

LEVERAGING MULTI-MESSENGER ASTROPHYSICS FOR DARK MATTER SEARCHES

By

Daniel Nicholas Salazar-Gallegos

A DISSERTATION

Submitted to  
Michigan State University  
in partial fulfillment of the requirements  
for the degree of

Physics—Doctor of Philosophy  
Computational Mathematics in Science and Engineering—Dual Major

Today

## **ABSTRACT**

I did Dark Matter with HAWC and IceCube. I also used Graph Neural Networks

Copyright by  
DANIEL NICHOLAS SALAZAR-GALLEGOS  
Today

## **ACKNOWLEDGMENTS**

I love my friends. Thanks to everyone that helped me figure this out. Amazing thanks to the people at LANL who supported me. Eames, etc Dinner Parties Jenny and her child Kaydince Kirsten, Pat, Andrea Family. You're so far but so critical to my formation. Unconditional love. Roommate

## TABLE OF CONTENTS

LIST OF TABLES . . . . .	vi
LIST OF FIGURES . . . . .	vii
LIST OF ABBREVIATIONS . . . . .	ix
CHAPTER 1 INTRODUCTION . . . . .	1
CHAPTER 2 DARK MATTER IN THE COSMOS . . . . .	2
2.1 Introduction . . . . .	2
2.2 Dark Matter Basics . . . . .	3
2.3 Evidence for Dark Matter . . . . .	4
2.4 Searching for Dark Matter . . . . .	11
2.5 Sources for Indirect Dark Matter Searches . . . . .	16
2.6 Multi-Messenger Dark Matter . . . . .	18
CHAPTER 3 MULTIMESSENGER ASTROPHYSICS: DETECTING HIGH ENERGY NEUTRAL MESSENGERS . . . . .	21
3.1 Introduction . . . . .	21
3.2 Charged Particles in a Medium . . . . .	21
3.3 Photons ( $\gamma$ ) . . . . .	22
3.4 Neutrinos ( $\nu$ ) . . . . .	22
3.5 Opportunities to Combine for Dark Matter . . . . .	22
CHAPTER 4 HIGH ALTITUDE WATER CHERENKOV (HAWC) OBSERVATORY . . . . .	24
4.1 The Detector . . . . .	24
4.2 Events Reconstruction and Data Acquisition . . . . .	24
4.3 Remote Monitoring . . . . .	24
CHAPTER 5 ICECUBE NEUTRINO OBSERVATORY . . . . .	25
5.1 The Detector . . . . .	25
5.2 Events Reconstruction and Data Acquisition . . . . .	25
5.3 Northern Test Site . . . . .	25
CHAPTER 6 COMPUTATIONAL TECHNIQUES IN PARTICLE ASTROPHYSICS . . . . .	26
6.1 Neural Networks for Gamma/Hadron Separation . . . . .	26
6.2 Parallel Computing for Dark Matter Analyses . . . . .	26
CHAPTER 7 GLORY DUCK . . . . .	27
CHAPTER 8 NU DUCK . . . . .	28
BIBLIOGRAPHY . . . . .	29

## **LIST OF TABLES**

Proof I know how to include

## LIST OF FIGURES

<p>Figure 2.1 Rotation curve fit to NGC 6503 from [7]. Dashed line is the contribution from visible matter. Dotted curves are from gas. Dash-dot curves are from dark matter (DM). Solid line is the composite contribution from all matter and DM sources. Data are indicated with bold dots with error bars. Data agree strongly with matter + DM composite prediction . . . . .</p> <p>Figure 2.2 Light from distant galaxy is bent in unique ways depending on the distribution of mass between the galaxy and Earth. Yellow dashed lines indicate where the light would have gone if the matter were not present [8]. . . . .</p> <p>Figure 2.3 (left) Optical image of galactic cluster 1E0657-558. (right) X-ray image of the cluster with redder meaning hotter and higher baryon density. (both) Green contours are reconstruction of gravity contours from lensing. White rings are the best fit mass maxima at 68.3%, 95.5%, and 99.7% confidence. The maxima of the clusters are clearly separated from x-ray maxima. [9] . . . . .</p> <p>Figure 2.4 Plank CMB sky. Sky map features small variations in temperature in primordial light. These anisotropies can be used to make inferences about the universe's energy budget. [10] . . . . .</p> <p>Figure 2.5 Observed Cosmic Microwave Background power spectrum as a function of multipole momentfrom Plank [10]. Blue line is best fit model from <math>\Lambda</math>CDM. Red points and lines are data and error respectively. . . . .</p> <p>Figure 2.6 Predicted power spectra of CMB for different <math>\Omega_m h^2</math> values. (left) Low <math>\Omega_m h^2</math> increases the prominence of first and second peaks. (middle) <math>\Omega_m h^2</math> is most similar to the observed power spectrum. The second and third peaks are similar in height. (right) <math>\Omega_m h^2</math> is large which suppresses the first peak and raises the prominence of the third peak. . . . .</p> <p>Figure 2.7 The Standard Model (SM) of particle physics. Figure taken from <a href="http://www.quantumdiaries.org/2014/03/14/the-standard-model-a-beautiful-but-flawed-theory/">http://www.quantumdiaries.org/2014/03/14/the-standard-model-a-beautiful-but-flawed-theory/</a> . . . . .</p> <p>Figure 2.8 Simplified Feynman diagram demonstrating with different ways DM can interact with SM particles. The 'X's refer to the DM particles whereas the SM refer to fundamental particles in the SM. The large circle in the center indicates the vertex of interaction and is purposely left vague. The colored arrows refer to different directions of time as well as their respective labels. The arrows indicate the initial and final state of the DM -SM interaction in time. . . . .</p>	<p>5</p> <p>7</p> <p>8</p> <p>9</p> <p>9</p> <p>10</p> <p>12</p> <p>13</p>
---	--

Figure 2.9 A single jet event in ATLAS detector from 2017 [16]. Jet momentum observed to be 1.9 TeV. Missing transverse momentum observed to be 1.9 TeV as the initial momentum of the event was 0. Implied MET is shown as a red dashed line in event display. . . . .	14
Figure 2.10 More detailed pseudo-Feynman diagram of particle cascade from dark matter annihilation into 2 quarks. The quarks hadronize and down to stable particles like $\gamma$ or the anti-proton ( $p^-$ ). Diagram pulled from ICRC 2021 presentation on DM annihilation search [17]. . . . .	15
Figure 2.11 Dark Matter (DM) decay spectrum for $b\bar{b}$ initial state and $\gamma$ final state. Redder spectra are for larger DM masses. Bluer spectra are light DM masses. $x$ is a unitless factor defined as the ratio of the mass of DM, $m_\chi$ , and the final state particle energy $E_\gamma$ . Figure from [19]. . . . .	17
Figure 2.12 Different dark matter density profiles compared. Some models produce very large densities at small r [20]. . . . .	18
Figure 2.13 The Milky Way Galaxy in photons ( $\gamma$ ) and neutrinos ( $\nu$ ) [22]. Galactic center is at $l=0^\circ$ and is the brightest region in all panels. (top) An Optical color image of the Milky Way galaxy seen from Earth. Clouds of gas and dust obscure some of the light from stars. (2nd down) Integrated flux of $\gamma$ -rays observed by the Fermi-LAT telescope [23]. (middle) Expected neutrino emmision that corresponds with Fermi-LAT observations. (2nd up) Expected neutrino emmision profile after considering detector systematics of IceCube. (bottom) Observed neutrino emmision from region of the galactic plane. Substantial neutrino emmision is detected. . . . .	19
Figure 2.14 Dark Matter annihilation spectra for different final state particle and standard model annihilation channels [24]. Photons (red), $e^\pm$ (green), $\bar{p}$ (blue), $\nu$ (black). . . . .	20
Figure 3.1 TODO: Show a nuclear reactor with cherenkov radiation[NEEDS A SOURCE][FACT CHECK THIS] . . . . .	22
Figure 3.2 TODO: absorption spectrum of liquid and solid water[NEEDS A SOURCE][FACT CHECK THIS] . . . . .	23

## LIST OF ABBREVIATIONS

<b>MSU</b>	Michigan State University
<b>LANL</b>	Los Alamos National Laboratory
<b>DM</b>	Dark Matter
<b>SM</b>	Standard Model
<b>HAWC</b>	High Altitude Water Cherenkov Observatory

## **CHAPTER 1**

### **INTRODUCTION**

Is the text not rendering right? Ah ok it knows im basically drafting the doc still

## CHAPTER 2

### DARK MATTER IN THE COSMOS

#### 2.1 Introduction

The dark matter problem can be summarized in part by following thought experiment.

Let's say you're the teacher for an elementary school classroom. You take them on a field trip to your local science museum and among exhibits is one for mass and weight. The exhibit has a gigantic scale, and you come up with a fun problem for your class.

You ask your class, "What is the total weight of the classroom? Give your best estimation to me in 30 minutes, and then we'll check your guess on the scale. If your guess is within 10% of the right answer, we will stop for ice cream on the way back"

The students are ecstatic to hear this, and they get to work. The solution is some variation of the following strategy. The students should give each other their weight or best guess if they do not know. Then, all they must do is add each student's weight and get a grand total for the class. The measurement on the giant scale should show the true weight of the class. When comparing the measured weight to your estimation, multiply the measurement by  $1.0 \pm 0.1$  in order to get the  $\pm 10\%$  tolerances for your estimation.

Two of your students, Sandra and Mario, return to you with a solution.

They say, "We weren't sure of everyone's weight. We used 65 lbs. for the people we didn't know and added everyone who does know. There are 30 of us, and we got 2,000 lbs.! That's a ton!"

You estimated 1,900 lbs. assuming the average weight of a student in your class was 60 lbs. So, you are pleased with Sandra's and Mario's answer. You instruct your students to all gather on the giant scale and read off the weight together. To all your surprise, the scale reads *10,000 lbs!* 10,000 is significantly more than a 10% error from 2,000. In fact, it is approximately 5 times more massive than either your or your students' estimates. You think to yourself and conclude there must be something wrong with the scale. You ask an employee to check the scale and verify it is calibrated well. They confirm that the scale is in working order. You weigh a couple of students individually to test that the scale is well calibrated. Sandra weighs 59 lbs., and Mario weighs 62

lbs., typical weights for their age. You then weigh each student individually and see that their weights individually do not deviate greatly from 60 lbs. So, where does all the extra weight come from?

This thought experiment serves as an analogy to the Dark Matter problem. The important substitution to make however is to replace the students with stars and the classroom with a galaxy, say the Milky Way. Individually the mass of stars is well measured and defined with the Sun as our nearest test case. However, when we set out to measure the mass of a collection of stars as large as galaxies, our well-motivated estimation is wildly incorrect. There simply is no way to account for this discrepancy except without some unseen, or dark, contribution to mass and matter in galaxies. I set out in my thesis to narrow the possibilities of what this Dark Matter could be.

This chapter is organized like the following... **TODO: Text should look like ... Chaper x has blah blah blah.**

## 2.2 Dark Matter Basics

Presently, the most compelling of cosmology that includes Dark Matter (DM) in order to explain a variety of observations is  $\Lambda$  Cold Dark Matter, or  $\Lambda$ CDM. I present the evidence supporting  $\Lambda$ CDM in 2.3, yet discuss the conclusions of the  $\Lambda$ CDM model here. According to  $\Lambda$ CDM fit to observations on the Cosmic Microwave Background (CMB), DM is 26.8% of the universe's current energy budget Baryonic matter, stuff like atoms, gas, and stars, contributes to 4.9% of the universe's current energy budget [1, 2, 3].

DM is dark; it does not interact readily with light at any wavelength. DM also does not interact noticeably with the other standard model forces (Strong and Weak) at a rate that is readily observed [3]. DM is cold, which is to say that the average velocity of DM is below relativistic speeds [1]. 'Hot' DM would not likely manifest the dense structures we observe like galaxies, and instead would produce much more diffuse galaxies than what is observed [3, 1]. DM is old; it played a critical role in the formation of the universe and the structures within it [1, 2].

Observations of DM have so far been only gravitational. The parameter space available to what DM could be therefore is extremely broad. The broad DM parameter space is iteratively tested in

DM searches by supposing a hypothesis that has not yet been ruled out and performing observations to test them. When the observations yield a null result, the parameter space is further constrained. I present some approaches for DM searches in Section 2.4.

### 2.3 Evidence for Dark Matter

Dark Matter (DM) has been a looming problem in physics for almost 100 years. Anomalies have been observed in galactic dynamics as early as 1933 when Fritz Zwicky noticed unusually large velocity dispersions in the Coma cluster. Zwicky's measurement was the first recorded to use the Virial theorem to measure the mass fraction of visible and invisible matter in celestial bodies [4]. From Zwicky in [5], "*If this would be confirmed, we would get the surprising result that dark matter is present in much greater amount than luminous matter.*" Zwicky's and others' observation did not instigate a crisis in astrophysics because the measurements did not entirely conflict with their understanding of galaxies [4]. In 1978, Rubin, Ford, and Norbert measured rotation curves for ten spiral galaxies [6]. Rubin et. al.'s 1978 publication presented a major challenge to the conventional understanding of galaxies that could no longer be accredited to measurement uncertainties. Evidence has been mounting ever since for this exotic form of matter. The following subsections sample some of the compelling evidence supporting DM.

#### 2.3.1 First Clues: Stellar Velocities

Zwicky, and later Rubin, measured the stellar velocities of various galaxies to estimate their virial mass. The Virial Theorem upon which these observations are interpreted is written as

$$2T + V = 0. \quad (2.1)$$

Where  $T$  is the kinetic energy and  $V$  is the potential energy in a self-gravitating system. The potential was defined as the classical Newton's law of gravity from stars and gas contained in the observed galaxies

$$V = -\frac{1}{2} \sum_i \sum_{j \neq i} \frac{m_i m_j}{r_{ij}}. \quad (2.2)$$

Zwicky et. al. measured just the velocities of stars apparent in optical wavelengths [5]. Rubin et. al. added by measuring the velocity of the hydrogen gas via the 21 cm emission line of Hydrogen

[6]. The velocities of the stars and gas are used to infer the total mass of galaxies and galaxy clusters via Eq. (2.1). An inferred mass is also made from the luminosity of the selected sources. The two inferences are compared to each other as a luminosity to mass ratio and typically yields [1]

$$\frac{M}{L} \sim 400 \frac{M_{\odot}}{L_{\odot}} \quad (2.3)$$

$M_{\odot}$  and  $L_{\odot}$  referring to stellar mass and stellar luminosity, respectively. These ratios clearly indicate a discrepancy in apparent light and mass from stars and gas and their velocities.

Rubin et.al. [6] demonstrated that the discrepancy was unlikely to be an under-estimation of the mass of the stars and gas. The inferred "dark" mass was up to 5 times more than the luminous mass. This dark mass also needed to extend well beyond the extent of the luminous matter.

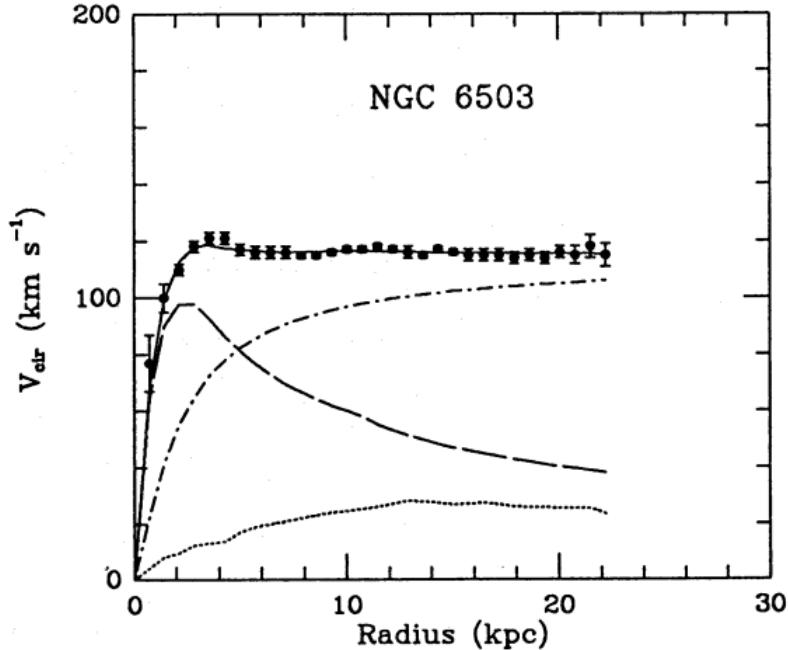


Figure 2.1 Rotation curve fit to NGC 6503 from [7]. Dashed line is the contribution from visible matter. Dotted curves are from gas. Dash-dot curves are from dark matter (DM). Solid line is the composite contribution from all matter and DM sources. Data are indicated with bold dots with error bars. Data agree strongly with matter + DM composite prediction

Fig. 2.1: features one of many observations made on the stellar velocities within galaxies. The measured rotation curves mostly feature a flattening of velocities at higher radius which is not expected if the gravity was only coming from gas and luminous matter. The extension of

the flat velocity region also indicates that the DM is distributed far from the center of the galaxy. Modern velocity measurements include significantly larger objects, galactic clusters, and smaller objects, dwarf galaxies. Yet, measurements along this regime are leveraging the Virial theorem with Newtonian potential energies. We know Newtonian gravity is not a comprehensive description of gravity. New observational techniques have been developed since 1978, and those are discussed in the following sections.

### 2.3.2 Evidence for Dark Matter: Gravitational Lensing

Modern evidence for dark matter comes from new avenues beyond stellar velocities. Gravitational lensing from DM is a new channel from general relativity. General relativity predicts aberrations in light caused by massive objects. In recent decades we have been able to measure the lensing effects from compact objects and DM haloes. Fig. 2.2 shows how different compact bodies change the final image of a faraway galaxy resulting from gravitational lensing. Gravitational lensing developed our understanding of dark matter in two important ways.

Gravitational lensing provides additional compelling evidence for DM. The observation of two merging galactic clusters in 2006, shown in Fig. 2.3, provided a compelling argument for DM outside the Standard Model. These clusters merged recently in astrophysical time scales. Their recent merger separated the stars and galaxies are separated from the intergalactic gas. For these clusters, the hot, intergalactic gas is responsible for most of the baryonic mass in the systems [4]. The hot gas is observed from the cluster's x-ray emmision. Two observations of the clusters were made independently of each other. The first was the lensing of light around the galaxies due to their gravitational influences. When celestial bodies are large enough, the gravity they exert bends space and time itself. These bending effects light and will deflect light an analogous way to how lenses will bend light. With a sufficient understanding of light sources behind a celestial body, we can reconstruct the contours of the gravitational lenses. The gradient of the contours then indicates how dense the matter is and where it is.

The x-ray emission can then be observed from the clusters. Since these galaxies are mostly gas and are merging, then the gas should be getting hotter. If they are merging, the x-ray emissions

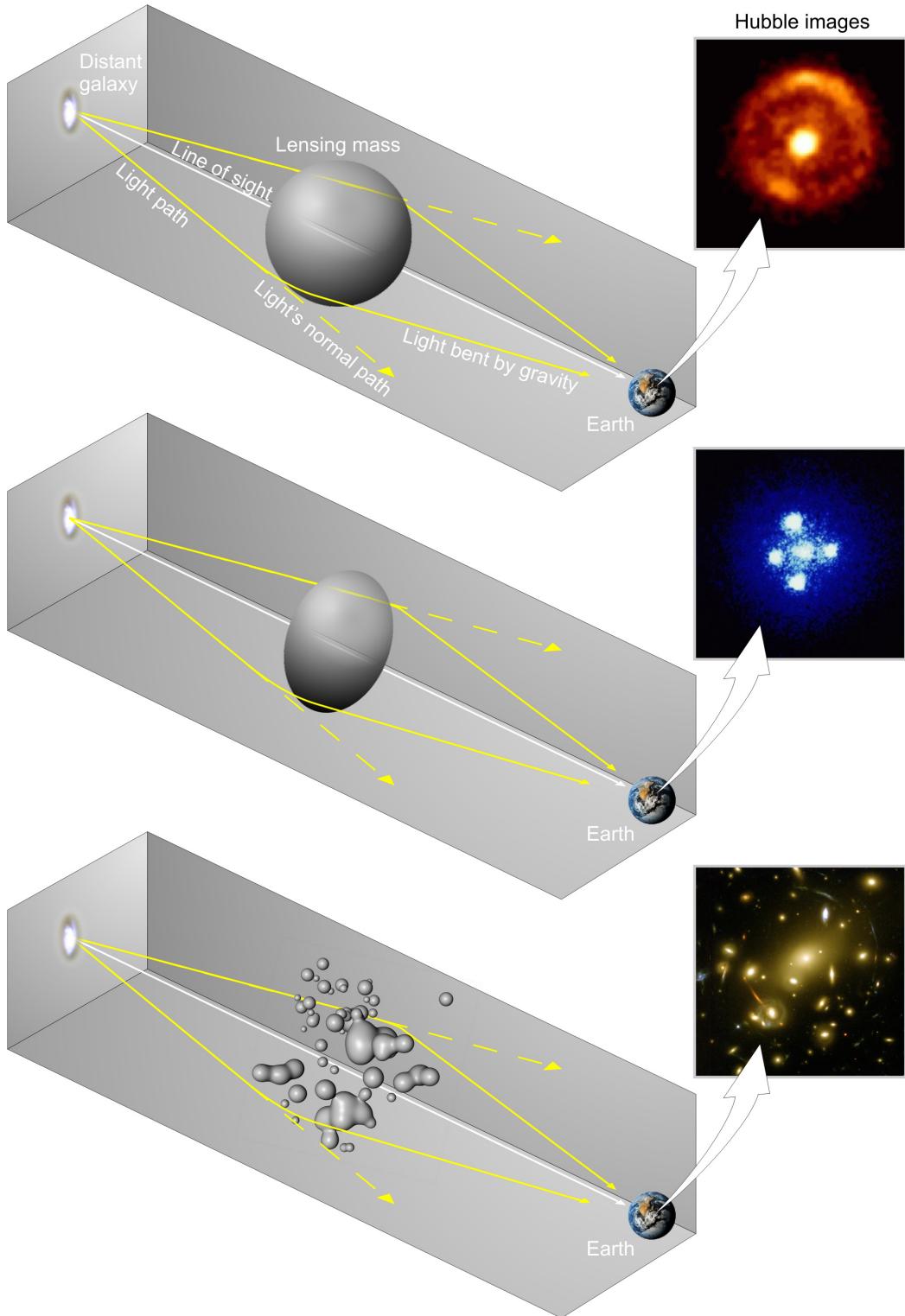


Figure 2.2 Light from distant galaxy is bent in unique ways depending on the distribution of mass between the galaxy and Earth. Yellow dashed lines indicate where the light would have gone if the matter were not present [8].

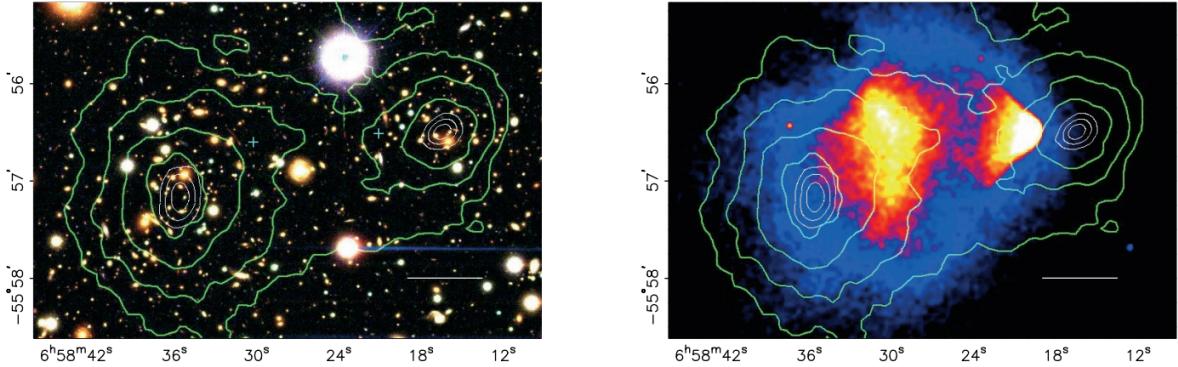


Figure 2.3 (left) Optical image of galactic cluster 1E0657-558. (right) X-ray image of the cluster with redder meaning hotter and higher baryon density. (both) Green contours are reconstruction of gravity contours from lensing. White rings are the best fit mass maxima at 68.3%, 95.5%, and 99.7% confidence. The maxima of the clusters are clearly separated from x-ray maxima. [9]

should be the strongest where the gas is mostly moving through each other. Hence, X-ray emission maps out where the gas is in the merging galaxy cluster.

The lensing and x-ray observations were done on the Bullet cluster featured on Fig. 2.3. The x-ray emmisions does not align with the gravitational countours from lensing. The incongruence in mass density and baryon density suggests that there is a lot of matter somewhere that does not interact with light. Moreover, this dark matter is can not be baryonic [9]. The Bullet Cluster measurement did not really tell us what DM is exactly, but it did give the clue that DM also does not interact with itself very strongly. If DM did interact strongly with itself, then it would have been more aligned with the x-ray emmision [9]. There have been follow-up studies of galaxy clusters with similar results. The Bullet Cluster and others like it provide a strong case against something possibly amiss in our gravitational theories.

### 2.3.3 Evidence for Dark Matter: Cosmic Microwave Background

The Cosmic Microwave Background (CMB) is the primordial light from the early universe when Hydrogen atoms formed from the free electron and proton soup in the early universe. The CMB is the earliest light we can observe; released when the universe was about 380,000 years old. Then we look at how the simulated universes look like compared to what we see. Fig. 2.4 is the most recent CMB image from the Plank observatory [10]. Redder regions indicate a slightly hotter

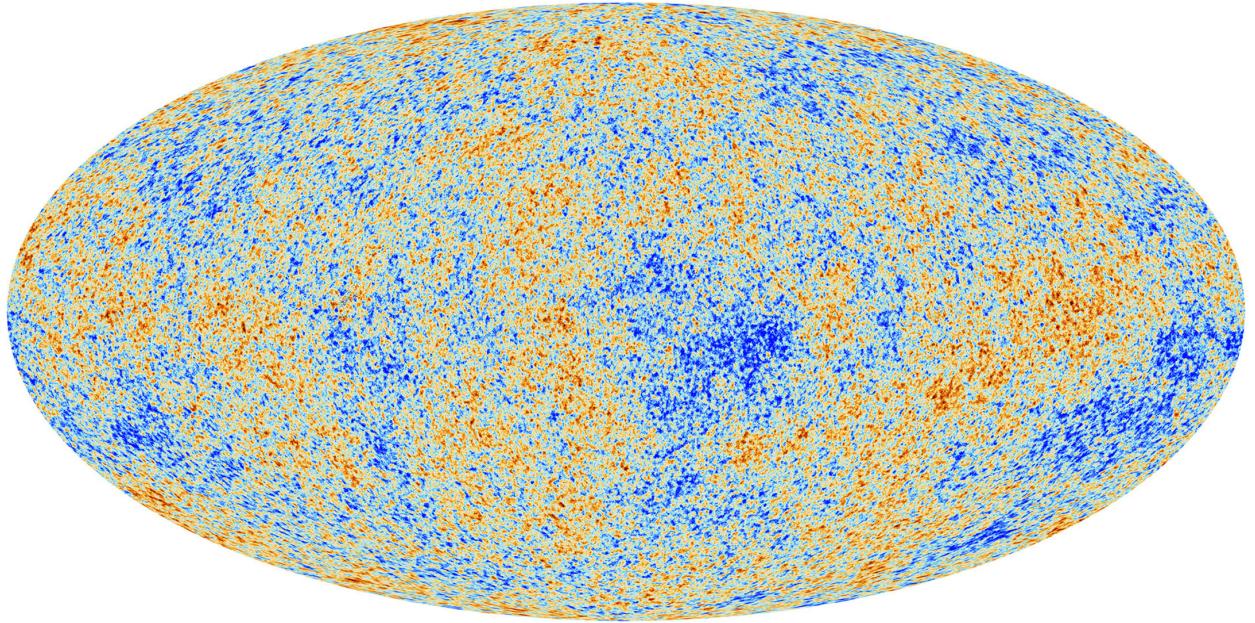


Figure 2.4 Plank CMB sky. Sky map features small variations in temperature in primordial light. These anisotropies can be used to make inferences about the universe's energy budget. [10]

region of the early universe and blue indicates colder.

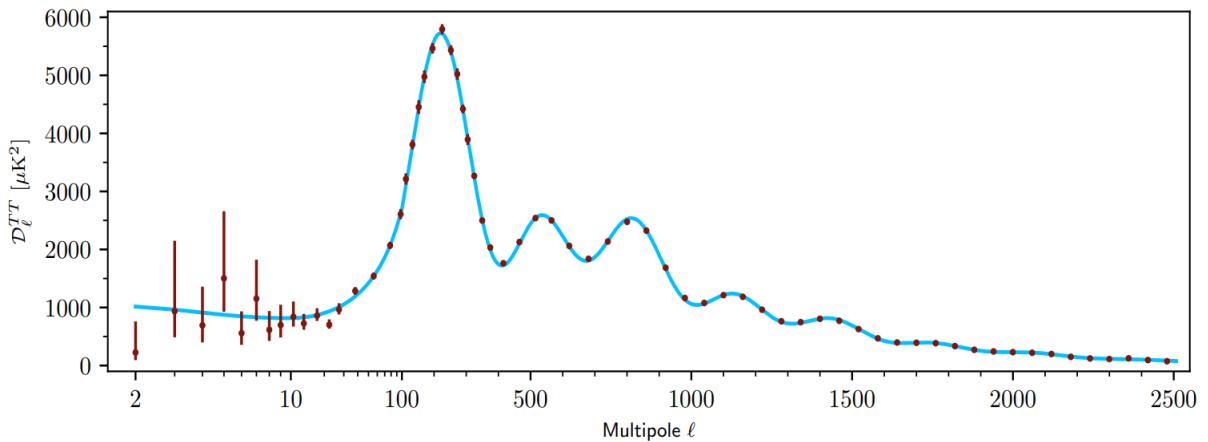


Figure 2.5 Observed Cosmic Microwave Background power spectrum as a function of multipole momentfrom Plank [10]. Blue line is best fit model from  $\Lambda$ CDM. Red points and lines are data and error respectively.

The Cosmic Microwave Background shows that the universe had DM in it from an incredibly early stage. To measure the DM, Dark Energy, and matter fractions of the universe from the CMB, the image is deconstructed into a power spectrum, which shows the amplitude of the fluctuations as a function of spherical multipole moments.  $\Lambda$ CDM provides the best fit to the power spectra of the

CDM as shown in Fig. 2.5. The CMB power spectrum is very sensitive to the fraction of each energy contribution in the early universe. Low  $l$  modes are dominated by variations in gravitational potential. Intermediate  $l$  emerge from oscillations in photon-baryon fluid from competing baryon pressures and gravity. High  $l$  is a damped region from the diffusion of photons during electron-proton recombination. [1]

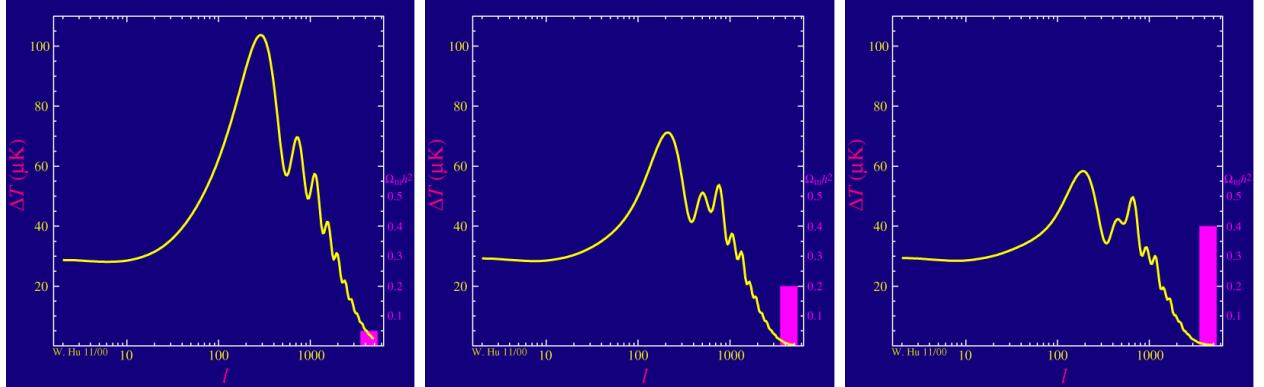


Figure 2.6 Predicted power spectra of CMB for different  $\Omega_m h^2$  values. (left) Low  $\Omega_m h^2$  increases the prominence of first and second peaks. (middle)  $\Omega_m h^2$  is most similar to the observed power spectrum. The second and third peaks are similar in height. (right)  $\Omega_m h^2$  is large which suppresses the first peak and raises the prominence of the third peak.

The harmonics would look very different for a universe with less DM. Fig. 2.6 shows the differences expected in the power spectrum for different matter fractions of the universe's energy budget. The observations fit well with the  $\Lambda$ CDM model and the derived fractions are as follows. The matter fraction:  $\Omega_m = 0.3153$ ; and the baryon fraction:  $\Omega_b = 0.04936$  [10]. These findings do rely however on a few assumptions and the precision of the Hubble constant,  $H_0$ .  $H_0$  especially has seen a growing tension in recent decades that continues to deepened with observatories like the James Webb Telescope [11, 12]

Overall these observations form a compelling body of research in favor of dark matter. However, these observations really only confirm that DM is there. It takes another leap of theory and experimentation to make observations of DM that are non-gravitational in nature. One hypothesis is the Weakly Interacting Massive Particle DM. This DM candidate theory is discussed further in the next section and is the hypothesis to this thesis.

## 2.4 Searching for Dark Matter

**TODO:** Start section that needs adoption

First, micro-lensing observations, or the lack of them, of our Milky Way halo resulted in a conspicuous absence of massive astrophysical compact halo objects (MACHOs). The hypothesis was that 'dark matter' could be accounted for by sufficiently dim compact objects. Such objects include things like planets, brown dwarves, black holes, or neutron stars. **TODO:** End section that needs adoption

Whenever these objects passed in front of a large luminous source, such as the Large Magellenic Clouds, a variation in light should be observed [4]. The MACHO and EROS collaborations performed this observation and did not find a substantial contribution to the DM Milky Way halo from MACHOs. They measured that MACHOs of mass range 0.15 to 0.9  $M_{\odot}$  contributes to an upper limit of 8% of the DM halo mass [13]. There remains many options available to what Dark Matter could be. For a particle dark matter hypothesis, we assume that DM interacts in some way, even if very weakly, with the Standard Model (SM), see Section 2.4. The current status of the SM does not have a viable DM candidate. When looking at the standard model, we can immediately exclude any charged particle. This is because charged particles interact with light. If DM is charged, it would be immediately visible if it had similar charge to many SM particles. Specifically this will rule out the following charged, fundamental particles:  $e, \mu, \tau, W, u, d, s, c, t, b$  and their corresponding antiparticles. Recalling from earlier that DM must be long lived and stable over the age of the universe, this would exclude all SM particles with decay half-lives at or shorter than the age of the universe. The lifetime constraint additionally eliminates the  $Z$  and  $H$  bosons. Finally, the candidate DM needs to be somewhat massive. Recall from Section 2.2 that DM is cold or not relativistic through the universe. This eliminates the remaining SM particles:  $\nu_{e,\mu,\tau}, g, \gamma$  as DM candidates. Because there are no DM candidates within the SM, the DM problem strongly hints to physics beyond the SM (BSM).

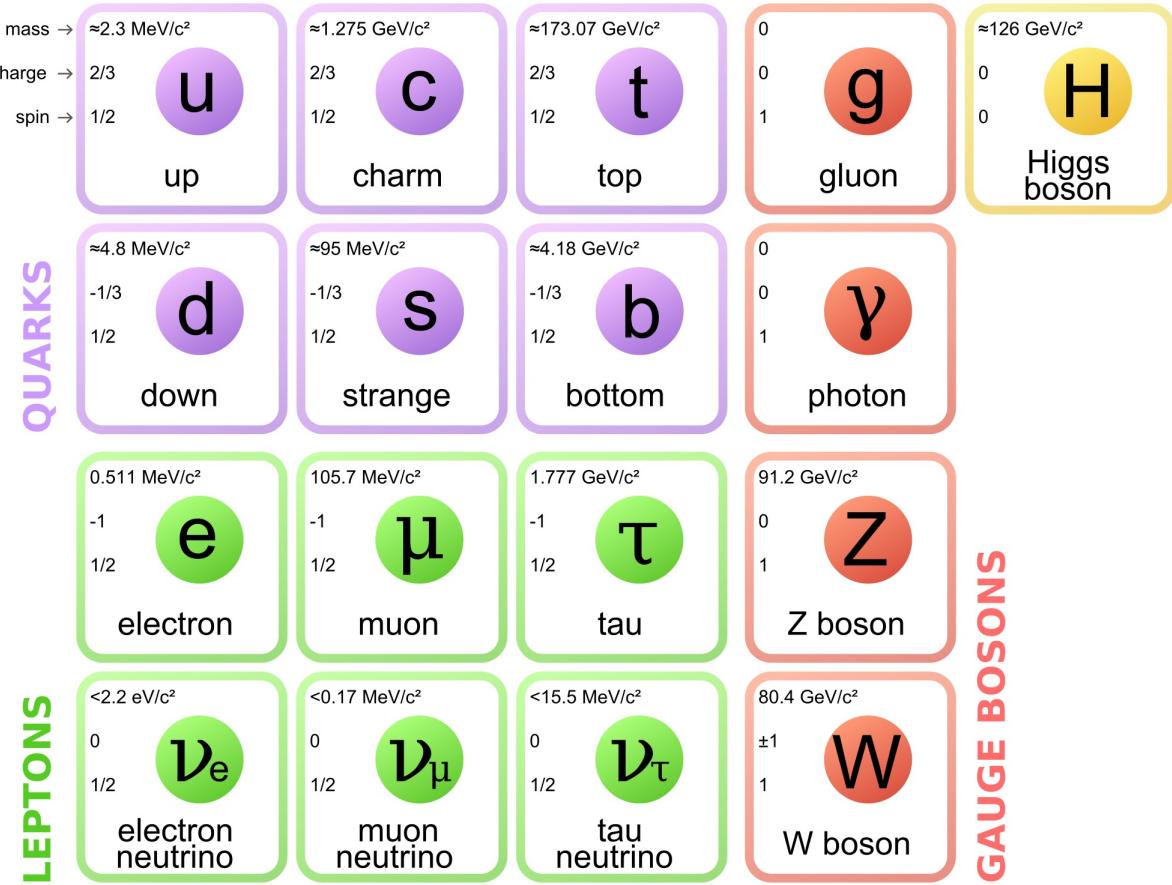


Figure 2.7 The Standard Model (SM) of particle physics. Figure taken from <http://www.quantumdiaries.org/2014/03/14/the-standard-model-a-beautiful-but-flawed-theory/>

### 2.4.1 Shake it, Break it, Make it

When considering DM that couples in some way with the SM, the interactions are roughly demonstrated by interaction demonstrated in Fig. 2.8. The figure is a simplified Feynman diagram where the arrow of time represents the interaction modes of: **Shake it, Break it, Make it**.

**Shake it** refers to the direct detection of dark matter. Direct detection interactions start with a free DM particle and some SM particle. The DM and SM interact under some elastic or inelastic collision and recoil away from each other. The DM remains in the dark sector and imparts some momentum onto the SM particle. The hope is that the momentum imparted onto the SM particle is sufficiently high enough to pick up with highly sensitive instruments. Because we cannot create the DM in the lab, a direct detection experiment must wait until DM is incident on the detector. Most direct detection experiments are therefore placed in low-background environments with inert

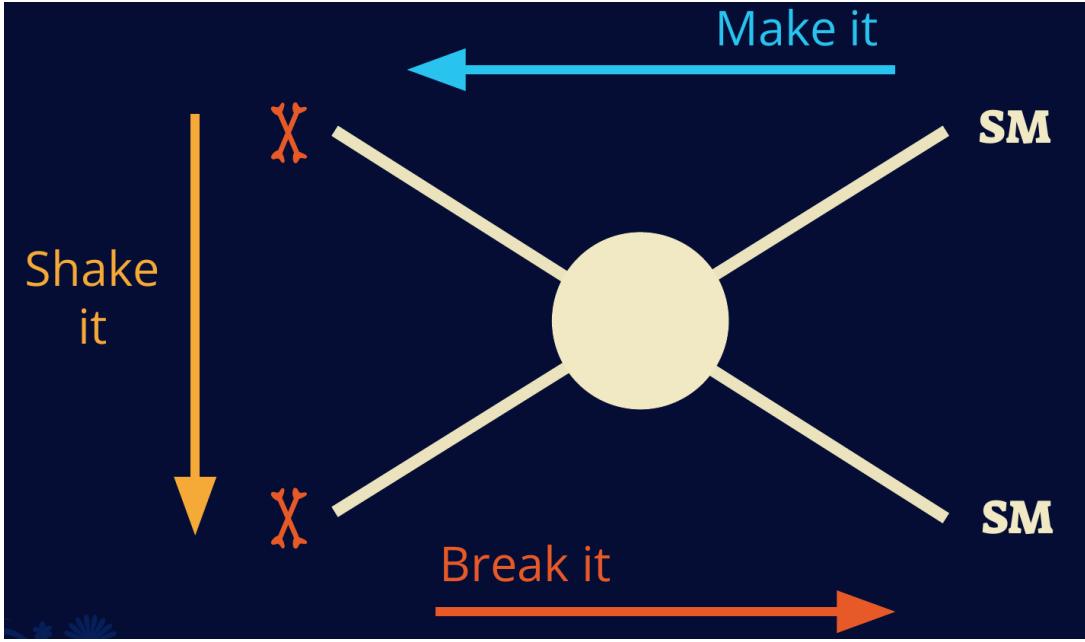


Figure 2.8 Simplified Feynman diagram demonstrating different ways DM can interact with SM particles. The 'X's refer to the DM particles whereas the SM refer to fundamental particles in the SM. The large circle in the center indicates the vertex of interaction and is purposely left vague. The colored arrows refer to different directions of time as well as their respective labels. The arrows indicate the initial and final state of the DM -SM interaction in time.

detection media like the noble gas Xenon. [14]

**Make it** refers to the production of DM from SM initial states. The experiment starts with particles in the SM. These SM particles are accelerated to incredibly high energies and then collided with each other. In the confluence of energy, DM hopefully emerges as a byproduct of the SM annihilation. Often it is the collider experiments that are able to generate energies high enough to probe DM production. These experiments include the world-wide collaborations ATLAS and CMS at CERN where protons are collided together at extreme energies. The DM searches however are complex. DM likely does not interact with the detectors and lives long enough to escape the detection apparatus of CERN's colliders. This means any DM production experiment searches for an excess of events with missing momentum or energy in the events. An example event with missing transverse momentum is shown in Fig. 2.9. The missing momentum with no particle tracks implies a neutral particle carried the energy out of the detector. However, there are other neutral particles in the SM, like neutrons or neutrinos, so any analysis have to account for SM signatures of missing

momentum. [15]

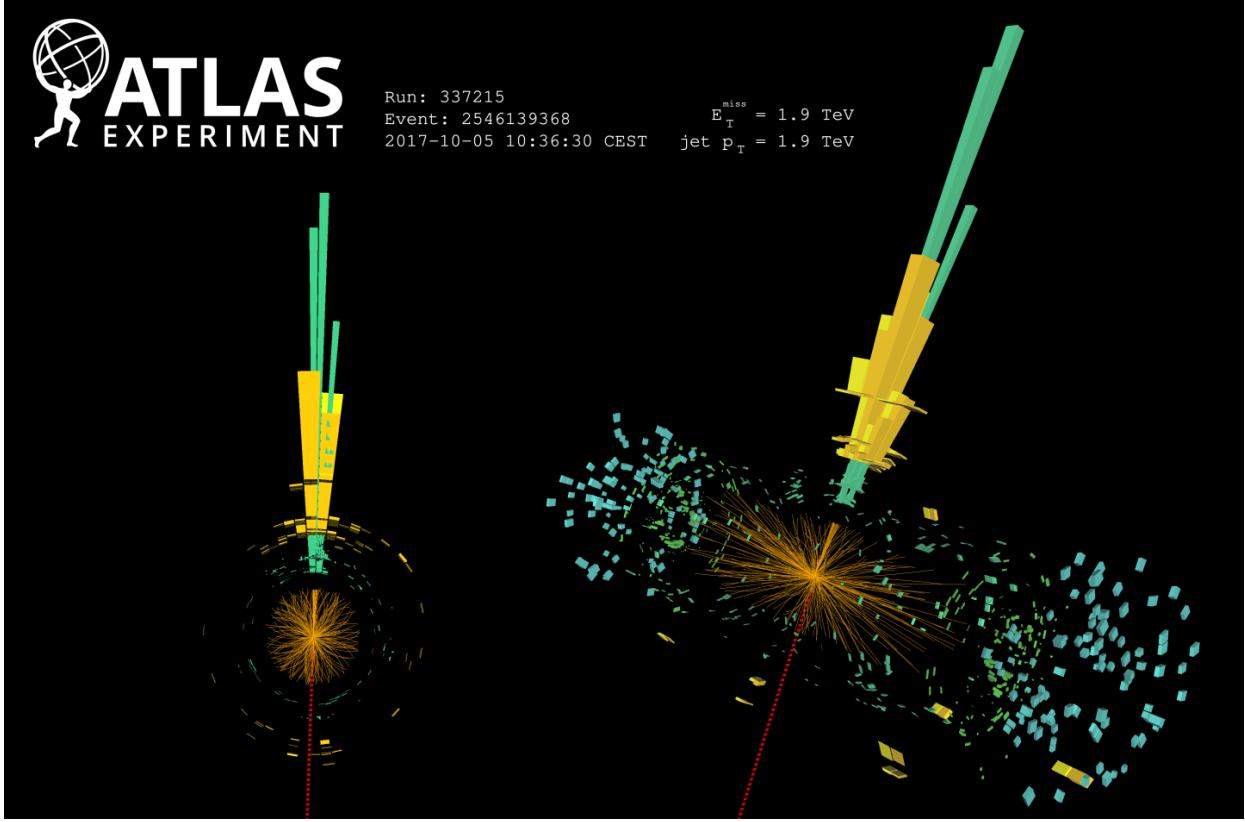


Figure 2.9 A single jet event in ATLAS detector from 2017 [16]. Jet momentum observed to be 1.9 TeV. Missing transverse momentum observed to be 1.9 TeV as the initial momentum of the event was 0. Implied MET is shown as a red dashed line in event display.

#### 2.4.2 Break it: Standard Model Signatures of Indirect Dark Matter Searches

**Break it** refers to the creation of SM particles from the dark sector, and it is the primary focus of this thesis. The interaction begins with DM or in the dark sector. The hypothesis is that this DM will either annihilate with itself or decay and produce a SM byproduct. This method is often referred to the Indirect Detection of DM because we have no lab to directly control or manipulate the DM. Therefore most DM primary observations will be performed from observations of known DM densities among the astrophysical sources. The strength is that we have the whole of the universe and its 13.6 billion year lifespan to use as the detector or particle accelerator. Additionally, locations of dark matter are also well understood since it was astrophysical observations that presented the problem of DM in the first place.

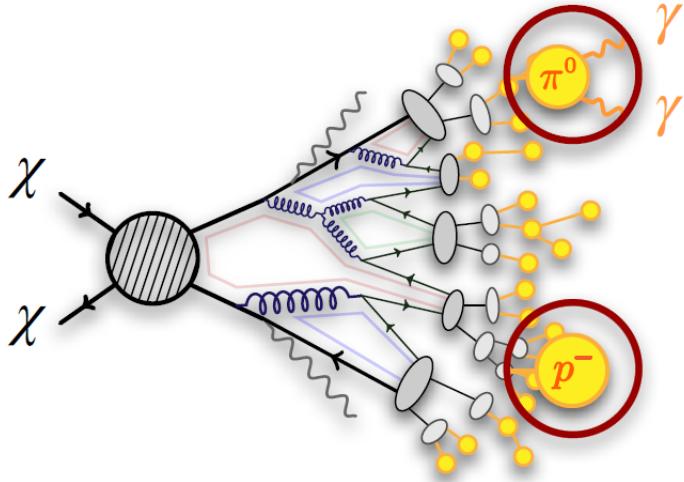


Figure 2.10 More detailed pseudo-Feynman diagram of particle cascade from dark matter annihilation into 2 quarks. The quarks hadronize and down to stable particles like  $\gamma$  or the anti-proton ( $p^-$ ). Diagram pulled from ICRC 2021 presentation on DM annihilation search [17].

However, anything can happen in the universe. There are many difficult to deconvolve backgrounds when searching for DM. One prominent example is the galactic center. There's a lot of DM there since the Milky Way definitely has a lot of DM. But any signal coming from there is hard to parse apart from the extreme environment of our supermassive black hole, Sagittarius A\* [18]. Despite the challenges, any DM model that yields evidence in the other observation two methods, **Shake it or Make it** must be corroborated with indirect observations of the known DM sources. Without corroborating evidence, DM observation in the lab is hard-pressed to demonstrate that it is the model contributing to the DM seen at the universal scale.

In the case of WIMP DM, signals are typically described in terms of primary SM particles produced from a DM decay or annihilation. The SM initial state particles are then simulated to stable final states such as the  $\gamma$ ,  $\nu$ ,  $p$ , or  $e$  which can traverse galactic lengths to reach Earth.

Fig. 2.10 shows the quagmire of SM particles that emerges from SM initial states that are not stable [17]. There are many different particles with varying energies that can be produced in such an interaction. For any arbitrary DM source and stable SM particle, the SM flux from DM annihilating

to some neutral particle in the SM,  $\phi$ , from a region in the sky is described by the following

$$\frac{d\Phi_\phi}{dE_\phi} = \frac{\langle\sigma v\rangle}{8\pi m_\chi^2} \frac{dN_\phi}{dE_\phi} \times \int_{\text{source}} d\Omega \int_{l.o.s} \rho_\chi^2 dl(r, \theta') \quad (2.4)$$

In Eq. (2.4),  $\langle\sigma v\rangle$  is the velocity-weighted annihilation cross-section of DM to the SM.  $m_\chi$  refers to the mass of DM, noted with greek letter  $\chi$ .  $\frac{dN_\phi}{dE_\phi}$  is the N particle flux weighted by the particle energy. An example is provided in Fig. 2.11 for the  $\gamma$  final state. The integrated terms are performed over the solid angle,  $d\Omega$ , and line of sight, l.o.s.  $\rho$  is the density of DM for a location  $(r, \theta')$  in the sky. The terms left of the '×' are often referred to as the particle physics component. The terms on the right are referred to as the astrophysical component. For decaying DM, the equation changes to...

$$\frac{d\Phi_\phi}{dE_\phi} = \frac{1}{4\pi\tau m_\chi} \frac{dN_\phi}{dE_\phi} \times \int_{\text{source}} d\Omega \int_{l.o.s} \rho_\chi dl(r, \theta') \quad (2.5)$$

In Eq. (2.5),  $\tau$  is the decay lifetime of the DM. Just as in Eq. (2.4), the left and right terms are the particle physics and the astrophysical components respectively. The integrated astrophysical component of Eq. (2.4) is often called the J-Factor. Whereas the integrated astrophysical component of Eq. (2.5) is often called the D-Factor.

Exact DM  $\text{DM} \rightarrow \text{SM}$  branching ratios are not known, so it is usually assumed to go 100% into a SM particle/anti-particle. Additionally, when a DM annihilation or decay produces one of the neutral, long-lived SM particles ( $\nu$  or  $\gamma$ ), the particle can be traced back to a DM source. For DM above GeV energies, there are very few SM processes that can produce particles with such a high energy. Seeing such a signal would almost certainly be an indication of the presence of dark matter. The universe fortunately provides us with the largest volume and lifetime ever for a particle physics experiment.

## 2.5 Sources for Indirect Dark Matter Searches

We of course have to know where to look. Thankfully, we have a good idea of where. The first detection of DM relied on optical observations. Since then, we've developed new techniques to find DM dense regions. As described in Section 2.3.1, many DM dense regions were through observing galactic rotation curves. Our Milky Way galaxy is among DM dense regions discovered,

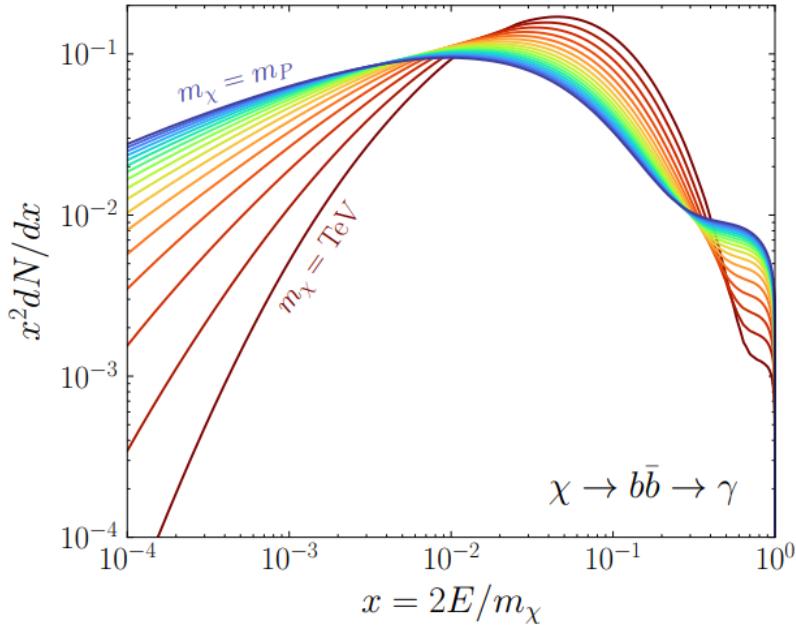


Figure 2.11 Dark Matter (DM) decay spectrum for  $b\bar{b}$  initial state and  $\gamma$  final state. Redder spectra are for larger DM masses. Bluer spectra are light DM masses.  $x$  is a unitless factor defined as the ratio of the mass of DM,  $m_\chi$ , and the final state particle energy  $E_\gamma$ . Figure from [19].

and it is the largest nearby DM dense region to look at. Additionally, the DM halo surrounding the Milky Way is somewhat clumpy [18]. There are regions in the DM halo of the Milky Way that have more DM than others they have captured gas over time. In some cases these sub-haloes were dense enough to host stars. These apparent sub galaxies are known was dwarf spheroidal galaxies and are the main sources studied in this thesis. Each source type comes with different trade offs. Galactic Center studies will be very sensitive to the assume distribution of DM. The central DM density can very substantially as demonstrated in Fig. 2.12. At small  $r$ , the differences in DM densities can be 3-4 orders of magnitude.

Dwarf Spheroidal Galaxies (dSph's) studies suffer from uncertainties in the DM density less than the galactic center studies. This is mostly from their diminutive size being smaller than the angular resolution of most  $\gamma$ -ray observatories [18]. The DM content dSph's are typically determined with the virial theorem, Eq. (2.1), and are usually majority DM [18] in mass. DSph's tend to be ideal sources to look at for DM searches. Their environments are fairly quiet with little

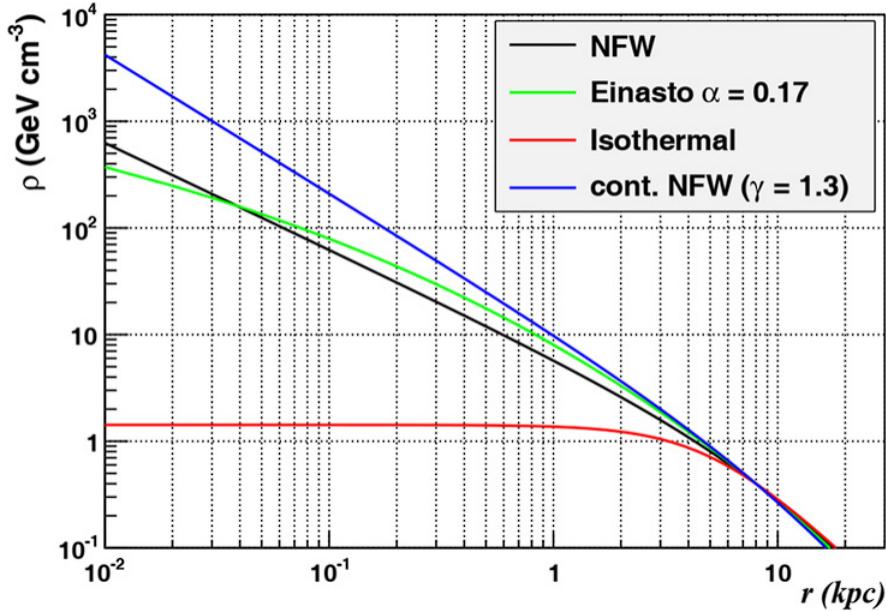


Figure 2.12 Different dark matter density profiles compared. Some models produce very large densities at small  $r$  [20].

astrophysical backgrounds. Unlike the galactic center, the most active components of dSph's are the stars within them versus a violent accretion disc around a black hole. All this together means that dSph's are among the best sources to look at for indirect DM searches. dSph's are the targets of focus for this thesis.

## 2.6 Multi-Messenger Dark Matter

Astrophysics entered a new phase in the past few decades that leverages our increasing sensitivity to SM channels and general relativity (GR). Up until the 21st century, astrophysical observations were done with photons ( $\gamma$ ) only. Astrophysics with this 'messenger' is fairly mature now. Novel observations of the universe have since only adjusted the sensitivity of the wavelength of light that's observed. Gems like the CMB [10], and more have ultimately been observations of different wavelengths of light. Multi-messenger astrophysics proposes using other SM particles such the  $p^{+/-}$ , or  $\nu$  or gravitation waves predicted by general relativity.

The experiments LIGO had a revolutionary discovery in 2016 with the first detection of a binary black hole merger [21]. This opened the collective imagination entirely to observing the universe

through gravitational waves. There's also been a surge of interest in the neutrino ( $\nu$ ) sector. IceCube demonstrated that we are sensitive to neutrinos in regions that correlate with significant photon emmission like the galactic plane [22]. Neutrinos, like gravitational waves and light, travel mostly unimpeded from their source to our observatories. This makes pointing to the oringinating source of the these messengers much easier than it is for cosmic rays that are almost always deflected from their source.

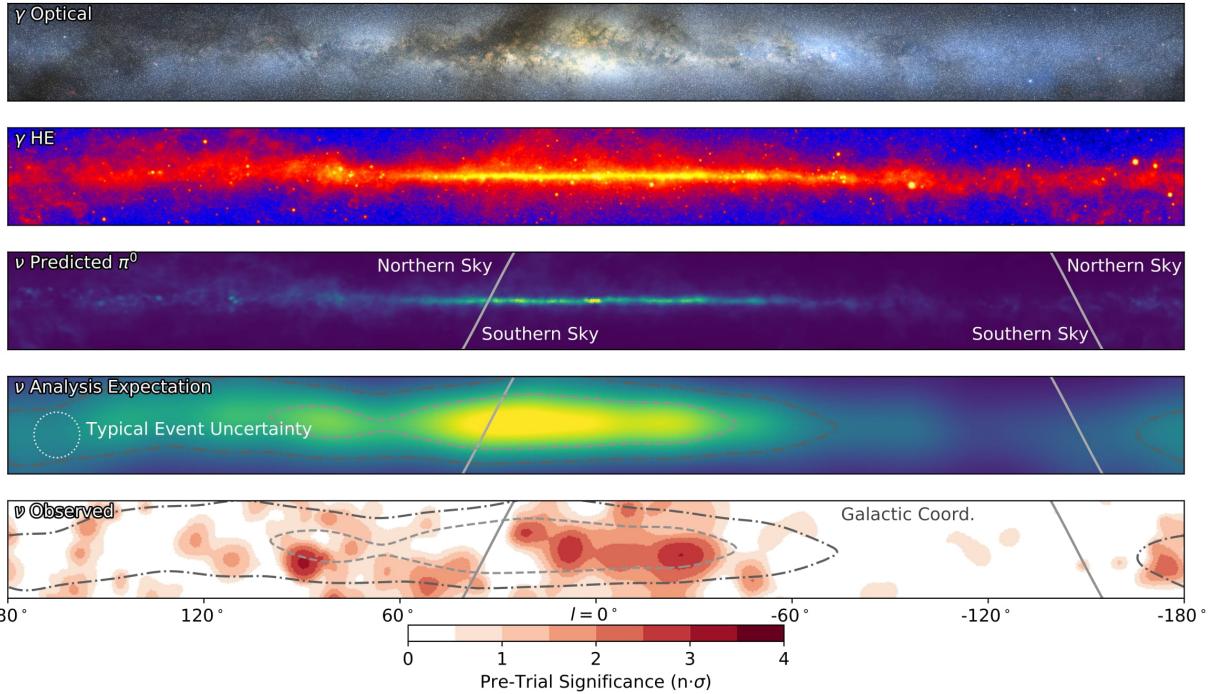


Figure 2.13 The Milky Way Galaxy in photons ( $\gamma$ ) and neutrinos ( $\nu$ ) [22]. Galactic center is at  $l=0^\circ$ and is the brightest region in all panels. (top) An Optical color image of the Milky Way galaxy seen from Earth. Clouds of gas and dust obscure some of the light from stars. (2nd down) Integrated flux of  $\gamma$ -rays observed by the Fermi-LAT telescope [23]. (middle) Expected neutrino emmision that corresponds with Fermi-LAT observations. (2nd up) Expected neutrino emmision profile after considering detector systematics of IceCube. (bottom) Observed neutrino emmision from region of the galactic plane. Substantial neutrino emmision is detected.

The recent result from IceCube, shown in Fig. 2.13, proves that we can make obervations under different messenger regimes. The top two panels are the appearance of the galactic plane to different wavelengths of light. Some sources are more apparent in some panels, while others are not. The IceCube collaboration recently published a groundbreaking result of the Milky Way in neutrinos.

This new channel is potentially very powerful because neutrinos are readily able to penetrate see through gas and dust in the Milky Way. This new image also refines our understanding of how high energy particles are accelerated since the fit to IceCube data prefers one standard model process over the other [22].

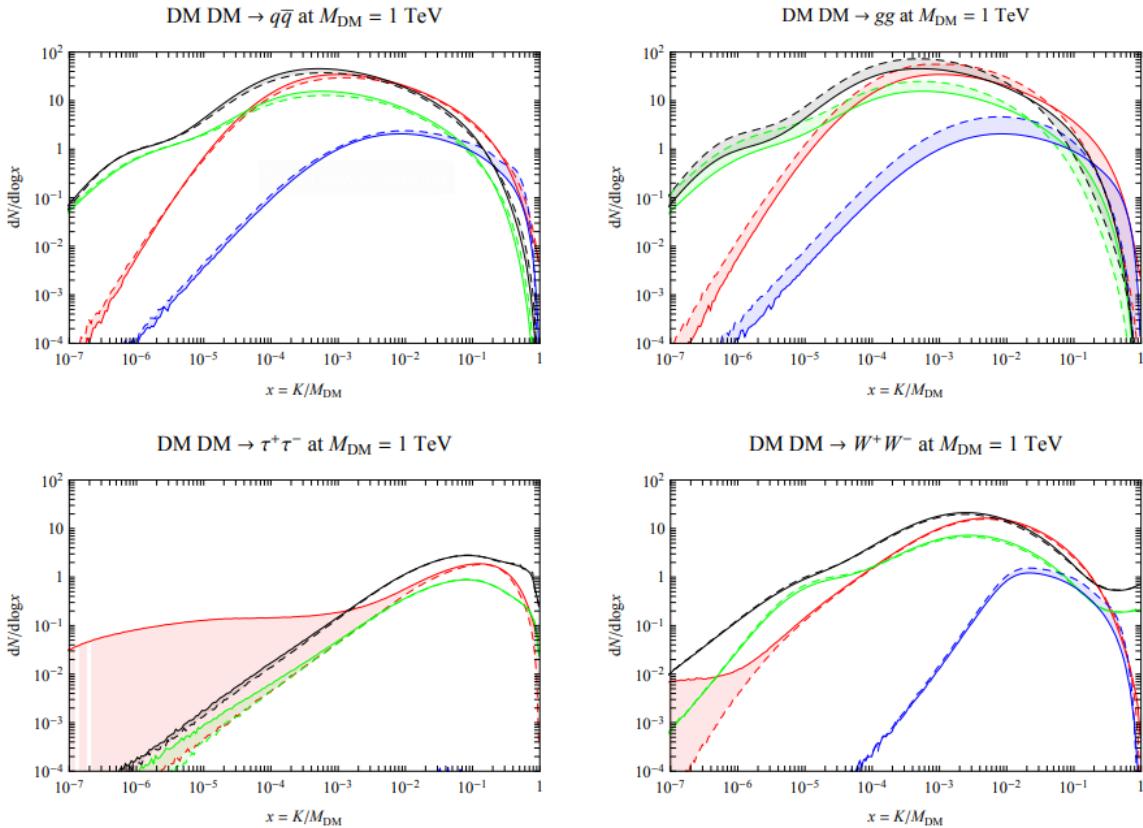


Figure 2.14 Dark Matter annihilation spectra for different final state particle and standard model annihilation channels [24]. Photons (red),  $e^\pm$  (green),  $\bar{p}$  (blue),  $\nu$  (black).

Exposing our observations to more cosmic messengers greatly increases our sensitivity to rare processes. In the case of DM, Fig. 2.14, there are many SM particles produced in a dark matter annihilation. Among the final state fluxes are gammas and neutrinos. Charged particles are also produced however they would not likely make it to Earth since they will be deflected by magnetic fields between the source and Earth. This means observatories that can see the neutral messengers are especially good for DM searches and for combining data for a multi-messenger DM search.

## CHAPTER 3

### MULTIMESSENGER ASTROPHYSICS: DETECTING HIGH ENERGY NEUTRAL MESSENGERS

#### 3.1 Introduction

Before the 20th century, all astrophysics observations were optical in nature. We literally only saw things with highly magnified optical observations. Then we discovered cosmic rays. cosmic rays are charged particles, typically naked protons or H+. This was seen by Victor Hess in 19???. Around the same time we discovered neutrinos from beta decay. Sometime around 1950 we started to build neutrino detectors which were mostly sensitive to neutrinos from the sun. Finally, it was theorized that compact objects like black holes and neutron stars would create waves in space-time when they experience mergers or collisions.

In the 21st century, we have developed new observation techniques and detectors that are no only sensitive to these four messengers - photons ( TODO: photon), neutrinos ( TODO: nu), Cosmic Rays (CR), and Gravitational Wave (WV) - we're collect high energy versions of these events. For the standad model particles, we're now sensitive to all messengers above the MeV eneryg range. Additionally, the GW's were sensitive to are in the stellar mass black hole region and above within our galactic neighborhood. This means were becoming sensitive to the fundamental physics occuring within the universe and we can rely on the universe as a TeV+ particle accelerator. We also have the abaility to correlate high energy events across messengers and gain new insights on the processes that occur in our universe.

This thesis focuses on very high energy (VHE) gamma rays and neutrinos. These can both be observed through the water cherenkov detection technique altho not exclusively. Methods on how to detect and observe these neutral messengers are discussed Section 3.3 and Section 3.4

#### 3.2 Charged Particles in a Medium

For high enery gamma-rays and neutrinos, we can exploit the same effect that charged particles have with water. This effect is known as Cherenkov radiation. Cherenkov Radiation occurs when a charged particle, usually electrons ( $e$ ) or muons ( $\mu$ ), traverse a medium, like water, faster than the

speed of light in that medium. This is similar to sonic boom where an object moves through air faster than the speed of sound in air. Cherenkov radiation can therefore be thought of as an 'optic boom'. Many astro-particle physics experiments will use water as the medium as because water has a unique set of properties ideal for charged particle tracking.



Figure 3.1 TODO: Show a nuclear reactor with cherenkov radiation[NEEDS A SOURCE][FACT CHECK THIS]

The frequency of light emitted due to cherenkov radiation follows the equation:

$$\text{INSERTCherenkovwavelengthcalcHERE.} \quad (3.1)$$

The absorption spectra is shown in the following figure:

### 3.3 Photons ( $\gamma$ )

### 3.4 Neutrinos ( $\nu$ )

### 3.5 Opportunities to Combine for Dark Matter



Figure 3.2 TODO: absorption spectrum of liquid and solid water[NEEDS A SOURCE][FACT CHECK THIS]

## **CHAPTER 4**

### **HIGH ALTITUDE WATER CHERENKOV (HAWC) OBSERVATORY**

#### **4.1 The Detector**

#### **4.2 Events Reconstruction and Data Acquisition**

##### **4.2.1 G/H Discrimination**

##### **4.2.2 Angle**

##### **4.2.3 Energy**

#### **4.3 Remote Monitoring**

##### **4.3.1 ATHENA Database**

##### **4.3.2 HOMER**

## **CHAPTER 5**

### **ICECUBE NEUTRINO OBSERVATORY**

#### **5.1 The Detector**

#### **5.2 Events Reconstruction and Data Acquisition**

##### **5.2.1 Angle**

##### **5.2.2 Energy**

#### **5.3 Northern Test Site**

##### **5.3.1 PIgeon remote dark rate testing**

##### **5.3.2 Bulkhead Construction**

## **CHAPTER 6**

### **COMPUTATIONAL TECHNIQUES IN PARTICLE ASTROPHYSICS**

**6.1 Neural Networks for Gamma/Hadron Separation**

**6.2 Parallel Computing for Dark Matter Analyses**

**CHAPTER 7**

**GLORY DUCK**

## **CHAPTER 8**

### **NU DUCK**

## BIBLIOGRAPHY

- [1] Anne M. Green. “Dark matter in astrophysics/cosmology”. In: *SciPost Phys. Lect. Notes* (2022), p. 37. doi: [10.21468/SciPostPhysLectNotes.37](https://doi.org/10.21468/SciPostPhysLectNotes.37). URL: <https://scipost.org/10.21468/SciPostPhysLectNotes.37>.
- [2] Bing-Lin Young. “A survey of dark matter and related topics in cosmology”. In: *Frontiers of Physics* 12 (Oct. 2016). doi: <https://doi.org/10.1007/s11467-016-0583-4>. URL: <https://doi.org/10.1007/s11467-016-0583-4>.
- [3] Gianfranco Bertone, Dan Hooper, and Joseph Silk. “Particle dark matter: evidence, candidates and constraints”. In: *Physics Reports* 405.5 (2005), pp. 279–390. issn: 0370-1573. doi: <https://doi.org/10.1016/j.physrep.2004.08.031>. URL: <https://www.sciencedirect.com/science/article/pii/S0370157304003515>.
- [4] Gianfranco Bertone and Dan Hooper. “History of dark matter”. In: *Rev. Mod. Phys.* 90 (4 Aug. 2018), p. 045002. doi: [10.1103/RevModPhys.90.045002](https://doi.org/10.1103/RevModPhys.90.045002). URL: <https://link.aps.org/doi/10.1103/RevModPhys.90.045002>.
- [5] Fritz Zwicky. “The Redshift of Extragalactic Nebulae”. In: *Helvetica Physica Acta* 6. (1933), pp. 110–127. doi: [10.5169/seals-110267](https://doi.org/10.5169/seals-110267).
- [6] Vera C. Rubin and Jr. Ford W. Kent. “Rotation of the Andromeda Nebula from a Spectroscopic Survey of Emission Regions”. In: *ApJ* 159 (Feb. 1970), p. 379. doi: [10.1086/150317](https://doi.org/10.1086/150317).
- [7] K. G. Begeman, A. H. Broeils, and R. H. Sanders. “Extended rotation curves of spiral galaxies: dark haloes and modified dynamics”. In: *Monthly Notices of the Royal Astronomical Society* 249.3 (Apr. 1991), pp. 523–537. issn: 0035-8711. doi: [10.1093/mnras/249.3.523](https://doi.org/10.1093/mnras/249.3.523). eprint: <https://academic.oup.com/mnras/article-pdf/249/3/523/18160929/mnras249-0523.pdf>. URL: <https://doi.org/10.1093/mnras/249.3.523>.
- [8] *Different types of gravitational lenses*. website. Feb. 2004. URL: <https://esahubble.org/images/heic0404b/>.
- [9] Douglas Clowe et al. “A Direct Empirical Proof of the Existence of Dark Matter”. In: *apjl* 648.2 (Sept. 2006), pp. L109–L113. doi: [10.1086/508162](https://doi.org/10.1086/508162). arXiv: [astro-ph/0608407 \[astro-ph\]](https://arxiv.org/abs/astro-ph/0608407).
- [10] Planck Collaboration and N. et. al. Aghanim. “Planck 2018 results I. Overview and the cosmological legacy of Planck”. In: *A&A* 641 (2020). doi: [10.1051/0004-6361/201833880](https://doi.org/10.1051/0004-6361/201833880). URL: <https://doi.org/10.1051/0004-6361/201833880>.
- [11] Wenlong Yuan et al. “A First Look at Cepheids in a Type Ia Supernova Host with JWST”. in: *The Astrophysical Journal Letters* 940.1 (Nov. 2022). doi: [10.3847/2041-8213/ac9b27](https://doi.org/10.3847/2041-8213/ac9b27).

URL: <https://dx.doi.org/10.3847/2041-8213/ac9b27>.

- [12] Wendy L. Freedman. “Measurements of the Hubble Constant: Tensions in Perspective”. In: *The Astrophysical Journal* 919.1 (Sept. 2021), p. 16. doi: [10.3847/1538-4357/ac0e95](https://doi.org/10.3847/1538-4357/ac0e95). URL: <https://dx.doi.org/10.3847/1538-4357/ac0e95>.
- [13] P. Tisserand et al. “Limits on the Macho content of the Galactic Halo from the EROS-2 Survey of the Magellanic Clouds”. In: *A&A* 469.2 (2007), pp. 387–404. doi: [10.1051/0004-6361:20066017](https://doi.org/10.1051/0004-6361:20066017). URL: <https://doi.org/10.1051/0004-6361:20066017>.
- [14] Jodi Cooley. “Dark Matter direct detection of classical WIMPs”. In: *SciPost Phys. Lect. Notes* (2022), p. 55. doi: [10.21468/SciPostPhysLectNotes.55](https://doi.org/10.21468/SciPostPhysLectNotes.55). URL: <https://scipost.org/10.21468/SciPostPhysLectNotes.55>.
- [15] “Search for new phenomena in events with an energetic jet and missing transverse momentum in  $pp$  collisions at  $\sqrt{s} = 13$  TeV with the ATLAS detector”. In: *Phys. Rev. D* 103 (11 July 2021), p. 112006. doi: [10.1103/PhysRevD.103.112006](https://doi.org/10.1103/PhysRevD.103.112006). URL: <https://link.aps.org/doi/10.1103/PhysRevD.103.112006>.
- [16] *Jetting into the dark side: a precision search for dark matter*. website. July 2020. URL: <https://atlas.cern/updates/briefing/precision-search-dark-matter>.
- [17] Celine Armand et. al. “Combined dark matter searches towards dwarf spheroidal galaxies with Fermi-LAT, HAWC, H.E.S.S., MAGIC, and VERITAS”. in: *Proceedings of Science*. Vol. 395. Mar. 2022. doi: <https://doi.org/10.22323/1.395.0528>.
- [18] Tracy R. Slatyer. “Les Houches Lectures on Indirect Detection of Dark Matter”. In: *SciPost Phys. Lect. Notes* (2022), p. 53. doi: [10.21468/SciPostPhysLectNotes.53](https://doi.org/10.21468/SciPostPhysLectNotes.53). URL: <https://scipost.org/10.21468/SciPostPhysLectNotes.53>.
- [19] Christian W Bauer, Nicholas L. Rodd, and Bryan R. Webber. “Dark matter spectra from the electroweak to the Planck scale”. In: *Journal of High Energy Physics* 2021.1029-8479 (June 2021). doi: [https://doi.org/10.1007/JHEP06\(2021\)121](https://doi.org/10.1007/JHEP06(2021)121).
- [20] Riccardo Catena and Piero Ullio. “A novel determination of the local dark matter density”. In: *Journal of Cosmology and Astroparticle Physics* 2010.08 (Aug. 2010), p. 004. doi: [10.1088/1475-7516/2010/08/004](https://doi.org/10.1088/1475-7516/2010/08/004). URL: <https://dx.doi.org/10.1088/1475-7516/2010/08/004>.
- [21] B. P. Abbott et al. “Observation of Gravitational Waves from a Binary Black Hole Merger”. In: *Phys. Rev. Lett.* 116 (6 Feb. 2016), p. 061102. doi: [10.1103/PhysRevLett.116.061102](https://doi.org/10.1103/PhysRevLett.116.061102). URL: <https://link.aps.org/doi/10.1103/PhysRevLett.116.061102>.
- [22] R. Abbasi et. al. “Observation of high-energy neutrinos from the Galactic plane”. In: *Science* 380.6652 (June 2023), pp. 1338–1343.

- [23] NASA Goddard Space Flight Center. *Fermi’s 12-year view of the gamma-ray sky*. website. 2022. URL: <https://svs.gsfc.nasa.gov/14090>.
- [24] Marco Cirelli et al. “PPPC 4 DM ID: a poor particle physicist cookbook for dark matter indirect detection”. In: *Journal of Cosmology and Astroparticle Physics* 2011.03 (Mar. 2011), p. 051. DOI: [10.1088/1475-7516/2011/03/051](https://doi.org/10.1088/1475-7516/2011/03/051). URL: <https://dx.doi.org/10.1088/1475-7516/2011/03/051>.