

## Second Midterm Examination Solutions

1. We have  $\Gamma = (200 - 100)/(200 + 100) = 1/3$ , and since  $\Gamma$  is purely real, we have

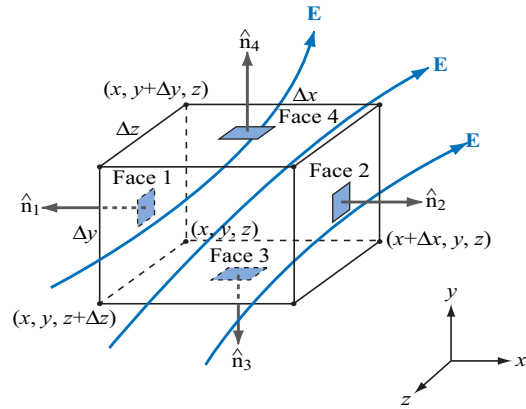
$$S = \frac{1 + |\Gamma|}{1 - |\Gamma|} = \frac{1 + \Gamma}{1 - \Gamma} = 2.$$

Because  $\Gamma$  is purely real and positive, the voltage has a maximum at the load and every distance  $-z = l = n\lambda/2$  away from the load, where  $n = 1, 2, 3, \dots$ . So, the first maximum away from the load is at  $l = \lambda/2$ . The minima occur in between the maxima when  $-z = l = (\lambda/4) + n\lambda/2$ , where  $n = 0, 1, 2, \dots$ . Hence, the first minimum occurs at  $l = \lambda/4$ .

2. Since the load impedance and the characteristic impedance are matched, there is only a forward-propagating wave, and the input impedance must therefore equal the load impedance, which is  $50 \Omega$ . Since the source impedance is also  $50 \Omega$ , half the power must be dissipated in the source impedance, and the other half must be dissipated in the load impedance. The power dissipated in the generator is  $P_g = (1/2) \text{Re}(VI^*)$ , and we have  $I^* = V^*/(Z_{\text{in}}^* + Z_{\text{source}}^*) = (100/100) \text{ A} = 1 \text{ A}$ . Hence, we have  $P_g = 25 \text{ W}$  and  $P_L = 25 \text{ W}$ .

3. a. The divergence is defined as the net flux of a vector field into a small volume in the limit as that volume tends to zero. The corresponding mathematical expression is

$$\nabla \cdot \mathbf{E} \equiv \lim_{\Delta v \rightarrow 0} \frac{\oint \mathbf{E} \cdot d\mathbf{S}}{\Delta v}.$$



- b. A picture of the flux fields passing through a small volume in Cartesian coordinates is given in the figure above, which is Ulaby et al.'s Fig. 3-21. If we consider the flux out of faces 1 and 2 in the figure to the right, we find that the flux out of face 1 is given by  $-E_x(x, y, z)\Delta y\Delta z$ , while the flux out of face 2 is given by  $E_x(x + \Delta x, y, z)\Delta y\Delta z \simeq E_x(x, y, z)\Delta y\Delta z + (\partial E_x/\partial x)\Delta x\Delta y\Delta z$ . Adding the contributions from faces 1 and 2, we obtain  $(\partial E_x/\partial x)\Delta x\Delta y\Delta z$ . Considering faces 3 and 4 yields analogously  $(\partial E_y/\partial y)\Delta x\Delta y\Delta z$ , and considering faces 5 and 6 yields analogously  $(\partial E_z/\partial z)\Delta x\Delta y\Delta z$ . Noting that  $\Delta v = \Delta x\Delta y\Delta z$  and applying the definition of the divergence, we conclude

$$\nabla \cdot \mathbf{E} = \frac{\partial E_x}{\partial x} + \frac{\partial E_y}{\partial y} + \frac{\partial E_z}{\partial z}.$$

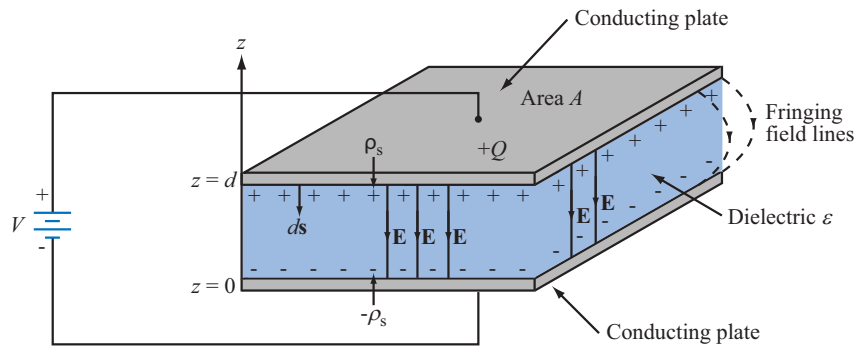
4. The geometry that we are considering is shown below, and we neglect fringing fields. We are going to suppose that we ramp up the voltage  $v$  from 0 to  $V$ . As that happens the charge  $q$  on the positive voltage end of the voltage source ramps up from 0 to  $Q$  coulombs. In order to move an increment of charge  $dq$  from the plate at the negative end of the voltage source to the positive end, we must do work, since the existing charges on the positive end will exert a force on the charge increment. This work is given by  $dW_e = v dq = (q/C)dq$ . Adding up all the work that is performed in charging the capacitor, we find

$$W_e = \int_0^Q \frac{q}{C} dq = \frac{1}{2} \frac{Q^2}{C}.$$

Using  $C = Q/V$ , we also have  $W_e = (1/2)CV^2$ . For the parallel plate capacitor, we have  $C = \epsilon A/d$ , where  $A$  is the area of the plate and  $d$  is the separation between the plates. We also have  $V = Ed$ . Substituting for  $C$  and  $V$  in the expression for  $W_e$ , we obtain

$$W_e = \frac{1}{2} \frac{\epsilon A}{d} (Ed)^2 = \frac{1}{2} \epsilon E^2 (Ad).$$

Noting that  $Ad$  equals the volume between the plates, we conclude that  $(1/2)\epsilon E^2$  equals the energy density.



5. a. The equations of motion are:  $d\mathbf{x}/dt = \mathbf{u}$  and  $d\mathbf{u}/dt = (q/m)(\mathbf{E} + \mathbf{u} \times \mathbf{B})$ , where  $x$  and  $u$  are the position and velocity of the charge.
- b. If  $B = 0$ , we may choose the direction in which the electric field points for convenience, and we will pick the  $z$ -direction. We may also pick the initial starting point of the particle to be  $(0, 0, 0)$ . The equations of motion become  $dx/dt = u_x$ ,  $dy/dt = u_y$ ,  $dz/dt = u_z$ ,  $du_x/dt = 0$ ,  $du_y/dt = 0$ , and  $du_z/dt = qE/m$ . The motion in the  $x$ ,  $y$ , and  $z$  directions is uncoupled. Since there is no change in the  $x$ - or  $y$ -velocities and we are starting from rest, we find  $u_x = 0$ ,  $u_y = 0$ ,  $x = 0$ , and  $y = 0$ . So, there is no motion at all in the  $x$ - and  $y$ -directions. Since  $E$  is constant, we find  $u_z = (qE/m)t$  and  $z = (qE/m)(t^2/2)$ , which is what we want to show. This result corresponds to a constant acceleration in the  $z$ -direction.