I dentical Portides $\Psi(\vec{r},\vec{r}_3) = \pm \Psi(\vec{r}_3,\vec{r}_3)$ +: bosons -: fermions symmetric antisymmetric e.g. Calculating $\langle (x_1 - x_2)^2 \rangle$ Distinguishable Particles $\Psi(x_1, x_2) = \Psi_a(x_1) \Psi_b(x_2)$ $\langle (x_1-x_2)^2 \rangle = \langle \chi_1^2 \rangle + \langle \chi_2^2 \rangle - 2 \langle \chi_1 \chi_2 \rangle$ $\langle x^2 \rangle = \langle \Psi | \chi^2 | \Psi \rangle$ = $\langle \Psi_a(x_1)\Psi_b(x_2)|\chi_1^2|\Psi_a(x_1)\Psi_b(x_2)\rangle$ $= \langle \Psi_{a}(x_{1}) | \chi_{1}^{2} | \Psi_{b}(x_{1}) \rangle \langle \Psi_{b}(x_{2}) | \Psi_{b}(x_{2}) \rangle$ $= \langle x_{s} \rangle$ $\langle \chi_2^2 \rangle = \langle \chi^2 \rangle_{\underline{L}}$ \(\text{x, x_2} \) = \(\frac{\frac{1}{a(\text{x_1})\frac{1}{b(\text{x_2})}}{\text{x_1 \text{x_2}}} \frac{\frac{1}{a(\text{x_1})\frac{1}{b(\text{x_2})}}{\text{x_1 \text{x_2}}} \] = < \times > $= \langle (x_1 - x_2)^2 \rangle = \langle x^2 \rangle_a + \langle x^2 \rangle_b - 2 \langle x \rangle_b$

$$\begin{split} & \psi_{\pm}(x_{1},x_{2}) = \frac{1}{12} \left[\psi_{a}(x_{1})\psi_{b}(x_{2}) \pm \psi_{a}(x_{2})\psi_{b}(x_{1}) \right] \\ & \langle \chi_{1}^{2} \rangle = \frac{1}{2} \langle \psi_{a}(x_{1})\psi_{b}(x_{2}) \pm \psi_{a}(x_{2})\psi_{b}(x_{1}) \right] \langle \chi_{1}^{2} | \psi_{a}(x_{1})\psi_{b}(x_{2}) \pm \psi_{a}(x_{2})\psi_{b}(x_{1}) \rangle \\ & = \frac{1}{2} \langle \psi_{a}(x_{1})\psi_{b}(x_{2}) | \chi_{1}^{2} | \psi_{a}(x_{1})\psi_{b}(x_{2}) \rangle \\ & \leq \frac{1}{2} \langle \psi_{a}(x_{2})\psi_{b}(x_{1}) | \chi_{1}^{2} | \psi_{a}(x_{1})\psi_{b}(x_{2}) \rangle \\ & = \frac{1}{2} \langle \psi_{a}(x_{2})\psi_{b}(x_{1}) | \chi_{1}^{2} | \psi_{a}(x_{1})\psi_{b}(x_{2}) \rangle \\ & = \frac{1}{2} \langle \chi^{2} \rangle_{b} + \langle \chi^{2} \rangle_{b} \right] \\ & \langle \chi_{1}^{2} \rangle_{c} = \frac{1}{2} \left[\langle \chi^{2} \rangle_{b} + \langle \chi^{2} \rangle_{b} \right] \\ & \langle \chi_{1}^{2} \rangle_{c} = \frac{1}{2} \left[\langle \chi^{2} \rangle_{b} + \langle \chi^{2} \rangle_{b} \right] \\ & \langle \chi_{1}^{2} \rangle_{c} = \frac{1}{2} \left[\langle \chi^{2} \rangle_{b} + \langle \chi^{2} \rangle_{b} \right] \\ & \langle \chi_{1}^{2} \rangle_{c} = \frac{1}{2} \left[\langle \chi^{2} \rangle_{b} + \langle \chi^{2} \rangle_{b} \right] \\ & \langle \chi_{1}^{2} \rangle_{c} = \frac{1}{2} \left[\langle \chi^{2} \rangle_{b} + \langle \chi^{2} \rangle_{b} \right] \\ & \langle \chi_{1}^{2} \rangle_{c} = \frac{1}{2} \left[\langle \chi^{2} \rangle_{b} + \langle \chi^{2} \rangle_{b} \right] \\ & \langle \chi_{1}^{2} \rangle_{c} = \frac{1}{2} \left[\langle \chi^{2} \rangle_{c} + \langle \chi^{2} \rangle_{c} \right] \\ & \langle \chi_{1}^{2} \rangle_{c} = \frac{1}{2} \left[\langle \chi^{2} \rangle_{c} + \langle \chi^{2} \rangle_{c} \right] \\ & \langle \chi_{1}^{2} \rangle_{c} = \frac{1}{2} \left[\langle \chi^{2} \rangle_{c} + \langle \chi^{2} \rangle_{c} \right] \\ & \langle \chi_{1}^{2} \rangle_{c} = \frac{1}{2} \left[\langle \chi^{2} \rangle_{c} + \langle \chi^{2} \rangle_{c} \right] \\ & \langle \chi_{1}^{2} \rangle_{c} = \frac{1}{2} \left[\langle \chi^{2} \rangle_{c} + \langle \chi^{2} \rangle_{c} \right] \\ & \langle \chi_{1}^{2} \rangle_{c} = \frac{1}{2} \left[\langle \chi^{2} \rangle_{c} + \langle \chi^{2} \rangle_{c} \right] \\ & \langle \chi_{1}^{2} \rangle_{c} = \frac{1}{2} \left[\langle \chi^{2} \rangle_{c} + \langle \chi^{2} \rangle_{c} \right] \\ & \langle \chi_{1}^{2} \rangle_{c} = \frac{1}{2} \left[\langle \chi^{2} \rangle_{c} + \langle \chi^{2} \rangle_{c} \right] \\ & \langle \chi_{1}^{2} \rangle_{c} = \frac{1}{2} \left[\langle \chi^{2} \rangle_{c} + \langle \chi^{2} \rangle_{c} \right] \\ & \langle \chi_{1}^{2} \rangle_{c} = \frac{1}{2} \left[\langle \chi^{2} \rangle_{c} + \langle \chi^{2} \rangle_{c} \right] \\ & \langle \chi_{1}^{2} \rangle_{c} = \frac{1}{2} \left[\langle \chi^{2} \rangle_{c} + \langle \chi^{2} \rangle_{c} \right] \\ & \langle \chi_{1}^{2} \rangle_{c} = \frac{1}{2} \left[\langle \chi^{2} \rangle_{c} + \langle \chi^{2} \rangle_{c} \right] \\ & \langle \chi_{1}^{2} \rangle_{c} = \frac{1}{2} \left[\langle \chi^{2} \rangle_{c} + \langle \chi^{2} \rangle_{c} \right] \\ & \langle \chi_{1}^{2} \rangle_{c} = \frac{1}{2} \left[\langle \chi^{2} \rangle_{c} + \langle \chi^{2} \rangle_{c} \right] \\ & \langle \chi_{1}^{2} \rangle_{c} = \frac{1}{2} \left[\langle \chi^{2} \rangle_{c} + \langle \chi^{2} \rangle_{c} \right] \\ & \langle \chi_{1}^{2} \rangle_{c} = \frac{1}{2} \left[\langle \chi^{2} \rangle_{c} + \langle \chi^{2} \rangle_{c} \right] \\ & \langle \chi_{1$$

e.g. Hydrogen molecule F E E F lectrons are in the middle,

they will pull the protons inverd

and hold molecule together

If electrons are on the outsider protons will push each other apart I for electorons to end up in the middle, they never to act like busons their spatial navekunction $\Psi(r_1, v_2)$ is symmetric Overall varefunction is $\psi(r_1, r_2) \chi(s_1, s_2)$ We can have symmetric Ψ If we have an antisymmetric X So that overall wavefunction is antisymmetric as befitting fermions X(S, S2) 15. one of four possibilities symmetric $\int_{\sqrt{2}}^{1} (fl+lf)$ 痘(アレーレア) "singlet state" is antisymmetric -, electrons in a Hz molecule ar in the singlet state 5=0 m=0