

Properties of the lowest mass Fe-core collapse supernovae

H. Umeda & K. Ishiguro (Univ. of Tokyo)



We calculate progenitors for Fe-core collapse SNe in a fine grid to investigate properties of low mass CCSNe. Treating the explosion energy as a free parameter, we attempt to constrain both the explosion energy and the ejected ^{56}Ni mass using observational constraints on neutron star mass and on chemical abundances in metal-poor stars.

Lowest Mass CCSN: What is ... ?

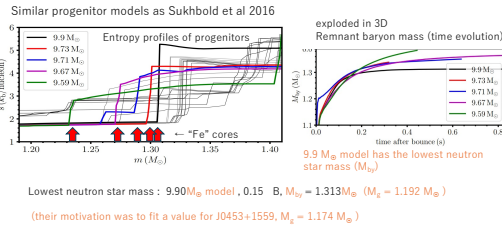
- the mass ? the neutron star mass for it ? the explosion energy ?
 - ^{56}Ni mass (brightness of the SN) ? the nucleosynthesis ?
 - the observed properties of light curves (different from ECSN) ?
- If it can explode in 1D simulation (as an ECSN)?

Previous work 1

Sukhbold et al 2016, ApJ : 9 to 120 M_{\odot} CCSN
Exploded in 1D using "calibrated central engines"
For the lowest mass models, an engine for ECSNe is used
The lowest mass
9 M_{\odot} , 0.006 M_{\odot} (^{56}Ni mass), 0.11 B (explosion energy), 1.35 M_{\odot} (remnant mass),
cf. (since Müller et al 2024 found the lowest remnant mass for 9.9 M_{\odot} model)
10 M_{\odot} , 0.031 M_{\odot} (^{56}Ni mass), 0.60 B (explosion energy), 1.45 M_{\odot} (remnant mass),

Previous work 2

Müller et al 2025, PhRvL :
The minimum NS mass in neutrino-driven SN



Method

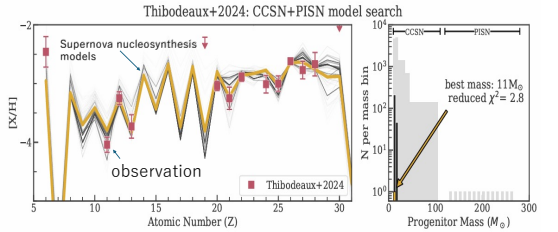
We use the HOSHI stellar evolution code (Takahashi, Yoshida, Umeda 2018)
Explosion: 1D PPM code with a thermal bomb (same as our previous work)
However, Mass cut is determined gravitationally, i.e., self-consistently, by inserting explosion energy deeper inside than before.

Previously we had two parameters (explosion energy and mass cut), but in this work just one parameter (explosion energy).

We constrain explosion energy by neutron star mass, and nucleosynthesis

For example, the abundance pattern of a metal poor star J1010+2358 is used, since this star might have an abundance pattern of the lowest mass CCSN.

Thibodeaux et al. (2024) \Rightarrow new abundance with non-LTE abundance



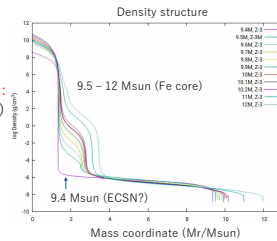
\Rightarrow The lowest mass CCSN model fits best to J1010+2358.

Our progenitor models

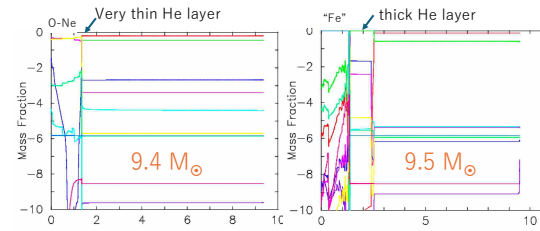
We consider $10^{-3} Z_{\odot}$ models because we would like to compare with a metal poor star J1010+2358.

The lowest mass for Fe-core collapse :
9.5 M_{\odot} (9.4 M_{\odot} forms an ONeMg core)

c.f., Z_{\odot} model : 9.6 M_{\odot} (lowest)
(9.5 M_{\odot} forms an ONeMg core)



Our progenitor models (Internal abundance)



Results: Explosion energy and NS mass for the exploding lowest mass model

| $10^{-3} Z_{\odot}$, 10.1 M_{\odot} | E _{ex} (B) | M _{ns} (baryon mass) | ejected ^{56}Ni mass (M_{\odot}) |
|--|---------------------|-------------------------------|---|
| | 0.089 | 1.482 | 0 |
| | 0.15 | 1.350 | 0.023 |
| | 0.176 | 1.330 | 0.040 |
| | 0.20 | 1.328 | 0.043 |
| | 0.22 | < 1.29 | |

| E _{ex} (B) | M _{ns} (remnant mass) | ejected ^{56}Ni mass |
|---------------------|--------------------------------|-------------------------------|
| 0.089 | 1.482 | 0 |
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Explosion energy should be smaller than 0.22B for M_{ns} > 1.3 Msun to avoid very small NS small

M \sim 10.0 Msun models require small explosion energies E_{ex} < 0.1 B, (which is consistent with Müller's 3D simulation results) and difficult to eject significant amount of ^{56}Ni to explain J1010+2358 abundance (though they may exist in nature)

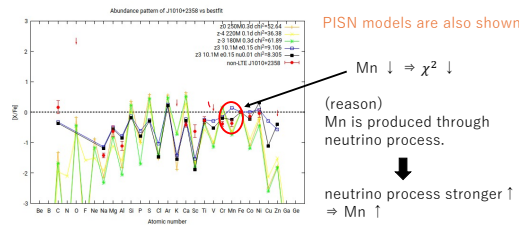
Artificially set mass-cut and ejected ^{56}Ni mass

Summary and references

- Best fit to a VMP star J1010+2358 is obtained for a metal poor CCSN model with initial mass 10.1 M_{\odot} , E_{exp}=0.15B, M(^{56}Ni)=0.048 M_{\odot} , and weak neutrino process.
- Low mass CCSNe have small E_{exp} and ^{56}Ni ejection.
- \sim 9.5-10.0 M_{\odot} can be dark SNe II with small E_{exp} < 0.1B without ^{56}Ni ejection.
- It will be interesting to study these models further (nucleosynthesis and light curves) with detailed 2D & 3D simulations.

Comparison with J1010+2358

Ishiguro Master's thesis (2025)



PISN models are also shown

Mn $\downarrow \Rightarrow \chi^2 \downarrow$
(reason)
Mn is produced through neutrino process.
 \downarrow
neutrino process stronger $\uparrow \Rightarrow$ Mn \uparrow

Fitting and χ^2 values: (Ishiguro Master's thesis 2025)

Here M_{cut} is also used as a fitting parameter

RED: Best Fit

- We also take the strength of neutrino nucleosynthesis process as a parameter,
nu=1: relatively strong
nu=0.01: basically off

| Model | Best fit : Mass cut and ejected ^{56}Ni mass (M_{\odot}) | χ^2 |
|---------------------------------|---|----------|
| 10.1 M_{\odot} e0.089 nu=1 | 1.3229 | 13.02 |
| 10.1 M_{\odot} e0.089 nu=0.01 | 1.3114 | 10.71 |
| 10.1 M_{\odot} e0.15 nu=1 | 1.3229 (0.044) | 9.106 |
| 10.1 M_{\odot} e0.15 nu=0.01 | 1.3157 (0.048) | 8.305 |
| 10.1 M_{\odot} e0.22 nu=0.01 | 1.2370 | 38.14 |
| 11 M_{\odot} e0.092 nu=1 | 1.0942 | 41.22 |
| 11 M_{\odot} e0.092 nu=0.01 | 1.0942 | 27.92 |
| 11 M_{\odot} e0.17 nu=1 | 1.0942 | 36.83 |
| 11 M_{\odot} e1.07 nu=1 | 1.2026 | 17.20 |
| 13 M_{\odot} e0.34 nu=1 | 1.4288 | 15.22 |

Discussions:

Best fit model for J1010+2358

- 10.1 M_{\odot} , 0.15 B, NS mass 1.3157 M_{\odot} , ^{56}Ni 0.048 M_{\odot}
- The NS mass is smaller and ^{56}Ni mass is larger than the "self-consistent" model, M_{rem} = 1.350 and ^{56}Ni = 0.023 M_{\odot}
- This suggests that the explosion energy should be slightly larger than 0.15B.
- For this best fit-model, the baryon mass 1.3157 M_{\odot} corresponds to the gravitational mass 1.196 M_{\odot} , which is close to the observational lower limit M_g = 1.174 M_{\odot} for J0453+1559 (though this particular NS may need binary interaction model to explain).

- We plan to apply these low mass progenitor models to 2D and 3D explosion simulations (just started).

| Upper limit of Explosion energy | Z _g model (mass) | Upper E _{ex} (B) |
|---------------------------------|-----------------------------|---------------------------|
| | 9.7 | 0.115 |
| | 10.0 | 0.62 |
| | 12.0 | 0.43 |
| | 13.0 | 0.69 |
| | 15.0 | 0.77 |
| | 18.0 | >>> (>1.04) |

This energy seems to be smaller than conventional wisdom
 \sim E_{ex} of typical SN II should be around 1B

- This work (to be submitted)
- Sukhbold et al. 2016, ApJ 821, 38
- Müller et al. 2025, PhRvL 134, 1403
- Takahashi et al. 2018, ApJ, 857, 111
- Thibodeaux et al. 2024, OJAp 7, id66