

Probability

if there are Ω states of a system
and every state is equally likely
then probability of being in one state
is $P = \frac{1}{\Omega}$

"accessible states" : states that system
could be in

microstate: complete description of state of the system

e.g. 3 coins,

HHH	HTH	TTH	TTT	$\Omega = 8$ microstates
HHT	HTT	THT	TTT	$P = \frac{1}{8}$

in many cases in physics, the accessible microstates
of a system are equally likely

Macrostate: a partial description

e.g. "exactly one head",
"first coin is tails"

multiplicity of a macrostate is

of microstates which satisfy the macrostate condition

e.g. 3 coins
"exactly one head" $\Omega = 3$
"first coin is tails" $\Omega = 4$

HTT
THT
<u>TTT</u>
TTH
TTT

probability of a macrostate, if all microstates are
equally likely,

$$P(\text{macro}) = \frac{\Omega(\text{macro})}{\Omega_{\text{all}}}$$

Ω_{all} : total # of accessible microstates

e.g. P of getting 2H after flipping 10 coins?

$$\Omega_{\text{all}} = 2^{10} = 1024 \quad \Omega = \binom{10}{2} = \frac{10!}{8!2!} = \frac{10 \cdot 9}{2} = 45$$

$$P = \frac{45}{1024} = 4.4\%$$

Useful Models for Statistical Mechanics

- "toy" models
- easy mathematically
- demonstrate interesting physical properties

1) Paramagnet

- small magnetic dipoles $\vec{\mu}$, independent
classical: $\uparrow \downarrow \leftarrow \rightarrow \uparrow \uparrow \leftarrow \searrow \searrow$

in quantum mechanics: $\uparrow\downarrow\uparrow\uparrow\downarrow\downarrow\uparrow\downarrow$

$\vec{\mu} \cdot \frac{\hbar}{2} = \mu_z = \pm \mu_0$
for spin- $\frac{1}{2}$ particle.

if N dipoles, $\Omega_{\text{all}} = 2^N$
 N_{\uparrow} point up, N_{\downarrow} point down
 $N = N_{\uparrow} + N_{\downarrow}$

Now place dipoles in an external magnetic field $\vec{B} = -B_z \hat{z}$
 dipoles want to align with the field \Downarrow

$\vec{B} \downarrow$ $\uparrow \mu$ high μ_B $\downarrow \mu$ low μ_B $U = -\vec{\mu} \cdot \vec{B}$
 $U = +\mu_B B$ $U = -\mu_B B$

$$\begin{aligned} \text{total PE} \quad U &= N_{\uparrow} (+\mu_B) + N_{\downarrow} (-\mu_B) \\ &= \mu_B (N_{\uparrow} - N_{\downarrow}) \\ &= \mu_B (2N_{\uparrow} - N) \end{aligned}$$

$$U = \underbrace{2\mu_0 B}_{\text{choose units of energy where } 2\mu_0 B = 1} N_T - \mu_0 B N$$

$$\rightarrow U = N_T$$

How many values can U take?

$U = 0, 1, \dots, N$ $N+1$ possibilities
 $N+1$ energy macrostates

eg. $U=2$
 $N=4$ $\uparrow\uparrow\downarrow\downarrow$ $\uparrow\downarrow\uparrow\downarrow$ \leftarrow two microstates
in the $U=2$
macrostate

multiplicity? $\Omega(U) = \binom{N}{U}$

$$\Omega(N_T) = \frac{N!}{N_T!(N-N_T)!} = \frac{N!}{N_T!N_b!}$$

isolated paramagnet (no energy exchange with environment)

U will be constant (conservation of energy)

accessible microstates are $N_T = 0$

Paramagnet jumps from one microstate to the next as energy sloshes around inside

at any moment probability of any one accessible microstate

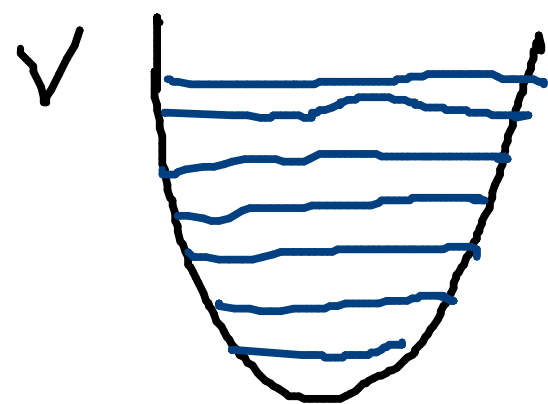
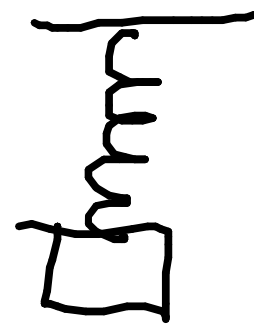
$$P = \frac{1}{\Omega(v)}$$

4 $N=10$ $\Sigma(U) = \binom{10}{2} = 45$

$U=2$ P_{off} 1111111111 is $\frac{1}{45}$

Einstein Solid

quantum harmonic oscillator

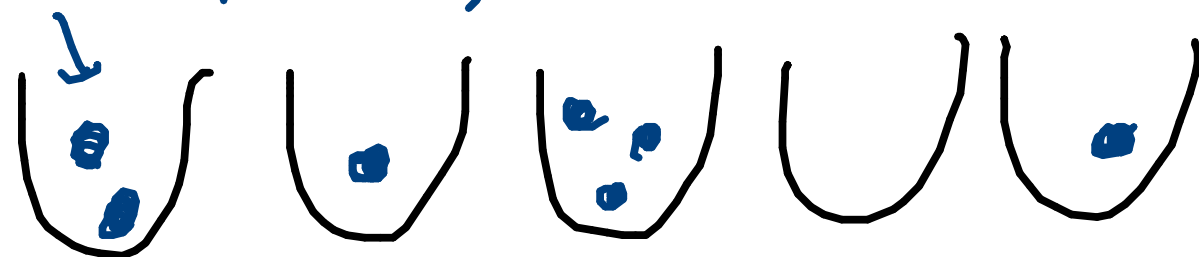


$$E = (q + \frac{1}{2}) h f \leftarrow \text{frequency of oscillator (fixed)}$$

$q = 0, 1, 2, 3, \dots$

consider N harmonic oscillators

quanta of energy



$$U = q = q_1 + q_2 + q_3 + q_4 + q_5$$

$2 \quad 1 \quad 3 \quad 0 \quad 1$

$$\begin{aligned} U &= \sum_{i=1}^N E_i \\ &= \sum_{i=1}^N (q_i + \frac{1}{2}) h f \\ &= h f \sum_{i=1}^N q_i + \frac{1}{2} h f N \\ &= 1 \quad \text{in same units} \end{aligned}$$

\uparrow baseline