

I am absolutely starving.

The balance we're interested in is:

$$L_{fusion} \approx L_{rad/conv} \quad (1)$$

$$+ \text{HE} \rightarrow \text{main sequence} \quad (2)$$

For stars of  $M \leq M_{\odot}$ , supported by  $P_{gas}$ , pp chain fusion, and  $\kappa \sim \kappa_{ff}$ . And if  $\gamma$  carry energy out  $L_{rad} \propto \frac{M^{5.5}}{\sqrt{R}}$ . Estimating  $L_{fusion} = \epsilon_c M$ . In the case of pp fusion,  $\epsilon_{pp} \propto \rho T_c^{4/5}$  where  $kT_c \sim \frac{GM\mu m_p}{R}$ .

$$L_{fusion} \propto \frac{M^{6/5}}{R^{7.5}} \quad (3)$$

If you change the  $L$  of the star,  $T$  and  $\rho$  must change to accommodate. In steady state:

$$L_{fusion} \propto \frac{M^{6/5}}{R^{7.5}} \propto \frac{M^{5.5}}{\sqrt{R}} \quad (4)$$

$$R \propto M^{1/7} \quad (5)$$

$$T_c \propto M^{6/7} \quad (6)$$

$$L \propto M^{5.4} \quad (7)$$

$$L \propto T_{eff}^4 \propto 4\pi R^2 \sigma T^4, \text{ but } R \text{ dep is so weak it's estimated to be constant} \quad (8)$$

$$M \uparrow, T_c \uparrow \sim M^{6/7} \quad (9)$$

$$\epsilon_{pp} \propto T^{4.5}, \epsilon_{CNO} \propto T^{20} \quad (10)$$

Stars that are more massive than the sun are very  $T$  dependent.

For  $M \geq M_{\odot}$ :  $\kappa \sim \kappa_T \sim \text{constant}$ .  $\gamma$  still dominate  $E$  transport. CNO cycle is dominant mechanism and  $P_{gas}$  dominates.  $L_{rad} \propto M^3$ .

$$L_{fusion} = \int \epsilon_{CNO} dM_r \quad (11)$$

$$\sim \epsilon_{CNO} M, \epsilon_{CNO} \propto \rho T_c^{20} \quad (12)$$

$$\epsilon_{CNO} \propto \frac{M^{21}}{R^{23}} \quad (13)$$

$$L_{fusion} \sim L_{rad} \quad (14)$$

$$\frac{M^{22}}{R^{23}} \propto M^3 \quad (15)$$

$$R \propto M^{19/23} \propto M^{.8} \quad (16)$$

$$T_c \propto M^{.2} \quad (17)$$

$$L = 4\pi R^2 \sigma T_{eff}^4 \quad (18)$$

$$L \propto R^2 T_{eff}^4 \quad (19)$$

$$R^2 \propto M^{1.6} \propto L^{1/2} \quad (20)$$

$$L \propto M^3 \quad (21)$$

$$\propto L^{1/2} T_{eff}^4 \quad (22)$$

$$L^{1/2} \propto T_{eff}^4 \quad (23)$$

$$L \propto T_{eff}^8 \quad (24)$$

$$T_{eff} \propto M^{3/8} \quad (25)$$

A huge change in  $L$  corresponds to a small change in  $T$ . Only for stars with masses slightly more than the sun.

For  $M \geq 1.2M_\odot$ , they have convective cores and  $\gamma$  transport energy on outer part of star. Reverse of our sun. Convection sets in if  $\frac{ds}{dr} < 0$ .  $\frac{ds}{dr}$  is implied by  $\gamma$  transport of energy. You can then use Radiative diffusion equation to see if  $\frac{ds}{dr} < 0$ . i.e. Convection sets in if  $\frac{d \ln T}{d \ln P} > \frac{\gamma-1}{\gamma}$ .  $\gamma$  is the one for photons, not of the particles convecting.

$$\frac{d \ln T}{d \ln P} = \frac{1}{4} \frac{P}{P_{rad}} \frac{L}{L_{edd}} \frac{L_r/L}{M_r/M}, L_{edd} = \frac{4\pi G M_c}{\kappa} \quad (26)$$

$P$  is the total pressure. This tells us convection sets in if  $\frac{1}{4} \frac{P}{P_{rad}} \frac{L}{L_{edd}} \frac{L_r/L}{M_r/M} > \frac{2}{5}$ . Recall for CNO,  $\epsilon \propto \rho T^\beta$ . At almost all  $r$ ,  $L_r \approx L$ . For CNO-dominated stars, only 1% of star's mass fuses. TINY! It's this enormous flux that originates so close to the core that it drives convection.  $\frac{M_r}{M} < \frac{5}{8} \frac{P}{P_{rad}} \frac{L}{L_{edd}}$ , then convection sets in. We're interested where  $P \approx P_{rad}$ . We want to know  $\frac{P}{P_{rad}}$  and  $\frac{L}{L_{edd}}$ .  $L \propto M^3$  and  $L \propto M$  so  $\frac{L}{L_{edd}} = 4.5 \times 10^{-5} \left( \frac{M}{M_\odot} \right)^2$ .

In the sun,  $r \sim 0 \rightarrow .5R_\odot$ ,  $\frac{P_{gas}}{P_{rad}} \sim 3000$ .

$$P_{gas} \propto \rho T \propto \frac{M^2}{R^4} \quad (27)$$

$$P_{rad} = \frac{1}{3} a T^4 \propto \frac{M^4}{R^4} \quad (28)$$

$$\frac{P_{gas}}{P_{rad}} \propto M^{-2} \quad (29)$$

$$\frac{P_{gas}}{P_{rad}} \approx 3000 \left( \frac{M}{M_\odot} \right)^{-2} \quad (30)$$

Convection sets in if  $\frac{M_r}{M} \leq 0.1$ .

Lifetime of MS star  $\approx \frac{E_{nuc}}{L}$ .

$$L = L_\odot \left( \frac{M}{M_\odot} \right)^{3.5} \quad (31)$$

$$E_{nuc} = N_p E \quad (32)$$

$$\approx .1 \frac{M}{m_p} \cdot 7 \text{ MeV} \quad (33)$$

$$\frac{E_{nuc}}{L} \approx 10^{10} \left( \frac{M}{M_\odot} \right)^{-2.5} \quad (34)$$

ANY star with a mass less than  $.85M_{\odot}$  is still fusing after 13.7 billion years. For  $M \sim 30M_{\odot}$ ,  $t_{MS}$  is about  $10^6$  years. Massive stars live and die where they are born.