

Hyperbolic plane \mathbb{H}^2

This Jupyter notebook illustrates some differential geometry capabilities of SageMath on the example of the hyperbolic plane. The corresponding tools have been developed within the [SageManifolds \(https://sagemanifolds.obspm.fr\)](https://sagemanifolds.obspm.fr) project.

A version of SageMath at least equal to 7.5 is required to run this notebook:

```
In [1]: version()
```

```
Out[1]: 'SageMath version 9.2, Release Date: 2020-10-24'
```

First we set up the notebook to display mathematical objects using LaTeX formatting:

```
In [2]: %display latex
```

We also tell Maxima, which is used by SageMath for simplifications of symbolic expressions, that all computations involve real variables:

```
In [3]: maxima_calculus.eval("domain: real;")
```

```
Out[3]: real
```

We declare \mathbb{H}^2 as a 2-dimensional differentiable manifold:

```
In [4]: H2 = Manifold(2, 'H2', latex_name=r'\mathbb{H}^2', start_index=1)
print(H2)
H2
```

```
2-dimensional differentiable manifold H2
```

```
Out[4]:  $\mathbb{H}^2$ 
```

We shall introduce charts on \mathbb{H}^2 that are related to various models of the hyperbolic plane as submanifolds of \mathbb{R}^3 . Therefore, we start by declaring \mathbb{R}^3 as a 3-dimensional manifold equipped with a global chart: the chart of Cartesian coordinates (X, Y, Z) :

```
In [5]: R3 = Manifold(3, 'R3', latex_name=r'\mathbb{R}^3', start_index=1)
X3.<X,Y,Z> = R3.chart()
X3
```

```
Out[5]:  $(\mathbb{R}^3, (X, Y, Z))$ 
```

Hyperboloid model

The first chart we introduce is related to the **hyperboloid model of \mathbb{H}^2** , namely to the representation of \mathbb{H}^2 as the upper sheet ($Z > 0$) of the hyperboloid of two sheets defined in \mathbb{R}^3 by the equation $X^2 + Y^2 - Z^2 = -1$:

```
In [6]: X_hyp.<X,Y> = H2.chart()
X_hyp
```

```
Out[6]:  $(\mathbb{H}^2, (X, Y))$ 
```

The corresponding embedding of \mathbb{H}^2 in \mathbb{R}^3 is

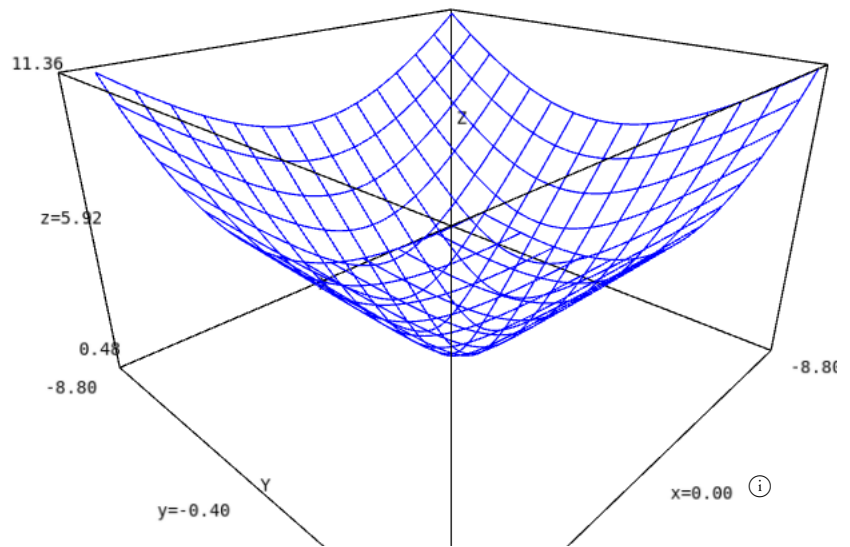
Hyperbolic plane (SageMath 9.2)

```
In [7]: Phi1 = H2.diff_map(R3, [X, Y, sqrt(1+X^2+Y^2)], name='Phi_1', latex_name=r'\Phi_1')
Phi1.display()
```

```
Out[7]:  $\Phi_1 : \mathbb{H}^2 \longrightarrow \mathbb{R}^3$ 
 $(X, Y) \longmapsto (X, Y, Z) = (X, Y, \sqrt{X^2 + Y^2 + 1})$ 
```

By plotting the chart $(\mathbb{H}^2, (X, Y))$ in terms of the Cartesian coordinates of \mathbb{R}^3 , we get a graphical view of $\Phi_1(\mathbb{H}^2)$:

```
In [8]: show(X_hyp.plot(X3, mapping=Phi1, number_values=15, color='blue'),
          aspect_ratio=1, figsize=7)
```



A second chart is obtained from the polar coordinates (r, φ) associated with (X, Y) . Contrary to (X, Y) , the polar chart is not defined on the whole \mathbb{H}^2 , but on the complement U of the segment $\{Y = 0, x \geq 0\}$:

```
In [9]: U = H2.open_subset('U', coord_def={X_hyp: (Y!=0, X<0)})
print(U)
```

Open subset U of the 2-dimensional differentiable manifold H2

Note that $(y!=0, x<0)$ stands for $y \neq 0$ OR $x < 0$; the condition $y \neq 0$ AND $x < 0$ would have been written $[y!=0, x<0]$ instead.

```
In [10]: X_pol.<r,ph> = U.chart(r'r:(0,+oo) ph:(0,2*pi):\varphi')
X_pol
```

```
Out[10]: (U, (r, \varphi))
```

```
In [11]: X_pol.coord_range()
```

```
Out[11]: r : (0, +\infty); \varphi : (0, 2\pi)
```

We specify the transition map between the charts $(U, (r, \varphi))$ and $(\mathbb{H}^2, (X, Y))$ as $X = r \cos \varphi, Y = r \sin \varphi$:

```
In [12]: pol_to_hyp = X_pol.transition_map(X_hyp, [r*cos(ph), r*sin(ph)])
pol_to_hyp
```

```
Out[12]: (U, (r, \varphi)) \rightarrow (U, (X, Y))
```

In [13]: `pol_to_hyp.display()`

Out[13]:
$$\begin{cases} X &= r \cos(\varphi) \\ Y &= r \sin(\varphi) \end{cases}$$

In [14]: `pol_to_hyp.set_inverse(sqrt(X^2+Y^2), atan2(Y, X))`

Check of the inverse coordinate transformation:

```
r == r *passed*
ph == arctan2(r*sin(ph), r*cos(ph)) **failed**
X == X *passed*
Y == Y *passed*
```

NB: a failed report can reflect a mere lack of simplification.

In [15]: `pol_to_hyp.inverse().display()`

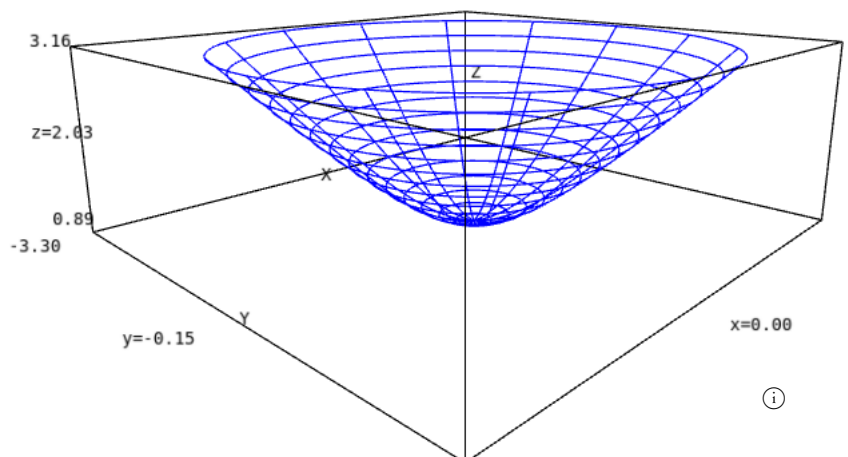
Out[15]:
$$\begin{cases} r &= \sqrt{X^2 + Y^2} \\ \varphi &= \arctan(Y, X) \end{cases}$$

The restriction of the embedding Φ_1 to U has then two coordinate expressions:

In [16]: `Phi1.restrict(U).display()`

Out[16]:
$$\begin{aligned} \Phi_1 : U &\longrightarrow \mathbb{R}^3 \\ (X, Y) &\longmapsto (X, Y, Z) = (X, Y, \sqrt{X^2 + Y^2 + 1}) \\ (r, \varphi) &\longmapsto (X, Y, Z) = (r \cos(\varphi), r \sin(\varphi), \sqrt{r^2 + 1}) \end{aligned}$$

In [17]: `graph_hyp = X_pol.plot(X3, mapping=Phi1.restrict(U), number_values=15, ranges={r: (0, 3)},
color='blue')
show(graph_hyp, aspect_ratio=1, figsize=7)`



In [18]: `Phi1._coord_expression`

Out[18]:
$$\left\{ \left((\mathbb{H}^2, (X, Y)), (\mathbb{R}^3, (X, Y, Z)) \right) : (X, Y, \sqrt{X^2 + Y^2 + 1}) \right\}$$

Metric and curvature

The metric on \mathbb{H}^2 is that induced by the Minkowsky metric on \mathbb{R}^3 :

$$\eta = dX \otimes dX + dY \otimes dY - dZ \otimes dZ$$

Hyperbolic plane (SageMath 9.2)

```
In [19]: eta = R3.lorentzian_metric('eta', latex_name=r'\eta')
eta[1,1] = 1 ; eta[2,2] = 1 ; eta[3,3] = -1
eta.display()
```

Out[19]: $\eta = dX \otimes dX + dY \otimes dY - dZ \otimes dZ$

```
In [20]: g = H2.metric('g')
g.set( Phil.pullback(eta) )
g.display()
```

Out[20]: $g = \left(\frac{Y^2 + 1}{X^2 + Y^2 + 1} \right) dX \otimes dX + \left(-\frac{XY}{X^2 + Y^2 + 1} \right) dX \otimes dY + \left(-\frac{XY}{X^2 + Y^2 + 1} \right) dY \otimes dX + \left(\frac{X^2}{X^2 + Y^2 + 1} \right) dY \otimes dY$

The expression of the metric tensor in terms of the polar coordinates is

```
In [21]: g.display(X_pol.frame(), X_pol)
```

Out[21]: $g = \left(\frac{1}{r^2 + 1} \right) dr \otimes dr + r^2 d\varphi \otimes d\varphi$

The Riemann curvature tensor associated with g is

```
In [22]: Riem = g.riemann()
print(Riem)
```

Tensor field Riem(g) of type (1,3) on the 2-dimensional differentiable manifold H2

```
In [23]: Riem.display(X_pol.frame(), X_pol)
```

Out[23]: $Riem(g) = -r^2 \frac{\partial}{\partial r} \otimes d\varphi \otimes dr \otimes d\varphi + r^2 \frac{\partial}{\partial r} \otimes d\varphi \otimes d\varphi \otimes dr + \left(\frac{1}{r^2 + 1} \right) \frac{\partial}{\partial \varphi} \otimes dr \otimes dr \otimes d\varphi + \left(-\frac{2r}{r^2 + 1} \right) \frac{\partial}{\partial \varphi} \otimes dr \otimes d\varphi \otimes dr$

The Ricci tensor and the Ricci scalar:

```
In [24]: Ric = g.ricci()
print(Ric)
```

Field of symmetric bilinear forms Ric(g) on the 2-dimensional differentiable manifold H2

```
In [25]: Ric.display(X_pol.frame(), X_pol)
```

Out[25]: $Ric(g) = \left(-\frac{1}{r^2 + 1} \right) dr \otimes dr - r^2 d\varphi \otimes d\varphi$

```
In [26]: Rscal = g.ricci_scalar()
print(Rscal)
```

Scalar field r(g) on the 2-dimensional differentiable manifold H2

```
In [27]: Rscal.display()
```

Out[27]: $r(g): \mathbb{H}^2 \rightarrow \mathbb{R}$
 $(X, Y) \mapsto -2$
on $U: (r, \varphi) \mapsto -2$

Hence we recover the fact that (\mathbb{H}^2, g) is a space of **constant negative curvature**.

Hyperbolic plane (SageMath 9.2)

In dimension 2, the Riemann curvature tensor is entirely determined by the Ricci scalar R according to

$$R^i_{jlk} = \frac{R}{2} (\delta^i_k g_{jl} - \delta^i_l g_{jk})$$

Let us check this formula here, under the form $R^i_{jlk} = -R g_{j[k} \delta^i_{l]}$:

```
In [28]: delta = H2.tangent_identity_field()
Riem == - Rscal*(g*delta).antisymmetrize(2,3) # 2,3 = last positions of the type-(1,
3) tensor g*delta
```

Out[28]: True

Similarly the relation $\text{Ric} = (R/2) g$ must hold:

```
In [29]: Ric == (Rscal/2)*g
```

Out[29]: True

Poincaré disk model

The Poincaré disk model of \mathbb{H}^2 is obtained by stereographic projection from the point $S = (0, 0, -1)$ of the hyperboloid model to the plane $Z = 0$. The radial coordinate R of the image of a point of polar coordinate (r, φ) is

$$R = \frac{r}{1 + \sqrt{1 + r^2}}.$$

Hence we define the Poincaré disk chart on \mathbb{H}^2 by

```
In [30]: X_Pdisk.<R,ph> = U.chart(r'R:(0,1) ph:(0,2*pi):\varphi')
X_Pdisk
```

Out[30]: $(U, (R, \varphi))$

```
In [31]: X_Pdisk.coord_range()
```

Out[31]: $R : (0, 1); \quad \varphi : (0, 2\pi)$

and relate it to the hyperboloid polar chart by

```
In [32]: pol_to_Pdisk = X_pol.transition_map(X_Pdisk, [r/(1+sqrt(1+r^2)), ph])
pol_to_Pdisk
```

Out[32]: $(U, (r, \varphi)) \rightarrow (U, (R, \varphi))$

```
In [33]: pol_to_Pdisk.display()
```

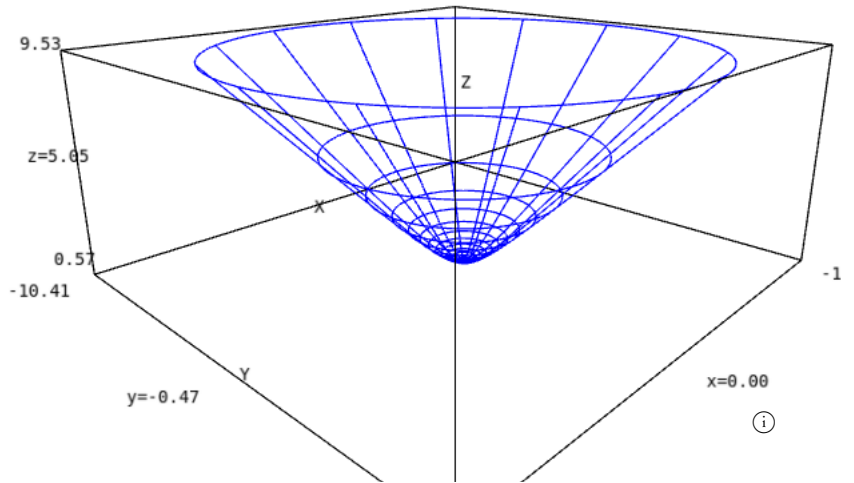
Out[33]:
$$\begin{cases} R &= \frac{r}{\sqrt{r^2+1}+1} \\ \varphi &= \varphi \end{cases}$$

```
In [34]: pol_to_Pdisk.set_inverse(2*R/(1-R^2), ph)
pol_to_Pdisk.inverse().display()
```

Out[34]:
$$\begin{cases} r &= -\frac{2R}{R^2-1} \\ \varphi &= \varphi \end{cases}$$

A view of the Poincaré disk chart via the embedding Φ_1 :

```
In [35]: show(X_Pdisk.plot(X3, mapping=Phil.restrict(U), ranges={R: (0,0.9)}, color='blue',
                        number_values=15),
          aspect_ratio=1, figsize=7)
```



The expression of the metric tensor in terms of coordinates (R, φ) :

```
In [36]: g.display(X_Pdisk.frame(), X_Pdisk)
```

```
Out[36]: 
$$g = \left( \frac{4}{R^4 - 2R^2 + 1} \right) dR \otimes dR + \left( \frac{4R^2}{R^4 - 2R^2 + 1} \right) d\varphi \otimes d\varphi$$

```

We may factorize each metric component:

```
In [37]: for i in [1,2]:
          g[X_Pdisk.frame(), i, i, X_Pdisk].factor()
          g.display(X_Pdisk.frame(), X_Pdisk)
```

```
Out[37]: 
$$g = \frac{4}{(R+1)^2(R-1)^2} dR \otimes dR + \frac{4R^2}{(R+1)^2(R-1)^2} d\varphi \otimes d\varphi$$

```

Cartesian coordinates on the Poincaré disk

Let us introduce Cartesian coordinates (u, v) on the Poincaré disk; since the latter has a unit radius, this amounts to define the following chart on \mathbb{H}^2 :

```
In [38]: X_Pdisk_cart.<u,v> = H2.chart('u: (-1,1) v: (-1,1)')
          X_Pdisk_cart.add_restrictions(u^2+v^2 < 1)
          X_Pdisk_cart
```

```
Out[38]:  $(\mathbb{H}^2, (u, v))$ 
```

On U , the Cartesian coordinates (u, v) are related to the polar coordinates (R, φ) by the standard formulas:

```
In [39]: Pdisk_to_Pdisk_cart = X_Pdisk.transition_map(X_Pdisk_cart, [R*cos(ph), R*sin(ph)])
          Pdisk_to_Pdisk_cart
```

```
Out[39]:  $(U, (R, \varphi)) \rightarrow (U, (u, v))$ 
```

```
In [40]: Pdisk_to_Pdisk_cart.display()
```

```
Out[40]: { u  =  R cos(φ)
          { v  =  R sin(φ)
```

```
In [41]: Pdisk_to_Pdisk_cart.set_inverse(sqrt(u^2+v^2), atan2(v, u))
Pdisk_to_Pdisk_cart.inverse().display()
```

Check of the inverse coordinate transformation:

```
R == R *passed*
ph == arctan2(R*sin(ph), R*cos(ph)) **failed**
u == u *passed*
v == v *passed*
```

NB: a failed report can reflect a mere lack of simplification.

```
Out[41]: { R  =  sqrt(u^2 + v^2)
          { φ  =  arctan(v, u)
```

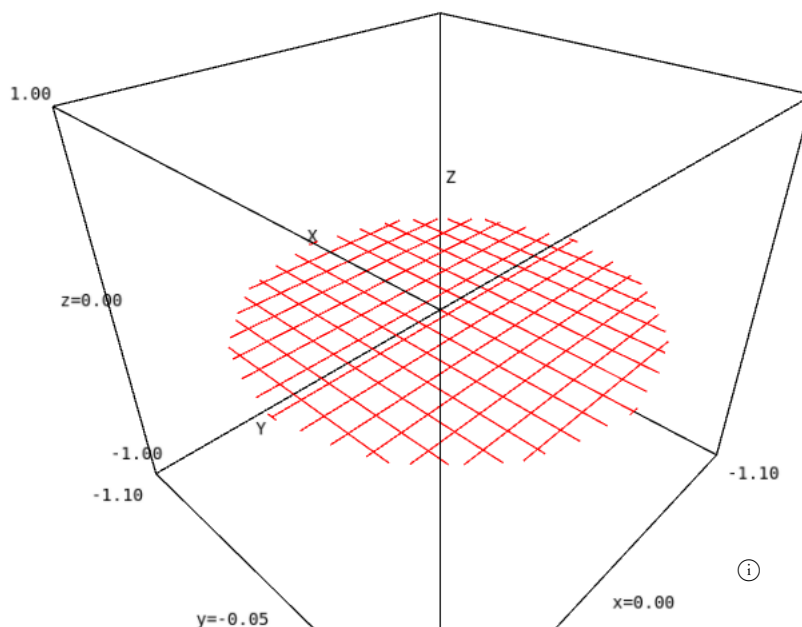
The embedding of \mathbb{H}^2 in \mathbb{R}^3 associated with the Poincaré disk model is naturally defined as

```
In [42]: Phi2 = H2.diff_map(R3, {(X_Pdisk_cart, X3): [u, v, 0]},
                             name='Phi_2', latex_name=r'\Phi_2')
Phi2.display()
```

```
Out[42]: Φ2 : ℍ2    → ℝ3
          (u, v) ↦ (X, Y, Z) = (u, v, 0)
```

Let us use it to draw the Poincaré disk in \mathbb{R}^3 :

```
In [43]: graph_disk_uv = X_Pdisk_cart.plot(X3, mapping=Phi2, number_values=15)
show(graph_disk_uv, figsize=7)
```



On U , the change of coordinates $(r, \varphi) \rightarrow (u, v)$ is obtained by combining the changes $(r, \varphi) \rightarrow (R, \varphi)$ and $(R, \varphi) \rightarrow (u, v)$:

```
In [44]: pol_to_Pdisk_cart = Pdisk_to_Pdisk_cart * pol_to_Pdisk
pol_to_Pdisk_cart
```

```
Out[44]: (U, (r, φ)) → (U, (u, v))
```

```
In [45]: pol_to_Pdisk_cart.display()
```

$$\text{Out[45]: } \begin{cases} u &= \frac{r \cos(\varphi)}{\sqrt{r^2+1}+1} \\ v &= \frac{r \sin(\varphi)}{\sqrt{r^2+1}+1} \end{cases}$$

Still on U , the change of coordinates $(X, Y) \rightarrow (u, v)$ is obtained by combining the changes $(X, Y) \rightarrow (r, \varphi)$ with $(r, \varphi) \rightarrow (u, v)$:

```
In [46]: hyp_to_Pdisk_cart_U = pol_to_Pdisk_cart * pol_to_hyp.inverse()
hyp_to_Pdisk_cart_U
```

```
Out[46]: (U, (X, Y)) -> (U, (u, v))
```

```
In [47]: hyp_to_Pdisk_cart_U.display()
```

$$\text{Out[47]: } \begin{cases} u &= \frac{X}{\sqrt{X^2+Y^2+1}+1} \\ v &= \frac{Y}{\sqrt{X^2+Y^2+1}+1} \end{cases}$$

We use the above expression to extend the change of coordinates $(X, Y) \rightarrow (u, v)$ from U to the whole manifold \mathbb{H}^2 :

```
In [48]: hyp_to_Pdisk_cart = X_hyp.transition_map(X_Pdisk_cart, hyp_to_Pdisk_cart_U(X,Y))
hyp_to_Pdisk_cart
```

```
Out[48]: (\mathbb{H}^2, (X, Y)) -> (\mathbb{H}^2, (u, v))
```

```
In [49]: hyp_to_Pdisk_cart.display()
```

$$\text{Out[49]: } \begin{cases} u &= \frac{X}{\sqrt{X^2+Y^2+1}+1} \\ v &= \frac{Y}{\sqrt{X^2+Y^2+1}+1} \end{cases}$$

```
In [50]: hyp_to_Pdisk_cart.set_inverse(2*u/(1-u^2-v^2), 2*v/(1-u^2-v^2))
hyp_to_Pdisk_cart.inverse().display()
```

Check of the inverse coordinate transformation:

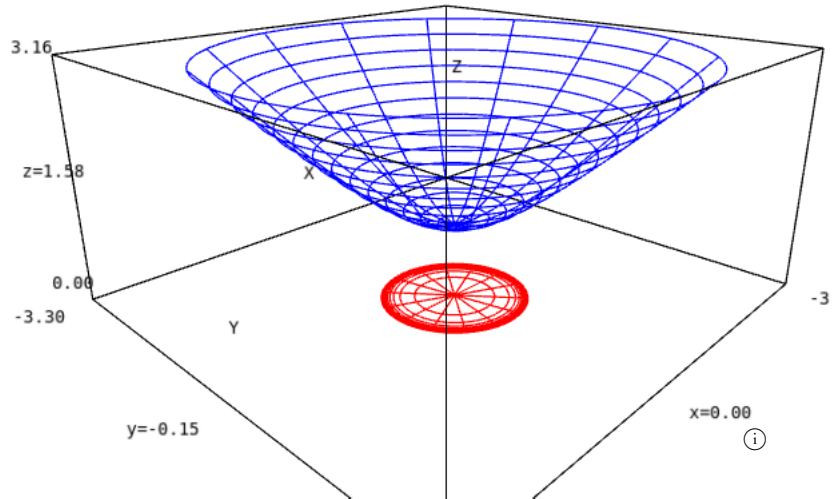
```
X == X *passed*
Y == Y *passed*
u == -2*u*abs(u^2 + v^2 - 1)/(u^4 + 2*u^2*v^2 + v^4 + (u^2 + v^2 - 1)*abs(u^2 + v^2 - 1) - 1) **failed**
v == -2*v*abs(u^2 + v^2 - 1)/(u^4 + 2*u^2*v^2 + v^4 + (u^2 + v^2 - 1)*abs(u^2 + v^2 - 1) - 1) **failed**
```

NB: a failed report can reflect a mere lack of simplification.

$$\text{Out[50]: } \begin{cases} X &= -\frac{2u}{u^2+v^2-1} \\ Y &= -\frac{2v}{u^2+v^2-1} \end{cases}$$

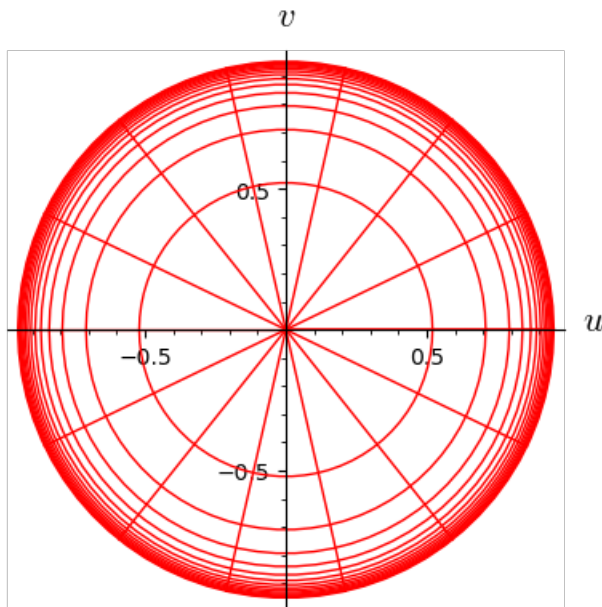
Hyperbolic plane (SageMath 9.2)

```
In [51]: graph_Pdisk = X_pol.plot(X3, mapping=Phi2.restrict(U), ranges={r: (0, 20)}, number_val
        lues=15,
        label_axes=False)
        show(graph_hyp + graph_Pdisk, aspect_ratio=1, figsize=7)
```



```
In [52]: X_pol.plot(X_Pdisk_cart, ranges={r: (0, 20)}, number_values=15)
```

Out[52]:



Metric tensor in Poincaré disk coordinates (u, v)

From now on, we are using the Poincaré disk chart $(\mathbb{H}^2, (u, v))$ as the default one on \mathbb{H}^2 :

```
In [53]: H2.set_default_chart(X_Pdisk_cart)
        H2.set_default_frame(X_Pdisk_cart.frame())
```

In [54]: `g.display(X_hyp.frame())`

Out[54]:
$$g = \left(\frac{u^4 + v^4 + 2(u^2 + 1)v^2 - 2u^2 + 1}{u^4 + v^4 + 2(u^2 + 1)v^2 + 2u^2 + 1} \right) dX \otimes dX + \left(-\frac{4uv}{u^4 + v^4 + 2(u^2 + 1)v^2 + 2u^2 + 1} \right) dX \otimes dY + \left(-\frac{4uv}{u^4 + v^4 + 2(u^2 + 1)v^2 + 2u^2 + 1} \right) dY \otimes dX + \left(\frac{u^4 + v^4 + 2(u^2 - 1)v^2 + 2u^2 + 1}{u^4 + v^4 + 2(u^2 + 1)v^2 + 2u^2 + 1} \right) dY \otimes dY$$

In [55]: `g.display()`

Out[55]:
$$g = \left(\frac{4}{u^4 + v^4 + 2(u^2 - 1)v^2 - 2u^2 + 1} \right) du \otimes du + \left(\frac{4}{u^4 + v^4 + 2(u^2 - 1)v^2 - 2u^2 + 1} \right) dv \otimes dv$$

In [56]: `g[1,1].factor() ; g[2,2].factor()
g.display()`

Out[56]:
$$g = \frac{4}{(u^2 + v^2 - 1)^2} du \otimes du + \frac{4}{(u^2 + v^2 - 1)^2} dv \otimes dv$$

Hemispherical model

The **hemispherical model** of \mathbb{H}^2 is obtained by the inverse stereographic projection from the point $S = (0, 0, -1)$ of the Poincaré disk to the unit sphere $X^2 + Y^2 + Z^2 = 1$. This induces a spherical coordinate chart on U :

In [57]: `X_spher.<th,ph> = U.chart(r'th:(0,pi/2):\theta ph:(0,2*pi):\varphi')
X_spher`

Out[57]: $(U, (\theta, \varphi))$

From the stereographic projection from S , we obtain that

$$\sin \theta = \frac{2R}{1 + R^2}$$

Hence the transition map:

In [58]: `Pdisk_to_spher = X_Pdisk.transition_map(X_spher, [arcsin(2*R/(1+R^2)), ph])
Pdisk_to_spher`

Out[58]: $(U, (R, \varphi)) \rightarrow (U, (\theta, \varphi))$

In [59]: `Pdisk_to_spher.display()`

Out[59]:
$$\begin{cases} \theta &= \arcsin\left(\frac{2R}{R^2+1}\right) \\ \varphi &= \varphi \end{cases}$$

In [60]: `Pdisk_to_spher.set_inverse(sin(th)/(1+cos(th)), ph)
Pdisk_to_spher.inverse().display()`

Out[60]:
$$\begin{cases} R &= \frac{\sin(\theta)}{\cos(\theta)+1} \\ \varphi &= \varphi \end{cases}$$

In the spherical coordinates (θ, φ) , the metric takes the following form:

In [61]: `g.display(X_spher.frame(), X_spher)`

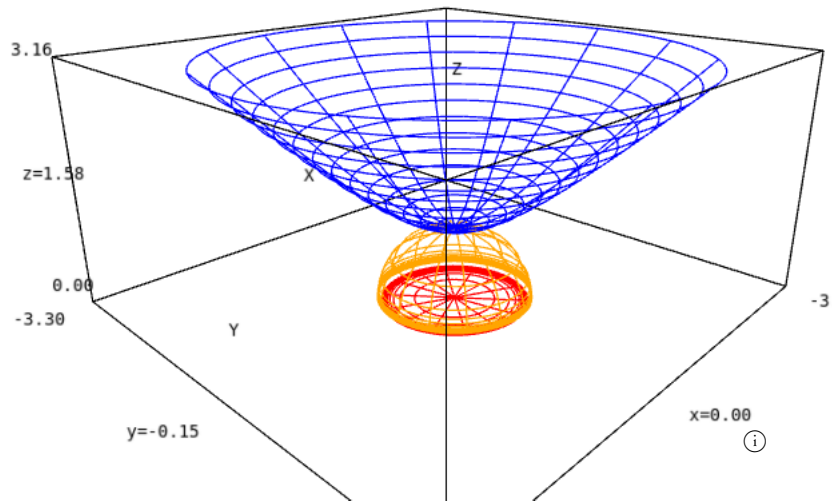
Out[61]:
$$g = \frac{1}{\cos^2(\theta)} d\theta \otimes d\theta + \frac{\sin^2(\theta)}{\cos^2(\theta)} d\varphi \otimes d\varphi$$

The embedding of \mathbb{H}^2 in \mathbb{R}^3 associated with the hemispherical model is naturally:

```
In [62]: Phi3 = H2.diff_map(R3, {(X_spher, X3): [sin(th)*cos(ph), sin(th)*sin(ph), cos(th)]},
name='Phi_3', latex_name=r'\Phi_3')
Phi3.display()
```

```
Out[62]:  $\Phi_3 : \mathbb{H}^2 \longrightarrow \mathbb{R}^3$ 
on  $U : (R, \varphi) \longmapsto (X, Y, Z) = \left( \frac{2R \cos(\varphi)}{R^2+1}, \frac{2R \sin(\varphi)}{R^2+1}, -\frac{R^2-1}{R^2+1} \right)$ 
on  $U : (\theta, \varphi) \longmapsto (X, Y, Z) = (\cos(\varphi) \sin(\theta), \sin(\varphi) \sin(\theta), \cos(\theta))$ 
```

```
In [63]: graph_spher = X_pol.plot(X3, mapping=Phi3, ranges={r: (0, 20)}, number_values=15,
color='orange', label_axes=False)
show(graph_hyp + graph_Pdisk + graph_spher, aspect_ratio=1,
figsize=7)
```



Poincaré half-plane model

The **Poincaré half-plane model** of \mathbb{H}^2 is obtained by stereographic projection from the point $W = (-1, 0, 0)$ of the hemispherical model to the plane $X = 1$. This induces a new coordinate chart on \mathbb{H}^2 by setting $(x, y) = (Y, Z)$ in the plane $X = 1$:

```
In [64]: X_hplane.<x,y> = H2.chart('x y:(0,+oo)')
X_hplane
```

```
Out[64]:  $(\mathbb{H}^2, (x, y))$ 
```

The coordinate transformation $(\theta, \varphi) \rightarrow (x, y)$ is easily deduced from the stereographic projection from the point W :

```
In [65]: spher_to_hplane = X_spher.transition_map(X_hplane, [2*sin(th)*sin(ph)/(1+sin(th)*cos
(ph)),
2*cos(th)/(1+sin(th)*cos(ph))])
spher_to_hplane
```

```
Out[65]:  $(U, (\theta, \varphi)) \rightarrow (U, (x, y))$ 
```

```
In [66]: spher_to_hplane.display()
```

```
Out[66]: 
$$\begin{cases} x &= \frac{2 \sin(\varphi) \sin(\theta)}{\cos(\varphi) \sin(\theta)+1} \\ y &= \frac{2 \cos(\theta)}{\cos(\varphi) \sin(\theta)+1} \end{cases}$$

```

```
In [67]: Pdisk_to_hplane = spher_to_hplane * Pdisk_to_spher
Pdisk_to_hplane
```

```
Out[67]: (U, (R, ϕ)) → (U, (x, y))
```

```
In [68]: Pdisk_to_hplane.display()
```

```
Out[68]: { x = 4 R sin(ϕ) / (R^2 + 2 R cos(ϕ) + 1)
          y = -2 (R^2 - 1) / (R^2 + 2 R cos(ϕ) + 1) }
```

```
In [69]: Pdisk_cart_to_hplane_U = Pdisk_to_hplane * Pdisk_to_Pdisk_cart.inverse()
Pdisk_cart_to_hplane_U
```

```
Out[69]: (U, (u, v)) → (U, (x, y))
```

```
In [70]: Pdisk_cart_to_hplane_U.display()
```

```
Out[70]: { x = 4 v / (u^2 + v^2 + 2 u + 1)
          y = -2 (u^2 + v^2 - 1) / (u^2 + v^2 + 2 u + 1) }
```

Let us use the above formula to define the transition map $(u, v) \rightarrow (x, y)$ on the whole manifold \mathbb{H}^2 (and not only on U):

```
In [71]: Pdisk_cart_to_hplane = X_Pdisk_cart.transition_map(X_hplane, Pdisk_cart_to_hplane_U
(u, v))
Pdisk_cart_to_hplane
```

```
Out[71]: (ℍ^2, (u, v)) → (ℍ^2, (x, y))
```

```
In [72]: Pdisk_cart_to_hplane.display()
```

```
Out[72]: { x = 4 v / (u^2 + v^2 + 2 u + 1)
          y = -2 (u^2 + v^2 - 1) / (u^2 + v^2 + 2 u + 1) }
```

```
In [73]: Pdisk_cart_to_hplane.set_inverse((4 - x^2 - y^2) / (x^2 + (2 + y)^2), 4 * x / (x^2 + (2 + y)^2))
Pdisk_cart_to_hplane.inverse().display()
```

```
Out[73]: { u = -(x^2 + y^2 - 4) / (x^2 + (y + 2)^2)
          v = 4 x / (x^2 + (y + 2)^2) }
```

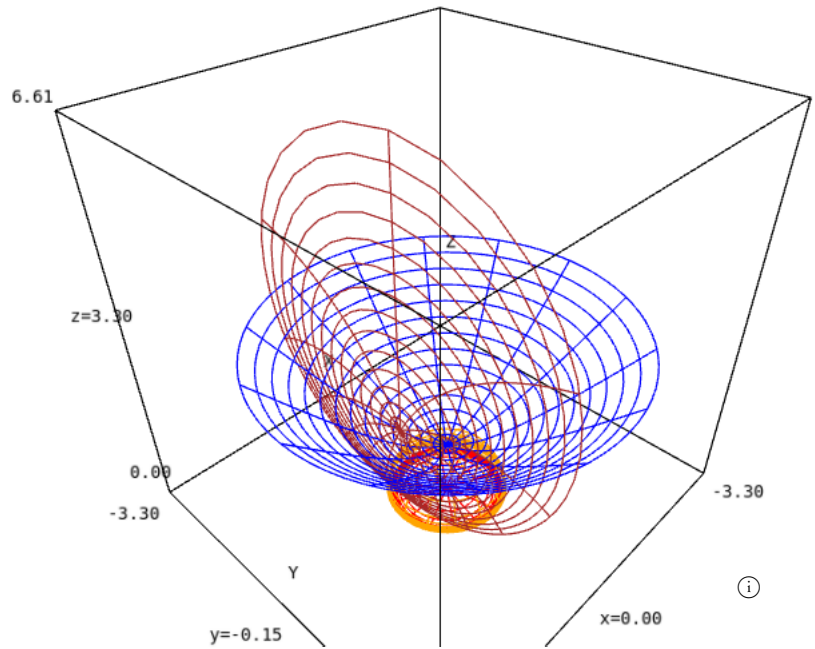
Since the coordinates (x, y) correspond to (Y, Z) in the plane $X = 1$, the embedding of \mathbb{H}^2 in \mathbb{R}^3 naturally associated with the Poincaré half-plane model is

```
In [74]: Phi4 = H2.diff_map(R3, {(X_hplane, X3): [1, x, y]}, name='Phi_4', latex_name=r'\Phi_4')
Phi4.display()
```

```
Out[74]: Φ4 : ℍ2 → ℝ3
          (u, v) ↦ (X, Y, Z) = (1, 4 v / (u^2 + v^2 + 2 u + 1), -2 (u^2 + v^2 - 1) / (u^2 + v^2 + 2 u + 1))
          (x, y) ↦ (X, Y, Z) = (1, x, y)
```

Hyperbolic plane (SageMath 9.2)

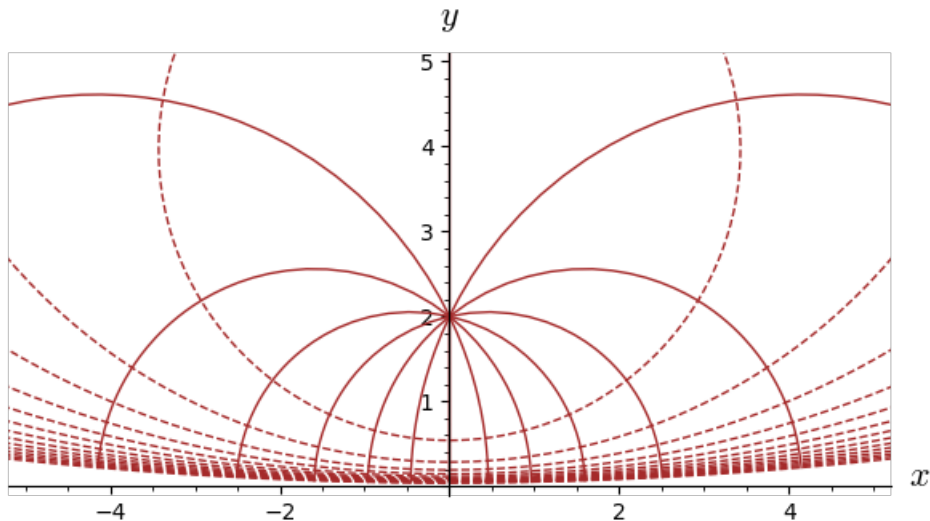
```
In [75]: graph_hplane = X_pol.plot(X3, mapping=Phi4.restrict(U), ranges={r: (0, 1.5)},
                                     number_values=15, color='brown', label_axes=False)
show(graph_hyp + graph_Pdisk + graph_spher + graph_hplane,
     aspect_ratio=1, figsize=8)
```



Let us draw the grid of the hyperboloidal coordinates (r, φ) in terms of the half-plane coordinates (x, y) :

```
In [76]: pol_to_hplane = Pdisk_to_hplane * pol_to_Pdisk
```

```
In [77]: show(X_pol.plot(X_hplane, ranges={r: (0,24)}, style={r: '-', ph: '--'}, number_values
=15,
           plot_points=200, color='brown'), xmin=-5, xmax=5, ymin=0, ymax=5,
           aspect_ratio=1)
```



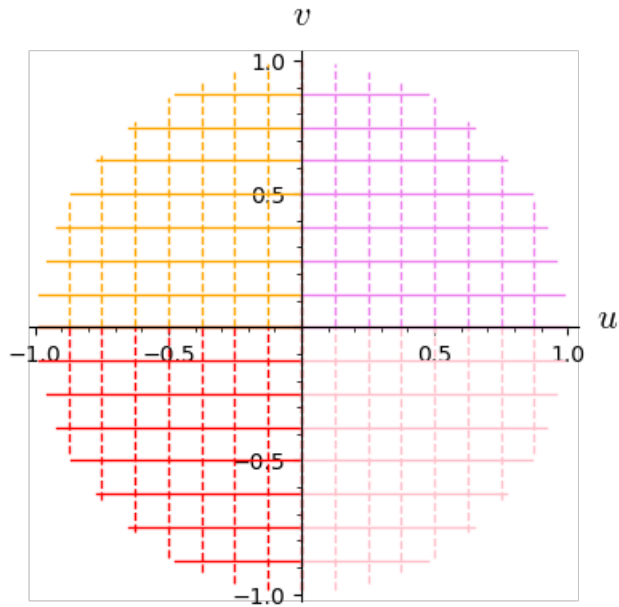
The solid curves are those along which r varies while φ is kept constant. Conversely, the dashed curves are those along which φ varies, while r is kept constant. We notice that the former curves are arcs of circles orthogonal to the half-plane boundary $y = 0$, hence they are geodesics of (\mathbb{H}^2, g) . This is not surprising since they correspond to the intersections of the hyperboloid with planes through the origin (namely the plane $\varphi = \text{const}$). The point $(x, y) = (0, 2)$ corresponds to $r = 0$.

We may also depict the Poincaré disk coordinates (u, v) in terms of the half-plane coordinates (x, y) :

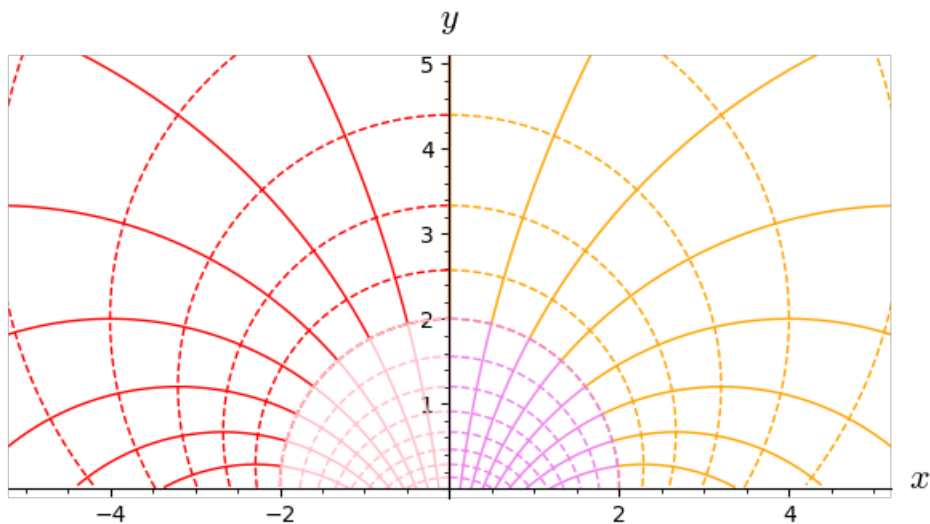
Hyperbolic plane (SageMath 9.2)

```
In [78]: X_Pdisk_cart.plot(ranges={u: (-1, 0), v: (-1, 0)},
                        style={u: '-', v: '--'}) + \
X_Pdisk_cart.plot(ranges={u: (-1, 0), v: (0., 1)},
                        style={u: '-', v: '--'}, color='orange') + \
X_Pdisk_cart.plot(ranges={u: (0, 1), v: (-1, 0)},
                        style={u: '-', v: '--'}, color='pink') + \
X_Pdisk_cart.plot(ranges={u: (0, 1), v: (0, 1)},
                        style={u: '-', v: '--'}, color='violet')
```

Out[78]:



```
In [79]: show(X_Pdisk_cart.plot(X_hplane, ranges={u: (-1, 0), v: (-1, 0)},
                        style={u: '-', v: '--'}) + \
X_Pdisk_cart.plot(X_hplane, ranges={u: (-1, 0), v: (0, 1)},
                        style={u: '-', v: '--'}, color='orange') + \
X_Pdisk_cart.plot(X_hplane, ranges={u: (0, 1), v: (-1, 0)},
                        style={u: '-', v: '--'}, color='pink') + \
X_Pdisk_cart.plot(X_hplane, ranges={u: (0, 1), v: (0, 1)},
                        style={u: '-', v: '--'}, color='violet'),
xmin=-5, xmax=5, ymin=0, ymax=5, aspect_ratio=1)
```



The expression of the metric tensor in the half-plane coordinates (x, y) is

```
In [80]: g.display(X_hplane.frame(), X_hplane)
```

Out[80]: $g = \frac{1}{y^2}dx \otimes dx + \frac{1}{y^2}dy \otimes dy$

Summary

9 charts have been defined on \mathbb{H}^2 :

```
In [81]: H2.atlas()
```

```
Out[81]: [(H^2, (X, Y)), (U, (X, Y)), (U, (r, phi)), (U, (R, phi)), (H^2, (u, v)), (U, (u, v)), (U, (theta, phi)), (H^2, (x, y)), (U,
```

There are actually 6 main charts, the other ones being subcharts:

```
In [82]: H2.top_charts()
```

```
Out[82]: [(H^2, (X, Y)), (U, (r, phi)), (U, (R, phi)), (H^2, (u, v)), (U, (theta, phi)), (H^2, (x, y))]
```

The expression of the metric tensor in each of these charts is

```
In [83]: for chart in H2.top_charts():
          show(g.display(chart.frame(), chart))
```

$$g = \left(\frac{Y^2 + 1}{X^2 + Y^2 + 1} \right) dX \otimes dX + \left(-\frac{XY}{X^2 + Y^2 + 1} \right) dX \otimes dY + \left(-\frac{XY}{X^2 + Y^2 + 1} \right) dY \otimes dX + \left(\frac{1}{X^2 + Y^2 + 1} \right) dY \otimes dY$$

$$g = \left(\frac{1}{r^2 + 1} \right) dr \otimes dr + r^2 d\varphi \otimes d\varphi$$

$$g = \frac{4}{(R+1)^2(R-1)^2} dR \otimes dR + \frac{4R^2}{(R+1)^2(R-1)^2} d\varphi \otimes d\varphi$$

$$g = \frac{4}{(u^2 + v^2 - 1)^2} du \otimes du + \frac{4}{(u^2 + v^2 - 1)^2} dv \otimes dv$$

$$g = \frac{1}{\cos(\theta)^2} d\theta \otimes d\theta + \frac{\sin(\theta)^2}{\cos(\theta)^2} d\varphi \otimes d\varphi$$

$$g = \frac{1}{y^2} dx \otimes dx + \frac{1}{y^2} dy \otimes dy$$

For each of these charts, the non-vanishing (and non-redundant w.r.t. the symmetry on the last 2 indices) **Christoffel symbols of g** are

```
In [84]: for chart in H2.top_charts():
          show(chart)
          show(g.christoffel_symbols_display(chart=chart))
```

$(\mathbb{H}^2, (X, Y))$

$$\Gamma^X_{XX} = -\frac{XY^2+X}{X^2+Y^2+1}$$

$$\Gamma^X_{XY} = \frac{X^2Y}{X^2+Y^2+1}$$

$$\Gamma^X_{YY} = -\frac{X^3+X}{X^2+Y^2+1}$$

$$\Gamma^Y_{XX} = -\frac{Y^3+Y}{X^2+Y^2+1}$$

$$\Gamma^Y_{XY} = \frac{XY^2}{X^2+Y^2+1}$$

$$\Gamma^Y_{YY} = -\frac{(X^2+1)Y}{X^2+Y^2+1}$$

$(U, (r, \varphi))$

$$\Gamma^r_{rr} = -\frac{r}{r^2+1}$$

$$\Gamma^r_{\varphi\varphi} = -r^3 - r$$

$$\Gamma^\varphi_{r\varphi} = \frac{1}{r}$$

$(U, (R, \varphi))$

$$\Gamma^R_{RR} = -\frac{2R}{R^2-1}$$

$$\Gamma^R_{\varphi\varphi} = \frac{R^3+R}{R^2-1}$$

$$\Gamma^\varphi_{R\varphi} = -\frac{R^2+1}{R^3-R}$$

$(\mathbb{H}^2, (u, v))$

$$\Gamma^u_{uu} = -\frac{2u}{u^2+v^2-1}$$

$$\Gamma^u_{uv} = -\frac{2v}{u^2+v^2-1}$$

$$\Gamma^u_{vv} = \frac{2u}{u^2+v^2-1}$$

$$\Gamma^v_{uu} = \frac{2v}{u^2+v^2-1}$$

$$\Gamma^v_{uv} = -\frac{2u}{u^2+v^2-1}$$

$$\Gamma^v_{vv} = -\frac{2v}{u^2+v^2-1}$$

$(U, (\theta, \varphi))$

$$\Gamma^\theta_{\theta\theta} = \frac{\sin(\theta)}{\cos(\theta)}$$

$$\Gamma^\theta_{\varphi\varphi} = -\frac{\sin(\theta)}{\cos(\theta)}$$

$$\Gamma^\varphi_{\theta\varphi} = \frac{1}{\cos(\theta) \sin(\theta)}$$

$(\mathbb{H}^2, (x, y))$

$$\Gamma^x_{xy} = -\frac{1}{y}$$

$$\Gamma^y_{xx} = \frac{1}{y}$$

$$\Gamma^y_{yy} = -\frac{1}{y}$$