

in $|a_1\rangle |a_2\rangle$ basis

$$\begin{pmatrix} 3 \\ 4 \end{pmatrix} \text{ shorthand for } 3|a_1\rangle + 4|a_2\rangle$$

$$\underbrace{\begin{pmatrix} 1 & 2 \\ 1 & 2 \end{pmatrix}} \begin{pmatrix} 3 \\ 4 \end{pmatrix} \text{ shorthand for } (1|a_1\rangle + 2|a_2\rangle)(3|a_1\rangle + 4|a_2\rangle)$$

$$|x_1\rangle = a|y_1\rangle + b|y_2\rangle$$

$$|x_2\rangle = c|y_1\rangle + d|y_2\rangle$$

\Rightarrow

$$|y_1\rangle = a'|x_1\rangle + b'|x_2\rangle$$

$$|y_2\rangle = c'|x_1\rangle + d'|x_2\rangle$$

A spin in a magnetic field

$$\vec{B} = B_0 \hat{z} + B_1 [\cos \omega t \hat{x} + \sin \omega t \hat{y}]$$

$$H_{\text{fl}} = \begin{pmatrix} \omega_0 & \omega_1 e^{-i\omega t} \\ \omega_1 e^{i\omega t} & -\omega_0 \end{pmatrix} \quad \omega_0 = \frac{e B_0}{m} \quad \omega_1 = \frac{e B_1}{m}$$

strength of fields
 ω : frequency of spinny field

We need to solve $i\hbar \frac{d}{dt} |\psi\rangle = H |\psi\rangle$

In fl basis

$$|\psi(t)\rangle = c_+(t) |\uparrow\rangle + c_-(t) |\downarrow\rangle = \begin{pmatrix} c_+ \\ c_- \end{pmatrix}$$

$$i\hbar \begin{pmatrix} \dot{c}_+ \\ \dot{c}_- \end{pmatrix} = \frac{\hbar}{2} \begin{pmatrix} \omega_0 & \omega_1 e^{-i\omega t} \\ \omega_1 e^{i\omega t} & -\omega_0 \end{pmatrix} \begin{pmatrix} c_+ \\ c_- \end{pmatrix}$$

$$i\hbar \dot{c}_+ = \frac{\hbar}{2} \omega_0 c_+ + \frac{\hbar}{2} \omega_1 e^{-i\omega t} c_-$$

$$i\hbar \dot{c}_- = \omega_1 e^{i\omega t} c_+ - \omega_0 c_-$$

Let $c_{\pm}(t) = \alpha_{\pm}(t) e^{\mp i\omega t/2}$

$$i\hbar \left[\dot{\alpha}_+ e^{-i\omega t/2} + \alpha_+ \left(-\frac{i\omega}{2}\right) e^{-i\omega t/2} \right] = \frac{\hbar}{2} \omega_0 \alpha_+ e^{-i\omega t/2} + \frac{\hbar \omega_1}{2} \alpha_- e^{-i\omega t/2} e^{-i\omega t}$$

$$i\hbar \dot{\alpha}_+ = i\hbar \left(\frac{\omega}{2} \right) \alpha_+ + \frac{\hbar \omega_0}{2} \alpha_+ + \frac{\hbar \omega_1}{2} \alpha_-$$

$$i\hbar \dot{\alpha}_+ = -\frac{\hbar \Delta\omega}{2} \alpha_+ + \frac{\hbar \omega_1}{2} \alpha_-$$

$$i\hbar \dot{\alpha}_- = \frac{\hbar \omega_1}{2} \alpha_+ + \frac{\hbar \Delta\omega}{2} \alpha_-$$

$\Delta\omega = \omega - \omega_0$
 \uparrow
 rotating freq.
 \uparrow
 strength of B_1
 spacing between energy levels
 $\hbar \omega_1$

$$i\hbar \frac{d}{dt} \begin{pmatrix} \alpha_+ \\ \alpha_- \end{pmatrix} = \frac{\hbar}{2} \begin{pmatrix} -\Delta\omega & \omega_1 \\ \omega_1 & \Delta\omega \end{pmatrix} \begin{pmatrix} \alpha_+ \\ \alpha_- \end{pmatrix}$$

$$i\hbar \frac{d}{dt} |\alpha\rangle = H' |\alpha\rangle$$

This is just like spin precession problem before
 $H = \frac{\hbar}{2} \begin{pmatrix} \omega_0 & \omega_1 \\ \omega_1 & -\omega_0 \end{pmatrix}$

if $|\alpha(0)\rangle = |\uparrow\rangle$

$$P_{\uparrow \rightarrow \downarrow}(t) = \frac{\omega_1^2}{(\Delta\omega)^2 + \omega_1^2} \sin^2 \left(\frac{\sqrt{\Delta\omega^2 + \omega_1^2}}{2} t \right)$$

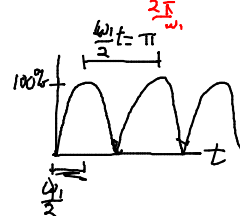
$|\psi\rangle = |\alpha\rangle \quad c_{\pm} = \alpha_{\pm} e^{\mp i\omega t/2}$

if $|\alpha(0)\rangle = |\uparrow\rangle = \begin{pmatrix} 1 \\ 0 \end{pmatrix} \quad \alpha_+ = 1 \quad \alpha_- = 0$

$$c_+ = 1 e^0 = 1 \quad c_- = 0$$

if $|\psi(0)\rangle = |\uparrow\rangle$

if $\Delta\omega \rightarrow 0$, then $P = \sin^2 \frac{\omega_1 t}{2}$



$$\begin{array}{l} |\uparrow\rangle \quad +\frac{\hbar\omega_0}{2} \\ \hline \hbar\omega_0 \\ \hline |\downarrow\rangle \quad -\frac{\hbar\omega_0}{2} \end{array} \quad \left[\begin{array}{l} \hbar\omega \\ \hline \hbar\omega \end{array} \right]$$

Choose rotation frequency to match energy jump required, & it will happen w/ 100% probability
Magnetic Resonance

Light-Matter Interactions similar

tune frequency of light to match

difference in energy levels

Function Spaces

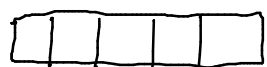
A vector space where the vectors are
complex functions of x, y, z, \dots


$$|\psi\rangle \equiv \psi(x)$$

In QM, we're particularly interested in
square-integrable functions

$$\int |\psi|^2 dx < \infty$$

Suppose a particle can be in one of 5 positions



$|1\rangle$: definitely in 1st position 

~~In general~~ $|1\rangle, |2\rangle, |3\rangle, |4\rangle, |5\rangle$ form a basis

$P_{|1\rangle \text{ on square 2}} = |\langle 2|1\rangle|^2 = 0$ orthogonal

In general

$$|\psi\rangle = \psi_1 |1\rangle + \psi_2 |2\rangle + \psi_3 |3\rangle + \psi_4 |4\rangle + \psi_5 |5\rangle$$

$$\langle 1|\psi\rangle = \psi_1$$

$$|\psi\rangle \equiv \begin{pmatrix} \psi_1 \\ \psi_2 \\ \psi_3 \\ \psi_4 \\ \psi_5 \end{pmatrix} \quad P(\psi \text{ is on square 1}) = |\langle 1|\psi\rangle|^2 = |\psi_1|^2$$

Now generalize to the number line

$$|3.257\dots m\rangle \quad |5m\rangle \text{ etc}$$

∞ -dimensional vector space

$$\langle x_1|x_2\rangle = \delta_{x_1, x_2}$$

$$|\psi\rangle = \int \psi(x) |x\rangle dx$$

$$\psi(x) = \langle x|\psi\rangle$$

e.g. $\psi(3m)$ = coefficient of term with $|x=3m\rangle$ in it
 $|x=3m\rangle$ = state where particle is at
position $x=3m$ w/ certainty.

Position operator has eigenstates $|x\rangle$

& if I measure $|\psi\rangle$ with position operator

I get $|x\rangle$ with probability $P = |\langle x|\psi\rangle|^2 = |\psi(x)|^2$