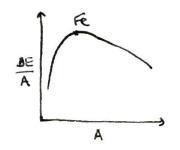
Let's recall the binding energies per nucleon of different elements:



So while heavier elements after iron home decreasing binding energy, elements lighter than iron have increasing binding energy with no nucleone the increase is also for more rapid than the decrease after iron.

Thus, much more energy can be released by fusion of two hight nuclei also a nearier nuclei than in the case of fission. Eg, let's consider the fusion of a deuteron and a hydrogen nucleus:  ${}^{2}H + {}^{1}H \rightarrow {}^{3}He + \Upsilon$ 

The  $\gamma$  is emitted since Helium is made in an excited state. The mass of  $^{2}$ H is  $3.34368 \times 10^{-27}$ kg  $^{1}$ H is  $1.67262 \times 10^{-27}$ kg  $^{3}$ He is  $5.00832 \times 10^{-27}$ kg

using  $E = D mc^2$  where D m is mass difference between left and right sides at the chemical formula, we find E = 4.4 MeV which is carried off by the r and the kinetic energy of the  $\frac{3}{2}$  He nucleus. Although this energy seems less than for fission, the energy per nucleon is much higher.

Fusion reactions also usually do not occur spontaneously because the nuclei repel each other due to coulomb interaction and must have enough energy to overcome the coulomb barrier to get close enough to puse.

An estimate of the required energy for 2 protons to tuse is:  $E = \frac{e^2}{4\pi \epsilon_0 R}$  where R is nuclear radius

if RN4 pm, then EN 350 keV
In order for the protons to have this average every associated with the degree of freedom corresponding to motion in the direction of the target proton, we would need a temperature T:

 $E = \frac{1}{2} KT$  where K is the Boltzman constant so  $T \sim 4 \times 10^9 K$  which is really hot.

In practice though, fusion can happen at lower temperatures such as in our Sun at  $1.3 \times 10^{7} \, \text{K}$  because:

- i) The temperature we determined was for the average energy per photor, but energies are distributed according to Movemell Botteman distribution, which has a tail, and this tail means that there are some particles whose energy is much larger than average.
- 2) It is not necessary for the incident probons to have sufficient energy to overcome the coulomb barrier. They are also get through by quantum turneling, provided the barrier height is not too high above the kinete energy of the incoming particle.

Fusion is the source of energy in stars and it works in cycles, a series at stages of turion in which the initial particles are proton but it subsequent stages tusion can happen between a turion product and a proton. At the end of the cycle, all intermediate turion products have disappeared, leaving only a stable tusion product, usually 4 He.

The most common such cycle is the proton cycle:

At the end of the cycle, there is no  $\frac{3}{2}$  He left. The overall process is:  $4 \mid H \rightarrow \frac{4}{2}H + 2e^{+} + 2V + 2V$ 

The total energy released by this is 24.7 MeV.

other cycles also occur within the sun:

This is called the carbon cycle

Two steps shown above are B decays of tusion products.

The initial carbon is produced by the tusion of 3 4 He nuclei.

The tusion process in the carbon cycle require more emergy to overcome the contomb barrier so is only likely at higher temperatures.

But at higher temperatures, this cycle is more likely than the proton cycle, because in the proton cycle we require 2 2 He to tuse which is unitely because the 3 He are produces from previous tusion processes and their density is low.

The sur is relatively cool for which the proton cycle dominates.

The produced Y-rows scatter around other charged particles in the sun, losing overgy at each scattering and eventually "themalize", i.e settle at our energy distribution which is the black-body distribution.