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PERSONAL PROGRAMMING PROJECT

Implementation of Iso-geometric Analysis (IGA) for Piezoelectric Material

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Abstract

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1 Introduction

Among all the numerical methods Finite Element Methods (FEM) are more popularly used to find approximate solutions of partial differential equations. FEM approximates the Computer Aided Drawing (CAD) geometry by discretizing it into smaller geometries called elements. Such geometrical approximations may create numerical errors and seriously effect the accuracy of the solution. Isogeometric analysis (IGA) is a technique to generate geometry using CAD concept of Non Uniform Rational B-Splines and analyse using its basis functions. !The pioneers of this technique are Tom Hughes and his group at The University of Texas at Austin!.

1.1 Advantages of IGA over FEA

- The exact representation of the geometry for analysis rules out the possibility of geometrical approximations.
- A huge amount of time involved in finite element modelling can be avoided.

2 Non Uniform Rational B-Splines

NURBS are very often used in computer-aided design(CAD), manufacturing (CAM) and engineering (CAE) due to its flexibility to represent complex geometries. NURBS curves and surfaces are considered as the generalization of B-Spline and Bezier curves and surfaces. A NURBS curve is defined by its order, control points and knot vectors.

2.1 Order

!The order of a NURBS curve defines the number of nearby control points that influence any given point on the curve. The curve is represented mathematically by a polynomial of degree one less than the order of the curve.(WIKI)! (Difference b/w degree and order)

2.2 Control points

!The control points determine the shape of the curve.[8] Typically, each point of the curve is computed by taking a weighted sum of a number of control points. The weight of each point varies according to the governing parameter. For a curve of degree d , the weight of any control point is only nonzero in $d+1$ intervals of the parameter space. Within those intervals, the weight changes according to a polynomial function (basis functions) of degree d . At the boundaries of the intervals, the basis functions go smoothly to zero, the smoothness being determined by the degree of the polynomial.(WIKI)!

2.3 Knot vector

The knot vector is a sequence of parameter values that determines where and how the control points affect the NURBS curve. The number of knots is always equal to the number of control points plus curve degree plus one (i.e. number of control points plus curve order). The knot vector divides the parametric space in the intervals mentioned before, usually referred to as knot spans. Each time the parameter value enters a new

knot span, a new control point becomes active, while an old control point is discarded. It follows that the values in the knot vector should be in nondecreasing order, so (0, 0, 1, 2, 3, 3) is valid while (0, 0, 2, 1, 3, 3) is not.(WIKI)!

2.4 NURBS Curve and its properties

The p^{th} degree NURBS curve is given by

$$C(u) = \frac{\sum_{i=0}^n N_{i,p}(u)w_i P_i}{\sum_{i=0}^n N_{i,p}(u)w_i} \quad (1)$$

refer NURBS book and write more info

2.5 NURBS Surface and Solids

NURBS Surfaces and solids are generated by the tensor product between NURBS curve basis functions. A NURBS surface with degree p in u direction and degree q in v direction is defined as

$$S(u) = \frac{\sum_{i=0}^n \sum_{j=0}^m N_{i,p}(u)N_{j,q}(v)w_{i,j}P_{i,j}}{\sum_{i=0}^n \sum_{j=0}^m N_{i,p}(u)N_{j,q}(v)w_{i,j}} \quad (2)$$

refer NURBS book and write more info

2.6 Derivatives of NURBS bivariate Basis Functions

The first partial derivatives of NURBS bivariate basis function are given by

$$\frac{\partial R_{i,j}^{p,q}}{\partial \xi} = \frac{N'_{i,p}N_{j,q}W - N_{i,p}N_{j,q}W'_{\xi}}{W^2}w_{i,j} \quad (3)$$

$$\frac{\partial R_{i,j}^{p,q}}{\partial \eta} = \frac{N_{i,p}N'_{j,q}W - N_{i,p}N_{j,q}W'_{\eta}}{W^2}w_{i,j} \quad (4)$$

where

$$W = \sum_{i=0}^n \sum_{j=0}^m N_{i,p}N_{j,q}w_{i,j} \quad (5)$$

$$W'_{\xi} = \sum_{i=0}^n \sum_{j=0}^m N'_{i,p}N_{j,q}w_{i,j} \quad (6)$$

$$W'_{\eta} = \sum_{i=0}^n \sum_{j=0}^m N_{i,p}N'_{j,q}w_{i,j} \quad (7)$$

3 Iso-geometric analysis

4 Implementation Procedure for IGA

This section describes step-by-step implementation of the Isogeometric analysis. A modified FEM code structure can be used to implement IGA. Similar to FEM, IGA can be divided into pre-processing, processing and post-processing stages.

4.1 Pre-processing Stage of the Analysis

This subsection mainly deals with NURBS based geometry creation, types of assembly arrays needed to assemble discretized geometries and how to deal with homogeneous and non-homogeneous boundary conditions on boundary defining control points due to their higher (C^{p-1}) continuity unlike C^0 continuity of FEM nodes.

4.1.1 Geometry Creation

As mentioned before in the section 2, the construction of NURBS discretized geometry requires parametric details such as control points, knot vectors and the order of the NURBS curve. Commercial softwares like "Rhino" can be used to extract parametric details of the complex NURBS geometry.

4.1.2 Assembly arrays

Assembly arrays are required to assemble local discretized geometries to global geometry. For IGA as knot vector and control points defines the geometry two assembly arrays (1) Control point assembly array and (2) Knot vector connectivity array are required.

1. Control point assembly array:

The degree of the NURBS curve determines the number of control points (n_{cp}^e) present in an IGA element. Considering a three-dimensional element (Ω^e)

$$n_{cp}^e = (p + 1)(q + 1)(r + 1)$$

The details of the control points are stored row wise in assembly array. ***This can be illustrated with an example below

2. Knot vector connectivity array:

The knot vector connectivity matrix is a row wise matrix with each row corresponds to respective element's knot spans. The columns corresponds to span ranges along ξ , η and ζ directions. ***This can be illustrated with an example below

4.1.3 Boundary Conditions

***Have to write

4.2 Processing Stage of the Analysis

In the processing stage of analysis it is required to compute global elemental stiffness matrix and global force vector and solve these for the solution field. To formulate these matrices it requires NURBS basis functions and their derivatives evaluation. Numerical integration scheme like Gauss-Legendre rule is employed to solve the volume and area integrals involved in forming the stiffness matrix and internal force vector. Numerical integration involves mapping elements from physical space to master space (which is also called unit domain). As NURBS basis functions are defined in parametric space as given in Eq. () it requires an additional mapping from physical space to parametric space ($\Omega_e \rightarrow \widetilde{\Omega}_e$). Later parametric space can be mapped on to master space ($(\widetilde{\Omega}_e \rightarrow \overline{\Omega}_e)$). This procedure is illustrated in Fig.().

1. Mapping from master space to parametric space:

Consider a discretized IGA surface which is defined in parametric space $\widetilde{\Omega}_e = [\xi_i, \xi_{i+1}] \otimes [\eta_i, \eta_{i+1}]$, Refer Fig.(). The NURBS basis functions and their derivatives are evaluated at ξ, η of the element $\widetilde{\Omega}_e$. These ξ, η co-ordinate values are calculated by a linear mapping as shown below

$$\xi = \frac{1}{2}[(\xi_{i+1} - \xi_i)\bar{\xi} + (\xi_{i+1} + \xi_i)] \quad (8)$$

$$\eta = \frac{1}{2}[(\eta_{i+1} - \eta_i)\bar{\eta} + (\eta_{i+1} + \eta_i)] \quad (9)$$

where $\bar{\xi}, \bar{\eta}$ are the integration points defined in master space

\mathbf{J}_2 matrix is defined as

$$\mathbf{J}_2 = \frac{\partial \xi}{\partial \bar{\xi}} \frac{\partial \eta}{\partial \bar{\eta}} \quad (10)$$

Determinant of \mathbf{J}_2 matrix is required in numerical integration scheme for linear mapping

2. Mapping from physical space to parametric space:

The Jacobian matrix \mathbf{J}_1 used to map from physical space to parametric space ($\Omega_e \rightarrow \widetilde{\Omega}_e$) is computed as:

$$\mathbf{J}_1 = \begin{bmatrix} \frac{\partial x}{\partial \xi} & \frac{\partial x}{\partial \eta} \\ \frac{\partial y}{\partial \xi} & \frac{\partial y}{\partial \eta} \end{bmatrix} \quad (11)$$

The components of the \mathbf{J}_1 matrix are calculated using Eq. (19).

$$\frac{\partial x}{\partial \xi} = \sum_{k=1}^{n_{cp}^e} \frac{\partial \mathbf{R}_k}{\partial \xi} x_i \quad \frac{\partial x}{\partial \eta} = \sum_{k=1}^{n_{cp}^e} \frac{\partial \mathbf{R}_k}{\partial \eta} x_i \quad (12)$$

$$\frac{\partial y}{\partial \xi} = \sum_{k=1}^{n_{cp}^e} \frac{\partial \mathbf{R}_k}{\partial \xi} y_i \quad \frac{\partial y}{\partial \eta} = \sum_{k=1}^{n_{cp}^e} \frac{\partial \mathbf{R}_k}{\partial \eta} y_i \quad (13)$$

The two different mappings described above can be illustrated with an example considering $\mathbf{F}(x, y)$ integrated over the physical space Ω

$$\begin{aligned}
\int_{\Omega} \mathbf{F}(x, y) d\Omega &= \sum_{e=1}^{nel} \int_{\Omega_e} \mathbf{F}(x, y) d\Omega \\
&= \sum_{e=1}^{nel} \int_{\tilde{\Omega}_e} \mathbf{F}(\xi, \eta) |\mathbf{J}_1| d\xi d\eta \\
&= \sum_{e=1}^{nel} \int_{\bar{\Omega}_e} \mathbf{F}(\bar{\xi}, \bar{\eta}) |\mathbf{J}_1| |\mathbf{J}_2| d\bar{\xi} d\bar{\eta} \\
&= \sum_{e=1}^{nel} \int_{-1}^1 \int_{-1}^1 \mathbf{F}(\bar{\xi}, \bar{\eta}) |\mathbf{J}_1| |\mathbf{J}_2| d\bar{\xi} d\bar{\eta} \\
&= \sum_{e=1}^{nel} \left[\sum_{i=1}^{n_{gp}^e} \mathbf{F}(\bar{\xi}_i, \bar{\eta}_i) gw_i |\mathbf{J}_1| |\mathbf{J}_2| d\bar{\xi} d\bar{\eta} \right]
\end{aligned}$$

where nel is the total number of elements and n_{gp}^e , gw_i denotes the number of Gauss points and their respective Gauss weights.

The formulated global stiffness matrix and force vector are solved using numerical scheme like Newton-Raphson method for the solution field.

4.3 Post-processing Stage of the Analysis

This section deals with visualization of the deformed geometry and how a displacement and stress fields are plotted.

1. Visualization of the deformed NURBS geometry:

Visualization of the deformed geometry will be done in the same manner as the visualization of the initial geometry. After determining the displacement field at control points they are added to the initial control points co-ordinates $[\mathbf{P}]$.

$$[\mathbf{P}_{new}] = [\mathbf{P}] + [\mathbf{U}]$$

$[\mathbf{P}^{new}]$ control points array is used to plot the deformed geometry.

2. Plotting of displacement and stress fields:

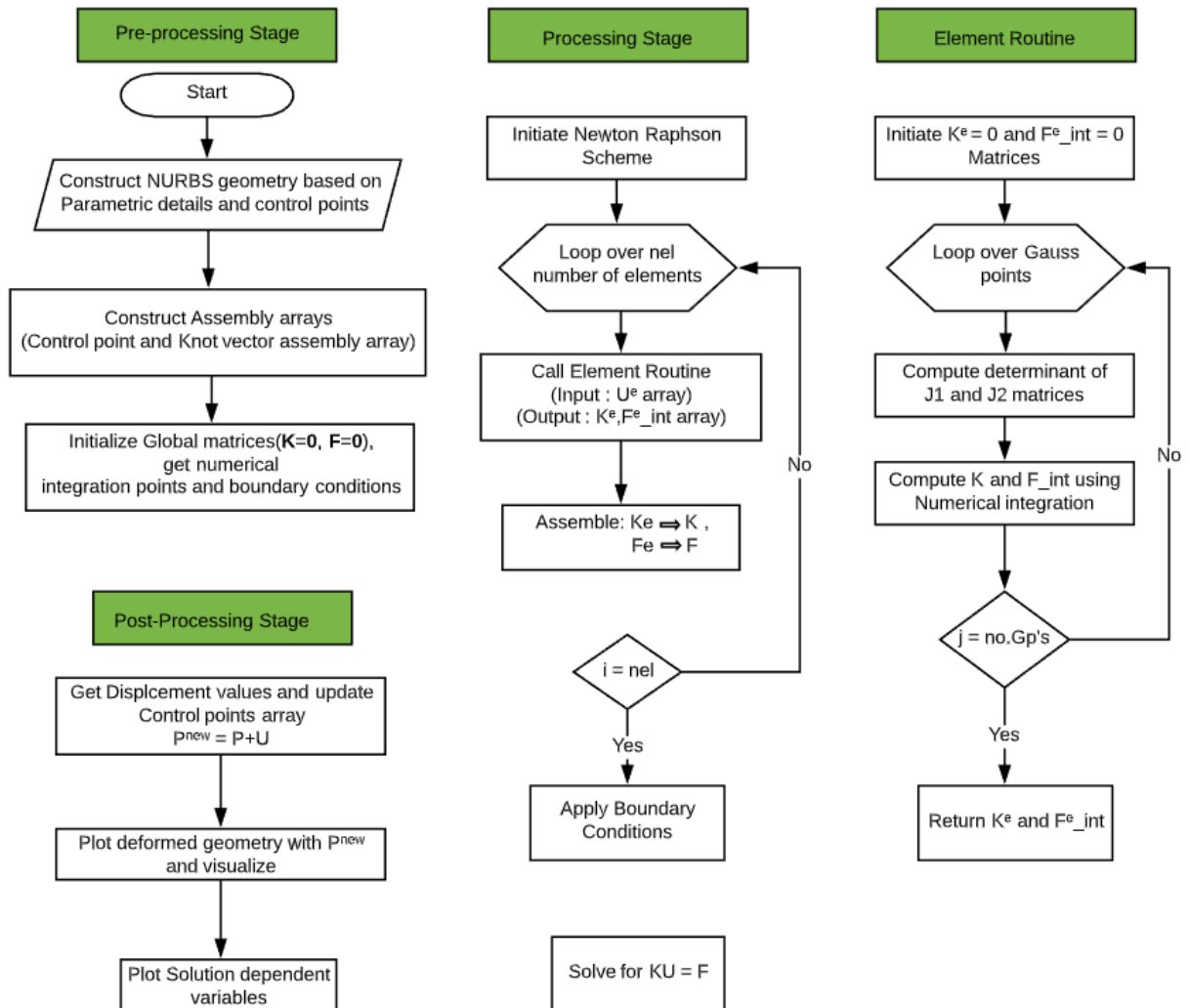
*** Have to write

5 Mechanical Case

5.1 Governing Equations

The governing equation for mechanical deformation is based on conservation of linear momentum which can be written as

$$\sigma_{ij,i} + b_j = 0 \tag{14}$$



where σ_{ij} and b_i is the Cauchy stress tensor and body force. Due to the static nature of the analysis the inertial term is not included in the eq().

Stress and strain are related by following constitutive equation

$$\sigma_{ij} = c_{ijkl}\epsilon_{kl} \quad (15)$$

The infinitesimal strain theory is adopted for the analysis in which displacements of the material particles are considered to be very small compared to the dimensions of the body under loading. Strain in a small strain setting can be written as

$$\epsilon_{ij} = \frac{1}{2}[u_{i,j} + u_{j,i}] \quad (16)$$

where u_i are the displacements in the body

5.2 Weak Formulation

Consider a domain Ω with Γ_u as prescribed displacements and Γ_t as traction boundary conditions. The domain boundary can be represented as $\Gamma = \Gamma_u \cup \Gamma_t$ and $\Gamma_u \cap \Gamma_t = \Phi$. By using the principle of virtual work the eq() can be written as

$$\delta W = \int_{\Omega} (\sigma_{ij,i} + b_j) \delta u_j dV = 0 \quad (17)$$

with, $u = u_o$ on Γ_u (essential boundary condition) and $\sigma_{ij}n_j = \bar{t}_j$ on Γ_t (natural boundary condition)

where δu_j is the virtual displacement field, n_j is unit normal to the surface

Applying integration by parts to the stress term under integral and by making use of conservation of angular momentum ($\sigma_{ij} = \sigma_{ji}$) and Gauss divergence theorem (converting volume integral to surface integral) we approach at the following equation

$$\delta W = \int_{\Omega} \sigma_{ij}\epsilon_{ij} d\Omega - \left[\int_{\Gamma} \bar{t}_j \delta u_j d\Gamma + \int_{\Omega} b_j \delta u_j d\Omega \right] \quad (18)$$

as an additional requirement δu_j must be zero at essential boundary conditions (Γ_u) for a unique solution. ***Additional data regarding detailed explanation of steps to derive weak form can be included.

5.3 IGA Formulation

The advantage of IGA over FEM formulation lies in its basis functions incorporation and its ability to capture the exact geometry. While the FEM uses lagrangian basis functions, IGA uses NURBS basis functions which are used to generate the geometry itself. As discussed in the previous sections a multidimensional NURBS basis function is represented by $R_{i,j,k}^{p,q,r}(\xi, \eta, \zeta) \rightarrow R_i$. The isogeometric element is represented by basis function R_i and control points P_i as

$$\mathbf{x}^e = \sum_{i=1}^{n_{cp}^e} R_i P_i \quad (19)$$

By Galerkin approach, the displacements and virtual displacements are given by

$$\mathbf{u}^e = \sum_{i=1}^{n_{cp}^e} R_i \mathbf{u}_i \quad \delta \mathbf{u}^e = \sum_{i=1}^{n_{cp}^e} R_i \delta \mathbf{u}_i \quad (20)$$

where \mathbf{u}_i and $\delta \mathbf{u}_i$ are values at i th control point. The strain displacement matrix \mathbf{B} is given by

$$\mathbf{B} = \begin{bmatrix} R_{1,x} & 0 & 0 & R_{2,x} & 0 & 0 & \dots & R_{n_{cp}^e,x} & 0 & 0 \\ 0 & R_{1,y} & 0 & 0 & R_{2,y} & 0 & \dots & 0 & R_{n_{cp}^e,y} & 0 \\ 0 & 0 & R_{1,z} & 0 & 0 & R_{2,z} & \dots & 0 & 0 & R_{n_{cp}^e,z} \\ R_{1,y} & R_{1,x} & 0 & R_{2,y} & R_{2,x} & 0 & \dots & R_{n_{cp}^e,y} & R_{n_{cp}^e,x} & 0 \\ 0 & R_{1,z} & R_{1,y} & 0 & R_{2,z} & R_{2,y} & \dots & 0 & R_{n_{cp}^e,z} & R_{n_{cp}^e,y} \\ R_{1,z} & 0 & R_{1,x} & R_{2,z} & 0 & R_{2,x} & \dots & R_{n_{cp}^e,z} & 0 & R_{n_{cp}^e,x} \end{bmatrix} \quad (21)$$

By substituting Eqs, (20) and (21) in Eq.(18) the weak form in matrix terms can be written as

$$\sum_{e=1}^{nel} \left[\left(\int_{\Omega_e} \mathbf{B}^T \mathbf{C} \mathbf{B} d\Omega \right) \right] \mathbf{u} = \int_{\Gamma_t^e} \mathbf{R}^T \cdot \mathbf{t} d\Gamma + \int_{\Omega_t^e} \mathbf{R}^T \cdot \mathbf{f} d\Omega \quad (22)$$

where \mathbf{R} is defined as
for the boundary Γ^e

$$\mathbf{R} = \begin{bmatrix} R_1(\xi, \eta) & 0 & R_2(\xi, \eta) & 0 & \dots & R_{n_{cp}^e}(\xi, \eta) & 0 \\ 0 & R_1(\xi, \eta) & 0 & R_2(\xi, \eta) & \dots & 0 & R_{n_{cp}^e}(\xi, \eta) \end{bmatrix} \quad (23)$$

for the domain Ω^e

$$\mathbf{R} = \begin{bmatrix} R_1(\xi, \eta, \zeta) & 0 & R_2(\xi, \eta, \zeta) & 0 & \dots & R_{n_{cp}^e}(\xi, \eta, \zeta) & 0 \\ 0 & R_1(\xi, \eta, \zeta) & 0 & R_2(\xi, \eta, \zeta) & \dots & 0 & R_{n_{cp}^e}(\xi, \eta, \zeta) \end{bmatrix} \quad (24)$$

Eq.(22) can be rewritten in a standard matrix form as

$$\sum_{e=1}^{nel} [\mathbf{K}^e \mathbf{U}^e = \mathbf{F}^e] \quad (25)$$

where \mathbf{K}^e is isogeometric element's stiffness matrix, \mathbf{U}^e is displacement vector and \mathbf{F}^e force vector

6 Piezoelectric Case

6.1 Governing Equations for Piezoelectric Materials

The coupled electro-mechanical interactions are governed by conservation of momentum and Gauss's law as:

$$\sigma_{ij,j} + f_i = 0 \quad (26)$$

$$D_{i,i} - q = 0 \quad (27)$$

Where, f_i is body force, q is electrical charge, σ_{ij} is Cauchy stress tensor and D_i is electrical displacement vector. The constitutive equations for electromechanical coupling are defined as

$$\sigma_{ij} = C_{ijkl}\varepsilon_{kl} - e_{kij}E_k \quad (28)$$

$$D_i = e_{ikl}\varepsilon_{kl} + \kappa_{ik}E_k \quad (29)$$

Where, C_{ijkl} , e_{ikj} and ε_{ij} are elastic, piezoelectric and dielectric material constants respectively. The Cauchy strain tensor is defined as:

$$\varepsilon_{ij} = \frac{1}{2}(u_{i,j} + u_{j,i}) \quad (30)$$

and electric field vector as:

$$E_i = -\phi_{,i} \quad (31)$$

6.2 Weak Formulation

Applying the principle of virtual work to the Eq.(26) and Eq. (27) we can write

$$\int_{\Omega} (\sigma_{ij,j} + f_i) \delta u_i d\Omega = 0 \quad (32)$$

$$\int_{\Omega} (D_{i,i} - q) \delta \phi d\Omega = 0 \quad (33)$$

with,

essential boundary conditions: $u = u_o$ on Γ_u and $\Phi = \Phi_0$ on Γ_{Φ}

natural boundary condition: $\sigma_{ij}n_j = \bar{t}_i$ on Γ_t and $D_i n_i = q_0$ on Γ_q

where δu_i and $\delta \phi$ are virtual or arbitrary displacement and potential fields.

Integrating Eq. (32) and Eq. (33) by parts and later applying Gauss divergence theorem and boundary conditions we approach at weak form

$$\int_{\Omega} \sigma_{ij} \delta \varepsilon_{ij} d\Omega - \left[\int_{\Gamma} \bar{t}_i \delta u_i d\Gamma + \int_{\Omega} f_i \delta u_i d\Omega \right] = 0 \quad (34)$$

$$\int_{\Omega} D_i \delta E_i d\Omega - \left[\int_{\Gamma} q \delta \phi d\Gamma + \int_{\Omega} q \delta \phi d\Omega \right] = 0 \quad (35)$$

6.3 IGA Formulation

By Galerkin approach, displacements, potentials and their virtual values are given by below equations

$$\mathbf{u}^e = \sum_{i=1}^{n_{cp}^e} R_i \mathbf{u}_i \quad \delta \mathbf{u}^e = \sum_{i=1}^{n_{cp}^e} R_i \delta \mathbf{u}_i \quad (36)$$

$$\Phi^e = \sum_{i=1}^{n_{cp}^e} R_i \Phi_i \quad \delta \Phi^e = \sum_{i=1}^{n_{cp}^e} R_i \delta \Phi_i \quad (37)$$

where \mathbf{u}_i , $\delta \mathbf{u}_i$, $\delta \Phi_i$ and Φ_i are values at i th control point.

By substituting Eqs. (36) and (37) in Eq.(34) and Eq.(35) the weak form in matrix notation can be written as

$$\begin{bmatrix} K_{MM} & K_{ME} \\ K_{EM} & K_{EE} \end{bmatrix} \begin{bmatrix} \mathbf{u} \\ \phi \end{bmatrix} = \begin{bmatrix} F_M \\ F_E \end{bmatrix} \quad (38)$$

Where,

$$K_{MM} = \int_{\Omega} \mathbf{B}_u^T \mathbf{C} \mathbf{B}_u d\Omega \quad (39)$$

$$K_{ME} = \int_{\Omega} \mathbf{B}_u^T \mathbf{e} \mathbf{B}_e d\Omega \quad (40)$$

$$K_{EM} = \int_{\Omega} \mathbf{B}_e^T \mathbf{e}^T \mathbf{B}_u d\Omega \quad (41)$$

$$K_{EE} = - \int_{\Omega} \mathbf{B}_e^T \kappa \mathbf{B}_e d\Omega \quad (42)$$

$$F_M = \int_{\Omega} \mathbf{R}_u^T \mathbf{f} d\Omega + \int_{\Gamma} \mathbf{R}_u^T \mathbf{t} d\Gamma \quad (43)$$

$$F_E = \int_{\Omega} \mathbf{R}_e^T q d\Omega + \int_{\Gamma} \mathbf{R}_e^T Q d\Gamma \quad (44)$$

where \mathbf{R} is defined as
for the boundary Γ

$$\mathbf{R}_u = \begin{bmatrix} R_1(\xi, \eta) & 0 & R_2(\xi, \eta) & 0 & \dots & R_{n_{cp}^e}(\xi, \eta) & 0 \\ 0 & R_1(\xi, \eta) & 0 & R_2(\xi, \eta) & \dots & 0 & R_{n_{cp}^e}(\xi, \eta) \end{bmatrix} \quad (45)$$

$$\mathbf{R}_e = [R_1(\xi, \eta) \quad R_2(\xi, \eta) \quad \dots \quad R_{n_{cp}^e}(\xi, \eta)] \quad (46)$$

for the domain Ω

$$\mathbf{R}_u = \begin{bmatrix} R_1(\xi, \eta, \zeta) & 0 & R_2(\xi, \eta, \zeta) & 0 & \dots & R_{n_{cp}^e}(\xi, \eta, \zeta) & 0 \\ 0 & R_1(\xi, \eta, \zeta) & 0 & R_2(\xi, \eta, \zeta) & \dots & 0 & R_{n_{cp}^e}(\xi, \eta, \zeta) \end{bmatrix} \quad (47)$$

$$\mathbf{R}_e = [R_1(\xi, \eta, \zeta) \quad R_2(\xi, \eta, \zeta) \quad \dots \quad R_{n_{cp}^e}(\xi, \eta, \zeta)] \quad (48)$$

B matix is given as

$$\mathbf{B}_u = \begin{bmatrix} R_{1,x} & 0 & 0 & R_{2,x} & 0 & 0 & \dots & R_{n_{cp}^e,x} & 0 & 0 \\ 0 & R_{1,y} & 0 & 0 & R_{2,y} & 0 & \dots & 0 & R_{n_{cp}^e,y} & 0 \\ 0 & 0 & R_{1,z} & 0 & 0 & R_{2,z} & \dots & 0 & 0 & R_{n_{cp}^e,z} \\ R_{1,y} & R_{1,x} & 0 & R_{2,y} & R_{2,x} & 0 & \dots & R_{n_{cp}^e,y} & R_{n_{cp}^e,x} & 0 \\ 0 & R_{1,z} & R_{1,y} & 0 & R_{2,z} & R_{2,y} & \dots & 0 & R_{n_{cp}^e,z} & R_{n_{cp}^e,y} \\ R_{1,z} & 0 & R_{1,x} & R_{2,z} & 0 & R_{2,x} & \dots & R_{n_{cp}^e,z} & 0 & R_{n_{cp}^e,x} \end{bmatrix} \quad (49)$$

$$\mathbf{B}_e = \begin{bmatrix} R_{1,x} & R_{2,x} & \dots & R_{n_{cp}^e,x} \\ R_{1,y} & R_{2,y} & \dots & R_{n_{cp}^e,y} \\ R_{1,z} & R_{2,z} & \dots & R_{n_{cp}^e,z} \end{bmatrix} \quad (50)$$

$$\varepsilon = \mathbf{B}_u \cdot \mathbf{u} \quad (51)$$

$$\mathbf{E} = -\mathbf{B}_e.\Phi \tag{52}$$

Eq. (38) can be solved using numerical methods like Newton Raphson scheme for displacements and potential solution field.