Phasor-Domain Equations

Transmission line equations — time domain:

$$-\frac{\partial v(z,t)}{\partial z} = R'i(z,t) + L'\frac{\partial i(z,t)}{\partial t}$$
$$-\frac{\partial i(z,t)}{\partial z} = G'v(z,t) + C'\frac{\partial v(z,t)}{\partial t}$$

Time domain → Phasor domain

$$v(z,t) = \text{Re} \left[\tilde{V}(z) \exp(j\omega t) \right], \quad i(z,t) = \text{Re} \left[\tilde{I}(z) \exp(j\omega t) \right]$$

Transmission line equations — phasor domain:



$$-\frac{d\tilde{V}(z)}{dz} = (R' + j\omega L')\tilde{I}(z)$$
$$-\frac{d\tilde{I}(z)}{dz} = (G' + j\omega C')\tilde{V}(z)$$

4.1

Phasor-Domain Equations

Advantages of Phasor-Domain Representation

- Reduction from partial differential equation to ordinary differential equation
 - so that it is easier to solve
- Allows generalization: $L' \to L'(\omega)$, $C' \to C'(\omega)$,

$$R' \to R'(\omega), \quad G' \to G'(\omega)$$

Advantages of Time-Domain Representation

• Allows study of transients



Phasor-Domain Equations

Second-Order Equations:

$$\frac{d^2\tilde{V}(z)}{dz^2} - \gamma^2 \tilde{V}(z) = 0, \quad \frac{d^2\tilde{I}(z)}{dz^2} - \gamma^2 \tilde{I}(z) = 0$$

with

$$\gamma = \alpha + j\beta = \sqrt{(R' + j\omega L')(G' + j\omega C')}$$

 $\gamma = \text{complex propagation constant}$

 $\alpha = \text{Re}(\gamma) = \text{attenuation constant (Np/m)}$

 $\beta = \text{Im}(\gamma) = \text{phase constant or wavenumber (rad/m)}$

- NOTES: We pick α and β so that both are positive
 - In a passive medium, α is always positive; it can be negative in an active medium



[An active medium, like a laser, has an energy source; a passive medium does not and must always lose energy.]

4.3

Phasor-Domain Equations

The second-order equations have the general solutions

$$\tilde{V}(z) = V_0^+ \exp(-\gamma z) + V_0^- \exp(\gamma z), \quad \tilde{I}(z) = I_0^+ \exp(-\gamma z) + I_0^- \exp(\gamma z)$$

- $\tilde{V}(z)$ and $\tilde{I}(z)$ are not independent
- + V_0^+ and V_0^- are arbitrary constants, while I_0^+ and I_0^- are not
- We may also write:

$$\tilde{I}(z) = \frac{V_0^+}{Z_0} \exp(-\gamma z) - \frac{V_0^-}{Z_0} \exp(\gamma z)$$

where

$$Z_0 = \sqrt{\frac{R' + j\omega L'}{G' + j\omega C'}}$$



is the characteristic impedance

NOTE: Z_0 (Ulaby et al.'s notation) $\rightarrow Z_C$ (Paul's notation)

Phasor-Domain Equations

Returning to the time domain:

We first write

$$V_0^+ = |V_0^+| \exp(j\phi^+), \quad V_0^- = |V_0^-| \exp(j\phi^-)$$

from which we obtain

$$v(z,t) = \operatorname{Re}\left[\tilde{V}(z)\exp(j\omega t)\right]$$

$$= \operatorname{Re}\left\{\left[V_0^+\exp(-\gamma z) + V_0^-\exp(\gamma z)\right]\exp(j\omega t)\right\}$$

$$= \operatorname{Re}\left\{\left|V_0^+\exp(j\phi^+)\exp(j\omega t)\exp(-(\alpha + j\beta)z)\right|\right\}$$

$$+\left|V_0^-\exp(j\phi^-)\exp(j\omega t)\exp((\alpha + j\beta)z)\right\}$$

$$= \left|V_0^+\exp(-\alpha z)\cos(\omega t - \beta z + \phi^+)\right| + \left|V_0^-\exp(\alpha z)\cos(\omega t + \beta z + \phi^-)\right|$$
forward propagation with attenuation backward propagation with attenuation



Phase velocity = $u_p = \omega/\beta$

4 5

Phasor-Domain Equations

Example: Ulaby et al. Exercise 2.4

Question: A two-wire air line has the following parameters: R' = 0.404 m Ω/m , $L' = 2.00 \mu H/m$, G' = 0, C' = 5.56 pF/m. For operation at 5 kHz, determine (a) the attenuation coefficient α , (b) the wavenumber β , (c) the phase velocity $u_{\rm p}$, and the characteristic impedance Z_0 .

Answer: $\omega = 2 \times 3.14159 \times (5 \times 10^3 \text{ s}^{-1}) = 3.14159 \times 10^4 \text{ s}^{-1}$. $R' + j\omega L' = (4.04000 \times 10^{-4} + j \times 3.14159 \times 10^4 \times 2.00000 \times 10^{-6}) \ \Omega/\text{m} = (4.04000 \times 10^{-4} + j \times 6.28318 \times 10^{-2}) \ \Omega/\text{m} = 6.28319 \times 10^{-2} \times \exp(j \times 1.56436) \ \Omega/\text{m}$. $G' + j\omega C' = j \times 1.74673 \times 10^{-7} \ \Omega^{-1}/\text{m} = 1.74673 \times 10^{-7} \times \exp(j \times 1.57080) \ \Omega^{-1}/\text{m}$. Note the small difference in phases! Six digits of accuracy are needed to keep three digits in the attenuation coefficient. $\gamma^2 = 1.09750 \times 10^{-8} \times \exp(j \times 3.13156) \ \text{m}^{-2}$, so that $\gamma = 1.04762 \times 10^{-4} \times \exp(j \times 3.13156) \ \text{m}^{-2} \times 1.04761 \times 10^{-4} \times 10^{$

Specializing to the case of no loss (R' = 0, G' = 0):

$$\alpha=0,\ \beta=\omega\sqrt{L'C'},\ Z_0=\sqrt{L'/C'},\ u_{\rm p}=1/\sqrt{L'C'}$$

We will later show that for any TEM transmission line:

$$L'C' = \mu\varepsilon$$

 $\mu = \text{magnetic permeability}$

 ε = electrical permittivity

- For any insulating material that would be used in a transmission line, $\mu = \mu_0 = 4\pi \times 10^{-7}$ H/m, where μ_0 is the vacuum permeability. — The value for μ_0 is exact and defines the relation between **B** and **H**.
- By contrast, the values for ε differ significantly for different materials. We write: $\varepsilon = \varepsilon_r \varepsilon_0$, where ε_r is referred to as the relative permittivity, and $\varepsilon_0 = 1/\mu_0 c^2 \simeq 8.854 \times 10^{-12}$ F/m is the vacuum permittivity



4.7

Lossless Transmission Line

We now have:

$$\beta = \omega \sqrt{\mu \varepsilon} = \omega \sqrt{\mu_0 \varepsilon_0 \varepsilon_r} = (\omega/c) \sqrt{\varepsilon_r}, \quad u_p = 1/\sqrt{\mu \varepsilon} = c/\sqrt{\varepsilon_r}$$

In optics, we define an index of refraction, $n = c / u_p$, so that $n = \sqrt{\varepsilon_p}$



The impedance relation:

$$Z_0 = \sqrt{L'/C'} = \sqrt{\mu_0/\varepsilon_0\varepsilon_r} K = \sqrt{\mu_0/\varepsilon_0} \left(K/\sqrt{\varepsilon_r} \right); \ 377 \left(K/\sqrt{\varepsilon_r} \right) \Omega$$

is a bit more complex. It involves a geometric factor *K*.

Table of Geometric Factor K

	Coaxial	Two wire	Parallel plane
K	$\frac{1}{2\pi}\ln(b/a)$	$\frac{1}{\pi} \ln \left[(d/2a) + \sqrt{(d/2a)^2 - 1} \right]$	$\frac{d}{w}$

The vacuum impedance of 377 Ω is a very important number



- Loads with lower impedance have large magnetic near fields
- Loads with higher impedance have large electric near fields

EMI properties are very different in the two cases!

4.9

Lossless Transmission Line

Voltage Reflection:

When $\alpha = 0$, our phasor relations become

$$\tilde{V}(z) = V_0^+ \exp(-j\beta z) + V_0^- \exp(j\beta z),$$

$$\tilde{I}(z) = \frac{V_0^+}{Z_0} \exp(-j\beta z) - \frac{V_0^-}{Z_0} \exp(j\beta z)$$

At the load, we have

$$Z_{\rm L} = \frac{\tilde{V}_{\rm L}}{\tilde{I}_{\rm L}}$$
, with $\tilde{V}_{\rm L} = V_0^+ + V_0^-$ and $\tilde{I}_{\rm L} = \frac{V_0^+}{Z_0} - \frac{V_0^-}{Z_0}$

which implies

$$V_0^- = \frac{Z_{\rm L} - Z_0}{Z_{\rm L} + Z_0} V_0^+ = \Gamma V_0^+$$



and defines the voltage reflection coefficient Γ

Load impedances and reflection coefficients are usually complex!

Example: Ulaby et al. Exercise 2.7

Question: A 50 Ω lossless transmission line is terminated in a load impedance $Z_L = (30 - j200) \Omega$. Calculate the voltage reflection coefficient at the load.

Answer:
$$\Gamma = (Z_L - Z_0) / (Z_L + Z_0) = (30 - j200 - 50) / (30 - j200 + 50) = (-20 - j200) / (80 - j200) = 201 \exp(-j1.67) / 215 \exp(-j1.19) = 0.93 \exp(-j0.48)$$
. NOTE: -0.48 rads $= -28^\circ$. Writing $\Gamma = |\Gamma| \exp(j\theta_r)$, we find $\theta_r = -0.48$ rads.



4.11

Lossless Transmission Line

Standing Waves

After using the relation $V_0^- = \Gamma V_0^+$ in the phasor equations,

$$\tilde{V}(z) = V_0^+ \Big[\exp(-j\beta z) + \Gamma \exp(j\beta z) \Big] = V_0^+ \Big[\exp(-j\beta z) + |\Gamma| \exp(j\theta_{\rm r} + j\beta z) \Big],$$

$$\tilde{I}(z) = \frac{V_0^+}{Z_0} \left[\exp(-j\beta z) - \Gamma \exp(j\beta z) \right] = \frac{V_0^+}{Z_0} \left[\exp(-j\beta z) - |\Gamma| \exp(j\theta_r + j\beta z) \right]$$

We then find

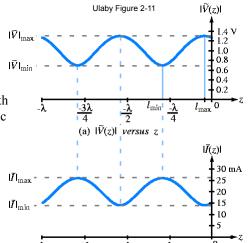
$$|\tilde{V}(z)| = |V_0^+| \left[1 + |\Gamma|^2 + 2|\Gamma|\cos(2\beta z + \theta_{\rm r}) \right]^{1/2},$$

so that the amplitude of the voltage varies sinusoidally with z.

- This pattern is called a *standing wave*.
- It comes from the interference of forward- and backward-propagating waves

Standing Waves

- The voltage and current maxima and minima are 180° out of phase
- The amplitude multiplies a cos(ωt) dependence with a complicated but periodic z-variation.
- The maxima are spaced λ/2 apart, and so are the minima.



(b) $|\tilde{I}(z)|$ versus z

Figure Parameters:



Figure 1 at a meter
$$|\Gamma| = 0.3$$
, $\theta_r = 30^\circ$ $Z_0 = 50 \Omega$ $|V_0^+| = 1.3 \text{ V}$

Lossless Transmission Line

Standing Waves

The total dependence of the standing wave voltage is

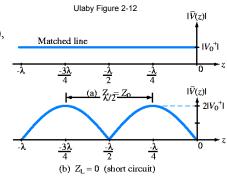
$$v(z,t) = |V_0^+| \left[1 + |\Gamma|^2 + 2|\Gamma| \cos(2\beta z + \theta_r) \right]^{1/2} \\ \times \cos \left\{ \omega t - \tan^{-1} \left[\frac{\sin(\beta z - \phi^+) - |\Gamma| \sin(\beta z + \theta_r + \phi^+)}{\cos(\beta z - \phi^+) + |\Gamma| \cos(\beta z + \theta_r + \phi^+)} \right] \right\}$$

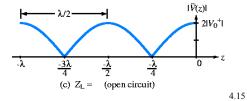


4.14

Standing Waves

- For a matched load $(Z_L = Z_0)$, there is no standing wave
- For a short-circuited load $(Z_{\rm L}=0)$ or an open-circuited load $(Z_{\rm L}=\infty)$, there are complete reflections and an oscillation depth of 100%
- With $|\Gamma| = 1$, there are points where the voltage is exactly zero, spaced $\lambda/2$ apart





OIVIDC AN HONORS UNIVERSITY

UNIVERSITY In Maryland

Lossless Transmission Line

Standing Waves

• We will designate the location of the maxima as $l_{\text{max}} = -z$, so that l_{max} is a positive number (since the load is at z = 0).

$$-z = l_{\text{max}} = \frac{\theta_{\text{r}} + 2n\pi}{2\beta} = \frac{\theta_{\text{r}}\lambda}{4\pi} + \frac{n\lambda}{2}$$

Only *n*-values that satisfy $l_{\text{max}} \ge 0$ are allowed

 The voltage standing wave ratio (VSWR or SWR) gives the oscillation depth. It is defined:

$$S = \frac{|\tilde{V}|_{\text{max}}}{|\tilde{V}|_{\text{min}}} = \frac{1 + |\Gamma|}{1 - |\Gamma|}$$



Example: Ulaby et al. Exercise 2.11

Question: A 140 Ω lossless line is terminated in a load impedance $Z_L = (280 + j \, 182) \, \Omega$. If $\lambda = 72$ cm, find (a) the reflection coefficient Γ , (b) the VSWR S, (c) the locations of the voltage maxima and minima.

Answer: (a) $\Gamma = (Z_L - Z_0) / (Z_L + Z_0) = (280 + j \cdot 182 - 140) / (280 + j \cdot 182 + 140) = 230 \exp(j \cdot 0.915) / 458 \exp(j \cdot 0.409) = 0.50 \exp(j \cdot 0.51)$. NOTE: 0.51 rads = 29°. (b) $S = (1 + |\Gamma|) / (1 - |\Gamma|) = (1 + 0.502) / (1 - 0.502) = 3.0$. (c) $l_{\text{max}} = (0.506 \times 72 / 4\pi + n \times 72 / 2) \text{ cm} = 2.9 \text{ cm}, 39 \text{ cm}, 75 \text{ cm}, ...;$ $l_{\text{min}} = (0.506 \times 72 / 4\pi + 18 + n \times 72 / 2) \text{ cm} = 21 \text{ cm}, 57 \text{ cm}, 93 \text{ cm}, ...$



4.17

Lossless Transmission Line

Input Impedance

With standing waves, the voltage-to-current ratio, which is referred to as the *input impedance*, varies as a function of position

$$Z_{\text{in}}(z) = \frac{\tilde{V}(z)}{\tilde{I}(z)} = \frac{V_0^+ \left[\exp(-j\beta z) + \Gamma \exp(j\beta z) \right]}{V_0^+ \left[\exp(-j\beta z) - \Gamma \exp(j\beta z) \right]} Z_0 = Z_0 \frac{\left[1 + \Gamma \exp(2j\beta z) \right]}{\left[1 - \Gamma \exp(2j\beta z) \right]},$$

Of particular interest is the input impedance at the generator, z = -l

$$\begin{split} Z_{\text{in}}(-l) &= Z_0 \left[\frac{1 + \Gamma \exp(-2j\beta l)}{1 - \Gamma \exp(-2j\beta l)} \right] \\ &= Z_0 \left(\frac{Z_{\text{L}} \cos \beta l + jZ_0 \sin \beta l}{Z_0 \cos \beta l + jZ_{\text{L}} \sin \beta l} \right) = Z_0 \left(\frac{Z_{\text{L}} + jZ_0 \tan \beta l}{Z_0 + jZ_{\text{L}} \tan \beta l} \right) \end{split}$$

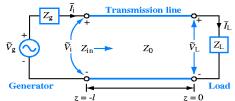


Input Impedance

From the standpoint of the generator circuit, the transmission line appears as an input impedance $Z_{\text{in}} = Z_{\text{in}}(-l)$, so that

$$\tilde{V_{\rm i}} = \tilde{I}_{\rm i} Z_{\rm in} = \frac{\tilde{V_{\rm g}} Z_{\rm in}}{Z_{\rm g} + Z_{\rm in}}$$

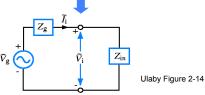
From the standpoint of the transmission line



 $\tilde{V}_{i} = \tilde{V}(-l) = V_{0}^{+}[\exp(j\beta l) + \Gamma \exp(-j\beta l)]$

Combining these relations...





4.19

Lossless Transmission Line

Input Impedance

...we conclude

$$V_0^+ = \left(\frac{\tilde{V}_{g} Z_{in}}{Z_{g} + Z_{in}}\right) \left(\frac{1}{\exp(j\beta l) + \Gamma \exp(-j\beta l)}\right)$$

Thus, we can now relate the wave parameters,

$$V_0^+, \quad V_0^- (= \Gamma V_0^+), \quad I_0^+ (= V_0^+ \, / \, Z_0), \quad I_0^- (= -V_0^- \, / \, Z_0 = -\Gamma V_0^+ \, / \, Z_0),$$

to the transmission line parameters

$$Z_{\rm g},\quad Z_{\rm L},\quad Z_0=\sqrt{L'/C'},\quad u_{\rm p}=1/\sqrt{L'C'},\quad l$$

and the input parameters



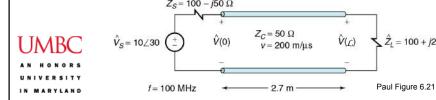
$$\tilde{V}_{\rm g}$$
, $f = \omega / 2\pi$

Paul Example 6.7 (extended)

Question: A line 2.7 m in length is excited by a 100 MHz source as shown in the figure. Determine the source and load voltages. Determine the voltages and the current everywhere in the transmission line

Answer: We will work in Ulaby et al.'s notation, and our first task is to translate Paul's problem specification into that notation:

$$\begin{split} \hat{V_S} &\to \tilde{V_g} = 10 \exp(j30^\circ) = 10 \exp(j\pi/6) = 10 \exp(j0.524) = (8.66 + 5.00 j) \text{ V} \\ \hat{V}(0) &\to \tilde{V_i} = \tilde{V}(-l); \quad \hat{V}(\mathcal{L}) \to \tilde{V_L} = \tilde{V}(0); \quad v \to u_p = 200 \text{ m} / \mu \text{ s} = 2 \times 10^8 \text{ m/s} \\ \hat{Z}_S &\to Z_g = 100 - j50 \,\Omega; \quad \hat{Z}_L \to Z_L = 100 + j200 \,\Omega; \quad Z_C \to Z_0 = 50 \,\Omega; \end{split}$$



Lossless Transmission Line

Paul Example 6.7

Answer (continued):*

(1) Find the reflection coefficient:

$$\Gamma = \frac{Z_{\rm L} - Z_0}{Z_{\rm L} + Z_0} = \frac{50 + j200}{150 + j200} = \frac{206 \exp(j1.326)}{250 \exp(j0.927)} = 0.825 \exp(j0.399)$$

(2) Find the propagation factor

$$\beta l = (2\pi f / u_p) l = (6.28 \times 10^8 / 2 \times 10^8) \times 2.7 = 8.48;$$

$$\exp(j\beta l) = \exp(j0.7\pi) = \exp(j2.20);$$

$$\exp(-j\beta l) = \exp(-j0.7\pi) = \exp(-j2.20);$$

$$\exp(-j2\beta l) = \exp(-j1.4\pi) = \exp(-j4.40) = \exp(j1.88)$$

UMBC

*NOTE: I am using MATLAB to calculate values. So, the calculations are good to 15 places, although I only report three.

4.22

4 21

Paul Example 6.7 Answer (continued):

(3) Find the input impedance:

$$\Gamma \exp(-j2\beta l) = 0.825 \exp[j(0.399 + 1.885)]$$

$$= 0.825 \exp(j2.283) = -0.539 + j0.624;$$

$$1 + \Gamma \exp(-j2\beta l) = 0.461 + j0.624;$$

$$1 - \Gamma \exp(-j2\beta l) = 1.539 - j0.624;$$

$$Z_{\text{in}} = Z_0 \frac{1 + \Gamma \exp(-j2\beta l)}{1 - \Gamma \exp(-j2\beta l)} = 50 \frac{0.461 + j0.624}{1.539 - j0.624} = 50 \frac{0.776 \exp(j0.935)}{1.661 \exp(-j0.385)}$$
$$= 23.35 \exp(j1.320) = 5.80 + j22.62$$



4.23

Lossless Transmission Line

Paul Example 6.7

Answer (continued):

(4) Find the input voltage:

$$Z_{g} + Z_{in} = (100 - j50) + (5.80 + j22.6)$$

$$= 106 - j27.4 = 109 \exp(-j0.253)$$

$$\tilde{V}_{i} = \tilde{V}_{g} \frac{Z_{in}}{Z_{g} + Z_{in}} = 10 \exp(j0.524) \frac{23.4 \exp(j1.320)}{109 \exp(-j0.253)}$$

$$= 2.14 \exp(j2.10) = 2.14 \angle 120^{\circ}$$



Paul Example 6.7 Answer (continued):

(5) Find the load voltage:

$$V_0^+ = \tilde{V}_i \frac{\exp(-j\beta l)}{1 + \Gamma \exp(-j2\beta l)} = 2.14 \exp(j2.10) \frac{\exp(-j2.20)}{0.776 \exp(j0.935)}$$
$$= 2.75 \exp(-j1.037)$$

$$1 + \Gamma = 1 + (0.760 + j0.320) = 1.79 \exp(j0.180)$$

$$\tilde{V}_L = V_0^+ (1 + \Gamma) = [2.75 \exp(-j1.037)][1.79 \exp(j0.180)]$$

= $4.93 \exp(-j0.857) = 4.93 \angle -49^\circ$



4.25

Lossless Transmission Line

Paul Example 6.7

Answer (continued):

(6) Write the time domain voltages:

$$v_{i}(t) = \text{Re}[\tilde{V}_{i} \exp(j\omega t)] = 2.14\cos(6.28 \times 10^{8} t + 120^{\circ})$$

 $v_{I}(t) = \text{Re}[\tilde{V}_{I} \exp(j\omega t)] = 4.93\cos(6.28 \times 10^{8} t - 49^{\circ})$

(7) Find the phasor domain voltage and current in the transmission line

$$\tilde{V}(z) = V_0^+ \exp(-j\beta z) + \Gamma V_0^+ \exp(j\beta z)$$

$$= 2.75 \exp(-j\beta z - j1.037) + 2.27 \exp(j\beta z - j0.638) \text{ (V)}$$

$$\tilde{I}(z) = \frac{V_0^+}{Z_0} \exp(-j\beta z) - \frac{\Gamma V_0^+}{Z_0} \exp(j\beta z)$$

$$= 55.1 \exp(-j\beta z - j1.037) + 45.4 \exp(j\beta z + j2.503) \text{ (mA)}$$



Paul Example 6.7 Answer (continued):

(8) Find the time domain voltage and current in the transmission line:

$$v(z,t) = \text{Re} \Big[2.75 \exp(j\omega t - j\beta z - j1.037) + 2.27 \exp(j\omega t + j\beta z - j0.638) \Big]$$

= 2.75 \cos(\omega t - \beta z - 1.037) + 2.27 \cos(\omega t + \beta z - 0.638) (V)

$$i(z,t) = \text{Re} \left[55.1 \exp(j\omega t - j\beta z - j1.037) + 45.4 \exp(j\omega t + j\beta z + j2.503) \right]$$

= 55.1 \cos(\omega t - \beta z - 1.037) + 45.4 \cos(\omega t + \beta z + 2.503) \text{ (mA)}

(9) Checks:

$$v(0,t) = 2.75\cos(\omega t - 1.037) + 2.27\cos(\omega t - 0.638)$$

$$= 4.93\cos(\omega t - 0.857) \text{ (V)}$$

$$v(-l,t) = 2.75\cos(\omega t + 8.48 - 1.037) + 2.27\cos(\omega t - 8.48 - 0.638)$$

$$= 2.14\cos(\omega t + 2.097) \text{ (V)}$$

