

Scattering states $E > 0$

$$x < -a: \psi(x) = A e^{ikx} + B e^{-ikx} \quad k = \frac{\sqrt{2mE}}{\hbar}$$

incident reflected

$$x > a: \psi(x) = F e^{ikx} \quad \text{transmitted}$$

$$T = \frac{|F|^2}{|A|^2} \quad R = \frac{|B|^2}{|A|^2}$$



$$-a < x < a: C \sin lx + D \cos lx \quad l = \frac{\sqrt{2m(E+V_0)}}{\hbar}$$

B.C. $\psi_-(-a) = \psi_+(-a) \Rightarrow A e^{-ika} + B e^{ika} = -C \sin la + D \cos la$

$$\psi'_-(-a) = \psi'_+(-a) \Rightarrow ik[A e^{-ika} - B e^{ika}] = l[C \cos la + D \sin la]$$

$$x = +a: C \sin la + D \cos la = F e^{ika}$$

$$l[C \cos la - D \sin la] = ik F e^{ika}$$

DO MATH

$$\frac{|A|^2}{|F|^2} = \frac{1}{T} = 1 + \frac{V_0^2}{4E(E+V_0)} \sin^2\left(\frac{2a}{\hbar} \sqrt{2m(E+V_0)}\right)$$

If $\sin^2 = 0$, then $T = 1$

for certain values of E , there is no reflection
the potential is "transparent"

To get transparency, $\frac{2a}{\hbar} \sqrt{2m(E+V_0)} = n\pi$

$$E_n + V_0 = \frac{n^2 \pi^2 \hbar^2}{2m(2a)^2}$$

distance above bottom of well

infinite square well
width = 2a

$$E_n = \frac{n^2 \pi^2 \hbar^2}{2m(\text{width})^2}$$

All energies can exist for scattering state,
but energies that match the allowed energies
of an ∞ \square well will have $T=1$ & $R=0$.

On other hand, if $\sin^2 = 1$

$$\frac{1}{T} = 1 + \frac{V_0^2}{4E(E+V_0)} \rightarrow T = \frac{4E(E+V_0)}{4E(E+V_0) + V_0^2}$$

as E goes up,
 $T \rightarrow 1$ anyway -

Chapter 3

Uncertainty Principle

$$\sigma_A \sigma_B \geq \left| \frac{1}{2i} \langle [A, B] \rangle \right|$$

$$\begin{aligned} \text{e.g. } [x, p] &= x \frac{\hbar}{i} \frac{d}{dx} f - \frac{\hbar}{i} \frac{d}{dx} x f \\ &= \frac{\hbar}{i} [x f' - (f + x f')] \\ &= \frac{\hbar}{i} (-f) \\ &= i\hbar \end{aligned}$$

$$\begin{aligned} \sigma_x \sigma_p &\geq \left| \frac{1}{2i} \langle i\hbar \rangle \right| \\ &\geq \left| \frac{i\hbar}{2i} \right| \end{aligned}$$

$$\sigma_x \sigma_p \geq \frac{\hbar}{2} \quad \text{Heisenberg Uncertainty Principle}$$

Compare $\sigma_{S_z} \sigma_{S_y}$ for $|↑\rangle$

$$\begin{aligned} \sigma_{S_z} \sigma_{S_y} &\geq \left| \frac{1}{2i} \langle [S_z, S_y] \rangle \right| \\ &\geq \left| \frac{1}{2i} \langle i\hbar (-S_x) \rangle \right| \\ &\geq \frac{\hbar}{2} |\langle S_x \rangle| \quad \langle S_x \rangle = 0 \\ &\geq 0 \end{aligned}$$

So we can have $\sigma_{S_z} = 0$

$\left(\begin{array}{ccc} \frac{1}{2} & \text{come out} & +\frac{\hbar}{2} \\ \frac{1}{2} & " & -\frac{\hbar}{2} \end{array} \right)$

$$\text{But } \sigma_x \sigma_p \geq \frac{\hbar}{2}$$

So x is never perfectly known
& p " " " "

Position & Momentum Eigenstates are not real states.

$$\Delta x \Delta p \geq \frac{\hbar}{2}$$

$$\begin{aligned} k &= \frac{2\pi}{\lambda} \\ \lambda &= \frac{h}{p} \quad \text{de Broglie} \\ k &= \frac{2\pi}{h} p \end{aligned}$$

position eigenstate at $x=x_0$ is $\delta(x-x_0)$
which is not a function
does not describe a real thing

momentum eigenstate is e^{ikx} $k = \frac{2\pi}{h} p$

$$\int_{-\infty}^{\infty} |\psi|^2 dx = \int_{-\infty}^{\infty} |e^{ikx}|^2 dx = \int_{-\infty}^{\infty} 1 dx = \infty,$$

non-normalizable

$$\Delta E \Delta t \geq \frac{\hbar}{2}$$