dnt2013 abundances. For each case the vertical scale is arbitrary. The XSPEC vnpw model was used with temperature of 1 keV. For N100dnt2013, H, Li, Be and B not shown because they are so small: they at 2.9×10^{-12} , 1.9×10^{-12} , 4.1×10^{-12} , and 1.2×10^{-12} respectively times the contribution of Fe.

Because the abundances of H and He are normally very large, they can contribute significantly to the X-rays spectrum. This is illustrated in Fig.2 below. I.e. H abundance needs to be 30 times solar for H to contribute equally to Fe. Because solar abundance of some elements is very low (e.g. Li) the abundance relative to solar to contribute equally as Fe is large ($\sim 10^{11}$).

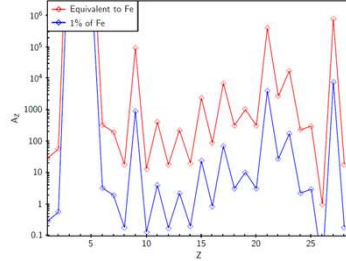


Figure 2. The abundance factors A_Z required to yield the same maximum contribution (red line) or to yield a contribution 1% as large (blue line) as the contribution from Fe to the 0.3 to 8 keV X-ray spectrum. The response matrices for the AstroSat SXT instrument was used and an equilibrium spectrum with temperature $kT=2$ keV and the abundance factor of Fe was set to 1 for the same-contribution line (or 0.01 for the 1% line). The abundance factors for Li, Be and B off scale, at values of 1.5×10^{11} , 7.8×10^{10} , and 1.1×10^{10} for equal contribution, with values 0.01 times as large for a 1% contribution.

We carried out a number of simulations in XSPEC simulating observed spectra using the AstroSat SXT response matrices. These were carried out for typical SNR fluxes, based on known type Ia SNRs with observed X-ray spectra, shown in Table 3.

Table 3. Measured temperatures and EM for type Ia SNRs^a

SNR ID	type	D(kpc)	EM(10^{50}cm^{-3})	norm	keV(kV)	Age(yr)	density(cm^{-3})	n_e (cm^{-3}) ^b
G536-2.2	Ia	7.8	0.022	3.62E-04	3.9	4600	0.096	5.41E+11
G299.2-2.9	Ia	5	0.029	9.70E-04	1.36	8900	0.022	2.45E+10
G358.5-0.9	Ia	30	1.8	3.76E-03	1.31	12800	2.3	4.00E+12
G313.4-2.3	Ia	3.8	0.096	4.91E-03	3.04	11000	0.317	3.62E+11
G252.7-0.1	ke ^c	7.5	0.11	1.64E-03	3.2	7600	1.66	1.50E+12
average:				2.32E-03	2.602			1.31E+12

^aValues are from Table 1 and 2 of Leahy et al. (2020), with references for EM, norm and kT given in Table 1.

^bValues for age and density are from Table 2 of Leahy et al. (2020), with n_e calculated using post-shock density of ρ_{sh} .

2 of the simulated X-ray spectra and their spectral fits are shown in Fig. 3. The results from a set of simulated spectra and fits to those spectra are shown in Table 4. These verify that the abundances A_Z derived from spectral fitting are quite uncertain, and the spectra are fit with significant changes in the H and He

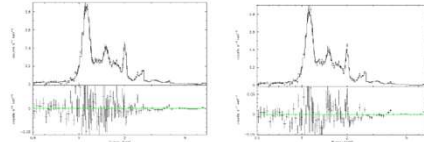


Figure 3. The HSN (high signal-to-noise, see text) simulated X-ray spectrum, which uses rmf, arf and background files for AstroSat SXT, and abundances given by column 8 in Table 2 and model fits. The other spectrum parameters are $N_H = 10^{21} \text{cm}^{-2}$, $kT=2.6$ keV, $n_e t = 1.31 \times 10^{12} \text{cm}^{-3} \text{s}$ and $norm=2.32 \times 10^{-2}$. The fit model on the left has original N100 abundances, whereas the fit model on the right has added H and He with $A_H = A_{He} = 1$. Although the fit with solar H and He is statistically poor (probability=1.6E-06, Table 4), it is not clearly different except for larger residuals by a factor of ~1.2.

Table 4. Fits to simulated spectra with added H and He^a

Spectrum	A_H	A_{He}	χ^2	probability	A_H	A_{He}	χ^2	probability
LSN	2.6E-96.75E-05	110.4	0.167	1.1	100.8	0.177		
	4.4	111.2	0.153	8.8	114.4	0.199		
	86.16	119.2	0.063	24.24	122.4	0.042		
	32.32	124.6	0.031	40.40	122.4	0.025		
MSN	2.6E-96.75E-05	90.6	0.662	1.1	105.6	0.258		
	1.41.4	113.1	0.127	1.5,1.8	119.2	0.063		
	2.2.2.2	124.7	0.033	2.6,2.6	128.5	0.018		
HSN	2.6E-96.75E-05	93.2	0.590	2.6E-16.1	106.8	0.233		
	1.7.5E-05	142.4	1.0E-03	1.1	176.1	1.6E-06		
	0.1.0.1	96.0	0.509	0.2,0.2	100.3	0.290		
	0.3,0.3	106.7	0.236	0.4,0.4	114.4	0.109		
	0.5,0.5	123.0	0.058					

^aLSN has exposure 10ks, norm 2.32×10^{-2} ; MSN has exposure 20ks, norm 6.96×10^{-3} ; HSN has exposure 20ks, norm 2.32×10^{-2} . All have $kT=2.6$ keV, $n_e t = 1.31 \times 10^{12}$ and N100 abundances from column 8 of Table 2. H and He are added to the fit model, as listed.

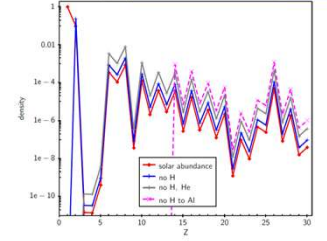


Figure 4. Element densities for different cases. The four cases shown are solar abundance, no H, no H and He, and no H through Al.

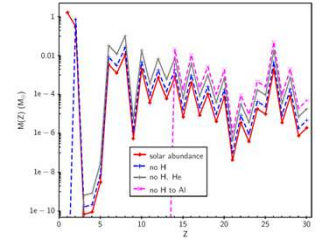


Figure 6. Masses for different cases. For calculation, the distance is 1 kpc, the emitting volume is a sphere of radius 2 pc, and the XSPEC norm is 0.1. The four cases shown are solar abundance, no H, no H and He, and no H through Al.

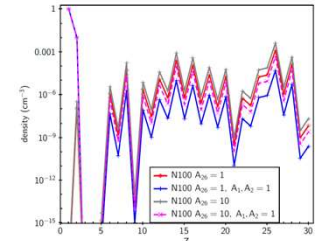


Figure 7. Element densities for N100 model (Seitzmuhl et al. 2013), with and without added H and He. For calculation, the distance is 1 kpc, the emitting volume is a sphere of radius 2 pc, and the XSPEC norm is 0.1, which yields XSPEC $EM_X = 1.2 \times 10^{50} \text{cm}^{-3}$. The four cases shown are N100 abundances normalized with $A_{Fe}=1$, N100 with $A_{Fe}=1$ and H and He added at solar abundances, N100 abundances normalized with $A_{Fe}=10$, and N100 with $A_{Fe}=10$ and H and He added at solar abundances.

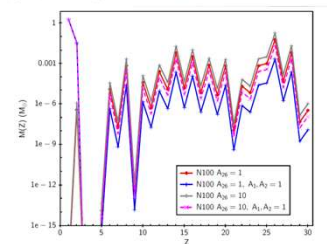


Figure 8. Masses for different cases. For calculation, the distance is 1 kpc, the emitting volume is a sphere of radius 2 pc, and the XSPEC norm is 0.1. The four cases shown are N100 abundances normalized with $A_{Fe}=1$, N100 with $A_{Fe}=1$ and H and He added at solar abundances, N100 abundances normalized with $A_{Fe}=10$, and N100 with $A_{Fe}=10$ and H and He added at solar abundances.

CONCLUSIONS/FUTURE WORK

For cases of non-solar abundances, with A_1 or A_2 different from 1, significant corrections are needed to the element densities to account for the difference between n_e and n_{eX} . n_e/n_{eX} can be much less than 1 as illustrated by the examples in Table 1. We define the electron density ratio r_e . The corrected densities are higher, by the factor $r_e^{-1/2}$, and the inferred element masses are corrected by the same factor ($r_e^{-1/2}$) thus can also be much larger. This implies that masses that may have been derived in the past using norm's from XSPEC spectral fitting on hydrogen poor plasmas, such as expected for Type Ia SNRs, may be significantly underestimated.

Reference: Leahy, Foster, Seitzmuhl 2024, ApJ