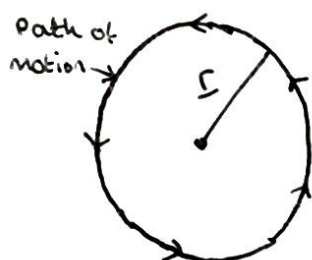


## Motion under Forces

If given a specified force  $\underline{F}$ , it may be possible to find the motion under this force using Newton's 2nd law. This is done by integrating twice:

$\underline{F} = m \cdot \underline{a}$  integrate both sides twice with respect to time to get the vector of motion  $\underline{r}$  from  $\underline{a}$ .

## Circular Motion



if  $|\underline{r}|$  remains constant through the object's motion, then the object is undergoing circular motion in two dimensions. To obtain the velocity, we start with:

$$\underline{r} \cdot \underline{r} = |\underline{r}|^2 \quad (\text{differentiating both sides with respect to } t:)$$

↑ differentiates to 0 since it is a constant

$$2\underline{r} \cdot \dot{\underline{r}} = 0$$

$$\underline{r} \cdot \dot{\underline{r}} = 0$$

↑  
this is velocity  $\underline{v}$

If  $\underline{v} \cdot \underline{r} = 0$ , then  $\underline{v}$  must be orthogonal to  $\underline{r}$ . Therefore velocity is always perpendicular to the radius. To find the acceleration, we start with:

$$\dot{\underline{r}} \cdot \dot{\underline{r}} = |\dot{\underline{r}}|^2 \quad (\text{differentiating b.s. w.r.t. } t:)$$

$$2\dot{\underline{r}} \cdot \ddot{\underline{r}} = 0$$

$$\dot{\underline{r}} \cdot \ddot{\underline{r}} = 0$$

↑  
this is acceleration

Therefore, acceleration is perpendicular to velocity  $\underline{v}$  and parallel (anti-parallel) to the radius  $\underline{r}$ .

We can actually derive the centripetal acceleration towards the centre from:

$$\underline{r} \cdot \dot{\underline{r}} = 0$$

velocity is perpendicular to radius. Differentiate both sides with respect to time:

$$\dot{\underline{r}} \cdot \dot{\underline{r}} + \underline{r} \cdot \ddot{\underline{r}} = 0$$

$$\underline{r} \cdot \ddot{\underline{r}} = -\dot{\underline{r}} \cdot \dot{\underline{r}} = -|\underline{v}|^2$$

but  $\underline{r}$  and  $\ddot{\underline{r}}$  are parallel so:

$$\begin{aligned}\underline{r} \cdot \ddot{\underline{r}} &= |\underline{r}| |\ddot{\underline{r}}| \cos 0 \\ &= |\underline{r}| |\ddot{\underline{r}}|\end{aligned}$$

but  $|\ddot{\underline{r}}| = \frac{\ddot{\underline{r}}}{\hat{\underline{r}}}$  note that  $\hat{\ddot{\underline{r}}}$  is just  $\hat{\underline{r}}$  since they are in the same direction

$$\therefore \frac{\ddot{\underline{r}}}{\hat{\underline{r}}} |\underline{r}| = -|\underline{v}|^2$$

$$\therefore \ddot{\underline{r}} = \frac{-|\underline{v}|^2}{|\underline{r}|} \hat{\underline{r}}$$

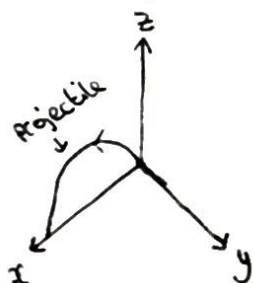
This is the centripetal acceleration towards the centre:

$$\underline{a} = \frac{-|\underline{v}|^2}{|\underline{r}|} \hat{\underline{r}}$$

Note that this has direction  $-\hat{\underline{r}}$  which is towards the centre.

## Projectile Motion

The force on a particle experiencing projectile motion is a constant given by  $m|g|$  where  $|g|$  is the strength of the gravitational vector.



If the projectile is moving in the  $x$ - $z$  plane, then  $g$  is given by:  $g = -|g|\hat{k}$ . This means the projectile experiences a constant force of  $F = -m|g|\hat{k}$ . Using Newton's 2nd Law, we get:

$$F = ma = -m|g|\hat{k} \quad \text{for force in the vertical direction}$$

$$\therefore a = -|g|\hat{k} \quad \text{integrating with respect to time:}$$

$$\begin{aligned} v &= -|g|t\hat{k} + \text{constant} \quad \left[ \begin{array}{l} \text{this constant is } v_0 \text{ since it is} \\ \text{the value when } t=0 \end{array} \right] \\ v &= -|g|t\hat{k} + v_0 \end{aligned} \quad \rightarrow \text{integrating again:}$$

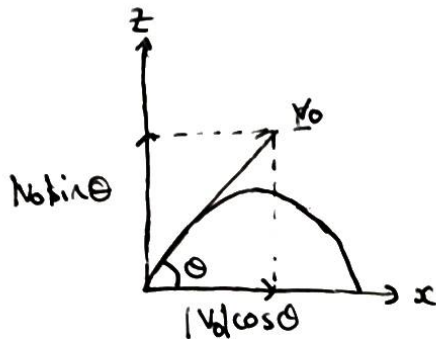
$$r = \frac{-|g|t^2\hat{k}}{2} + v_0t + \text{constant}$$

$$r = -\frac{1}{2}|g|t^2\hat{k} + v_0t + r_0 \quad \left[ \begin{array}{l} \text{this constant is } r_0 \text{ since it is} \\ \text{the value when } t=0 \end{array} \right]$$

$$\therefore \boxed{r(t) = r_0 + v_0t - \frac{1}{2}|g|t^2\hat{k}}$$

This is an equation for the position  $r$  of the particle as a function of time.

The equation of motion can be obtained by looking at the motion in the  $x$  and  $z$  directions:



$$\underline{v}_0 = v_0 \sin \theta + v_0 \cos \theta$$

We will take the initial position  $\underline{r}_0$  to be at the origin. We can now decompose the equation for  $\underline{r}(t)$  in the  $x$  and  $z$  directions.

$$\underline{r}(t) = \underline{r}_0 + \underline{v}_0 t - \frac{1}{2} g t^2 \hat{k}$$

$$\therefore r_x = |v_0| \cos \theta t \quad \Rightarrow \quad x = |v_0| \cos \theta t \quad (1)$$

$$r_z = |v_0| \sin \theta t - \frac{1}{2} g t^2 \hat{k} \quad \Rightarrow \quad z = |v_0| \sin \theta t - \frac{1}{2} g t^2 \quad (2)$$

For simplicity, we can write these as (1) and (2)

$$(2) \div (1): \quad \frac{z}{x} = \frac{|v_0| \sin \theta t}{|v_0| \cos \theta t} - \frac{\frac{1}{2} g t^2}{|v_0| \cos \theta t}$$

$$\frac{z}{x} = \tan \theta - \frac{1}{2} g t \frac{1}{|v_0| \cos \theta}$$

$$z = x \tan \theta - \frac{1}{2} g t \cdot \frac{|v_0| \cos \theta t}{|v_0| \cos \theta} \times \frac{|v_0| \cos \theta t}{|v_0| \cos \theta t}$$

$$z = x \tan \theta - \frac{1}{2} g \frac{x^2}{|v_0|^2 \cos^2 \theta}$$

This is the equation of a parabola in the form:  
 $z = ax^2 + bx + c$

## Time of flight

We can obtain the time of flight from the equation of the  $z$  component.

$$r_z = |V_0| \sin \theta t - \frac{1}{2} |g| t^2 \hat{k}$$

The projectile starts flight and finishes flight at  $r_z = 0$ :

$$|V_0| \sin \theta t - \frac{1}{2} |g| t^2 \hat{k} = 0$$

One solution is  $t=0$  (this is the start of flight)

Another solution is the time of flight:

$$\text{Time of flight} = \frac{2|V_0| \sin \theta}{|g|}$$

## Range of Projectile

This is the position of  $x$  at the end of the flight.

Substituting time of flight into the  $x$  equation:

$$r_x = |V_0| \cos \theta t \quad \text{where} \quad t = \frac{2|V_0| \sin \theta}{|g|}$$

$$\text{range of flight} = \frac{2|V_0|^2 \sin \theta \cos \theta}{|g|}$$

$$\therefore \text{range of flight} = \frac{|V_0|^2 \sin 2\theta}{|g|}$$

## Maximum Height

The maximum height is reached at half the time of flight.

Therefore, it is the  $z$  component when  $t = \frac{|V_0| \sin \theta}{|g|}$

Substituting this into the  $z$  equation gives:

$$\text{maximum height} = \frac{1}{2} \frac{|V_0|^2 \sin^2 \theta}{|g|}$$