

PH2202

Thermal physics

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Thermal physics deals with the topic of *temperature*. Temperature is a statistical property – thus, it makes no sense to talk of the temperature of one, two, or even a handful of particles.

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1 Kinetic Theory of Gases

1.1 The molecular picture of matter

Imagine looking into a container filled with steam, and magnifying by a factor of 10^{10} . A cubic metre might contain around 20 molecules, all of which are in constant motion, colliding with the walls and each other. Suppose that one of the walls is a piston. The molecules which collide with the piston and impart a force on it; in order to fix the piston in place, a counter force must be applied.

Definition 1.1 (Pressure). The force per unit area applied by a gas on the walls of its container is called the pressure of the gas.

Now provide the system with heat. We know that the temperature of the gas must increase – what this means is that the speeds of the molecules increase, on average.

Definition 1.2 (Temperature). The temperature of a gas is a measure of the average kinetic energy of the constituent particles.

Instead, consider an adiabatic container, which stops all flow of heat into and out of the gas. By compressing the gas with the piston, we observe that the temperature of the gas also rises.

Now, take away heat from the system. The temperature drops and the molecules tend to be close to each other. This is because of the dipolar attractive forces between the molecules

(which varies as the inverse cube of the distance of the dipoles, and is hence comparatively short range). On the other hand, they cannot get too close, since once the electron clouds of the molecules start to overlap, a repulsive force is introduced. At a certain point, we reach a condensed form of matter: liquid water.

Liquid water is very much incompressible, yet the molecules freely move and slide around, without any periodic arrangement. The molecules at the surface are attracted by like molecules inside; this cohesive force keeps the liquid condensed. This tendency of a liquid to minimize its surface area is related to the phenomenon of surface tension. Some molecules on the surface are energetic enough to escape this cohesive attraction and leave the liquid – this is called evaporation. Heating a liquid simply increases the average kinetic energy of the molecules, thus increasing the rate of evaporation. When these energetic molecules leave the liquid, the average kinetic energy of the liquid drops, hence it cools down. This is the phenomenon of latent heat.

When this happens in a closed container, the process of evaporation cannot go on indefinitely, since the air has a limited capacity for holding moisture. Condensation is the process where these airborne molecules return to the liquid. At a certain point, the rates of evaporation and condensation become equal, and we obtain a saturated vapour.

Return to the liquid, and take away even more heat. Now, the motion of the molecules decrease to a point where they occupy fixed positions. They are still in motion, but their movement is restricted around their mean position. This is the crystal state. The lower the temperature, the smaller the oscillations and vibrations.

1.2 Basic assumptions

1. Gases are made up of a large number of molecules, and all molecules of one gas are identical.
2. Molecules of a gas are always moving. The number of molecules per unit volume remains constant, i.e. the density remains constant.
3. Molecules behave as elastic spheres during collisions. Kinetic energy and momenta are conserved, and the collision time is negligible compared to the free path time.
4. No force acts on any molecule, except during collisions. Intermolecular forces are only short ranged. Between collisions, the molecules continue moving with uniform velocity in a straight line.
5. The entire gas is isotropic; for all molecules, all directions are the same.

Remark. The collisions between molecules can be modelled as the elastic collision of hard spheres. The repulsive forces between molecules, governed by the Lennard-Jones potential, varies as $1/r^{12}$, which is very short range and very powerful. In comparison, gravity is a long range force since it varies as $1/r^2$.

1.3 Ideal gases

For an ideal gas, we make a few more assumptions. The gas molecules have negligible size, so are essentially point masses. Also, there are no forces on the molecules except during collisions, so they have no potential.

No real gases are ideal. We may look at the limit where the temperature T is very high and the density (or n) is very low. Here, the kinetic energy far exceeds any potential energies, and the mean free path becomes very high.

We look at some absurdities of this model.

1. How do point masses collide?
2. If two gases of different temperatures are mixed, how do they exchange heat?

3. Without intermolecular forces, are there any phase changes?
4. How do we explain properties such as viscosity and thermal conductivity?

1.4 Pressure

Suppose that a volume dV , located at (r, θ, ϕ) , contains $n dV$ particles. If we consider a small flat, horizontal area ΔS , we can calculate the number of molecules moving towards ΔS , as

$$dN = \frac{n \cos \theta \Delta S}{4\pi r^2} dV.$$

Over a time Δt , we only consider the particles within the region $r = 0 \rightarrow c\Delta t$ above the xy plane. Integrating, we have

$$\int dN = \int_0^{c\Delta t} \int_0^{\pi/2} \int_0^{2\pi} \frac{n \cos \theta \Delta S}{4\pi r^2} r^2 \sin \theta d\phi d\theta dr.$$

Simplifying, we have

$$N = \frac{1}{2} n \Delta S \cdot c \Delta t \cdot \int_0^{\pi/2} \cos \theta \sin \theta d\theta = \frac{1}{4} n c \Delta S \Delta t.$$

Thus, the number of molecules hitting the wall per unit area per unit time is given by $nc/4$.

What if we have different molecules with different velocities? We can use this expression to conclude that if n_i molecules have velocity c_i , the average velocity is $\bar{c} = \sum n_i c_i / n$, $n = \sum n_i$, so

$$N = \frac{1}{4} n \bar{c}.$$

Now, each molecule can strike the walls of the container at some angle θ . For an elastic collision, the change in its momentum is $2mc \cos \theta$. Repeating the integration process, we write the momentum imparted as

$$\int 2mc \cos \theta \cdot \frac{n \cos \theta \Delta S}{4\pi r^2} dV = mnc^2 \Delta S \cdot \int_0^{\pi/2} \cos^2 \theta \sin \theta d\theta.$$

Simplifying, we have

$$\frac{1}{3} mnc^2 \Delta S \Delta t.$$

For a velocity distribution, we deal with the RMS velocity where $c_{rms}^2 = \sum n_i c_i^2 / n$. Thus the pressure, which is the momentum imparted per unit area per unit time, is given by

$$p = \frac{1}{3} \rho c_{rms}^2.$$

Note that $\rho = mn$ is the density of the gas. Now, with knowledge of Boyle's Law and Charles' Law, we are forced to conclude that the temperature T is linearly dependent on c_{rms}^2 .

For a volume V of gas, we see that

$$pV = \frac{1}{3} mnV c_{rms}^2 = \frac{1}{3} mN c_{rms}^2,$$

where $N = nV$ is the total number of molecules. Now, the average kinetic energy of these N molecules is

$$E = \sum \frac{1}{2} mn_i c_i^2 = \frac{1}{2} mN c_{rms}^2.$$

Combining these relations, we have

$$pV = \frac{2}{3} E.$$

Proposition 1.1 (Ideal gas law). *The ideal gas law gives a relation between the pressure p , the volume V , the temperature T , and the number of moles n of an ideal gas.*

$$pV = nRT.$$

Here, the constant of proportionality R is called the ideal gas constant, with value

$$R \approx 8.314 \text{ J mol}^{-1} \text{ K}^{-1}.$$