

Physics 112 Problem Set 1

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1. (a) $7^4 = 2401$
 (b) $\frac{7!}{4!} = 840$
 (c) $\frac{7!}{3!4!} = 140$
 (d) $\frac{10!}{4!6!} = 210$
2. (a) $P(5000) = \frac{10000!}{5000!5000!}(0.5)^{10000} = 0.00798$
 (b) $P(5100)/P(5000) = \frac{5000!5000!}{5100!4900!} = 0.135$
 (c) $P(6000)/P(5000) = \frac{5000!5000!}{6000!4000!} = 3.643 \times 10^{-88}$
 (d) $P(6)/P(5) = \frac{5!5!}{6!4!} = 0.833$
3. (a) The probability of a single sequence of N trials with n events occurring is $p^n(1-p)^{N-n}$, and there are $\binom{N}{n}$ sequences of length N with n events, so the total probability of n events occurring is

$$P(n) = \frac{N!}{n!(N-n)!} p^n (1-p)^{N-n}.$$

- (b) With $\lambda = Np$, we have

$$P(n) = \frac{N(n-1)(n-2)\dots(N-n+1)}{n!} \left(\frac{\lambda}{N}\right)^n \left(1 - \frac{\lambda}{N}\right)^{N-n}$$

$$P(n) = \frac{N^n + \mathcal{O}(N^{n-1})}{n!} \left(\frac{\lambda}{N}\right)^n \left(1 - \frac{\lambda}{N}\right)^{N-n}.$$

For large n , this is approximately

$$P(n) \approx \frac{\lambda^n}{n!} \left(1 - \frac{\lambda}{N}\right)^{N-n}.$$

Take the limit as $n \rightarrow \infty$, and recall $\lim_{x \rightarrow \infty} (1 + \frac{1}{x})^x = e$. Then

$$P(n) \approx \frac{\lambda^n}{n!} e^{-\lambda}.$$

4. (a)

$$\begin{aligned}
g(N, s) &= \frac{N!}{(N/2 + s)!(N/2 - s)!} \\
g(N, s) &= \frac{\sqrt{2\pi N} N^N e^{-N}}{\sqrt{2\pi(N/2 + s)} (N/2 + s)^{N/2+s} e^{-N/2-s} \sqrt{2\pi(N/2 - s)} (N/2 - s)^{N/2-s} e^{-N/2+s}} \\
g(N, s) &= \frac{\sqrt{N} N^N}{\sqrt{2\pi}} \frac{1}{N/2 \sqrt{(1 + 2s/N)(1 - 2s/N)}} \frac{1}{[N/2(1 + 2s/N)]^{N/2+s}} \frac{1}{[N/2(1 - 2s/N)]^{N/2-s}} \\
g(N, s) &= \frac{\sqrt{2} N^N}{\sqrt{\pi N}} \frac{1}{(N/2)^N} \frac{1}{[(1 + 2s/N)(1 - 2s/N)]^{N/2}} \frac{1}{(1 + 2s/N)^{1/2+s}} \frac{1}{(1 - 2s/N)^{1/2-s}} \\
g(N, s) &= 2^N \sqrt{\frac{2}{\pi N}} \left(1 - 4 \frac{s^2}{N^2}\right)^{-N/2} \frac{(1 - 2s/N)^{s-1/2}}{(1 + 2s/N)^{s+1/2}}
\end{aligned}$$

(b) Call the expression from (a) $g_1(N, s)$ and the expression from (b) $g_2(N, s)$. Then

- $g_1(10, 1)/g_2(10, 1) = 1.019$
- $g_1(1000, 100)/g_2(1000, 100) = 0.891$
- $g_1(1000, 10)/g_2(1000, 10) = 1.000$.

(c) These are both approximations to the true multiplicity function, so it's expected that they give similar results. I expect the first function, g_1 , to be more accurate, since it only used the assumption that $N \gg 1$, from the Stirling approximation. The function derived in class, g_2 , also used this assumption, but added the additional assumption that $s \ll N$, from the Taylor expansion of $\ln(1 - 2s/N)$. This is why we see better agreement between the two functions for $g(1000, 10)$ than for $g(1000, 100)$, because the $s \ll N$ assumption is more accurate in the first case. The main limitation of the function in (a), then, is just the requirement that N is large.

5. (a)

$$\begin{aligned}
\langle m(s) \rangle &= \sum_{m=-R}^R m P(m, s) \\
&= \sum_{m=-R}^R m \left[\frac{R+m+1}{2R} P(m+1, s-1) + \frac{R-m+1}{2R} P(m-1, s-1) \right].
\end{aligned}$$

Reindexing the sums,

$$\langle m(s) \rangle = \sum_{m'=-R+1}^{R+1} (m' - 1) \left[\frac{R+m'}{2R} P(m', s-1) \right] + \sum_{m'=-R-1}^{R-1} (m' + 1) \left[\frac{R-m'}{2R} P(m', s-1) \right].$$

We know $-R \leq m' \leq R$, since there are only $2R$ fleas, so the $m' = R+1$ and $m' = -R-1$ terms are unphysical, and we drop them. Furthermore, we can add the term indexed by $m = -R$ to the first sum and the term indexed by $m' = R$ to the second sum, since the value of these terms is zero. So

$$\begin{aligned}
\langle m(s) \rangle &= \sum_{m'=-R}^R (m' - 1) \left[\frac{R + m'}{2R} P(m', s - 1) \right] + \sum_{m'=-R}^R (m' + 1) \left[\frac{R - m'}{2R} P(m', s - 1) \right] \\
&= \left[\frac{1}{2} \langle m(s - 1) \rangle + \frac{1}{2R} \langle m^2(s - 1) \rangle - \frac{1}{2} - \frac{1}{2R} \langle m(s - 1) \rangle \right] + \\
&\quad \left[\frac{1}{2} \langle m(s - 1) \rangle - \frac{1}{2R} \langle m^2(s - 1) \rangle + \frac{1}{2} - \frac{1}{2R} \langle m(s - 1) \rangle \right] \\
&= \left(1 - \frac{1}{R} \right) \langle m(s - 1) \rangle.
\end{aligned}$$

- (b) We will prove this by induction. At the beginning, the excess number of fleas on Rover is n , so $\langle m(0) \rangle = n = n(1 - 1/R)^0$. So the base case holds.

Assume $\langle m(k) \rangle = n(1 - 1/R)^k$. Then from the recurrence relation in (a), $\langle m(k + 1) \rangle = (1 - 1/R) \times n(1 - 1/R)^k = n(1 - 1/R)^{k+1}$.

So, by induction, $\langle m(s) \rangle = n(1 - 1/R)^s$ for all integers $s \geq 0$.

Figure 1: $\langle m(s) \rangle$ for $R = 10$

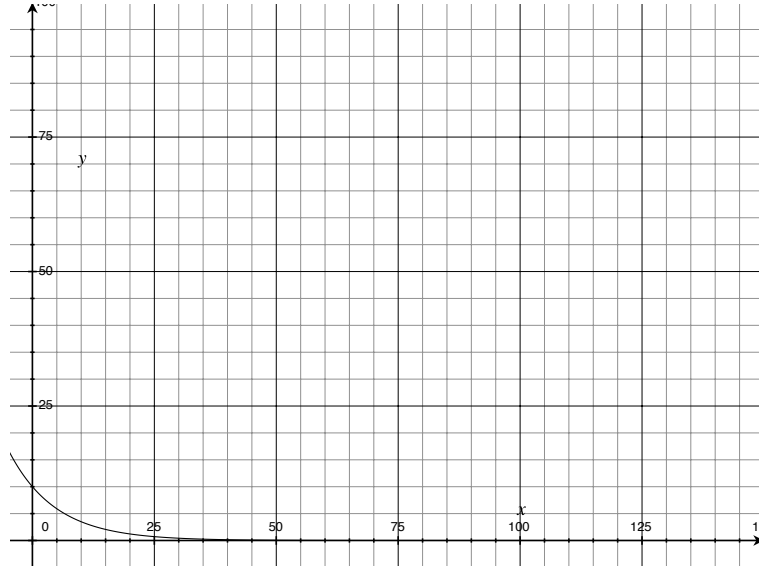
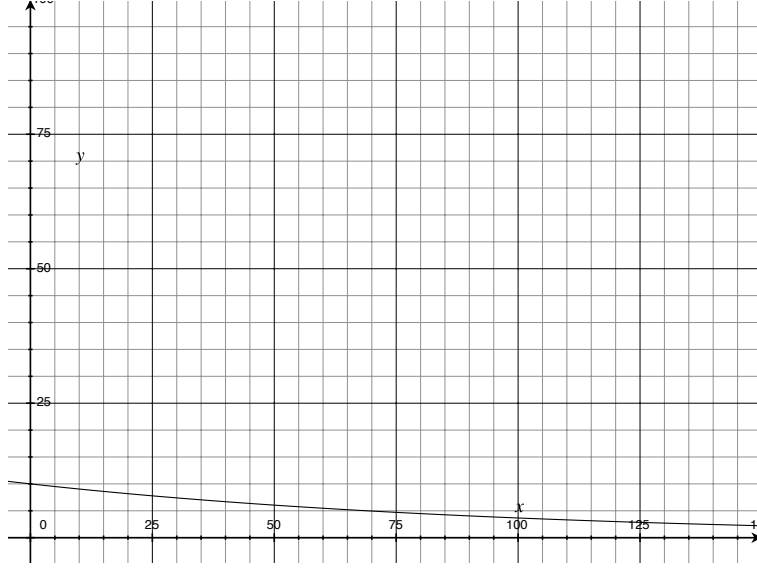


Figure 2: $\langle m(s) \rangle$ for $R = 100$



- (c) This behavior is roughly expected – we expect the fleas to be equalized between the two dogs ($\langle m(s) \rangle \rightarrow 0$ as $s \rightarrow \infty$), and the more fleas they are (larger R), the longer this will take.
- (d) For large R , $\langle m(s) \rangle = n \exp(s \ln(1 - 1/R)) \approx n \exp(-s/R)$. Making the analogy to RC circuits, the characteristic time scale of this exponential decay is R . That is, it takes about R steps for the flea excess to fall off appreciably (to $1/e$). The larger the number of fleas, the longer it will take to reach equilibrium, and this relationship is roughly linear in R .
6. (a) The average path length is 10^{-5} m, and the projection of a path of this length in the z direction is $10^{-5} \cos \theta$. So the average displacement in the z direction, integrating over all solid angles, is

$$\langle z \rangle = \frac{1}{4\pi} \int_0^\pi \int_0^{2\pi} 10^{-5} \cos \theta \sin \theta \, d\phi d\theta = 0,$$

as expected, since the particles are equally likely to travel in the $+z$ and $-z$ directions.

- (b) The same argument applies, so

$$\langle z^2 \rangle = \frac{1}{4\pi} \int_0^\pi \int_0^{2\pi} (10^{-5} \cos \theta)^2 \sin \theta \, d\phi d\theta = 3.333 \times 10^{-11},$$

thus $\sqrt{\langle z^2 \rangle} = 5.774 \times 10^{-6}$ m.

- (c) In 2 seconds, the ammonia molecules experience $N = 2 \times 10^7$ collisions. So we can think of the displacement in the z direction as the sum of 2×10^7 independent random variables, each with mean $\langle z \rangle = 0$ and standard deviation $\sigma = \sqrt{\langle z^2 \rangle - \langle z \rangle^2} = 5.774 \times 10^{-6}$ m.

So, the average displacement after 2 seconds is $N\langle z \rangle = 0$. The standard deviation is $\sqrt{N}\sigma = \sqrt{2 \times 10^7}(5.774 \times 10^{-6}) = 0.0258$ m.

- (d) The standard deviation after t seconds is $\sqrt{t \times 10^7}(5.774 \times 10^{-6})$. When the standard deviation of z is 6 m, approximately 32% of the molecules are further than 6 m from the origin (assuming a Gaussian distribution, which is reasonable given the large number of molecules). This occurs at

$$t = \left(\frac{6}{\sqrt{10^7}(5.774 \times 10^{-6})} \right)^2 = 108000 \text{ s},$$

or 300 hours.

7. (a) $P = \left(\frac{1}{44}\right)^{100000}$, so $\log_{10} P = -10000 \log_{10} 44 = -164345$, thus $P = 10^{-164345}$.
- (b) Typing 10 keys per second, each monkey has typed $10^{18} \times 10 = 10^{19}$ keys since the beginning of the universe. Each key (approximately) represents the start of another 10^5 character sequence. So the 10^{10} monkeys have had $10^{10} \times 10^{19} = 10^{29}$ opportunities to type Hamlet. The probability of any one random sequence resulting in Hamlet is $10^{-164345}$, from (a), so the probability that Hamlet has been typed is $10^{29} \times 10^{-164345} = 10^{-164316}$ (which is still extremely small).