

1 LEVERAGING MULTI-MESSENGER ASTROPHYSICS FOR DARK MATTER SEARCHES

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Today

**ABSTRACT**

3 I did Dark Matter with HAWC and IceCube. I also used Graph Neural Networks

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<sup>6</sup> Today

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8 I love my friends. Thanks to everyone that helped me figure this out. Amazing thanks to the people  
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**LIST OF ABBREVIATIONS**

- 407 **MSU** Michigan State University  
408 **LANL** Los Alamos National Laboratory  
409 **DM** Dark Matter  
410 **SM** Standard Model  
411 **HAWC** High Altitude Water Cherenkov Observatory  
412 **dSph** Dwarf Spheroidal Galaxy

## **CHAPTER 1**

### **INTRODUCTION**

414 Is the text not rendering right? Ah ok it knows im basically drafting the doc still

## CHAPTER 2

415

### DARK MATTER IN THE COSMOS

416 **2.1 Introduction**

417 The dark matter problem can be summarized in part by the following thought experiment.

418 Let us say you are the teacher for an elementary school classroom. You take them on a field  
419 trip to your local science museum and among the exhibits is one for mass and weight. The exhibit  
420 has a gigantic scale, and you come up with a fun problem for your class.

421 You ask your class, "What is the total weight of the classroom? Give your best estimation to  
422 me in 30 minutes, and then we'll check your guess on the scale. If your guess is within 10% of the  
423 right answer, we will stop for ice cream on the way back."

424 The students are ecstatic to hear this, and they get to work. The solution is some variation of  
425 the following strategy. The students should give each other their weight or best guess if they do  
426 not know. Then, all they must do is add each student's weight and get a grand total for the class.  
427 The measurement on the giant scale should show the true weight of the class. When comparing  
428 the measured weight to your estimation, multiply the measurement by  $1.0 \pm 0.1$  to get the  $\pm 10\%$   
429 tolerances for your estimation.

430 Two of your students, Sandra and Mario, return to you with a solution.

431 They say, "We weren't sure of everyone's weight. We used 65 lbs. for the people we didn't  
432 know and added everyone who does know. There are 30 of us, and we got 2,000 lbs.! That's a ton!"

433 You estimated 1,900 lbs. assuming the average weight of a student in your class was 60 lbs.  
434 So, you are pleased with Sandra's and Mario's answer. You instruct your students to all gather on  
435 the giant scale and read off the weight together. To all your surprise, the scale reads *10,000 lbs.*!  
436 10,000 is significantly more than a 10% error from 2,000. In fact, it is approximately 5 times more  
437 massive than either your or your students' estimates. You think to yourself and conclude there  
438 must be something wrong with the scale. You ask an employee to check the scale and verify it is  
439 well calibrated. They confirm that the scale is in working order. You weigh a couple of students  
440 individually to assess that the scale is well calibrated. Sandra weighs 59 lbs., and Mario weighs

441 62 lbs., typical weights for their age. You then weigh each student individually and see that their  
442 weights individually do not deviate greatly from 60 lbs. So, where does all the extra weight come  
443 from?

444 This thought experiment serves as an analogy to the Dark Matter problem. The important  
445 substitution to make however is to replace the students with stars and the classroom with a galaxy,  
446 say the Milky Way. Individually the mass of stars is well measured and defined with the Sun as our  
447 nearest test case. However, when we set out to measure the mass of a collection of stars as large as  
448 galaxies, our well-motivated estimation is wildly incorrect. There simply is no way to account for  
449 this discrepancy except without some unseen, or dark, contribution to mass and matter in galaxies.  
450 I set out in my thesis to narrow the possibilities of what this Dark Matter could be.

451 This chapter is organized like the following... **TODO: Text should look like ... Chapter x has**  
452 **blah blah blah.**

453 **2.2 Dark Matter Basics**

454 Presently, a more compelling theory of cosmology that includes Dark Matter (DM) in order  
455 to explain a variety of observations is  $\Lambda$  Cold Dark Matter, or  $\Lambda$ CDM. I present the evidence  
456 supporting  $\Lambda$ CDM in Section 2.3 yet discuss the conclusions of the  $\Lambda$ CDM model here. According  
457 to  $\Lambda$ CDM fit to observations on the Cosmic Microwave Background (CMB), DM is 26.8% of the  
458 universe's current energy budget. Baryonic matter, stuff like atoms, gas, and stars, contributes to  
459 4.9% of the universe's current energy budget [1, 2, 3].

460 DM is dark; it does not interact readily with light at any wavelength. DM also does not interact  
461 noticeably with the other standard model forces (Strong and Weak) at a rate that is readily observed  
462 [3]. DM is cold, which is to say that the average velocity of DM is below relativistic speeds [1].  
463 'Hot' DM would not likely manifest the dense structures we observe like galaxies, and instead  
464 would produce much more diffuse galaxies than what is observed [3, 1]. DM is old; it played a  
465 critical role in the formation of the universe and the structures within it [1, 2].

466 Observations of DM have so far been only gravitational. The parameter space available to what  
467 DM could be therefore is extremely broad. The broad DM parameter space is iteratively tested in

468 DM searches by supposing a hypothesis that has not yet been ruled out and performing observations  
469 to test them. When the observations yield a null result, the parameter space is constrained further.  
470 I present some approaches for DM searches in Section 2.4.

471 **2.3 Evidence for Dark Matter**

472 Dark Matter (DM) has been a looming problem in physics for almost 100 years. Anomalies  
473 have been observed by astrophysicists in galactic dynamics as early as 1933 when Fritz Zwicky  
474 noticed unusually large velocity dispersion in the Coma cluster. Zwicky's measurement was the  
475 first recorded to use the Virial theorem to measure the mass fraction of visible and invisible matter  
476 in celestial bodies [4]. From Zwicky in [5], "*If this would be confirmed, we would get the surprising*  
477 *result that dark matter is present in much greater amount than luminous matter.*" Zwicky's and  
478 others' observation did not instigate a crisis in astrophysics because the measurements did not  
479 entirely conflict with their understanding of galaxies [4]. In 1978, Rubin, Ford, and Norbert  
480 measured rotation curves for ten spiral galaxies [6]. Rubin et al.'s 1978 publication presented a  
481 major challenge to the conventional understanding of galaxies that could no longer be dismissed by  
482 measurement uncertainties. Evidence has been mounting ever since for this exotic form of matter.  
483 The following subsections provide three compelling pieces of evidence in support of the existence  
484 of DM.

485 **2.3.1 First Clues: Stellar Velocities**

486 Zwicky, and later Rubin, measured the stellar velocities of various galaxies to estimate their  
487 virial mass. The Virial Theorem upon which these observations are interpreted is written as

$$2T + V = 0. \quad (2.1)$$

488 Where  $T$  is the kinetic energy and  $V$  is the potential energy in a self-gravitating system. The  
489 classical Newton's law of gravity from stars and gas was used for gravitational potential modeled in  
490 the observed galaxies:

$$V = -\frac{1}{2} \sum_i \sum_{j \neq i} \frac{m_i m_j}{r_{ij}}. \quad (2.2)$$

491 Zwicky et al. measured just the velocities of stars apparent in optical wavelengths [5]. Rubin et al.  
 492 added by measuring the velocity of the hydrogen gas via the 21 cm emission line of Hydrogen [6].  
 493 The velocities of the stars and gas are used to infer the total mass of galaxies and galaxy clusters  
 494 via Equation (2.1). An inferred mass is obtained from the luminosity of the selected sources. The  
 495 two inferences are compared to each other as a luminosity to mass ratio which typically yielded [1]

$$\frac{M}{L} \sim 400 \frac{M_{\odot}}{L_{\odot}} \quad (2.3)$$

496  $M_{\odot}$  and  $L_{\odot}$  referring to stellar mass and stellar luminosity, respectively. These ratios clearly indicate  
 497 a discrepancy in apparent light and mass from stars and gas and their velocities.

498 Rubin et al. [6] demonstrated that the discrepancy was unlikely to be an underestimation of  
 499 the mass of the stars and gas. The inferred "dark" mass was up to 5 times more than the luminous  
 500 mass. This dark mass also needed to extend well beyond the extent of the luminous matter.

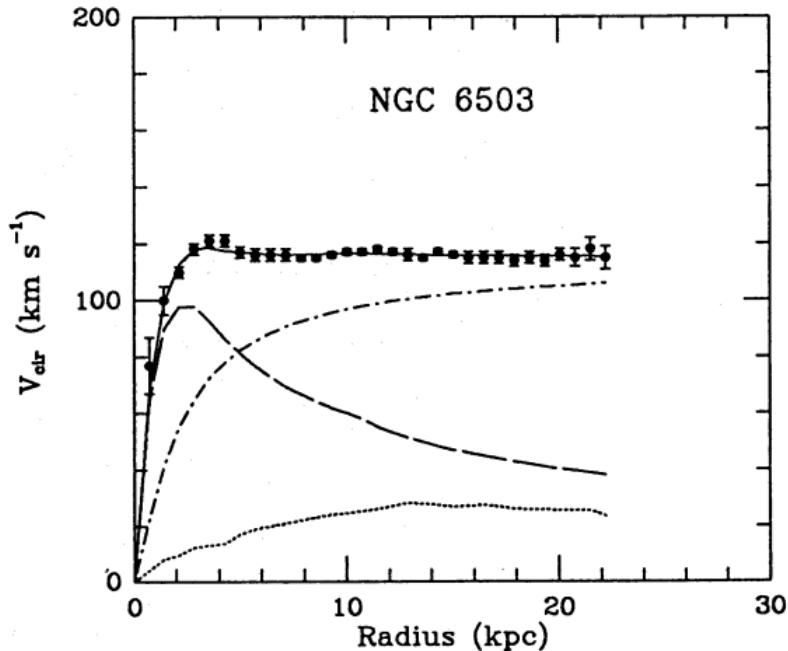


Figure 2.1 Rotation curve fit to NGC 6503 from [7]. Dashed line is the contribution from visible matter. Dotted curves are from gas. Dash-dot curves are from dark matter (DM). Solid line is the composite contribution from all matter and DM sources. Data are indicated with bold dots with error bars. Data agree strongly with a matter + DM composite prediction.

501 Figure 2.1 features one of many rotation curves plotted from the stellar velocities within galaxies.

502 The measured rotation curves mostly feature a flattening of velocities at higher radius which is not  
503 expected if the gravity was only coming from gas and luminous matter. The extension of the  
504 flat velocity region also indicates that the DM is distributed far from the center of the galaxy.  
505 Modern velocity measurements include significantly larger objects, galactic clusters, and smaller  
506 objects, dwarf galaxies. Yet, measurements along this regime are leveraging the Virial theorem  
507 with Newtonian potential energies. We know Newtonian gravity is not a comprehensive description  
508 of gravity. New observational techniques have been developed since 1978, and those are discussed  
509 in the following sections.

510 **2.3.2 Evidence for Dark Matter: Gravitational Lensing**

511 Modern evidence for dark matter comes from new avenues beyond stellar velocities. Grav-  
512 itational lensing from DM is a new channel from general relativity. General relativity predicts  
513 aberrations in light caused by massive objects. In recent decades we have been able to measure the  
514 lensing effects from compact objects and DM halos. Figure 2.2 shows how different massive ob-  
515 jects change the final image of a faraway galaxy resulting from gravitational lensing. Gravitational  
516 lensing developed our understanding of dark matter in two important ways.

517 Gravitational lensing provides additional compelling evidence for DM. The observation of two  
518 merging galactic clusters in 2006, shown in Figure 2.3, provided a compelling argument for DM  
519 outside the Standard Model. These clusters merged recently in astrophysical time scales. Galaxies  
520 and star cluster have mostly passed by each other as the likelihood of their collision is low. Therefore,  
521 these massive objects will mostly track the highest mass, dark and/or baryonic, density. Yet, the  
522 intergalactic gas is responsible for the majority of the baryonic mass in the systems [4]. These gas  
523 bodies will not phase through and will heat up as they collide together. The hot gas is located via  
524 x-ray emission from the cluster. Two observations of the clusters were performed independently of  
525 each other.

526 The first was the lensing of light around the galaxies due to their gravitational influences.  
527 When celestial bodies are large enough, the gravity they exert bends space and time itself. The  
528 warped space-time lenses light and will deflect in an analogous way to how glass lenses will bend

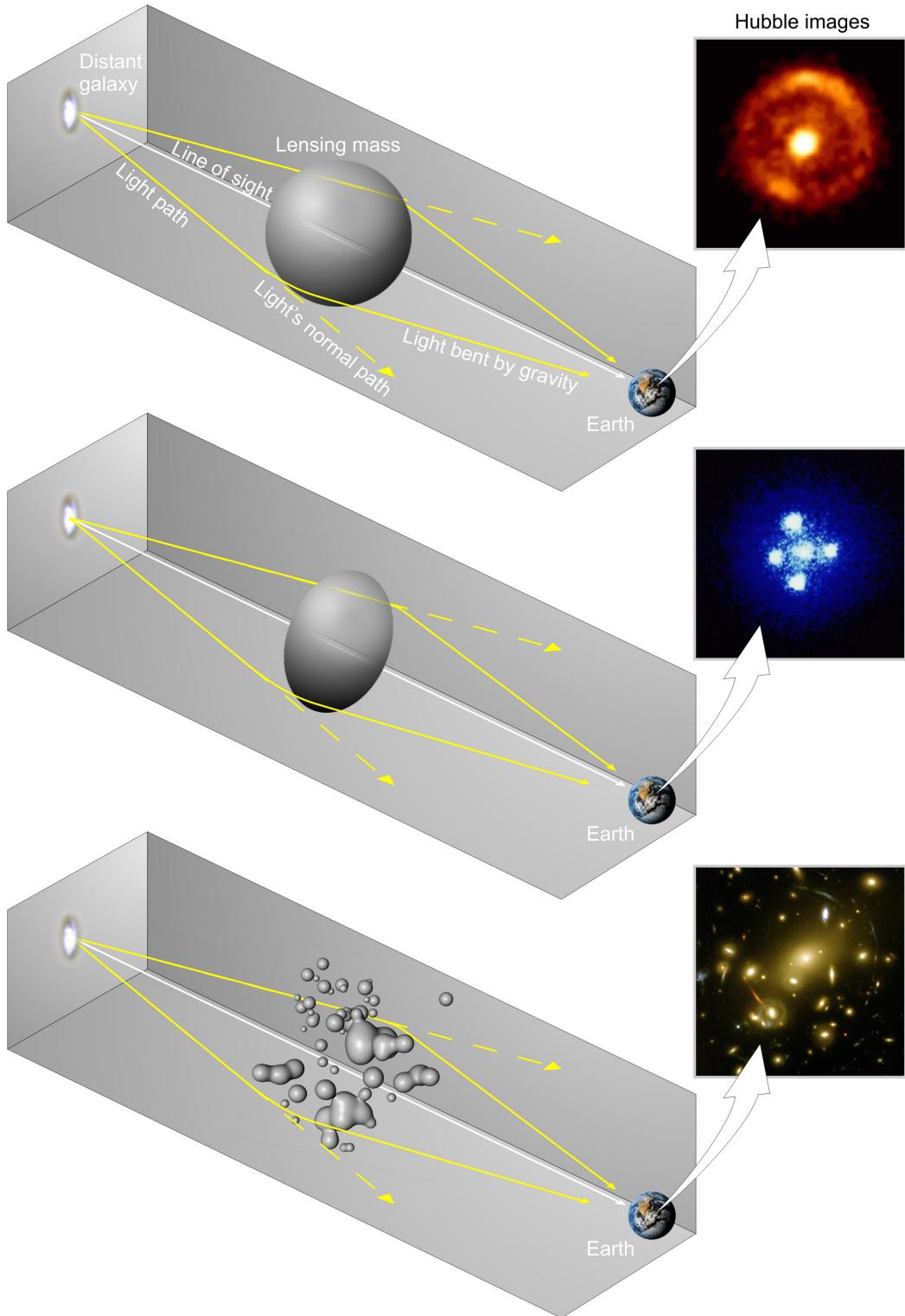


Figure 2.2 Light from distant galaxy is bent in unique ways depending on the distribution of mass between the galaxy and Earth. Yellow dashed lines indicate where the light would have gone if the matter were not present [8].

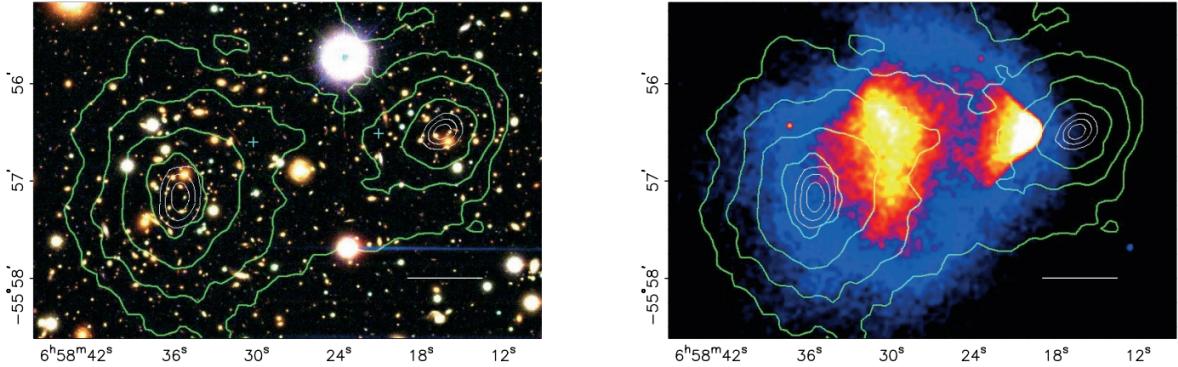


Figure 2.3 (left) Optical image of galactic cluster 1E0657-558. (right) X-ray image of the cluster with redder meaning hotter and higher baryon density. (both) Green contours are reconstruction of gravity contours from weak lensing. White rings are the best fit mass maxima at 68.3%, 95.5%, and 99.7% confidence. The matter maxima of the clusters are clearly separated from x-ray maxima. [9]

529 light, see Figure 2.2. With a sufficient understanding of light sources behind a massive object, we  
 530 can reconstruct the contours of the gravitational lenses. The gradient of the contours shown in  
 531 Figure 2.3 then indicates how dense the matter is and where it is.

532 The x-ray emission can then be observed from the clusters. Since these galaxies are mostly gas  
 533 and are merging, then the gas should be getting hotter. If they are merging, the x-ray emissions  
 534 should be the strongest where the gas is mostly moving through each other. Hence, X-ray emission  
 535 maps out where the gas is in the merging galaxy cluster.

536 The lensing and x-ray observations were done on the Bullet cluster featured on Figure 2.3.  
 537 The x-ray emissions do not align with the gravitational contours from lensing. The incongruence  
 538 in mass density and baryon density suggests that there is a lot of matter somewhere that does  
 539 not interact with light. Moreover, this dark matter cannot be baryonic [9]. The Bullet Cluster  
 540 measurement did not really tell us what DM is exactly, but it did give the clue that DM also does  
 541 not interact with itself very strongly. If DM did interact strongly with itself, then it would have been  
 542 more aligned with the x-ray emission [9]. There have been follow-up studies of galaxy clusters with  
 543 similar results. The Bullet Cluster and others like it provide a persuasive case against something  
 544 possibly amiss in our gravitational theories.

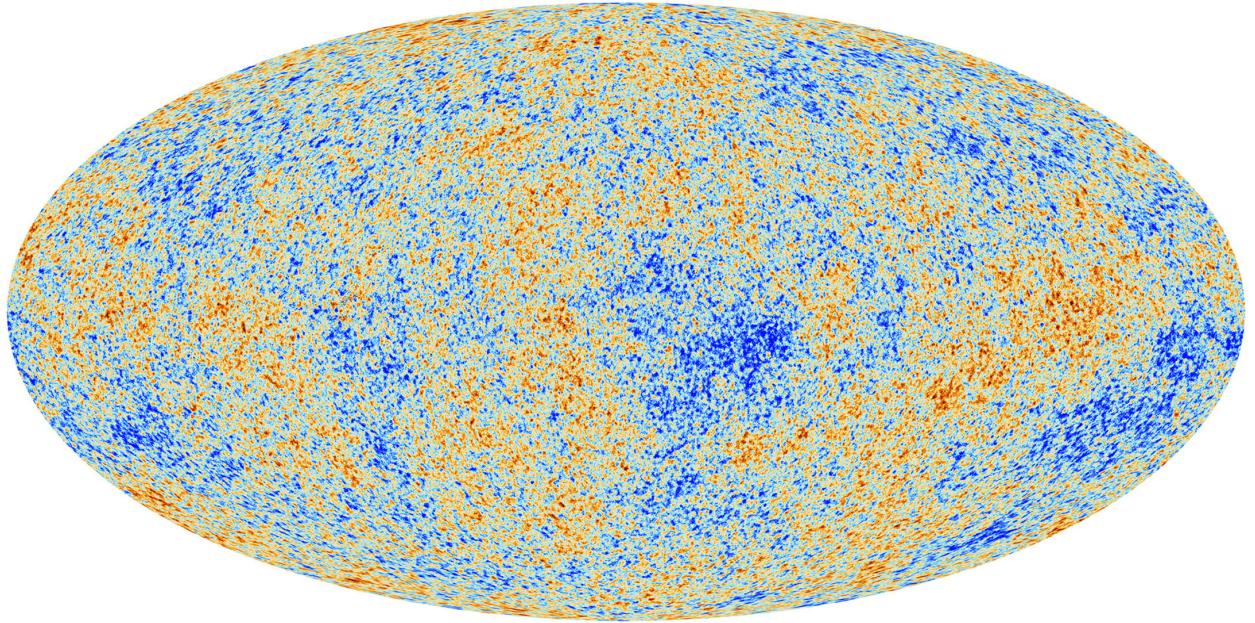


Figure 2.4 Plank CMB sky. Sky map features small variations in temperature in primordial light. These anisotropies are used to make inferences about the universe's energy budget and developmental history. [10]

545 **2.3.3 Evidence for Dark Matter: Cosmic Microwave Background**

546 The Cosmic Microwave Background (CMB) is the primordial light from the early universe  
547 when Hydrogen atoms formed from the free electron and proton soup in the early universe. The  
548 CMB is the earliest light we can observe; released when the universe was about 380,000 years old.  
549 Then we look at how the simulated universes look like compared to what we see. Figure 2.4 is the  
550 most recent CMB image from the Plank satellite after subtracting the average value and masking the  
551 galactic plane [10]. Redder regions indicate a slightly hotter region in the CMB, and blue indicates  
552 colder. The intensity variations are on the order of 1 in 1000 with respect to the average value.

553 The Cosmic Microwave Background shows that the universe had DM in it from an incredibly  
554 early stage. To measure the DM, Dark Energy, and matter fractions of the universe from the CMB,  
555 the image is analyzed into a power spectrum, which shows the amplitude of the fluctuations as  
556 a function of spherical multipole moments.  $\Lambda$ CDM provides the best fit to the power spectra of  
557 the CMB as shown in Figure 2.5. The CMB power spectrum is quite sensitive to the fraction  
558 of each energy contribution in the early universe. Low  $l$  modes are dominated by variations  
559 in gravitational potential. Intermediate  $l$  emerge from oscillations in photon-baryon fluid from

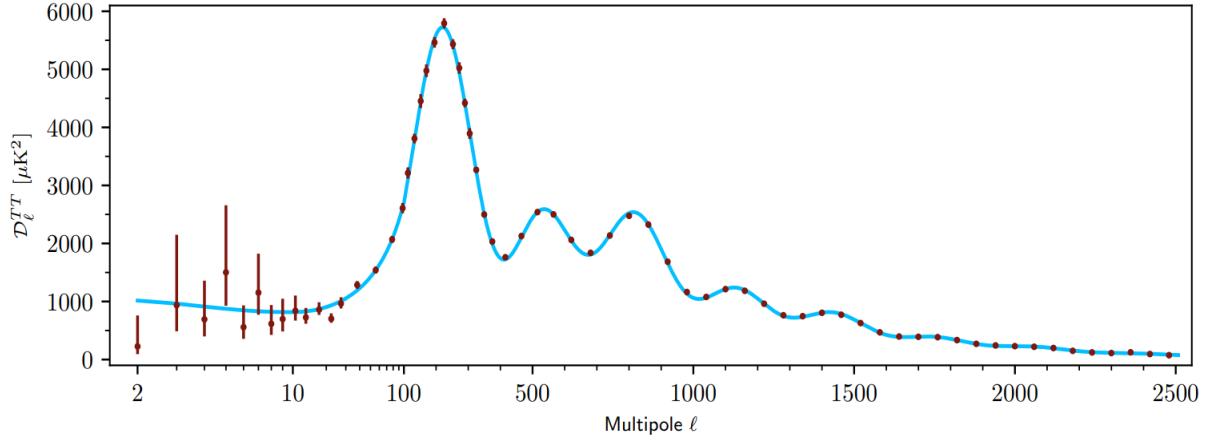


Figure 2.5 Observed Cosmic Microwave Background power spectrum as a function of multipole moment from Plank [10]. Blue line is best fit model from  $\Lambda$ CDM. Red points and lines are data and error, respectively.

560 competing baryon pressures and gravity. High  $l$  is a damped region from the diffusion of photons  
 561 during electron-proton recombination. [1]

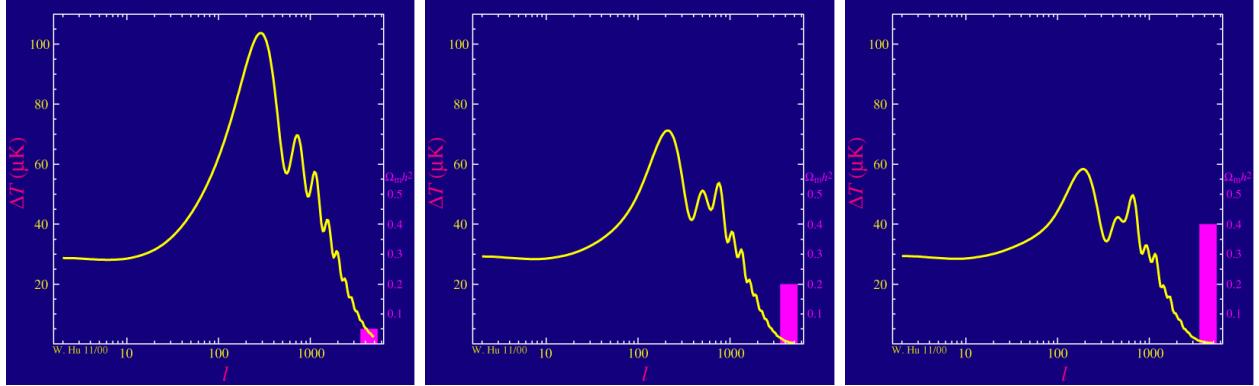


Figure 2.6 Predicted power spectra of CMB for different  $\Omega_m h^2$  values for fixed baryon density from [11]. (left) Low  $\Omega_m h^2$  increases the prominence of first and second peaks. (middle)  $\Omega_m h^2$  is most similar to the observed power spectrum. The second and third peaks are similar in height. (right)  $\Omega_m h^2$  is large which suppresses the first peak and raises the prominence of the third peak.

562 The harmonics would look quite different for a universe with less DM. Figure 2.6 demonstrates  
 563 the effect  $\Omega_m h^2$  has on the expected power spectrum for fixed baryon matter density. [11] Sweeping  
 564  $\Omega_m h^2$  in this way clearly shows the effect dark matter has on the CMB power spectrum. The  
 565 observations fit well with the  $\Lambda$ CDM model, and the derived fractions are as follows. The matter  
 566 fraction:  $\Omega_m = 0.3153$ ; and the baryon fraction:  $\Omega_b = 0.04936$  [10]. Plank's observations also  
 567 provide a measure of the Hubble constant,  $H_0$ .  $H_0$  especially has seen a growing tension in the

568 past decade that continues to deepened with observations from instruments like the James Webb  
569 Telescope [12, 13]. As Hubble tensions deepen, we may find that perhaps  $\Lambda$ **CDM**, despite its  
570 successes, is missing some critical physics.

571 Overall, the Newtonian motion of stars in galaxies, weak lensing from galactic clusters, and  
572 power spectra from primordial light form a compelling body of research in favor of dark matter.  
573 It takes another leap of theory and experimentation to make observations of DM that are non-  
574 gravitational in nature. In Section 2.3, the evidence for DM implies strongly that the DM is matter  
575 and not a lost parameter in the gravitational fields between massive objects. Finally, if we take one  
576 axiom: that this matter has quantum behavior, such as being described by some Bohr wavelength  
577 and abiding by some fermion or boson statistics; then we arrive at particle dark matter. One particle  
578 DM hypothesis is the Weakly Interacting Massive Particle (WIMP). This DM candidate theory is  
579 discussed further in the next section and is the focus of this thesis.

## 580 2.4 Searching for Dark Matter: Particle DM

581 Section 2.4 shows the Standard Model of particle physics and is currently the most accurate  
582 model for the dynamics of fundamental particles like electrons and photons. The current status  
583 of the SM does not have a viable DM candidate. When looking at the standard model, we can  
584 immediately exclude any charged particle because charged particles interact strongly with light.  
585 Specifically, this will rule out the following charged, fundamental particles:  $e, \mu, \tau, W, u, d, s, c, t, b$   
586 and their corresponding antiparticles. Recalling from Section 2.2 that DM must be long-lived and  
587 stable over the age of the universe which excludes all SM particles with decay half-lives at or shorter  
588 than the age of the universe. The lifetime constraint additionally eliminates the  $Z$  and  $H$  bosons.  
589 Finally, the candidate DM needs to be somewhat massive. Recall from Section 2.2 that DM is cold  
590 or not relativistic through the universe. This eliminates the remaining SM particles:  $\nu_{e,\mu,\tau}, g, \gamma$  as  
591 DM candidates. Because there are no DM candidates within the SM, the DM problem strongly  
592 hints to physics beyond the SM (BSM).

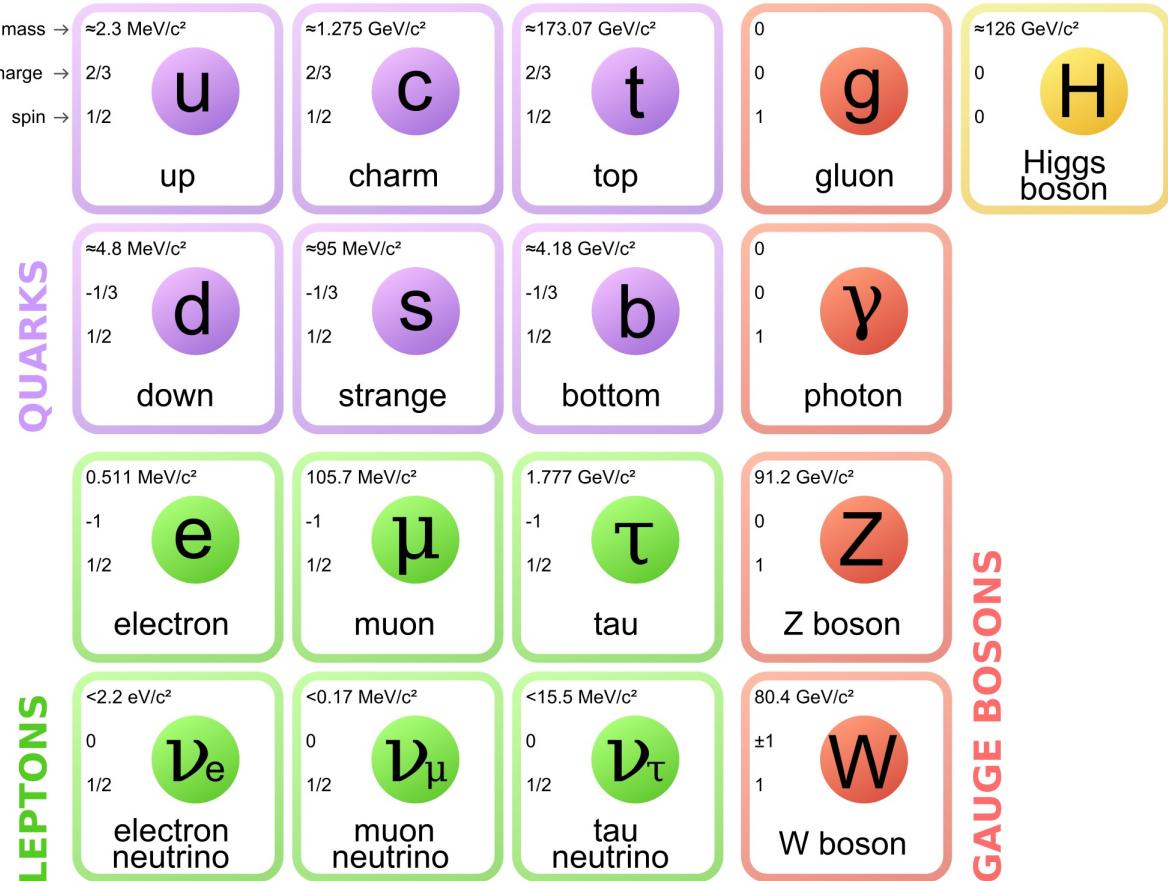


Figure 2.7 The Standard Model (SM) of particle physics. Figure taken from <http://www.quantumdiaries.org/2014/03/14/the-standard-model-a-beautiful-but-flawed-theory/>

### 593 2.4.1 Shake it, Break it, Make it

594 When considering DM that couples in some way with the SM, the interactions are roughly  
 595 demonstrated by interaction demonstrated in Figure 2.8. The figure is a simplified Feynman  
 596 diagram where the arrow of time represents the interaction modes of: **Shake it, Break it, Make it**.  
 597 **Shake it** refers to the direct detection of dark matter. Direct detection interactions start with  
 598 a free DM particle and an SM particle. The DM and SM interact via elastic or inelastic collision  
 599 and recoil away from each other. The DM remains in the dark sector and imparts some momentum  
 600 onto the SM particle. The hope is that the momentum imparted onto the SM particle is sufficiently  
 601 high enough to pick up with extremely sensitive instruments. Because we cannot create the DM in  
 602 the lab, a direct detection experiment must wait until DM is incident on the detector. Most direct  
 603 detection experiments are therefore placed in low-background environments with inert detection

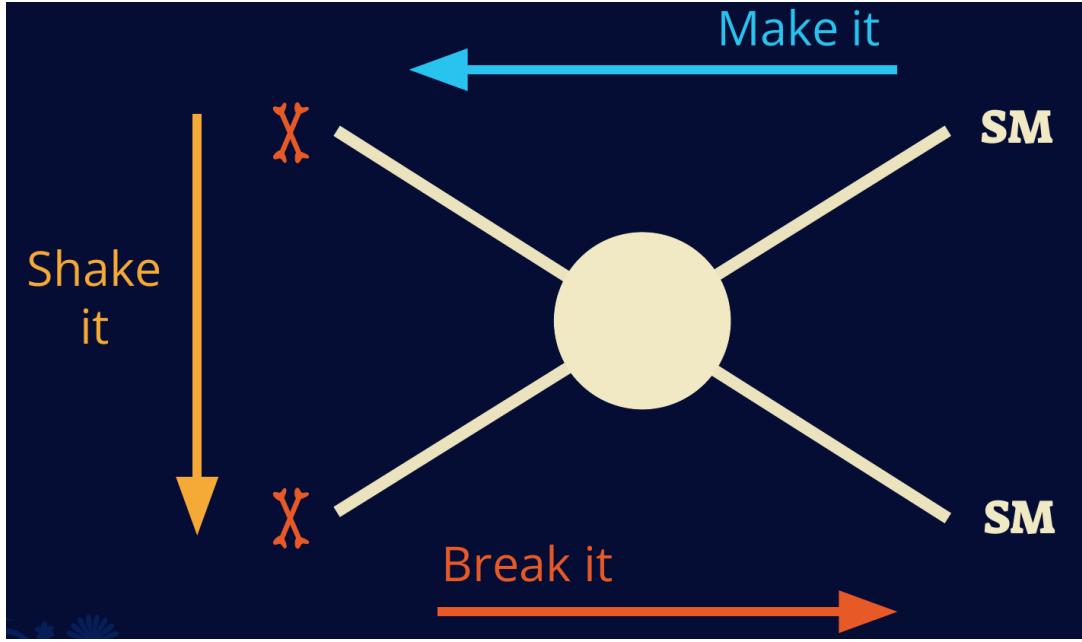


Figure 2.8 Simplified Feynman diagram demonstrating different ways DM can interact with SM particles. The 'X's refer to the DM particles whereas the SM refer to fundamental particles in the SM. The large circle in the center indicates the vertex of interaction and is purposely left vague. The colored arrows refer to different directions of time as well as their respective labels. The arrows indicate the initial and final state of the DM -SM interaction in time.

604 media like the noble gas Xenon. [14]

605 **Make it** refers to the production of DM from SM initial states. The experiment starts with  
 606 particles in the SM. These SM particles are accelerated to incredibly high energies and then collide  
 607 with each other. In the confluence of energy, DM hopefully emerges as a byproduct of the SM  
 608 annihilation. Often it is the collider experiments that are energetic enough to hypothetically produce  
 609 DM. These experiments include the world-wide collaborations ATLAS and CMS at CERN where  
 610 proton collide together at extreme energies. The DM searches, however, are complex. DM likely  
 611 does not interact with the detectors and lives long enough to escape the detection apparatus of  
 612 CERN's colliders. This means any DM production experiment searches for an excess of events  
 613 with missing momentum or energy in the events. An example event with missing transverse  
 614 momentum is shown in Figure 2.9. The missing momentum with no particle tracks implies a  
 615 neutral particle carried the energy out of the detector. However, there are other neutral particles  
 616 in the SM, like neutrons or neutrinos, so any analysis has to account for SM signatures of missing

617 momentum. [15]

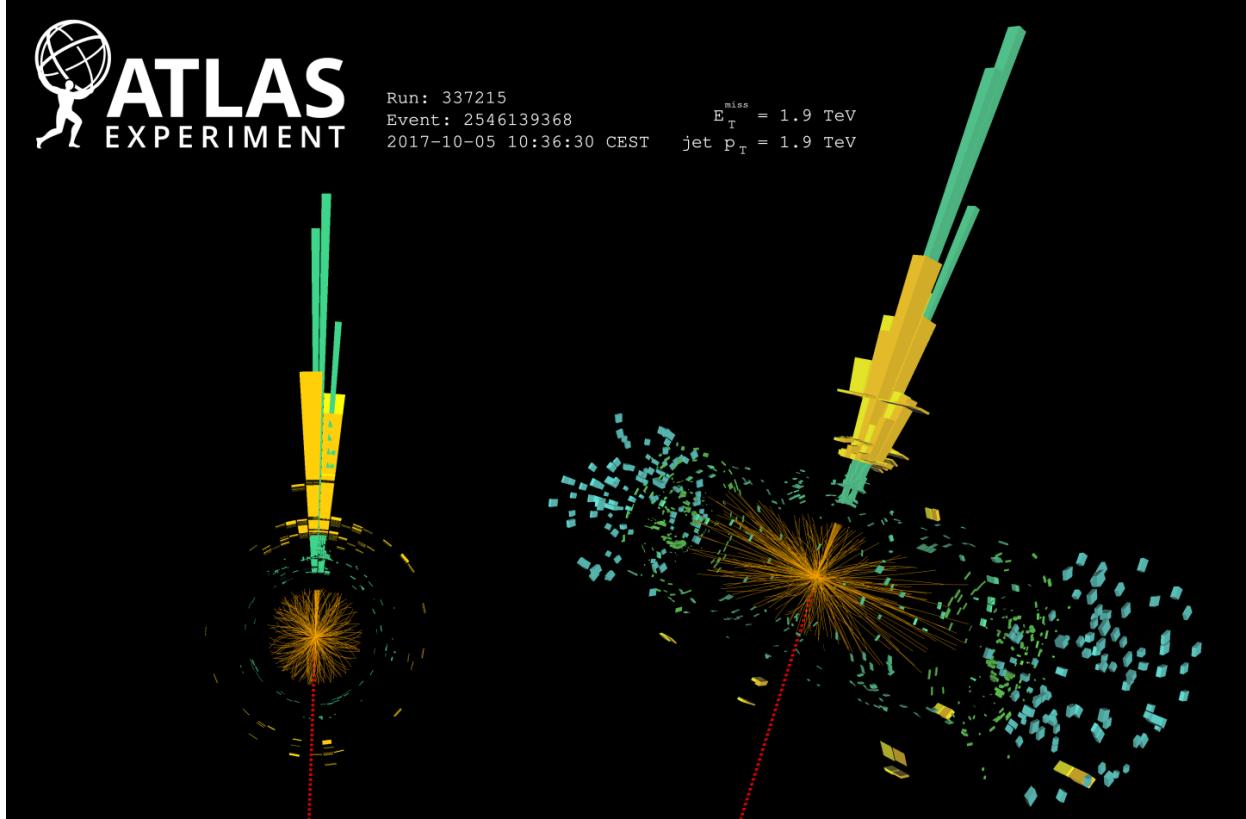


Figure 2.9 A single jet event in ATLAS detector from 2017 [16]. Jet momentum was observed to be 1.9 TeV. Missing transverse momentum was observed to be 1.9 TeV compared to the initial transverse momentum of the event was 0. Implied MET is traced by a red dashed line in event display.

#### 618 2.4.2 Break it: Standard Model Signatures of Dark Matter through Indirect Searches

619 **Break it** refers to the creation of SM particles from the dark sector, and it is the primary focus  
620 of this thesis. The interaction begins with DM or in the dark sector. The hypothesis is that this  
621 DM will either annihilate with itself or decay and produce an SM byproduct. This method is  
622 often referred to as the Indirect Detection of DM because we have no lab to directly control or  
623 manipulate the DM. Therefore, most indirect DM searches are performed using observations of  
624 known DM densities among the astrophysical sources. The strength is that we have the whole of the  
625 universe and its 13.6-billion-year lifespan to use as a detector and particle accelerator. Additionally,  
626 locations of dark matter are well cataloged since it was astrophysical observations that presented

627 the problem of DM in the first place.

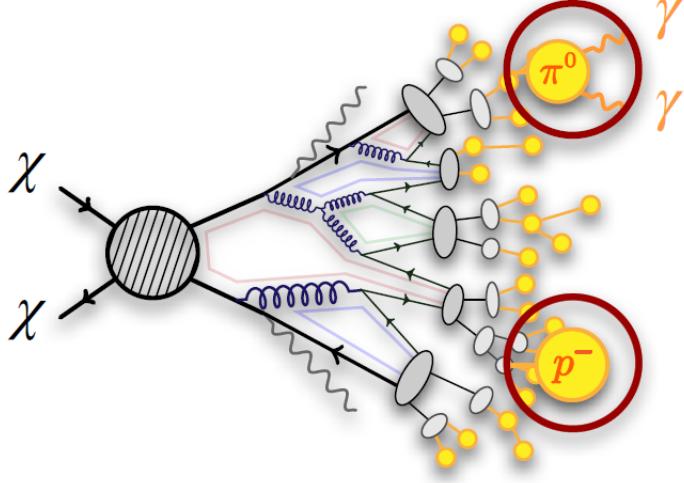


Figure 2.10 More detailed pseudo-Feynman diagram of particle cascade from dark matter annihilation into 2 quarks. The quarks hadronize and down to stable particles like  $\gamma$  or the anti-proton ( $p^-$ ). Diagram pulled from ICRC 2021 presentation on DM annihilation search [17].

628 However, anything can happen in the universe. There are many difficult to deconvolve back-  
629 grounds when searching for DM. One prominent example is the galactic center. We know the  
630 galactic center has a large DM content because of stellar kinematics in our Milky Way and DM halo  
631 simulations. Yet, any signal from the galactic center is challenging to parse apart from the extreme  
632 environment of our supermassive black hole, unresolved sources, and diffuse emission from the  
633 interstellar medium [18]. Despite the challenges, any DM model that yields evidence in the other  
634 two observation methods, **Shake it or Make it** must be corroborated with indirect observations of  
635 the known DM sources. Without corroborating evidence, DM observation in the lab is hard-pressed  
636 to demonstrate that it is the model contributing to the DM seen at the universal scale.

637 In the case of WIMP DM, signals are described in terms of primary SM particles produced  
638 from DM decay or annihilation. The SM initial state particles are then simulated down to stable  
639 final states such as the  $\gamma$ ,  $\nu$ ,  $p$ , or  $e$  which can traverse galactic lengths to reach Earth.

640 Figure 2.10 shows the quagmire of SM particles that emerges from SM initial states that are not  
641 stable [17]. There are many SM particles with varying energies that can be produced in such an

642 interaction. For any arbitrary DM source and stable SM particle, the SM flux from DM annihilating  
 643 to a neutral particle in the SM,  $\phi$ , from a region in the sky is described by the following.

$$\frac{d\Phi_\phi}{dE_\phi} = \frac{\langle\sigma v\rangle}{8\pi m_\chi^2} \frac{dN_\phi}{dE_\phi} \times \int_{\text{source}} d\Omega \int_{l.o.s} \rho_\chi^2 dl(r, \theta') \quad (2.4)$$

644 In Equation (7.1),  $\langle\sigma v\rangle$  is the velocity-weighted annihilation cross-section of DM to the SM.  $m_\chi$   
 645 refers to the mass of DM, noted with Greek letter  $\chi$ .  $\frac{dN_\phi}{dE_\phi}$  is the N particle flux weighted by the  
 646 particle energy. An example is provided in Figure 2.11 for the  $\gamma$  final state. The integrated terms  
 647 are performed over the solid angle,  $d\Omega$ , and line of sight, l.o.s.  $\rho$  is the density of DM for a  
 648 location  $(r, \theta')$  in the sky. The terms left of the ' $\times$ ' are often referred to as the particle physics  
 649 component. The terms on the right are referred to as the astrophysical component. For decaying  
 650 DM, the equation changes to...

$$\frac{d\Phi_\phi}{dE_\phi} = \frac{1}{4\pi\tau m_\chi} \frac{dN_\phi}{dE_\phi} \times \int_{\text{source}} d\Omega \int_{l.o.s} \rho_\chi dl(r, \theta') \quad (2.5)$$

651 In Equation (2.5),  $\tau$  is the decay lifetime of the DM. Just as in Equation (7.1), the left and right  
 652 terms are the particle physics and the astrophysical components respectively. The integrated  
 653 astrophysical component of Equation (7.1) is often called the J-Factor. Whereas the integrated  
 654 astrophysical component of Equation (2.5) is often called the D-Factor.

655     Exact DM  $\text{DM} \rightarrow \text{SM SM}$  branching ratios are not known, so it is usually assumed to go 100%  
 656 into an SM particle/anti-particle. Additionally, when a DM annihilation or decay produces one of  
 657 the neutral, long-lived SM particles ( $\nu$  or  $\gamma$ ), the particle is traced back to a DM source. For DM  
 658 above GeV energies, there are very few SM processes that can produce particles with such a high  
 659 energy. Seeing such a signal would almost certainly be an indication of the presence of dark matter.  
 660 Fortunately, the universe provides us with the largest volume and lifetime ever for a particle physics  
 661 experiment.

## 662 2.5 Sources for Indirect Dark Matter Searches

663     The first detection of DM relied on optical observations. Since then, we have developed new  
 664 techniques to find DM dense regions. As described in Section 2.3.1, many DM dense regions were  
 665 through observing galactic rotation curves. Our Milky Way galaxy is among DM dense regions

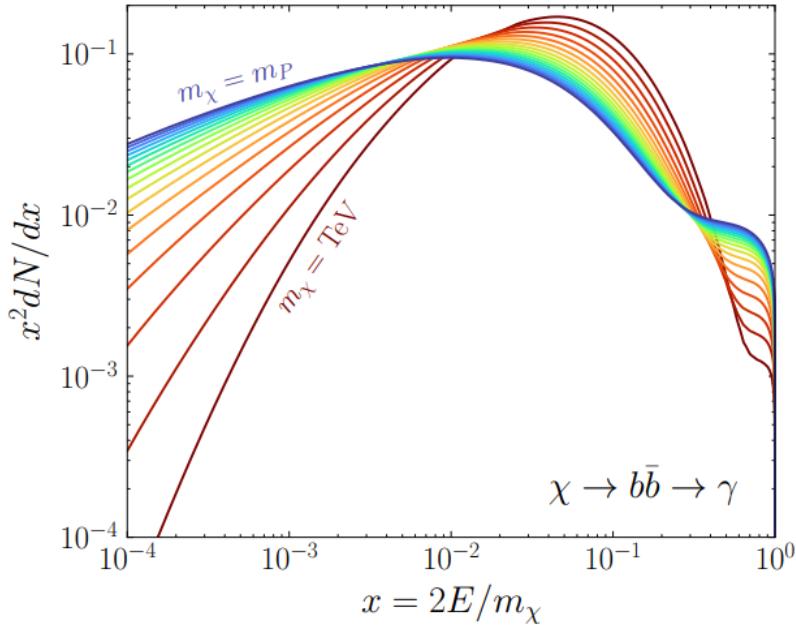


Figure 2.11 Dark Matter (DM) decay spectrum for  $b\bar{b}$  initial state and  $\gamma$  final state. Redder spectra are for larger DM masses. Bluer spectra are light DM masses.  $x$  is a unitless factor defined as the ratio of the mass of DM,  $m_\chi$ , and the final state particle energy  $E_\gamma$ . Figure from [19].

discovered, and it is the largest nearby DM dense region to look at. Additionally, the DM halo surrounding the Milky Way is clumpy [18]. There are regions in the DM halo of the Milky Way that have more DM than others that have captured gas over time. These sub-halos were dense enough collapse gas and form stars. These apparent sub galaxies are known as dwarf spheroidal galaxies and are the main sources studied in this thesis. Each source type comes with different trade-offs. Galactic Center studies will be very sensitive to the assumed distribution of DM. The central DM density can vary substantially as demonstrated in Figure 2.12. At distances close to the center of the galaxy, or small  $r$ , the differences in DM densities can be 3-4 orders of magnitude. Searches toward the galactic center will therefore be quite sensitive to the assumed DM distribution.

Searches toward Dwarf Spheroidal Galaxies (dSph's) suffer from uncertainties in the DM density less than the galactic center studies. This is mostly from their diminutive size being smaller than the angular resolution of most  $\gamma$ -ray observatories [18]. The DM content dSph's are typically determined with the Virial theorem, Equation (2.1), and are usually majority DM [18] in mass. DSph's tend to be ideal sources to look at for DM searches. Their environments are quiet with little

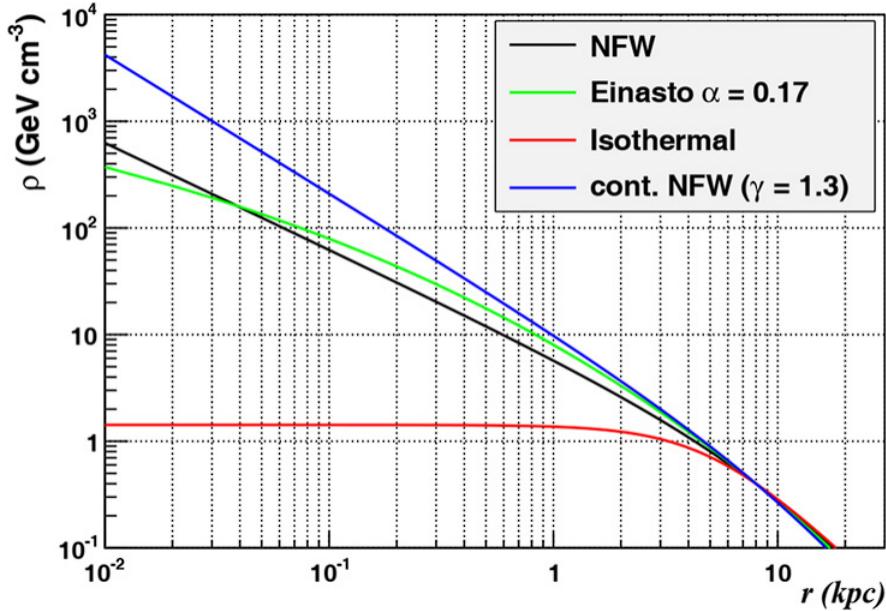


Figure 2.12 Different dark matter density profiles compared. Some models produce exceptionally large densities at small  $r$  [20].

680 astrophysical background. Unlike the galactic center, the most active components of dSph's are the  
 681 stars within them versus a violent accretion disc around a black hole. All this together means that  
 682 dSph's are among the best sources to look at for indirect DM searches. dSph's are the targets of  
 683 focus for this thesis.

## 684 2.6 Multi-Messenger Dark Matter

685 Astrophysics entered a new phase in the past few decades that leverages our increasing sensitivity  
 686 to SM channels and general relativity (GR). Up until the 21st century, astrophysical observations  
 687 were performed with photons ( $\gamma$ ) only. Astrophysics with this 'messenger' is fairly mature now.  
 688 Novel observations of the universe have since only adjusted the sensitivity of the wavelength of  
 689 light that is observed except at MeV energies. Gems like the CMB [10], and more have ultimately  
 690 been observations of different wavelengths of light. Multi-messenger astrophysics proposes using  
 691 other SM particles such the  $p^{+-}$ , or  $\nu$  or gravitation waves predicted by general relativity.

692 The experiment LIGO had a revolutionary discovery in 2016 with the first detection of a binary  
 693 black hole merger [21]. This opened the collective imagination to observing the universe through  
 694 gravitational waves. There has also been a surge of interest in the neutrino ( $\nu$ ) sector. IceCube

695 demonstrated that we are sensitive to neutrinos in regions that correlate with significant photon  
 696 emission like the galactic plane [22]. Neutrinos, like gravitational waves and light, travel mostly  
 697 unimpeded from their source to our observatories. This makes pointing to the originating source  
 698 of these messengers much easier than it is for cosmic rays which are deflected from their source by  
 699 magnetic fields.

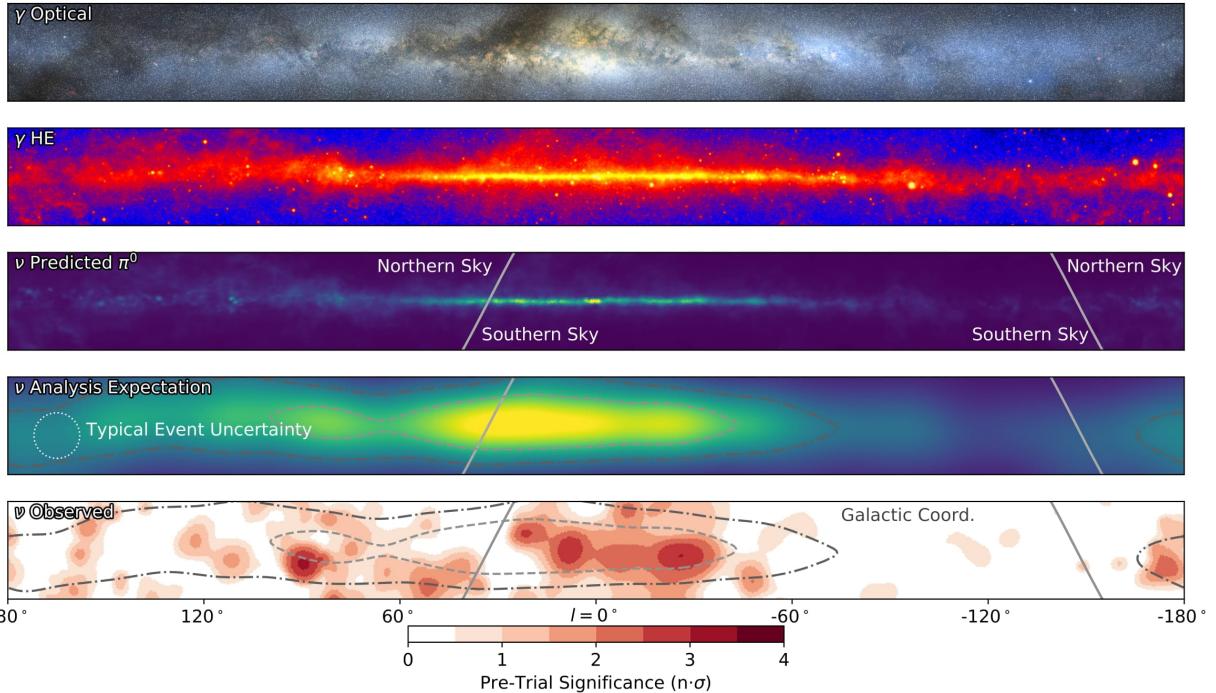


Figure 2.13 The Milky Way Galaxy in photons ( $\gamma$ ) and neutrinos ( $\nu$ ) [22]. The Galactic center is at  $l=0^\circ$  and is the brightest region in all panels. (top) An Optical color image of the Milky Way galaxy seen from Earth. Clouds of gas and dust obscure some light from stars. (2nd down) Integrated flux of  $\gamma$ -rays observed by the Fermi-LAT telescope [23]. (middle) Expected neutrino emission that corresponds with Fermi-LAT observations. (2nd up) Expected neutrino emission profile after considering detector systematics of IceCube. (bottom) Observed neutrino emission from region of the galactic plane. Substantial neutrino emission was detected.

700 The IceCube collaboration recently published a groundbreaking result of the Milky Way in  
 701 neutrinos. The recent result from IceCube, shown in Figure 2.13, proves that we can make  
 702 observations under different messenger regimes. The top two panels show the appearance of the  
 703 galactic plane to different wavelengths of light. Some sources are more apparent in some panels,  
 704 while others are not. This new channel is powerful because neutrinos are readily able to penetrate  
 705 through gas and dust in the Milky Way. This new image also refines our understanding of how high

706 energy particles are produced. For example, the fit to IceCube data prefers neutrino production  
 707 from the decay of  $\pi^0$  [22].

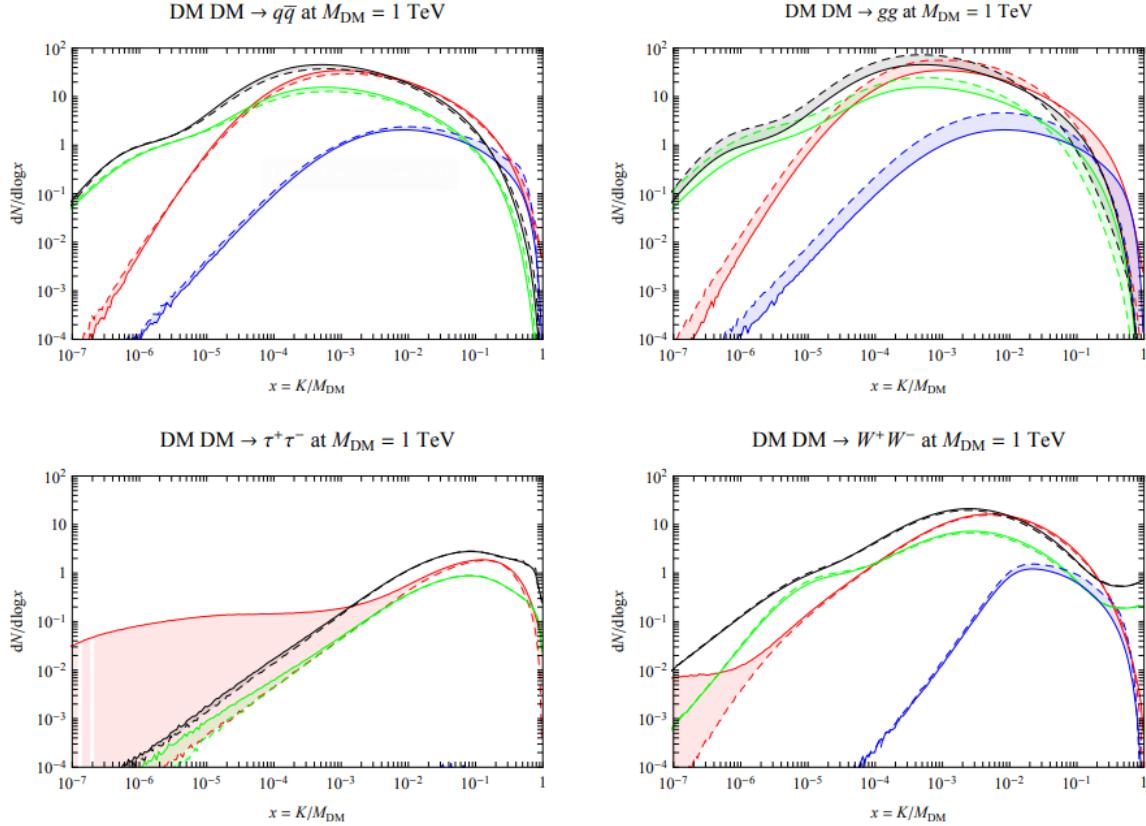


Figure 2.14 Dark Matter annihilation spectra for different final state particle and standard model annihilation channels [24]. Photons (red),  $e^\pm$  (green),  $\bar{p}$  (blue),  $\nu$  (black).

708 Exposing our observations to more cosmic messengers greatly increases our sensitivity to  
 709 rare processes. In the case of DM, Figure 2.14, there are many SM particles produced in DM  
 710 annihilation. Among the final state fluxes are gammas and neutrinos. Charged particles are also  
 711 produced however they would not likely make it to Earth since they will be deflected by magnetic  
 712 fields between the source and Earth. This means observatories that can see the neutral messengers  
 713 are especially good for DM searches and for combining data for a multi-messenger DM search.

## CHAPTER 3

714           **HIGH ALTITUDE WATER CHERENKOV (HAWC) OBSERVATORY**

715   **3.1 The Detector**

716   **3.2 Events Reconstruction and Data Acquisition**

717   **3.2.1 G/H Discrimination**

718   **3.2.2 Angle**

719   **3.2.3 Energy**

720   **3.3 Remote Monitoring**

721   **3.3.1 ATHENA Database**

722   **3.3.2 HOMER**

723

## CHAPTER 4

### ICECUBE NEUTRINO OBSERVATORY

724 **4.1 The Detector**

725 **4.2 Events Reconstruction and Data Acquisition**

726 **4.2.1 Angle**

727 **4.2.2 Energy**

728 **4.3 Northern Test Site**

729 **4.3.1 PIgeon remote dark rate testing**

730 **4.3.2 Bulkhead Construction**

## CHAPTER 5

### GLORY DUCK: MULTI-WAVELENGTH SEARCH FOR DARK MATTTER ANNIHILATION TOWARDS DWARF SPHEROIDAL GALAXIES

#### 5.1 Introduction

The field of astrophysics now has several instruments and observatories sensitive to high energy gamma-rays. The energy sensitivity for the modern gamma-ray program spans many orders of magnitude. Figure 5.1 demonstrates these comparable sensitivities across energies for the five experiments: Fermi-LAT, HAWC, HESS, MAGIC, and VERITAS.

Each of the five experiments featured in Figure 5.1 have independently searched for DM annihilation from dwarf spheroidal galaxies (dSph) and set limits on annihilation cross-section of WIMPs. Intriguingly, there are regions of substantial overlap in their energy sensitivities. This clearly motivates an analysis that combines data from these five. Each experiment has unique gamma-ray detection methods and their weaknesses and strengths can be leveraged with each other. The HAWC gamma-ray observatory is extensively introduced in chapter 3, so it is not introduced here. A brief description of the remaining experiments are in the following paragraphs.

The Large Area Telescope (LAT) is a pair conversion telescope mounted on the NASA Fermi satellite in orbit  $\sim$ 550 km above the Earth [26]. LAT's field of view covers about 20% of the whole sky, and it sweeps the whole sky approximately every 3 hours. LAT's gamma-ray energy sensitivity ranges from 20 MeV up to 1 TeV. Previous DM searches towards dSphs using Fermi-LAT are published in [27] and [28].

The High Energy Spectroscopic System (HESS), Major Atmospheric Gamma Imaging Cherenkov (MAGIC), and Very Energetic Radiation Imaging Telescope Array System (VERITAS) are arrays of Imaging Atmospheric Cherenkov Telescopes (IACT). These telescopes observe the Cherenkov light emitted from gamma-ray showers in the Earth's atmosphere. The field of view for these telescopes is no larger than  $5^\circ$  with energy sensitivities ranging from 30 GeV up to 100 TeV [29, 30, 31]. IACTs are able to make precise observations in selected regions of the sky, however can only be operated in ideal dark conditions. HESS's observations of the dwarves

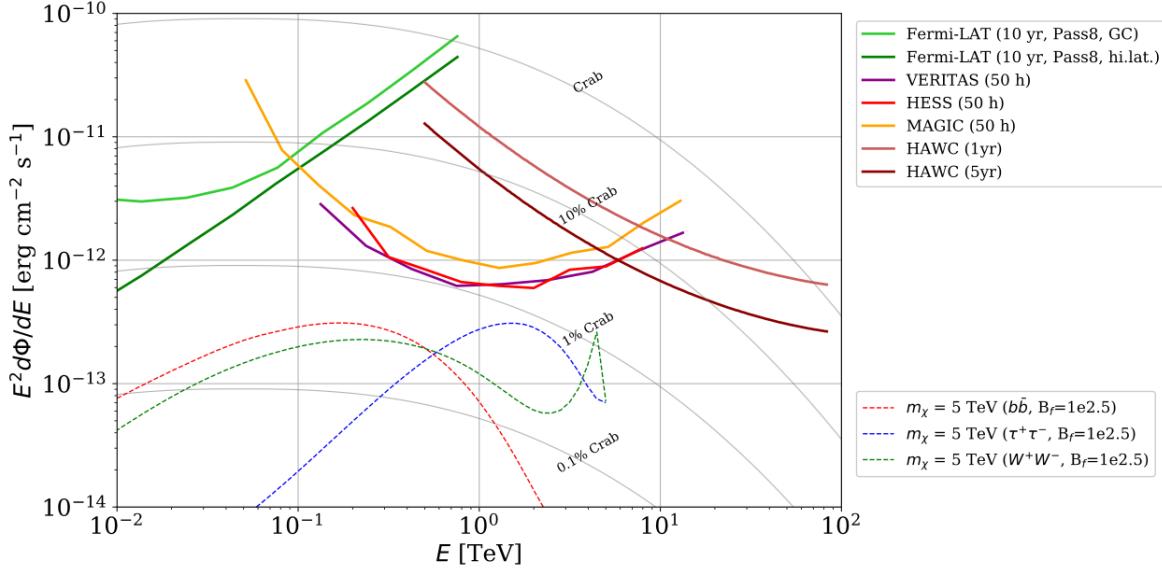


Figure 5.1 Sensitivities of five gamma-ray experiments compared to percentages of the Crab nebula's emission and dark matter annihilation. Solid lines present estimated sensitivities to power law spectra [FACT CHECK THIS] for each experiment. Green lines are Fermi-LAT sensitivities where lighter green is the sensitivity to the galactic center and dark green is its sensitivity to higher declinations. Orange, red, and purple solid lines represent the MAGIC, HESS, and VERITAS 50 hour sensitivities respectively. The maroon and brown lines are the HAWC 1 year and 5 year sensitivities. Across four decades of gamma-ray energy, these experiments have similar sensitivities on the order  $10^{-12}$  erg  $\text{cm}^{-2}\text{s}^{-1}$ . The dotted lines are estimated dark matter fluxes assuming  $m_\chi = 5$  TeV DM annihilating to bottom quarks (red), tau leptons (blue), and W bosons (green). Faded gray lines outline percentage flux of the Crab nebula. Figure is an augmented version of [25]

757 Sculptor and Carina were between January 2008 and December 2009. HESS's observations of  
 758 Coma Berenices were taken from 2010 to 2013, and Fornax was observed in 2010 [32, 33, 34].  
 759 MAGIC provided deep observations of Segue1 between 2011 and 2013 [35]. MAGIC also provides  
 760 data for three additional dwarves: Coma Berenices, Draco, and Ursa Major II where observations  
 761 were made in: January - June 2019 [36], March - September 2018 [36], and 2014 - 2016 [37]  
 762 respectively. VERITAS provided data for Boötes I, Draco, Segue 1, and Ursa Minor from 2009 to  
 763 2016 [38].

764 This chapter presents the Glory Duck analysis, the name given for the search for dark matter  
 765 annihilation from dSph by combining data from the five gamma-ray observatories: Fermi-LAT,  
 766 HAWC, HESS, MAGIC, and VERITAS. Specifically, the methods in analysis and modeling are  
 767 presented for the HAWC gamma-ray observatory. This work will be published in the Journal of

768 Cosmology and Astroparticle Physics and presented at the International Cosmic Ray Conference  
769 in 2019, 2021, and 2023 [39, 40, 41] and others.

770 **5.2 Dataset and Background**

771 This section enumerates the data analysis and background estimation methods used for HAWC's  
772 study of dSphs. Section 5.2.1 and Section 5.2.2 are most useful for fellow HAWC collaborators  
773 looking to replicate the Glory Duck analysis.

774 **5.2.1 Itemized HAWC files**

775 These files are only available withing HAWC's internal documentation and collaborators. They  
776 are not meant for public access, and are presented here so that HAWC collaborators can reproduce  
777 results accurately.

- 778 • Detector Response: `response_aerie_svn_27754_systematics_best_mc_test_noBr`  
779     `oadpulse\10pctlogchargesmearing_0.63qe_25kHzNoise_run5481_curvature`  
780     `0_index3.root`
- 781 • Data Map: `maps-20180119/liff/maptree_1024.root`
- 782 • Spectral Dictionary: `DM_CirrelliSpectrum_dict_gammas.npy`
- 783 • Analysis wiki: [https://private.hawc-observatory.org/wiki/index.php/Glory\\_Duck\\_Multi-Experiment\\_Dark\\_Matter\\_Search](https://private.hawc-observatory.org/wiki/index.php/Glory_Duck_Multi-Experiment_Dark_Matter_Search)

785 **5.2.2 Software Tools and Development**

786 This analysis was performed using HAL and 3ML [42, 43] in Python version 2. I built software  
787 to implement the *A Poor Particle Physicist Cookbook for Dark Matter Indirect Detection* (PPPC)  
788 [44] DM spectral model and dSphs spatial model from [45] for HAWC analysis. A NumPy version  
789 of this dictionary was made for both Py2 and Py3. The code base for creating this dictionary is  
790 linked on my GitLab sandbox:

- 791 • Py2: [Dictionary Generator \(Deprecated\)](#)

- 792     • Py3: [PPPC2Dict](#)

793     The analysis was performed using the  $f_{hit}$  framework as used and described in the HAWC Crab  
794     paper [42]. The Python2 NumPy dictionary file for gamma-ray final states is `dmCirSpecDict.npy`.  
795     The corresponding Python3 file is `DM_CirrelliSpectrum_dict_gammas.npy`. These files can  
796     also be used for decay channels and the PPPC describes how [44]. All other software used for data  
797     analysis, DM profile generation, and job submission to SLURM are also kept in my sandbox for  
798     [the Glory Duck](#) project.

799     **5.2.3 Data Set and Background Description**

800     The HAWC data maps used for this analysis contain 1017 days of data between runs 2104  
801     (2014-11-26) and 7476 (2017-12-20). They were generated from pass 4.0 reconstruction. The  
802     analysis is performed using the  $f_{hit}$  energy binning scheme with bins (1-9) similar to what was done  
803     for the Crab and previous HAWC dSph analysis [42, 46]. Bin 0 was excluded as it has substantial  
804     hadronic contamination and poor angular resolution.

805     This analysis was done on dSphs because of their large DM mass content relative to baryonic  
806     mass. We consider the following to estimate the background to this study.

- 807     • The dSphs' angular extent are small relative to HAWC's spatial resolution, so the analysis is  
808         not sensitive to large or small scale anisotropies.
- 809     • The dSphs used in this analysis are off the galactic plane and therefore not contaminated by  
810         diffuse emission from the galaxy.
- 811     • The dSphs are baryonically faint relative to their expected dark matter content and are not  
812         expected to contain high energy gamma-ray sources.

813     Therefor we make no additional assumptions on the background from our sources and use  
814     HAWC's standard direct integration method for background estimation [42]. The largest background  
815     under this consideration is from an isotropic flux of cosmic rays. The contamination of this hadronic  
816     flux is worse at lower energies where HAWC's gamma/hadron discrimination worse. It is possible

817 for gamma rays from DM annihilation to scatter in transit to HAWC via Inverse Compton Scattering  
818 (ICS). This was investigated and its impact on the flux is negligible. Supporting information on  
819 this is in Section 5.7.1

820 **5.3 Analysis**

821 The expected differential photon flux from DM-DM annihilation to standard model particles,  
822  $d\Phi_\gamma/dE_\gamma$ , over solid angle,  $\Omega$ , is described by the familiar equation.

$$\frac{d\Phi_\gamma}{dE_\gamma} = \frac{\langle\sigma v\rangle}{8\pi m_\chi^2} \frac{dN_\gamma}{dE_\gamma} \times \int_{\text{source}} d\Omega \int_{l.o.s} dl \rho_\chi^2 J(r, \theta') \quad (5.1)$$

823 Where  $\langle\sigma v\rangle$  is the velocity weighted annihilation cross-section.  $\frac{dN}{dE}$  is the expected differential  
824 number of photons produced at each energy per annihilation.  $m_\chi$  is the rest mass of the supposed  
825 DM particle.  $\rho_\chi$  is the DM density.  $J$  is the astrophysical J-factor and is defined as

$$J = \int d\Omega \int_{l.o.s} dl \rho_\chi^2(r, \theta') \quad (5.2)$$

826  $l$  is the distance to the source from Earth.  $r$  is the radial distance from the center of the source.  $\theta'$  is  
827 the half angle defining a cone containing the DM source. How each component is synthesized and  
828 considered for HAWC's analysis is presented in the following sections. Section 5.3.1 presents the  
829 particle physics model for DM annihilation. Section 5.3.2 presents the spatial distributions built  
830 for each dSph.

831 **5.3.1  $\frac{dN_\gamma}{dE_\gamma}$  - Particle Physics Component**

832 For these spectra, we import the PPPC with Electroweak (EW) corrections [44]. Public versions  
833 of the imported tables are provided by the [authors online](#). The spectrum is implemented as a model  
834 script in astromodels for 3ML. The EW corrections were previously not considered for HAWC and  
835 are significant for DM annihilating to EW coupled SM particles such as all leptons, and the  $\gamma$ ,  
836  $Z$ , and  $W$  bosons [46]. Figure 5.2 demonstrates the significance of EW corrections for W boson  
837 annihilation. Across EW SM channels, the gamma-ray spectra become harder than spectra without  
838 EW corrections. Tables from the PPPC were reformatted into Python NumPy dictionaries for  
839 collaboration-wide use. A class in astromodels was developed to include the EW correction from  
840 the PPPC and is aptly named `PPPCSpectra` within `DM_models.py`.

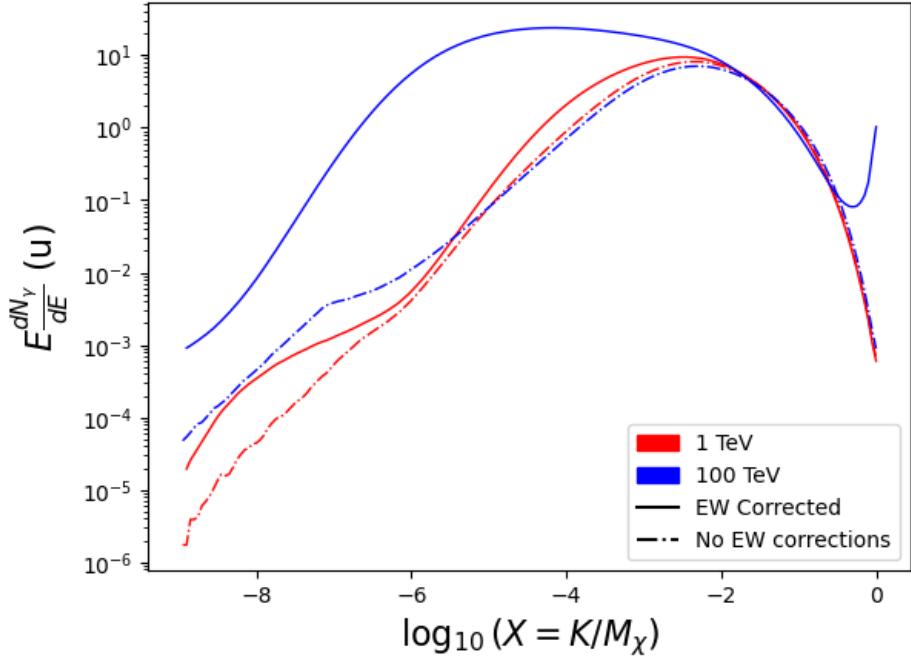


Figure 5.2 Effect of Electroweak (EW) corrections on expected DM annihilation spectrum for  $\chi\chi \rightarrow W^-W^+$ . Solid lines are spectral models that consider EW corrections. Dash-dot lines are spectral models without EW corrections. Red lines are models for  $M_\chi = 1$  TeV. Blue lines represent models for  $M_\chi = 100$  TeV. All models are sourced from the PPPC4DMID [44].

### 841 5.3.2 J- Astrophysical Component

842 The J-factor profiles for each source are imported from Geringer-Sameth (referred to with  $\mathcal{GS}$ )  
 843 [45].  $\mathcal{GS}$  fits the Zhao DM profile to the dSphs which has a DM density described as [47]

$$\rho(r) = \frac{\rho_0}{(r/R_s)^\gamma (1 + (r/R_s)^\alpha)^{(\beta-\gamma)}}. \quad (5.3)$$

844  $R_s$  is the scale radius and free parameter in the model.  $\gamma$  is the logarithmic slope in the region  
 845  $r \ll R_s$ .  $\beta$  is the logarithmic slope in the region  $r \gg R_s$ .  $\alpha$  is known as the sharpness of transition  
 846 where  $r \approx R_s$ . The classic Navarro-Frenk-White [48] (NFW) can be retrieved from Zhao by fixing  
 847  $(\alpha, \beta, \gamma) = (1, 3, 1)$ .

848  $\mathcal{GS}$  best fits were pulled from the publication as  $J(\theta)$ , where  $\theta$  is the angular separation from  
 849 the center of the source. HAWC requires maps in terms of  $\frac{dJ}{d\Omega}$ , so the conversion from the maps  
 850 was done in the following way...

851 First, convert the angular distances to solid angles

$$\Omega = 2 \cdot \pi(1 - \cos(\theta)) \quad (5.4)$$

852 which reduces with a small angle approximation to  $\pi\theta^2$ . Next, the central difference for both the  
853  $\Delta J$  and  $\Delta\Omega$  value were calculated from the discretized  $J(\theta)$  with the central difference stencil:

$$\Delta\phi_i = \phi_{i+1} - \phi_{i-1} \quad (5.5)$$

854 Where  $\phi$  is either  $\Omega$  or  $J$ . These were done separately in case the grid spacing in  $\theta$  was not uniform.  
855 Finally, these lists are divided so that we are left with an approximation of the  $dJ/d\Omega$  profile that  
856 is a function of  $\theta$ . Admittedly, this is an approximation method for the map which introduces small  
857 errors compared to the true profile estimate. This was checked as a systematic against the author's  
858 profiling of the spatial distribution and is documented in Section 5.8.1.

859 With  $\frac{dJ}{d\Omega}(\theta)$ , a map is generated, first by filling in the north-east quadrant of the map. This  
860 quadrant is then reflected twice, vertically then horizontally, to fill the full map. Maps are then  
861 normalized by dividing the discrete 2D integral of the map. The 2D integral was a simple height  
862 of bins, Newton's integral:

$$p^2 \cdot \sum_{i,j=0}^{N,M} \frac{dJ}{d\Omega}(\theta_{i,j}, \phi_{i,j}) \quad (5.6)$$

863 These maps are HEALpix maps with NSIDE 16384 and saved in the .fits format. The hyper fine  
864 resolution was selected to better preserve the total expected counts after integrating Eq. (7.1) with  
865 the detector response.

866 Another DM spatial distribution model from Bonnivard ( $\mathcal{B}$ ) [49] was used for the Glory Duck  
867 study. However, to save computational time, limits from  $\mathcal{GS}$  were scaled to  $\mathcal{B}$  instead of each  
868 experiment performing a full study a second time. How these models compare is demonstrated  
869 for each dSph in Figure 5.16 and Figure 5.17 Plots of these maps are provided for each source  
870 in chapter A Examples of the two most impactful dSphs derived from  $\mathcal{GS}$ , Segue1 and Coma  
871 Berenices are featured in Figure 5.3

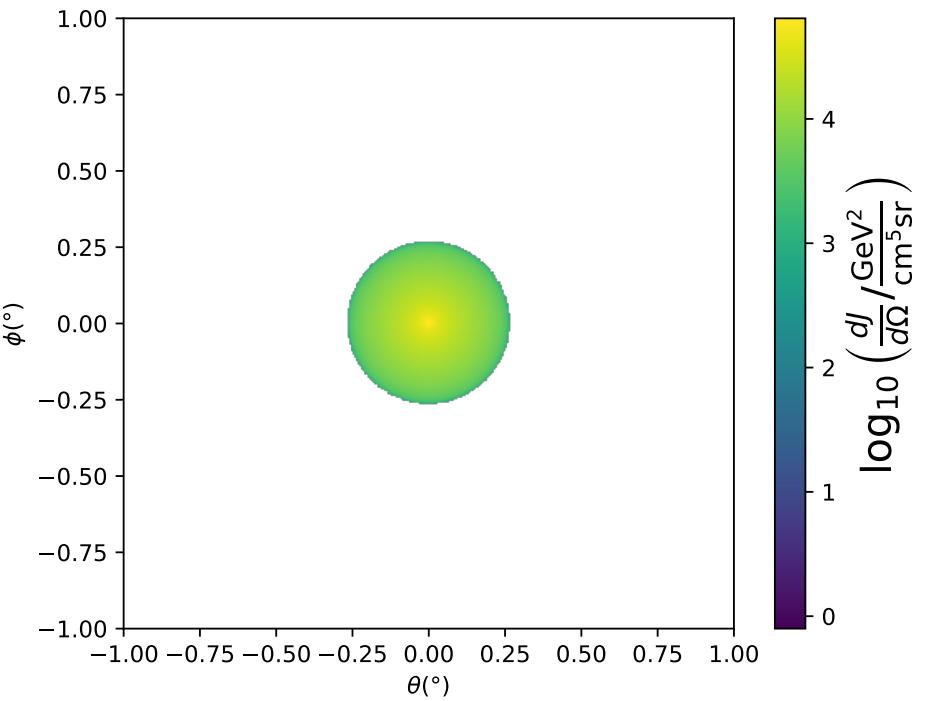
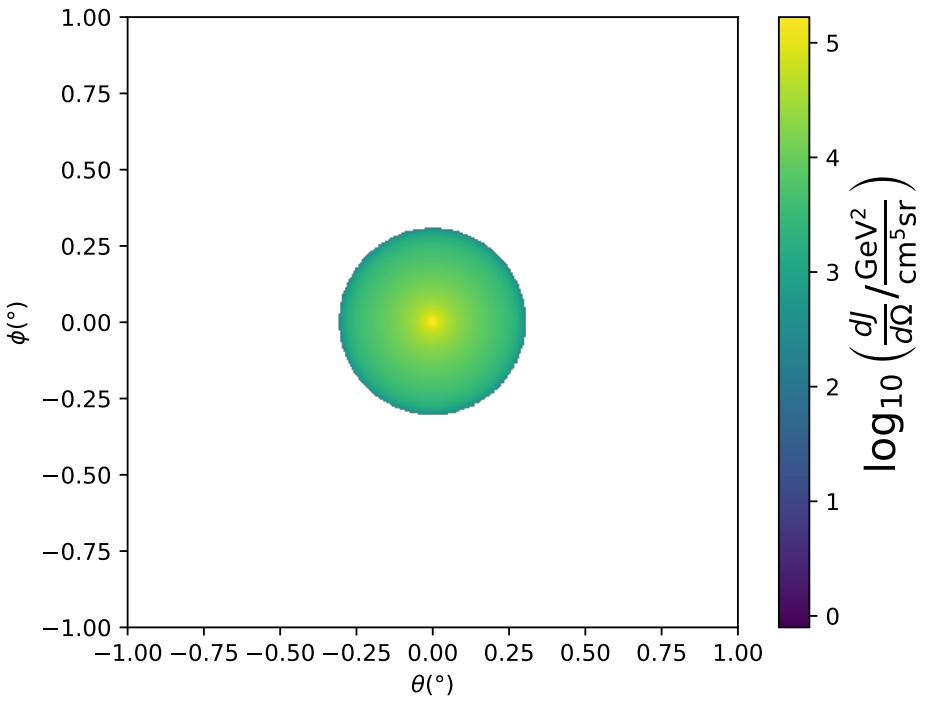


Figure 5.3  $\frac{dJ}{d\Omega}$  maps for Segue1 (top) and Coma Berenices (bottom). Origin is centered on the specific dwarf spheroidal galaxies (dSph). X and Y axes are the angular separation from the center of the dwarf. Profile is truncated at the scale radius. Plots of the remaining 11 dSph HAWC studied are linked in Fig. A.1.

872 **5.3.3 Source Selection and Annihilation Channels**

873 We use many of the dSphs presented in HAWC’s previous dSph DM search [46]. HAWC’s  
874 sources for Glory Duck include Boötes I, Coma Berenices, Canes Venatici I + II, Draco, Hercules,  
875 Leo I, II, + IV, Segue 1, Sextans, and Ursa Major I + II. A full description of all sources used  
876 in Glory Duck is found in Table 5.1. Triangulum II was excluded from the Glory Duck analysis  
877 because of large uncertainties in its  $J$  factor. Ursa Minor was excluded from HAWC’s contribution  
878 to the combination because the source extension model extended Ursa Minor beyond HAWC’s field  
879 of view. Ursa Minor was not expected to contribute significantly to the combined limit, so work  
880 was not invested in a solution to include Ursa Minor.

881 This analysis improves on the previous HAWC dSph paper [46] in the following ways. Pre-  
882 viously, the dSphs were treated and implemented as point sources. For this analysis, dSphs are  
883 modeled and treated as extended source. The impact of this change with respect to the upper limit  
884 is source dependent and is explored in Section 5.7.2. Previously, the particle physics model used for  
885 gamma-ray spectra from DM annihilation did not have EW corrections where the PPPC includes  
886 them. Finally, the gamma-ray ray dataset is much larger. The study performed here analyzes over  
887 1000 days of data compared to 507.

888 The SM annihilation channels probed for the Glory Duck combination include  $b\bar{b}$ ,  $e\bar{e}$ ,  $\mu\bar{\mu}$ ,  $\tau\bar{\tau}$ ,  
889  $t\bar{t}$ ,  $W^+W^-$ , and  $ZZ$ . A summary of all sources, with a description of each experiments’ sensitivity  
890 to the source, is provided in Table 5.2.

891 **5.4 Likelihood Methods**

892 **5.4.1 HAWC Likelihood**

893 For every analysis bin in energy,  $f_{hit}$  bins (1-9), and location, we can expect  $N$  signal events and  
894  $B$  background events. The expected number of excess signal events from dark matter annihilation,  
895  $S$ , is estimated by convolving Equation (7.1) with HAWC’s energy response and pixel point spread  
896 functions. I then compute the test statistic (TS) with the log-likelihood ratio test,

$$\text{TS}_{\max} = -2 \ln \left( \frac{\mathcal{L}_0}{\mathcal{L}^{\max}} \right) \quad (5.7)$$

Table 5.1 Summary of the relevant properties of the dSphs used in the present work. Column 1 lists the dSphs. Columns 2 and 3 present their heliocentric distance and galactic coordinates, respectively. Columns 4 and 5 report the  $J$ -factors of each source given from the  $\mathcal{GS}$  and  $\mathcal{B}$  independent studies and their estimated  $\pm 1\sigma$  uncertainties. The values  $\log_{10} J$  ( $\mathcal{GS}$  set) [45] correspond to the mean  $J$ -factor values for a source extension truncated at the outermost observed star. The values  $\log_{10} J$  ( $\mathcal{B}$  set) [49] are provided for a source extension at the tidal radius of each dSph. **Bolded sources are within HAWC's field of view and provided to the Glory Duck analysis.**

Name	Distance (kpc)	$l, b$ ( $^{\circ}$ )	$\log_{10} J$ ( $\mathcal{GS}$ set) $\log_{10}(\text{GeV}^2\text{cm}^{-5}\text{sr})$	$\log_{10} J$ ( $\mathcal{B}$ set) $\log_{10}(\text{GeV}^2\text{cm}^{-5}\text{sr})$
<b>Boötes I</b>	66	358.08, 69.62	$18.24^{+0.40}_{-0.37}$	$18.85^{+1.10}_{-0.61}$
<b>Canes Venatici I</b>	218	74.31, 79.82	$17.44^{+0.37}_{-0.28}$	$17.63^{+0.50}_{-0.20}$
<b>Canes Venatici II</b>	160	113.58, 82.70	$17.65^{+0.45}_{-0.43}$	$18.67^{+1.54}_{-0.97}$
Carina	105	260.11, -22.22	$17.92^{+0.19}_{-0.11}$	$18.02^{+0.36}_{-0.15}$
<b>Coma Berenices</b>	44	241.89, 83.61	$19.02^{+0.37}_{-0.41}$	$20.13^{+1.56}_{-1.08}$
<b>Draco</b>	76	86.37, 34.72	$19.05^{+0.22}_{-0.21}$	$19.42^{+0.92}_{-0.47}$
Fornax	147	237.10, -65.65	$17.84^{+0.11}_{-0.06}$	$17.85^{+0.11}_{-0.08}$
<b>Hercules</b>	132	28.73, 36.87	$16.86^{+0.74}_{-0.68}$	$17.70^{+1.08}_{-0.73}$
<b>Leo I</b>	254	225.99, 49.11	$17.84^{+0.20}_{-0.16}$	$17.93^{+0.65}_{-0.25}$
<b>Leo II</b>	233	220.17, 67.23	$17.97^{+0.20}_{-0.18}$	$18.11^{+0.71}_{-0.25}$
<b>Leo IV</b>	154	265.44, 56.51	$16.32^{+1.06}_{-1.70}$	$16.36^{+1.44}_{-1.65}$
Leo V	178	261.86, 58.54	$16.37^{+0.94}_{-0.87}$	$16.30^{+1.33}_{-1.16}$
Leo T	417	214.85, 43.66	$17.11^{+0.44}_{-0.39}$	$17.67^{+1.01}_{-0.56}$
Sculptor	86	287.53, -83.16	$18.57^{+0.07}_{-0.05}$	$18.63^{+0.14}_{-0.08}$
<b>Segue I</b>	23	220.48, 50.43	$19.36^{+0.32}_{-0.35}$	$17.52^{+2.54}_{-2.65}$
Segue II	35	149.43, -38.14	$16.21^{+1.06}_{-0.98}$	$19.50^{+1.82}_{-1.48}$
<b>Sextans</b>	86	243.50, 42.27	$17.92^{+0.35}_{-0.29}$	$18.04^{+0.50}_{-0.28}$
<b>Ursa Major I</b>	97	159.43, 54.41	$17.87^{+0.56}_{-0.33}$	$18.84^{+0.97}_{-0.43}$
<b>Ursa Major II</b>	32	152.46, 37.44	$19.42^{+0.44}_{-0.42}$	$20.60^{+1.46}_{-0.95}$
Ursa Minor	76	104.97, 44.80	$18.95^{+0.26}_{-0.18}$	$19.08^{+0.21}_{-0.13}$

Table 5.2 Summary of dSph observations by each experiment used in this work. A ‘-’ indicates the experiment did not observe the dSph for this study. For Fermi-LAT, the exposure at 1 GeV is given. For HAWC,  $|\Delta\theta|$  is the absolute difference between the source declination and HAWC latitude. HAWC is more sensitive to sources with smaller  $|\Delta\theta|$ . For IACTs, we show the zenith angle range, the total exposure, the energy range, the angular radius  $\theta$  of the signal or ON region, the ratio of exposures between the background-control (OFF) and signal (ON) regions ( $\tau$ ), and the significance of gamma-ray excess in standard deviations,  $\sigma$ .

Source name	Fermi-LAT	HAWC	H.E.S.S, MAGIC, VERITAS						
	Exposure ( $10^{11}$ s m $^2$ )	$ \Delta\theta $ ( $^\circ$ )	IACT	Zenith ( $^\circ$ )	Exposure (h)	Energy range (GeV)	$\theta$ ( $^\circ$ )	$\tau$	$S$ ( $\sigma$ )
Boötes I	2.6	4.5	VERITAS	15 – 30	14.0	100–41000	0.10	8.6	-1.0
Canes Venatici I	2.9	14.6	–	–	–	–	–	–	–
Canes Venatici II	2.9	15.3	–	–	–	–	–	–	–
Carina	3.1	–	H.E.S.S.	27 – 46	23.7	310 – 70000	0.10	18.0	-0.3
Coma Berenices	2.7	4.9	H.E.S.S.	47 – 49	11.4	550 – 70000	0.10	14.4	-0.4
			MAGIC	5 – 37	49.5	60 – 10000	0.17	1.0	–
			MAGIC	29 – 45	52.1	70 – 10000	0.22	1.0	–
Draco	3.8	38.1	VERITAS	25 – 40	49.8	120 – 70000	0.10	9.0	-1.0
Fornax	2.7	–	H.E.S.S.	11 – 25	6.8	230 – 70000	0.10	45.5	-1.5
Hercules	2.8	6.3	–	–	–	–	–	–	–
Leo I	2.5	6.7	–	–	–	–	–	–	–
Leo II	2.6	3.1	–	–	–	–	–	–	–
Leo IV	2.4	19.5	–	–	–	–	–	–	–
Leo V	2.4	–	–	–	–	–	–	–	–
Leo T	2.6	–	–	–	–	–	–	–	–
Sculptor	2.7	–	H.E.S.S.	10 – 46	11.8	200 – 70000	0.10	19.8	-2.2
Segue I	2.5	2.9	MAGIC	13 – 37	158.0	60 – 10000	0.12	1.0	-0.5
			VERITAS	15 – 35	92.0	80 – 50000	0.10	7.6	0.7
Segue II	2.7	–	–	–	–	–	–	–	–
Sextans	2.4	20.6	–	–	–	–	–	–	–
Ursa Major I	3.4	32.9	–	–	–	–	–	–	–
Ursa Major II	4.0	44.1	MAGIC	35 – 45	94.8	120 – 10000	0.30	1.0	-2.1
Ursa Minor	4.1	–	VERITAS	35 – 45	60.4	160 – 93000	0.10	8.4	-0.1

897 where  $\mathcal{L}_0$  is the null hypothesis, or no DM emission, likelihood.  $\mathcal{L}^{\max}$  is the best fit signal  
 898 hypothesis where  $\langle\sigma v\rangle$  maximizes the likelihood. We calculate the likelihood of each source and  
 899 model, assuming events are Poisson distributed, with

$$\mathcal{L} = \prod_i \frac{(B_i + S_i)^{N_i} e^{-(B_i + S_i)}}{N_i!} \quad (5.8)$$

901 where  $S_i$  is the sum of expected number of signal counts.  $B_i$  is the number of background counts  
 902 observed.  $N_i$  is the total number of counts.

903 I also calculate an upper limit on  $\langle\sigma v\rangle$  by calculating the 95% confidence level (CL). For the  
 904 CL, we define a parameter,  $TS_{95}$ , as

$$TS_{95} \equiv \sum_{\text{bins}} \left[ 2N \ln \left( 1 + \frac{\epsilon S_{\text{ref}}}{B} \right) - 2\epsilon S_{\text{ref}} \right] \quad (5.9)$$

905 where the expected signal counts from a dSph is scaled by  $\epsilon$ .  $S_{\text{ref}}$  is the expected number of excess  
 906 counts in a bin from DM emission from a dSph with a corresponding annihilation cross-section,  
 907  $\langle\sigma v\rangle$ . We scan  $\epsilon$  such that

$$2.71 = TS_{\max} - TS_{95} \quad (5.10)$$

908 HAWC's exclusive results are provided in Section 5.5.

### 909 5.4.2 Glory Duck Joint Likelihood

910 The joint likelihood for the 5-experiment combination was done similarly as Section 5.4.1. We  
 911 calculate upper limits on  $\langle\sigma v\rangle$  from the TS, Eq. (5.7), and define the likelihood ratio more generally

$$\lambda(\langle\sigma v\rangle | \mathcal{D}_{\text{dSphs}}) = \frac{\mathcal{L}(\langle\sigma v\rangle; \hat{\nu} | \mathcal{D}_{\text{dSphs}})}{\mathcal{L}(\widehat{\langle\sigma v\rangle}; \hat{\nu} | \mathcal{D}_{\text{dSphs}})} \quad (5.11)$$

912  $\mathcal{D}_{\text{dSphs}}$  is the totality of observations across experiments and dSphs.  $\nu$  are the nuisance parameters  
 913 which are the  $J$  factors in this study.  $\widehat{\langle\sigma v\rangle}$  and  $\hat{\nu}$  are the respective estimate that maximize  $\mathcal{L}$   
 914 globally. Finally,  $\hat{\nu}$  is the set of nuisance parameters that maximize  $\mathcal{L}$  for a fixed value of  $\langle\sigma v\rangle$ .

915 The *complete* joint likelihood,  $\mathcal{L}$  that encompasses all observations from all instruments and  
 916 dSphs can be factorized into *partial* functions for each dSph  $l$  (with  $\mathcal{L}_{\text{dSph},l}$ ) and its  $J$  factor ( $\mathcal{J}_l$ ):

$$\mathcal{L}(\langle\sigma v\rangle; \nu | \mathcal{D}_{\text{dSphs}}) = \prod_{l=1}^{N_{\text{dSphs}}} \mathcal{L}_{\text{dSph},l}(\langle\sigma v\rangle; J_l, \nu_l | \mathcal{D}_l) \times \mathcal{J}_l(J_l | J_{l,\text{obs}}, \sigma_{\log J_l}). \quad (5.12)$$

917 For this study,  $N_{\text{dSphs}} = 20$  is the number of dSphs studied.  $\mathcal{D}_l$  are the gamma-ray observations  
918 of dSph,  $l$ .  $\nu_l$  are the nuisance parameters modifying the gamma-ray observations of dSph,  $l$ ,  
919 but excludes  $\mathcal{J}_l$ .  $\mathcal{J}_l$  is the  $J$  factor for dSph,  $l$ , as defined in Equation (5.2), and it is a nuisance  
920 parameter whose value is unknown.  $\log_{10} J_{l,\text{obs}}$  and  $\sigma_{\log J_l}$  are obtained by fitting a log-normal  
921 function of  $J_{l,\text{obs}}$  to the posterior distribution of  $J_l$  [50].  $\log_{10} J_{l,\text{obs}}$  and  $\sigma_{\log J_l}$  values are provided  
922 in Table 5.1. The term  $\mathcal{J}_l$  constraining  $J_l$  is written as:

$$\mathcal{J}_l(J_l | J_{l,\text{obs}}, \sigma_{\log J_l}) = \frac{1}{\ln(10)J_{l,\text{obs}}\sqrt{2\pi}\sigma_{\log J_l}} \exp\left(-\frac{(\log_{10} J_l - \log_{10} J_{l,\text{obs}})^2}{2\sigma_{\log J_l}^2}\right). \quad (5.13)$$

923 Both the  $\mathcal{GS}$  and  $\mathcal{B}$ , displayed in Table 5.1, sets of  $J$  factors are used in this analysis. Equation (5.13)  
924 is also normalized, so it can also be interpreted as a probability density function (PDF) for  $J_{l,\text{obs}}$ .  
925 From Equation (7.1), we can also see that  $\langle\sigma v\rangle$  and  $J_l$  are degenerate when computing  $\mathcal{L}_{\text{dSph},l}$ .  
926 Therefore, as noted in [51], it is sufficient to compute  $\mathcal{L}_{\text{dSph},l}$  versus  $\langle\sigma v\rangle$  for a fixed value of  $J_l$ .  
927 We used  $J_{l,\text{obs}}(\mathcal{GS})$  reported in Tab. 5.1, in order to perform the profile of  $\mathcal{L}$  with respect to  $J_l$ .  
928 The degeneracy implies that for any  $J'_l \neq J_{l,\text{obs}}$  (in practice in our case we used  $J'_l = J_{l,\text{obs}}(\mathcal{B})$  to  
929 compute results from a different set of  $J$  factors):

$$\mathcal{L}_{\text{dSph},l}(\langle\sigma v\rangle; J'_l, \nu_l | \mathcal{D}_l) = \mathcal{L}_{\text{dSph},l}\left(\frac{J'_l}{J_{l,\text{obs}}} \langle\sigma v\rangle; J_{l,\text{obs}}, \nu_l | \mathcal{D}_l\right), \quad (5.14)$$

930 which is a straightforward rescaling operation that reduces the computational needs of the profiling  
931 operation since:

$$\mathcal{L}(\langle\sigma v\rangle; \hat{\nu} | \mathcal{D}_{\text{dSphs}}) = \prod_{l=1}^{N_{\text{dSphs}}} \max_{J_l} \left[ \mathcal{L}_{\text{dSph},l}(\langle\sigma v\rangle; J_l, \hat{\nu}_l | \mathcal{D}_l) \times \mathcal{J}_l(J_l | J_{l,\text{obs}}, \sigma_{\log J_l}) \right]. \quad (5.15)$$

932 In addition, Eq. (5.14) enables the combination of data from different gamma-ray instruments and  
933 observed dSphs via tabulated values of  $\mathcal{L}_{\text{dSph},l}$ , or equivalently of  $\lambda$  from Eq. (5.11) as was done in  
934 this work, versus  $\langle\sigma v\rangle$ .  $\mathcal{L}_{\text{dSph},l}$  is computed for a fixed value of  $J_l$  and profiled with respect to all  
935 instrumental nuisance parameters  $\nu_l$ , these nuisance parameters are discussed in more detail below.  
936 These values are produced by each detector independently and therefore there is no need to share  
937 sensitive low-level information used to produce them, such as event lists. Figure 5.4 illustrates the  
938 multi-instrument combination technique used in this study with a comparison of the upper limit

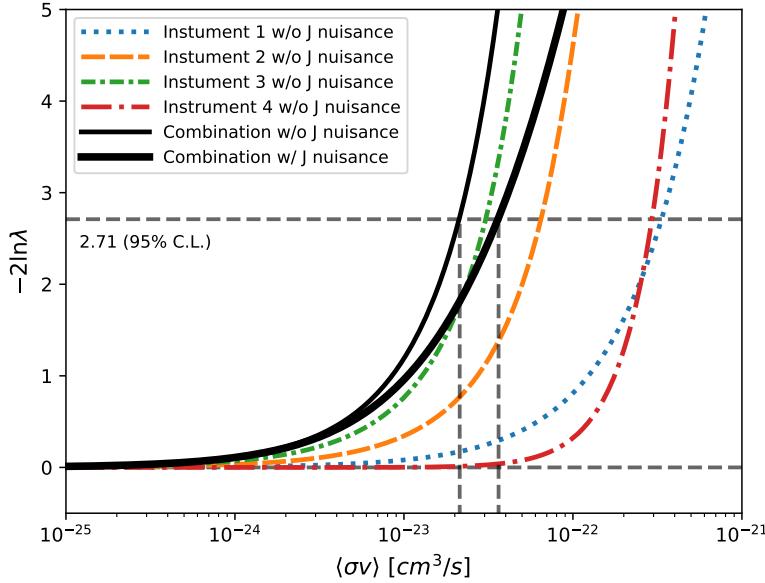


Figure 5.4 Illustration of the combination technique showing a comparison between  $-2 \ln \lambda$  provided by four instruments (colored lines) from the observation of the same dSph without any  $J$  nuisance and their sum, *i.e.* the resulting combined likelihood (thin black line). According to the test statistics of Equation (5.7), the intersection of the likelihood profiles with the line  $-2 \ln \lambda = 2.71$  indicates the 95% C.L. upper limit on  $\langle\sigma v\rangle$ . The combined likelihood (thin black line) shows a smaller value of upper limit on  $\langle\sigma v\rangle$  than those derived by individual instruments. We also show how the uncertainties on the  $J$  factor effects the combined likelihood and degrade the upper limit on  $\langle\sigma v\rangle$  (thick black line). All likelihood profiles are normalized so that the global minimum  $\widehat{\langle\sigma v\rangle}$  is 0. We note that each profile depends on the observational conditions in which a target object was observed. The sensitivity of a given instrument can be degraded and the upper limits less constraining if the observations are performed in non-optimal conditions such as a high zenith angle or a short exposure time.

939 on  $\langle\sigma v\rangle$  obtained from the combination of the observations of four experiments towards one dSph  
 940 versus the upper limit from individual instruments. It also shows graphically the effect of the  
 941  $J$ -factor uncertainty on the combined observations.

942 The *partial* joint likelihood function for gamma-ray observations of each dSph ( $\mathcal{L}_{dSph,l}$ ) is  
 943 written as the product of the likelihood terms describing the  $N_{exp,l}$  observations performed with  
 944 any of our observatories:

$$\mathcal{L}_{dSph,l} (\langle\sigma v\rangle; J_l, \nu_l | \mathcal{D}_l) = \prod_{k=1}^{N_{exp,l}} \mathcal{L}_{lk} (\langle\sigma v\rangle; J_l, \nu_{lk} | \mathcal{D}_{lk}), \quad (5.16)$$

945 where each  $\mathcal{L}_{lk}$  term refers to an observation of the  $l$ -th dSph with associated  $k$ -th instrument

946 responses.  $N_{\text{exp},l}$  varies from dSph to dSph and can be inferred from Table 5.2.

947     Each collaboration separately analyzes their data for  $\mathcal{D}_{lk}$  corresponding to dSph  $l$  and gamma-  
948     ray detector  $k$ , using as many common assumptions as possible in the analysis. HAWC's treatment  
949     was described earlier in Section 5.4.1 whereas the specifics of the remaining experiments is left to  
950     the publication. We compute the values for the likelihood functions  $\mathcal{L}_{lk}$  (see Eq. (5.16)) for a fixed  
951     value of  $J_l$  and profile over the rest of the nuisance parameters  $\nu_{lk}$ . Then, values of  $\lambda$  from Eq. (5.11)  
952     are computed as a function of  $\langle \sigma v \rangle$ , and shared using a common format. Results are computed for  
953     seven annihilation channels,  $W^+W^-$ ,  $ZZ$ ,  $b\bar{b}$ ,  $t\bar{t}$ ,  $e^+e^-$ ,  $\mu^+\mu^-$ , and  $\tau^+\tau^-$  over 62  $m_\chi$  values between  
954     5 GeV and 100 TeV provided in [44]. The  $\langle \sigma v \rangle$  range is defined between  $10^{-28}$  and  $10^{-18} \text{cm}^3 \cdot \text{s}^{-1}$ ,  
955     with 1001 logarithmically spaced values. The likelihood combination, i.e. Equation (5.12), and  
956     profile over the  $J$ -factor to compute the profile likelihood ratio  $\lambda$ , Equation (5.11), are carried out  
957     with two different public analysis software packages, namely `gLike` [52] and `LklCom` [53], that  
958     provide the same results [54].

959     As mentioned previously, each experiment computes the  $\mathcal{L}_{lk}$  from Equation (5.11) differently.  
960     The remainder of this section highlights the differences in this calculation across the experiments.  
961     Four experiments, namely *Fermi*-LAT, H.E.S.S., HAWC and MAGIC, use a binned likelihood to  
962     compute the  $\mathcal{L}_{lk}$ . For these experiments, for each observation  $\mathcal{D}_{lk}$  of a given dSph  $l$  carried out  
963     using a given gamma-ray detector  $k$ , the binned likelihood function is:

$$\mathcal{L}_{lk} (\langle \sigma v \rangle; J_l, \nu_{lk} | \mathcal{D}_{lk}) = \prod_{i=1}^{N_E} \prod_{j=1}^{N_P} \left[ \mathcal{P}(s_{lk,ij}(\langle \sigma v \rangle, J_l, \nu_{lk}) + b_{lk,ij}(\nu_{lk}) | N_{lk,ij}) \right] \times \mathcal{L}_{lk,\nu} (\nu_{lk} | \mathcal{D}_{\nu_{lk}}) \quad (5.17)$$

964     where  $N_E$  and  $N_P$  are the number of considered bins in reconstructed energy and arrival direction,  
965     respectively;  $\mathcal{P}$  represents a Poisson PDF for the number of gamma-ray candidate events  $N_{lk,ij}$   
966     observed in the  $i$ -th bin in energy and  $j$ -th bin in arrival direction, when the expected number is  
967     the sum of the expected mean number of signal events  $s_{ij}$  (produced by DM annihilation) and of  
968     background events  $b_{ij}$ ;  $\mathcal{L}_{lk,\nu}$  is the likelihood term for the extra  $\nu_{lk}$  nuisance parameters that vary  
969     from one instrument  $k$  to another. The expected counts for signal events  $s_{ij}$  for a given dSph  $l$  and

970 detector  $k$  is given by:

$$s_{ij}(\langle\sigma\nu\rangle, J) = \int_{E'_{\min,i}}^{E'_{\max,i}} dE' \int_{P'_{\min,j}}^{P'_{\max,j}} d\Omega' \int_0^\infty dE \int_{\Delta\Omega_{tot}} d\Omega \int_0^{T_{\text{obs}}} dt \frac{d^2\Phi(\langle\sigma\nu\rangle, J)}{dEd\Omega} \text{IRF}(E', P' | E, P, t) \quad (5.18)$$

971 where  $E'$  and  $E$  are the reconstructed and true energies,  $P'$  and  $P$  the reconstructed and true  
972 arrival directions;  $E'_{\min,i}$ ,  $P'_{\min,j}$ ,  $E'_{\max,i}$ , and  $P'_{\max,j}$  are their lower and upper limits of the  $i$ -th  
973 energy bin and the  $j$ -th arrival direction bin;  $T_{\text{obs}}$  is the (dead-time corrected) total observation  
974 time;  $t$  is the time along the observations;  $d^2\Phi/dEd\Omega$  is the DM flux in the source region (see  
975 Equation (7.1)); and  $\text{IRF}(E', P' | E, P, t)$  is the IRF, which can be factorized as the product of the  
976 effective collection area of the detector  $A_{\text{eff}}(E, P, t)$ , the PDFs for the energy estimator  $f_E(E' | E, t)$ ,  
977 and arrival direction  $f_P(P' | E, P, t)$  estimators. Note that for Fermi-LAT, HAWC, MAGIC, and  
978 VERITAS the effect of the finite angular resolution is taken into account through the convolution  
979 of  $d\Phi/dEd\Omega$  with  $f_P$  in Equation (5.18), whereas in the cases of H.E.S.S.  $f_P$  is approximated by a  
980 delta function. This approximation has been made in order to maintain compatibility of the result  
981 with what has been previously published. The difference introduced by this approximation is  $< 5\%$   
982 for all considered dSphs. A more comprehensive review of the differences between the analyses of  
983 different instruments can be found in [25].

984 **5.5 HAWC Results**

985 13 of the 20 dSphs considered for the Glory Duck analysis are within HAWC's field of view.  
986 These dSph are analyzed for emission from DM annihilation according to the likelihood method  
987 described in Section 5.4. The 13 likelihood profiles are then stacked to synthesize a combined  
988 limit on the dark matter cross-section,  $\langle\sigma\nu\rangle$ . This combination is done for the 7 SM annihilation  
989 channels used in the Glory Duck analysis. Figure 5.5 shows the combined limit for all annihilation  
990 channels with HAWC only observations. We also perform 300 studies of Poisson trials on the  
991 background. These trials are used to produce HAWC sensitivities with  $\pm 1\sigma$  and  $\pm 2\sigma$  uncertainty  
992 bands which were shared with the other collaborators for combination. The results on fitting to  
993 HAWC's Poisson trials of the DM hypothesis is shown in Figure 5.7 for all the DM annihilation  
994 channels studied for Glory Duck.

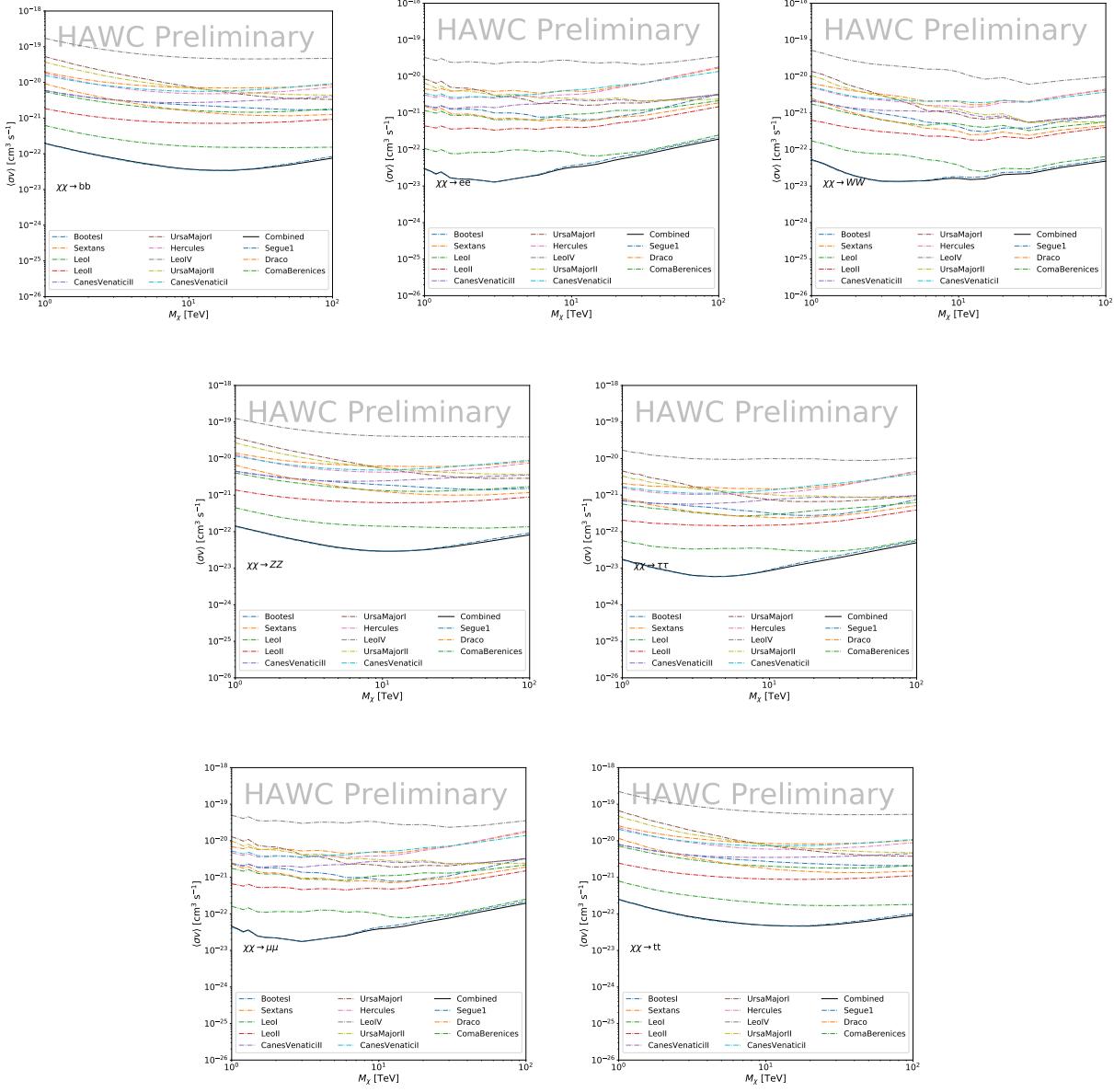


Figure 5.5

995 No DM was found in HAWC observations. HAWC's limits are dominated by the dSphs Segue1  
996 and Coma Berenices. The remaining 11 dSphs do not contribute significantly to the limit because  
997 they are at high zenith and/or have much smaller  $J$  factors. Even though some remaining dSphs  
998 have large  $J$  factors, they are towards the edge of HAWC's field of view where HAWC analysis is  
999 less sensitive.

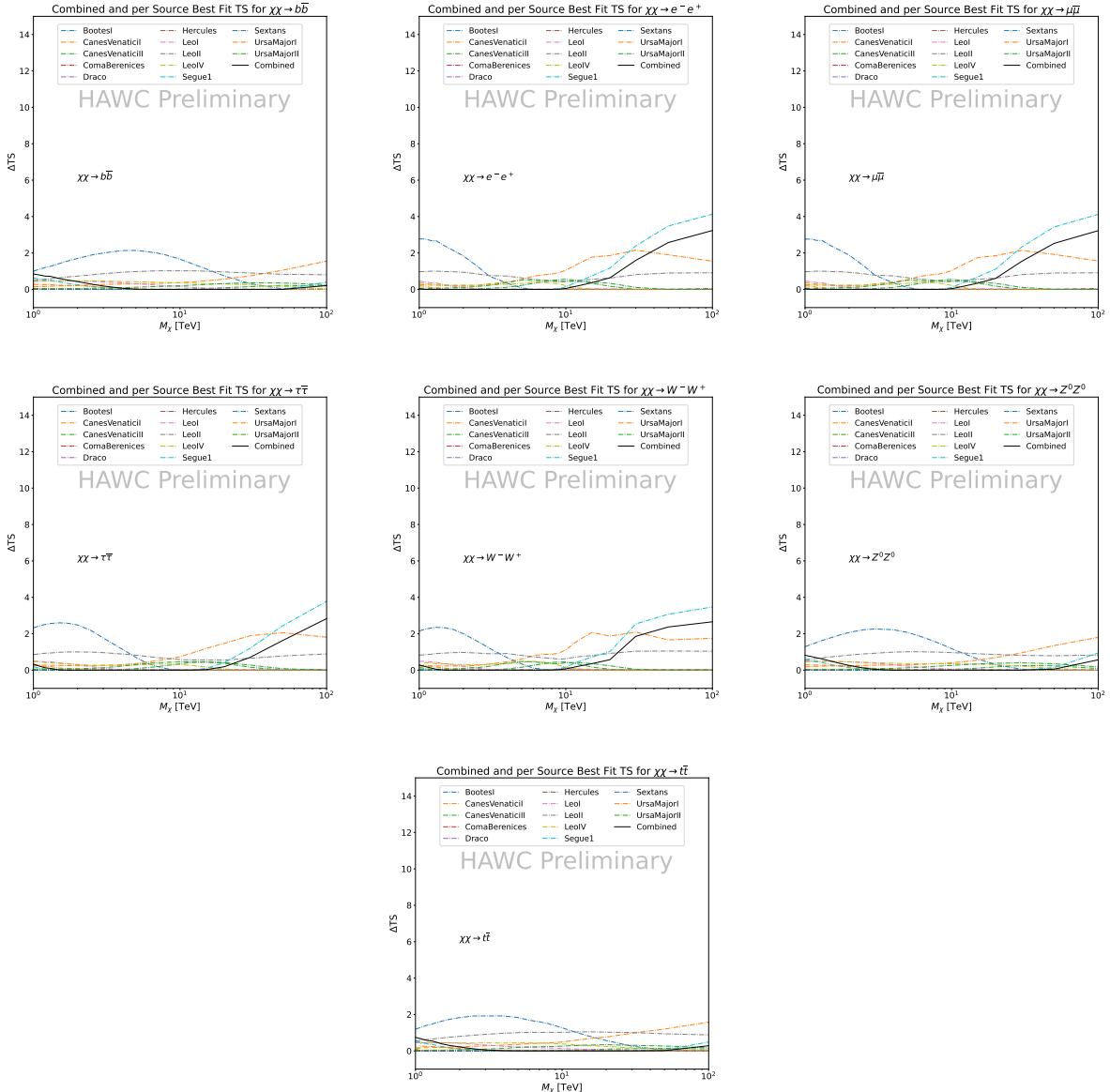


Figure 5.6 HAWC TS values for best fit  $\langle\sigma v\rangle$  versus  $m_\chi$  for seven SM annihilation channels with  $J$  factors from  $\mathcal{GS}$ . The solid black line shows the combined best fit TS values. The colored, dashed lines are the TS values for each of the 13 sources HAWC studied.

## 5.6 Glory Duck Combined Results

The crux of this analysis is that HAWC's results are combined with 4 other gamma-ray observatories: Fermi-LAT, H.E.S.S., MAGIC, and VERITAS. No significant DM emission was observed by any of the five instruments. We present the upper limits on  $\langle\sigma v\rangle$  assuming seven independent DM self annihilation channels, namely  $W^+W^-$ ,  $Z^+Z^-$ ,  $b\bar{b}$ ,  $t\bar{t}$ ,  $e^+e^-$ ,  $\mu^+\mu^-$ , and  $\tau^+\tau^-$ . The 68%

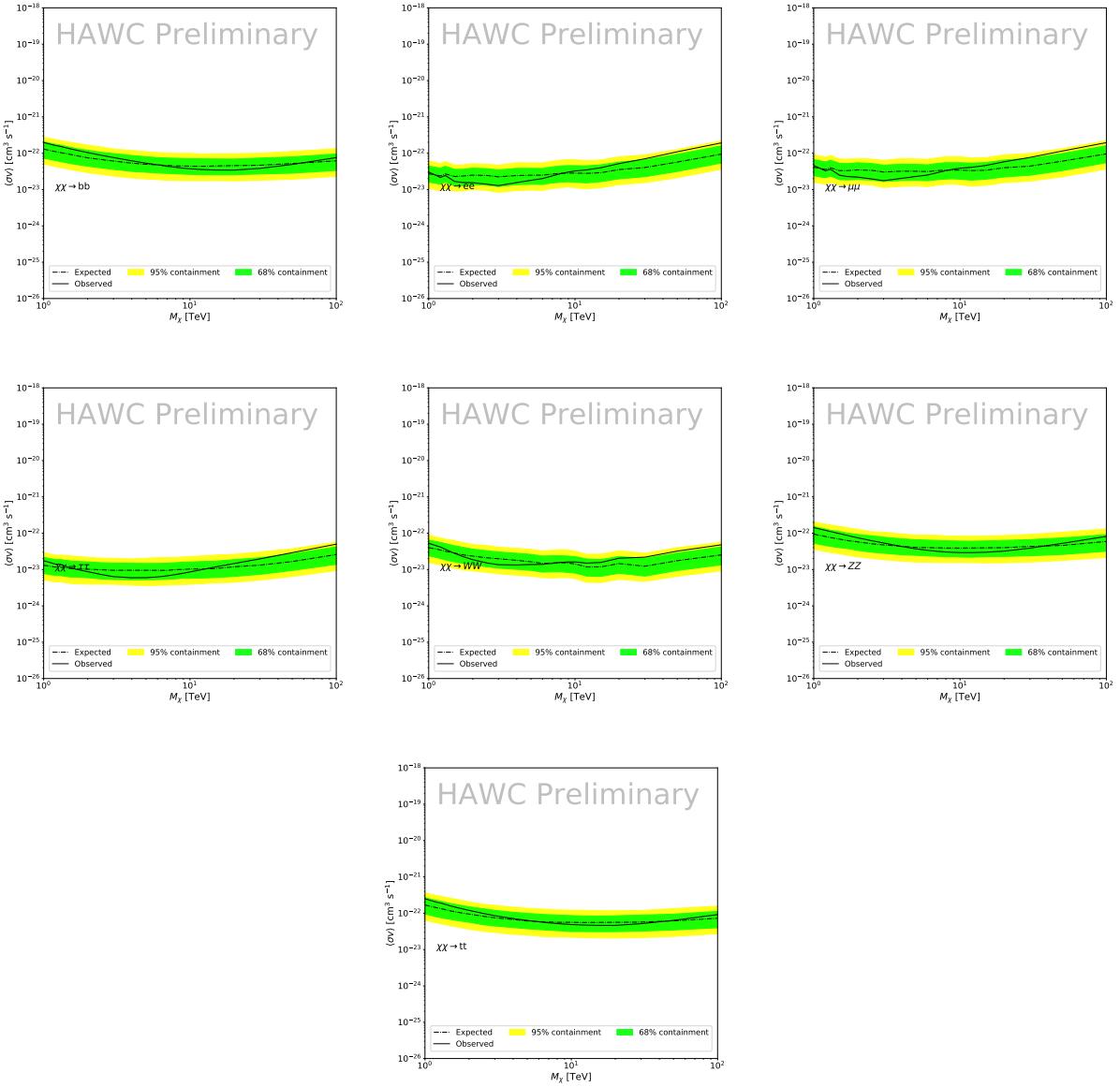


Figure 5.7 HAWC Brazil bands at 95% confidence level on  $\langle\sigma v\rangle$  versus DM mass for seven annihilation channels with  $J$ -factors from  $\mathcal{GS}$  [55]. The solid line represents the combined limit from 13 dSphs. The dashed line is the expected limit. The green band is the 68% containment. The yellow band is the 95% containment.

and 95% containment bands are produced from 300 Poisson realizations of the null hypothesis corresponding to each of the combined datasets. These 300 realizations are combined identically to dSph observations. The containment bands and the median are extracted from the distribution of resulting limits on the null hypothesis. These 300 realizations are obtained either by fast simulations of the OFF observations, for H.E.S.S., MAGIC, VERITAS, and HAWC, or taken from real

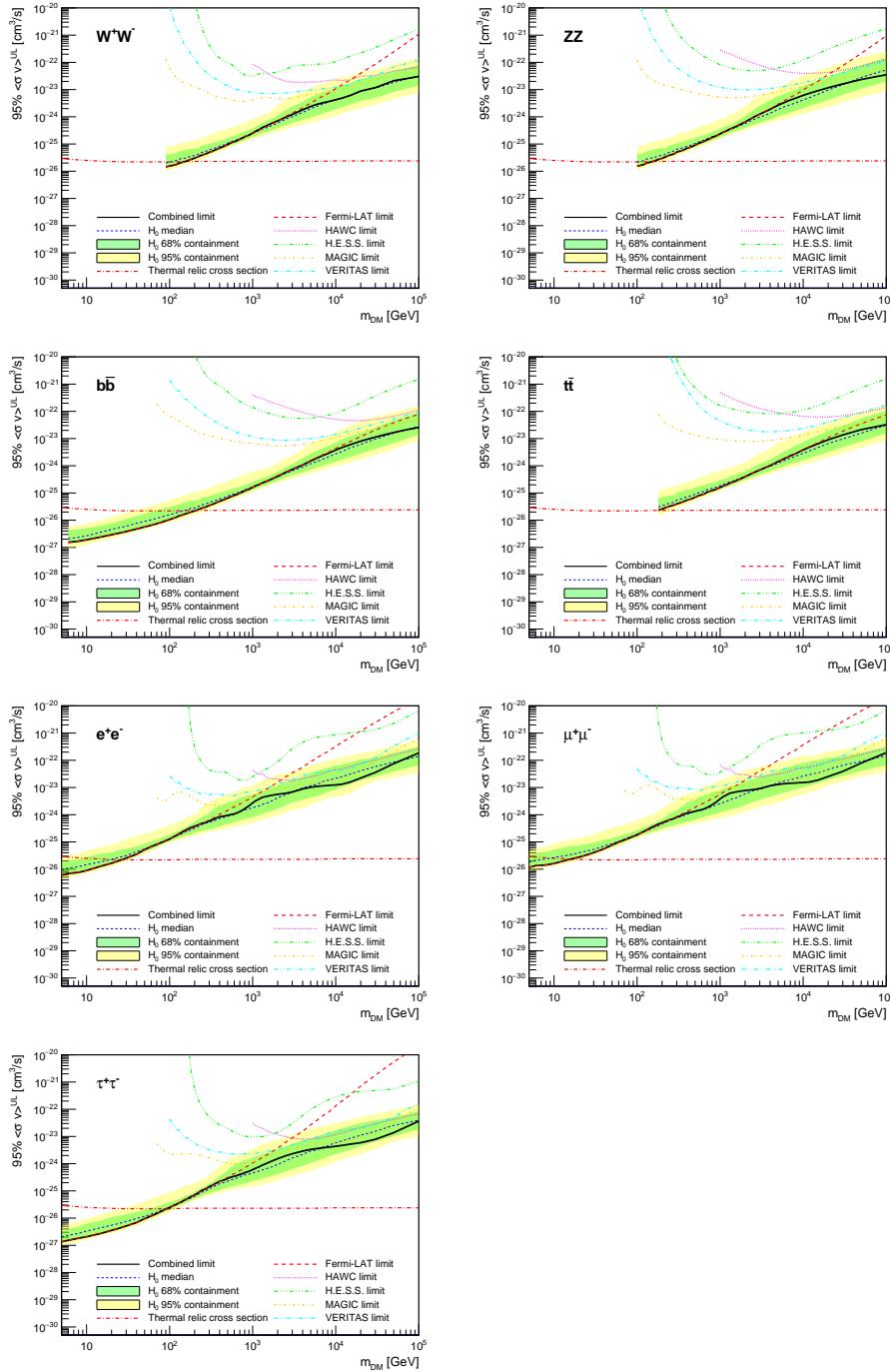


Figure 5.8 Upper limits at 95% confidence level on  $\langle\sigma v\rangle$  in function of the DM mass for eight annihilation channels, using the set of  $J$  factors from Ref. [55] ( $\mathcal{GS}$  set in Table 5.1). The black solid line represents the observed combined limit, the black dashed line is the median of the null hypothesis corresponding to the expected limit, while the green and yellow bands show the 68% and 95% containment bands. Combined upper limits for each individual detector are also indicated as solid, colored lines. The value of the thermal relic cross-section in function of the DM mass is given as the red dotted-dashed line [56].

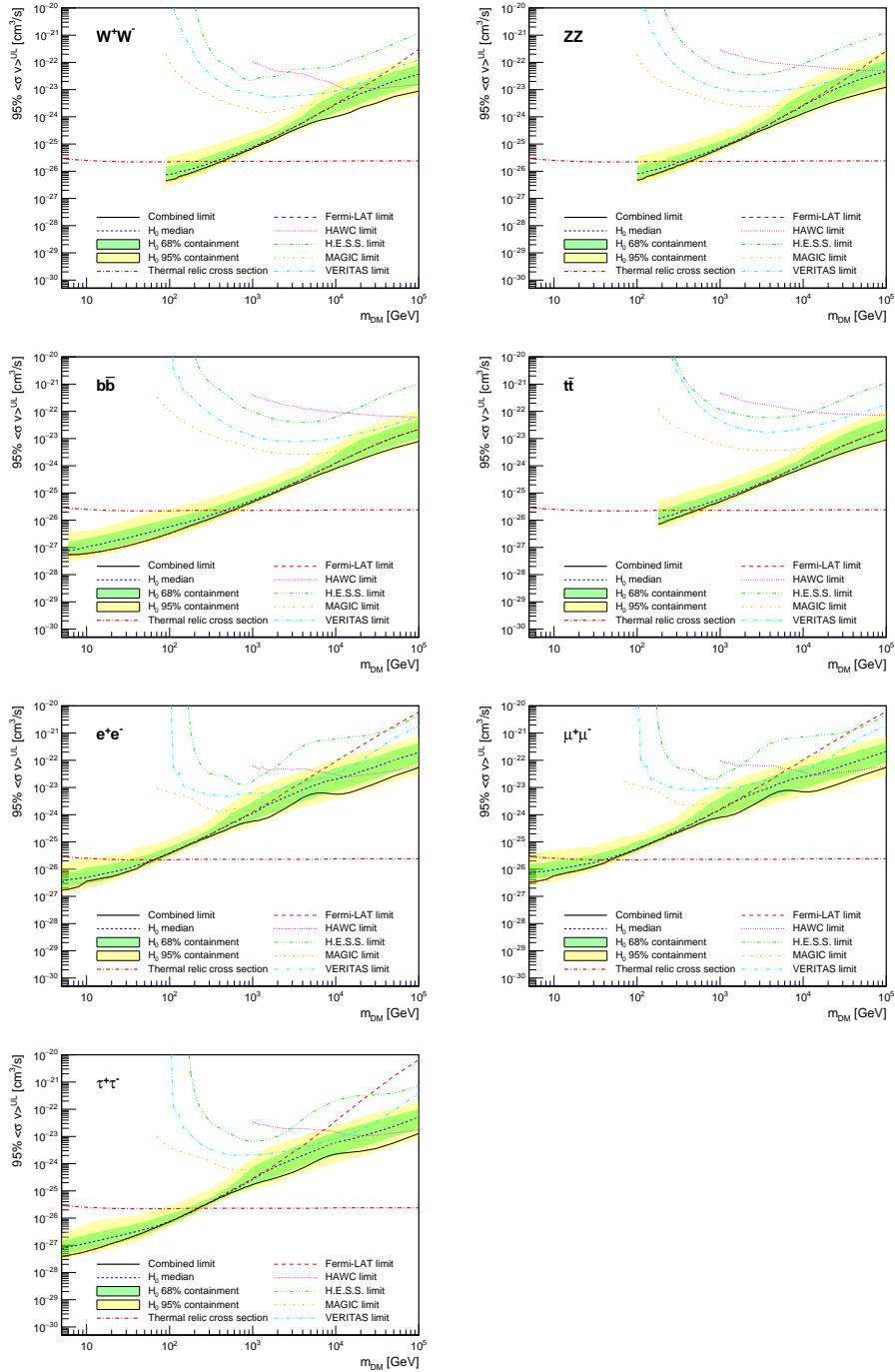


Figure 5.9 Same as Fig. 5.8, using the set of  $J$  factors from Ref. [49, 57] ( $\mathcal{B}$  set in Table 5.1).

1010 observations of empty fields of view in the case of Fermi-LAT [50, 58, 59].

1011 The obtained limits are shown in Figure 5.8 for the  $\mathcal{G}\mathcal{S}$  set of  $J$ -factors [55] and in Figure 5.9  
1012 for the  $\mathcal{B}$  set of  $J$ -factors [49, 57]. The combined limits are presented with their 68% and 95%  
1013 containment bands, and are expected to be close to the median limit when no signal is present.

1014 We observe agreement with the null hypothesis for all channels, within  $2\sigma$  standard deviations,  
1015 between the observed limits and the expectations given by the median limits. Limits obtained from  
1016 each detector are also indicated in the figures, where limits for all dSphs observed by the specific  
1017 instrument have been combined.

1018 Below  $\sim 300$  GeV, the *Fermi*-LAT dominates the DM limits for all annihilation channels. From  
1019  $\sim 300$  GeV to  $\sim 2$  TeV, *Fermi*-LAT continues to dominate for the hadronic and bosonic DM channels,  
1020 yet the IACTs (H.E.S.S., MAGIC, and VERITAS) and *Fermi*-LAT all contribute to the limit for  
1021 leptonic DM channels. For DM masses between  $\sim 2$  TeV to  $\sim 10$  TeV, the IACTs dominate leptonic  
1022 DM annihilation channels, whereas both the *Fermi*-LAT and the IACTs dominate bosonic and  
1023 hadronic DM annihilation channels. From  $\sim 10$  TeV to  $\sim 100$  TeV, both the IACTs and HAWC  
1024 contribute significantly to the leptonic DM limit. For hadronic and bosonic DM, the IACTs and  
1025 *Fermi*-LAT both contribute strongly.

1026 We notice that the limits computed using the  $\mathcal{B}$  set of  $J$ -factor are always better compared to the  
1027 ones calculated with the  $\mathcal{GS}$  set. For the  $W^+W^-$ ,  $Z^+Z^-$ ,  $b\bar{b}$ , and  $t\bar{t}$  channels, the ratio between the  
1028 limits computed with the two sets of  $J$ -factor is varying between a factor of  $\sim 3$  and  $\sim 5$  depending  
1029 on the energy, with the largest ratio around 10 TeV. For the channels  $e^+e^-$ ,  $\mu^+\mu^-$ , and  $\tau^+\tau^-$ , the  
1030 ratio lies between  $\sim 2$  to  $\sim 6$ , being maximum around 1 TeV. Examining Figure 5.16 and Figure 5.17  
1031 in Section 5.8, these differences are explained by the fact that the  $\mathcal{B}$  set provides higher  $J$ -factors  
1032 for the majority of the studied dSphs, with the notable exception of Segue I. The variation on the  
1033 ratio of the limits for the two sets is due to different dSph dominating the limits depending on the  
1034 energy. One set,  $\mathcal{B}$ , pushes the range of which thermal cross-section which can be excluded to  
1035 higher mass. This comparison demonstrates the magnitude of systematic uncertainties associated  
1036 with the choice of the  $J$ -factor calculation. The  $\mathcal{GS}$  and  $\mathcal{B}$  sets present a difference in the limits for  
1037 all channels of about This difference is explained, see Figure 5.16 and Figure 5.17, by the fact that  
1038 the  $\mathcal{B}$  set provides higher  $J$ -factors for all dSph except for Segue I.

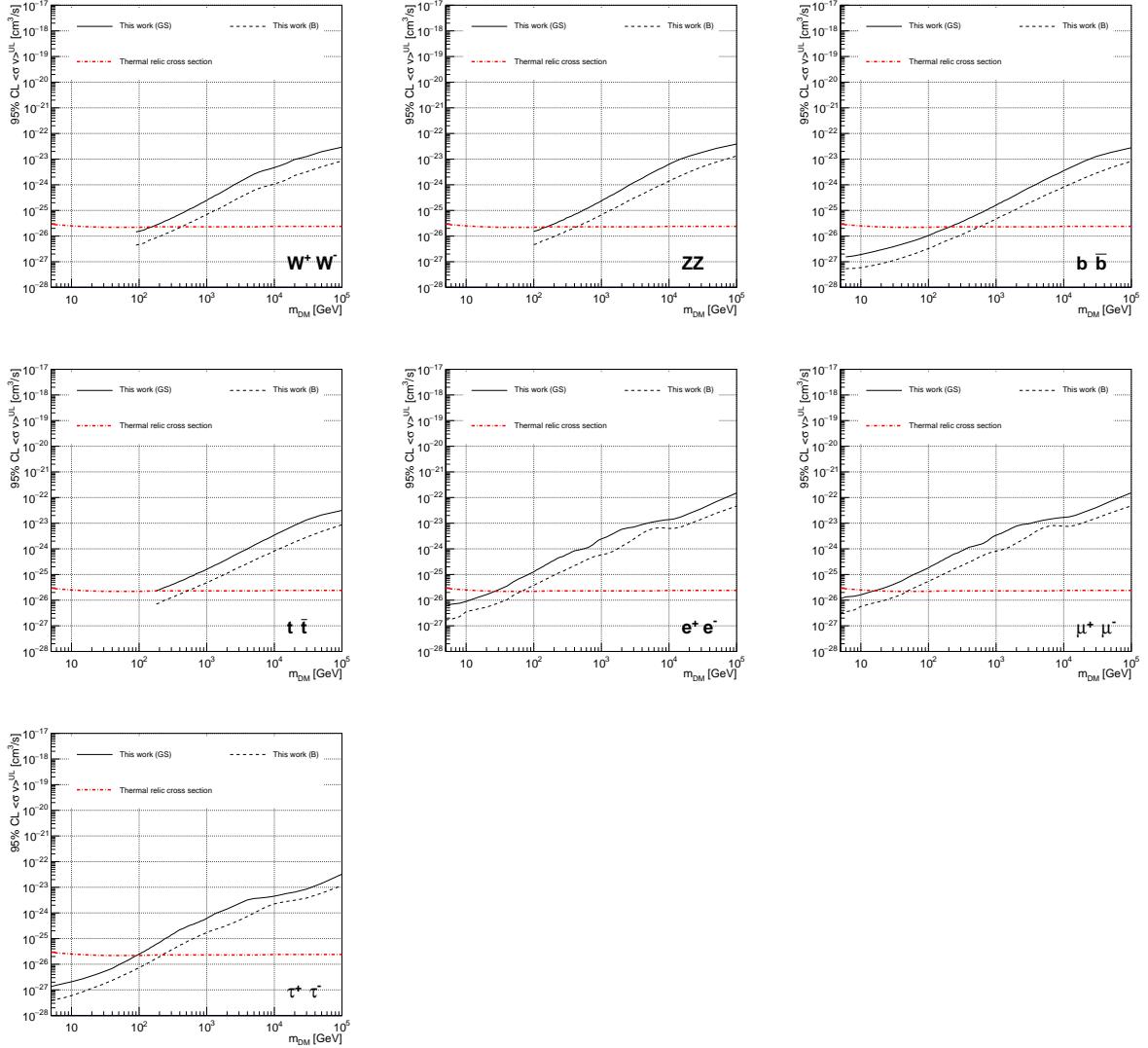


Figure 5.10 Comparisons of the combined limits at 95% confidence level for each of the eight annihilation channels when using the  $J$  factors from Ref. [55] ( $\mathcal{GS}$  set in Table 5.1), plain lines, and the  $J$  factor from Ref. [49, 57] ( $\mathcal{B}$  set in Table 5.1), dashed lines. The cross-section given by the thermal relic is also indicated [56].

## 1039 5.7 HAWC Systematics

### 1040 5.7.1 Inverse Compton Scattering

1041 The DM-DM annihilation channels produce many high energy electrons regardless of the  
 1042 primary annihilation channel. These high energy electrons can produce high energy gamma-rays  
 1043 through Inverse Compton Scattering (ICS). If this effect is strong, it would change the morphology  
 1044 of the source and increase the total expected gamma-ray counts from any source. The PPPC [44]

1045 provides tools in Mathematica for calculating the impact of ICS for an arbitrary location in the  
1046 sky for a specified annihilation channel. We calculated the change in gamma-ray counts for DM  
1047 annihilation to primary  $e\bar{e}$  for RA and Dec corresponding to Segue1 and Coma Berenices. These  
1048 dSphs were chosen because they are the strongest contributors to the limit.  $e\bar{e}$  was selected because  
1049 it would have the largest number of high energy electrons. The effect was found to be on the order  
1050 of  $10^{-7}$  on the gamma-ray spectrum. As a result, this systematic is not considered in our analysis.

### 1051 **5.7.2 Point Source Versus Extended Source Limits**

1052 The previous DM search toward dSph approximated the dSphs as point sources [46]. In  
1053 this analysis, the dSphs are implemented as extended with J-factor distributions following those  
1054 produced by [55]. The resolution of the cited map is much finer than HAWC's angular resolution.  
1055 The vast majority of the J-factor distribution is represented on the central HAWC pixel of the dSph  
1056 spatial map. However, the neighboring 8 pixels are not negligible and contribute to our limit.

1057 Figure 5.11 shows a substantial improvement to the limit for Segue1. Fig. 5.12 however showed  
1058 identical limits. These disparities are best explained by the relative difference in their J-Factors.  
1059 Both dSphs pass almost overhead the HAWC detector, however Segue1 has the larger J-Factor  
1060 between the two. Adjacent pixels to the central pixel will therefore contribute to the limits. This is  
1061 the case for other dSph that are closer to the zenith of the HAWC detector.

1062 Comparison plots for all sources and the combined limit can be found in the sandbox for the  
1063 Glory Duck project.

### 1064 **5.7.3 Impact of Pointing Systematic**

1065 During the analysis it was discovered that directional reconstruction of gamma-rays had a  
1066 systematic bias at large zenith angles. Slides describing this systematic can be found [here](#). Shown  
1067 on the presentation is dependence on the pointing systematic on declination. New spatial profiles  
1068 were generated for every dSph and limits were computed for the adjusted declination.

1069 Section 5.7.3 demonstrates the impact of this systematic for all DM annihilation channels  
1070 studied by HAWC. The impact is a tiny improvement, yet mostly identical, to the combined limits.

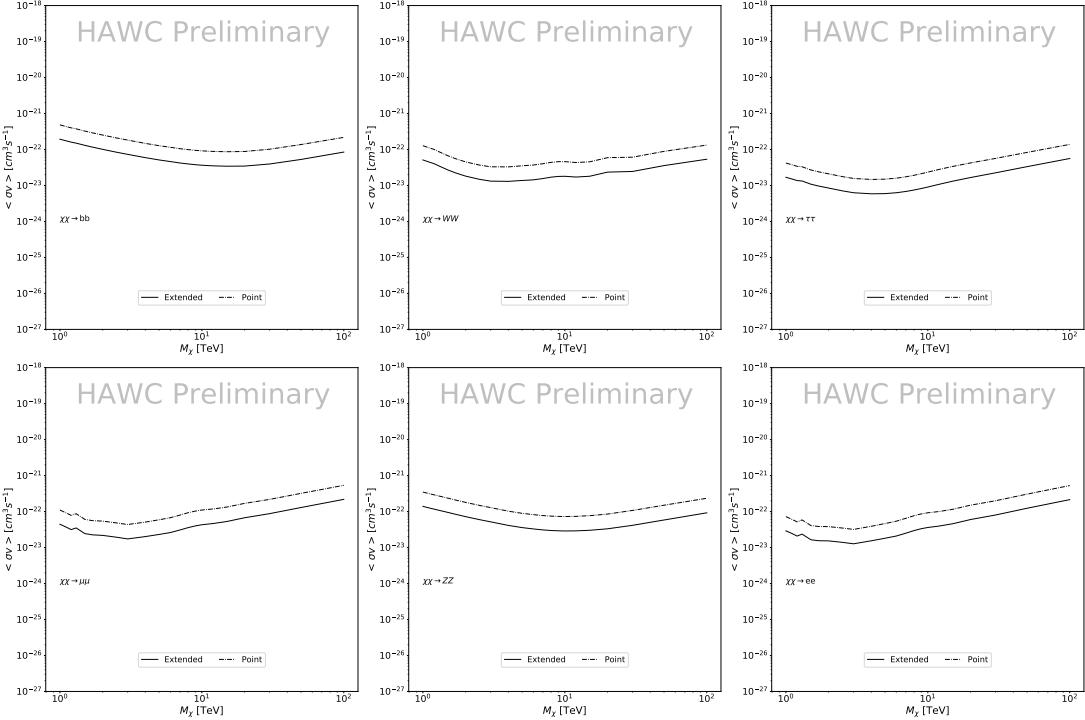


Figure 5.11 Comparisons of the combined limits at 95% confidence level for a point source analysis and extended source using [55]  $\mathcal{GS}$  J-factor distributions and PPPC [44] annihilation spectra. Shown are the limits for Segue1 which will have the most significant impact on the combined limit. 6 of the 7 DM annihilation channels are shown. Solid lines are extended source studies. Dashed lines are point source studies. Overall, the extended source analysis improves the limit by a factor of 2.

## 1071 5.8 J-factor distributions

### 1072 5.8.1 Numerical integration of $\mathcal{GS}$ maps

1073 It was discovered well after the HAWC analysis was completed that the published tables from  
 1074  $\mathcal{GS}$  [45] quoted median  $J$ -factors were computed in a non-trivial manner. The assumption myself  
 1075 and collaborators had been that the published tables represented the  $J$ -factor as a function of  $\theta$  for  
 1076 the best global fit model on a per-source basis. However, this is not the case. Instead, what is  
 1077 published are the best fit model for each dwarf that only considers stars up to the angular separation  
 1078  $\theta$ . Therefore, the model is changing for each value of  $\theta$  for each dwarf. Yet, the introduced features  
 1079 from unique models at each  $\theta$  are much smaller than the angular resolution of HAWC. It is not  
 1080 expected for these effects to impact the limits and TS greatly as a result.

1081 Median  $J$ -factor model profiles were provided by the authors. New maps were generated

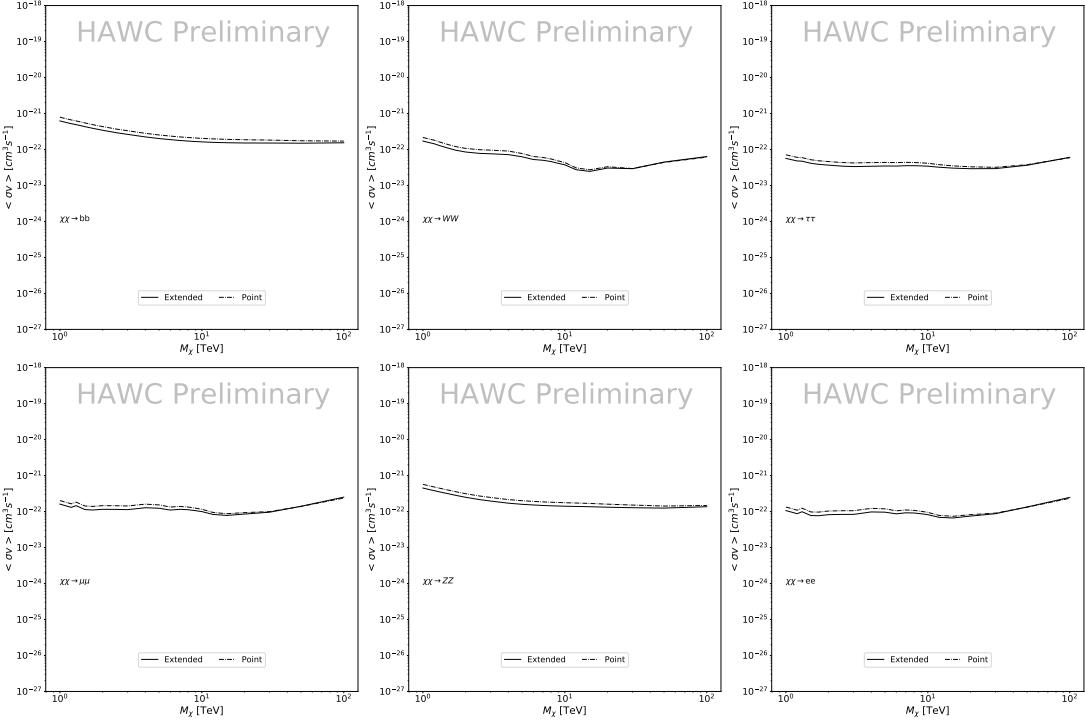


Figure 5.12 Same as Fig. 5.11 on Coma Berenices. This dSph also contributes significantly to the limit. The limits are identical in this case.

and analyzed for Segue1 and Coma Berenices. Figure 5.14 shows the differential between maps generated with the method from Section 5.8.1 and from the authors of [45]. These maps were reanalyzed for all SM DM annihilation channels. Upper limits for these channels are shown in Figure 5.15

From Figure 5.15, we can see that the impact of these model difference was no substantial. The observed impact was a fractional effect which is much smaller than the impact from selecting another DM spatial distribution model as was shown in Figure 5.10.

### 5.8.2 $\mathcal{GS}$ Versus $\mathcal{B}$ spatial models

We show in this appendix a comparison between the  $J$ -factors computed by Geringer-Sameth *et al.* [55] (the  $\mathcal{GS}$  set) and the ones computed by Bonnivard *et al.* [49, 57] (the  $\mathcal{B}$  set). The  $\mathcal{GS}$   $J$ -factors are computed through a Jeans analysis of the kinematic stellar data of the selected dSphs, assuming a dynamic equilibrium and a spherical symmetry for the dSphs. They adopted the generalized DM density distribution, known as Zhao-Hernquist, introduced by [47], carrying

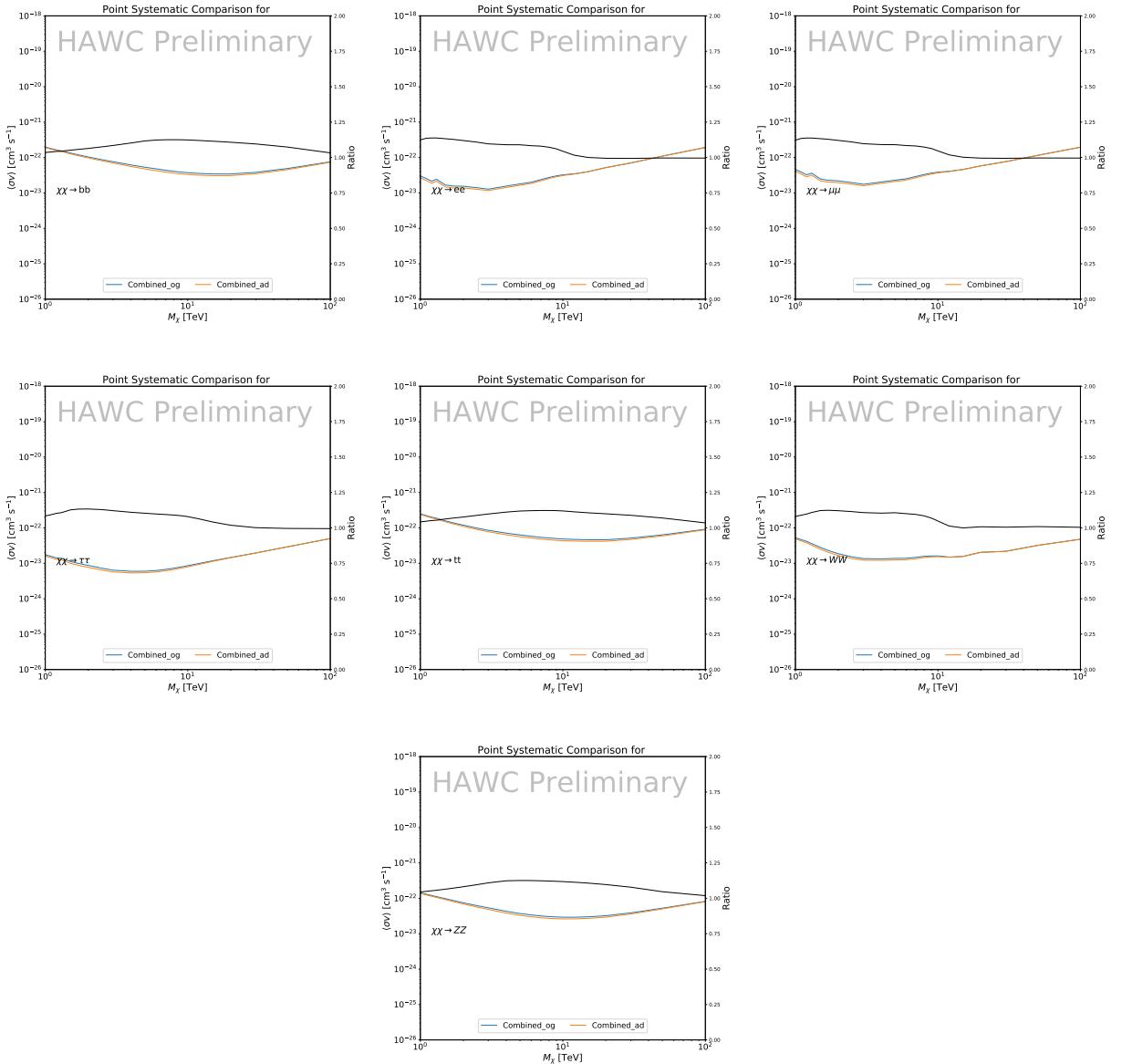


Figure 5.13 Comparison of combined limits when correcting for HAWC's pointing systematic. All DM annihilation channels are shown. The solid black line is the ratio between published limit to the declination corrected limit. The blue solid line or "Combined\_og" represented the limits computed for Glory Duck. The solid orange line or "Combined\_ad" represented the limits computed after correcting for the pointing systematic.

1095 three additional index parameters to describe the inner and outer slopes, and the break of the  
 1096 density profile. Such a profile parametrization allows the reduction of the theoretical bias from  
 1097 the choice of a specific radial dependency on the kinematic data. In other words, the increase of  
 1098 free parameters with the use of the Zhao-Hernquist profile allows a better description of the mass

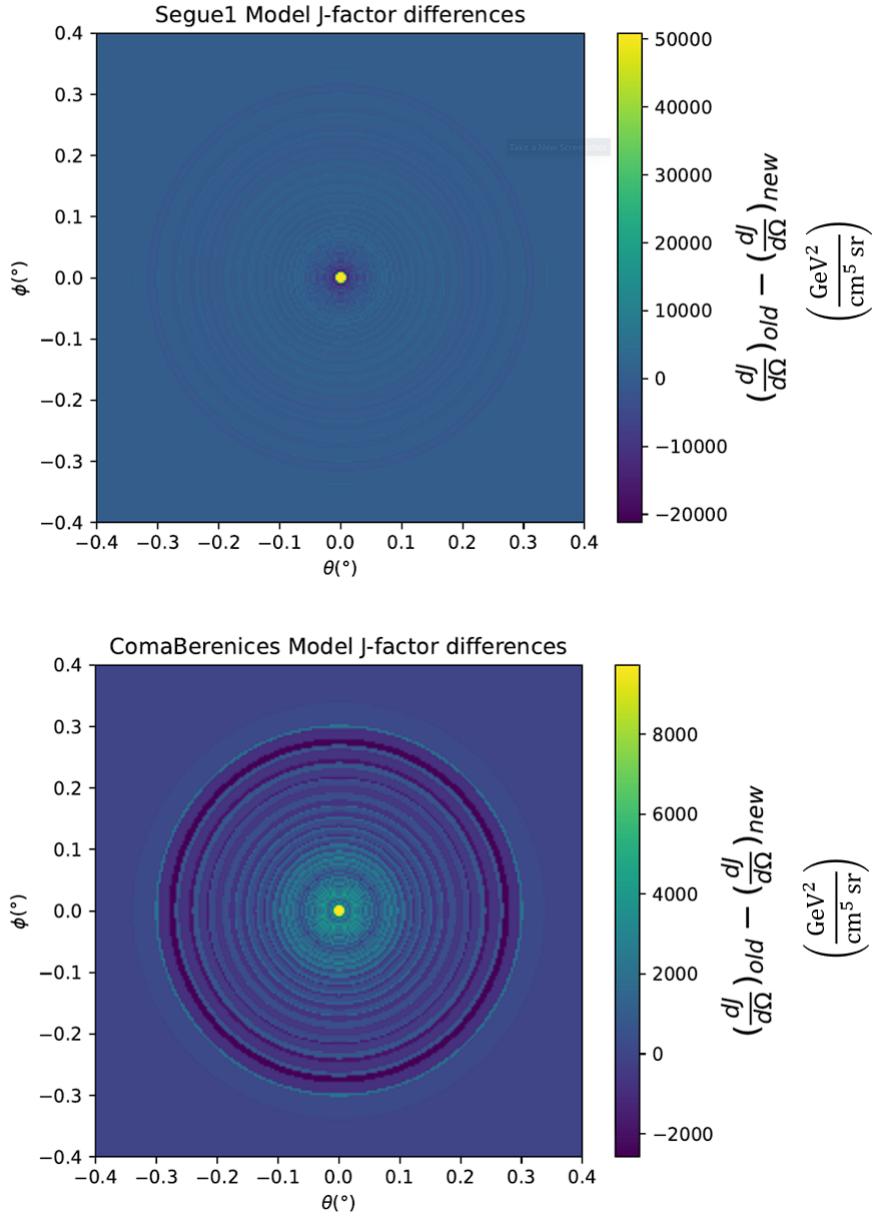


Figure 5.14 Differential map of  $dJ/\Omega$  from model built in Section 5.8.1 and profiles provided directly from authors. (Top) Differential from Segue1. (bottom) Differential from Coma Berenices. Note that their scales are not the same. Segue1 shows the deepest discrepancies which is congruent with its large uncertainties. Both models show anuli where unique models become apparent.

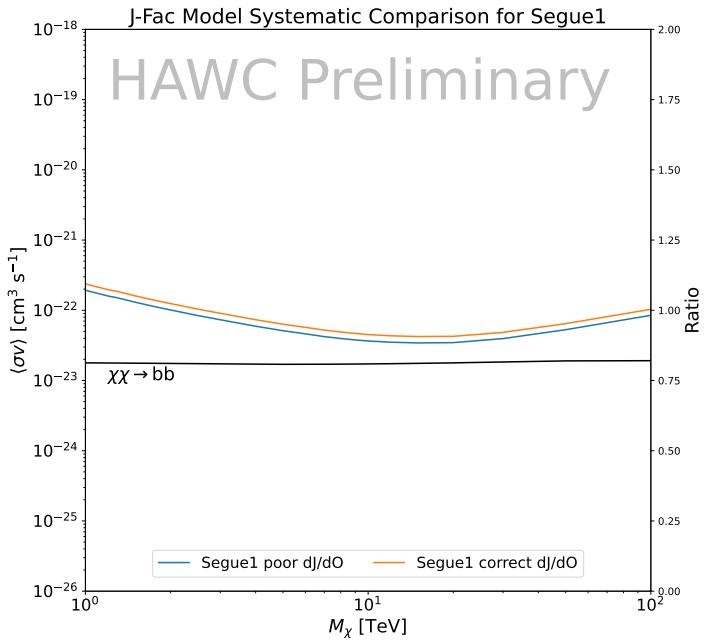
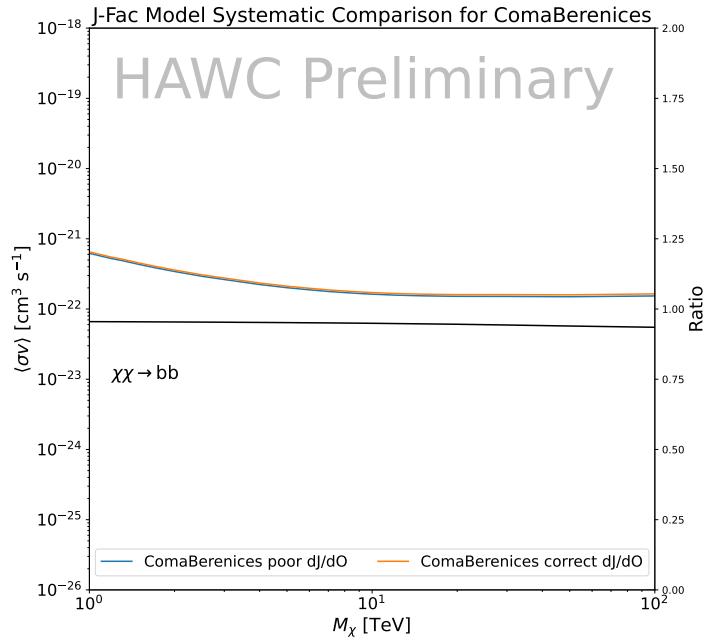


Figure 5.15 HAWC limits for Coma Berenices (top) and Segue1 (bottom) for two different map sets. Blue lines are limits calculated on maps with poor model representation. Orange lines are limits calculated on spatial profiles provided by the authors of [45]. Black line is the ratio of the poor spatial model limits to the corrected spatial models. The left y-axis measures  $\langle \sigma v \rangle$  for the blue and orange lines. The right y-axis measures the ratio and is unitless.

1099 density distribution of dark matter.

1100 In addition, a constant velocity anisotropy profile and a Plummer light profile [60] for the stellar  
1101 distribution were assumed. The velocity anisotropy profile depends on the radial and tangential  
1102 velocity dispersion. However, its determination remains challenging since only the line-of-sight  
1103 velocity dispersion can be derived from velocity measurements. Therefore, the parametrization of  
1104 the anisotropy profile is obtained from simulated halos (see [61] for more details). They provide the  
1105 values of the  $J$ -factors of regions extending to various angular radius up to the outermost member  
1106 star.

1107 The  $\mathcal{B}$   $J$ -factors were computed through a Jeans analysis taking into account the systematic  
1108 uncertainties induced by the DM profile parametrization, the radial velocity anisotropy profile, and  
1109 the triaxiality of the halo of the dwarf galaxies. They performed a more complete study of the dSph  
1110 kinematics and dynamics than  $\mathcal{GS}$  for the determination of the  $J$ -factor. Conservative values of the  
1111  $J$ -factors where obtained using an Einasto DM density profile [62], a realistic anisotropy profile  
1112 known as the Baes & Van Hese profile [63] which takes into account that the inner regions can be  
1113 significantly non-isotropic, and a Zhao-Hernquist light profile [47].

1114 For both sets,  $J$ -factor values are provided for all dSphs as a function of the radius of the  
1115 integration region [55, 49, 57]. Table 5.1 shows the heliocentric distance and Galactic coordinates  
1116 of the twenty dSphs, together with the two sets of  $J$ -factor values integrated up to the outermost  
1117 observed star for  $\mathcal{GS}$  and the tidal radius for  $\mathcal{B}$ . Both  $J$ -factor sets were derived through a Jeans  
1118 analysis based on the same kinematic data, except for Draco where the measurements of [64] have  
1119 been adopted in the computation of the  $\mathcal{B}$  value. The computations for producing the  $\mathcal{GS}$  and  $\mathcal{B}$   
1120 samples differ in the choice of the DM density, velocity anisotropy, and light profiles, for which the  
1121 set  $\mathcal{B}$  takes into account some sources of systematic uncertainties.

1122 Figure 5.16 and Figure 5.17 show the comparisons for the  $J$ -factor versus the angular radius  
1123 for each of the 20 dSphs used in this study. The uncertainties provided by the authors are also  
1124 indicated in the figures. For the  $\mathcal{GS}$  set, the computation stops at the angular radius corresponding  
1125 to the outermost observed star, while for the  $\mathcal{B}$  set, the computation stops at the angular radius

1126 corresponding to the tidal radius.

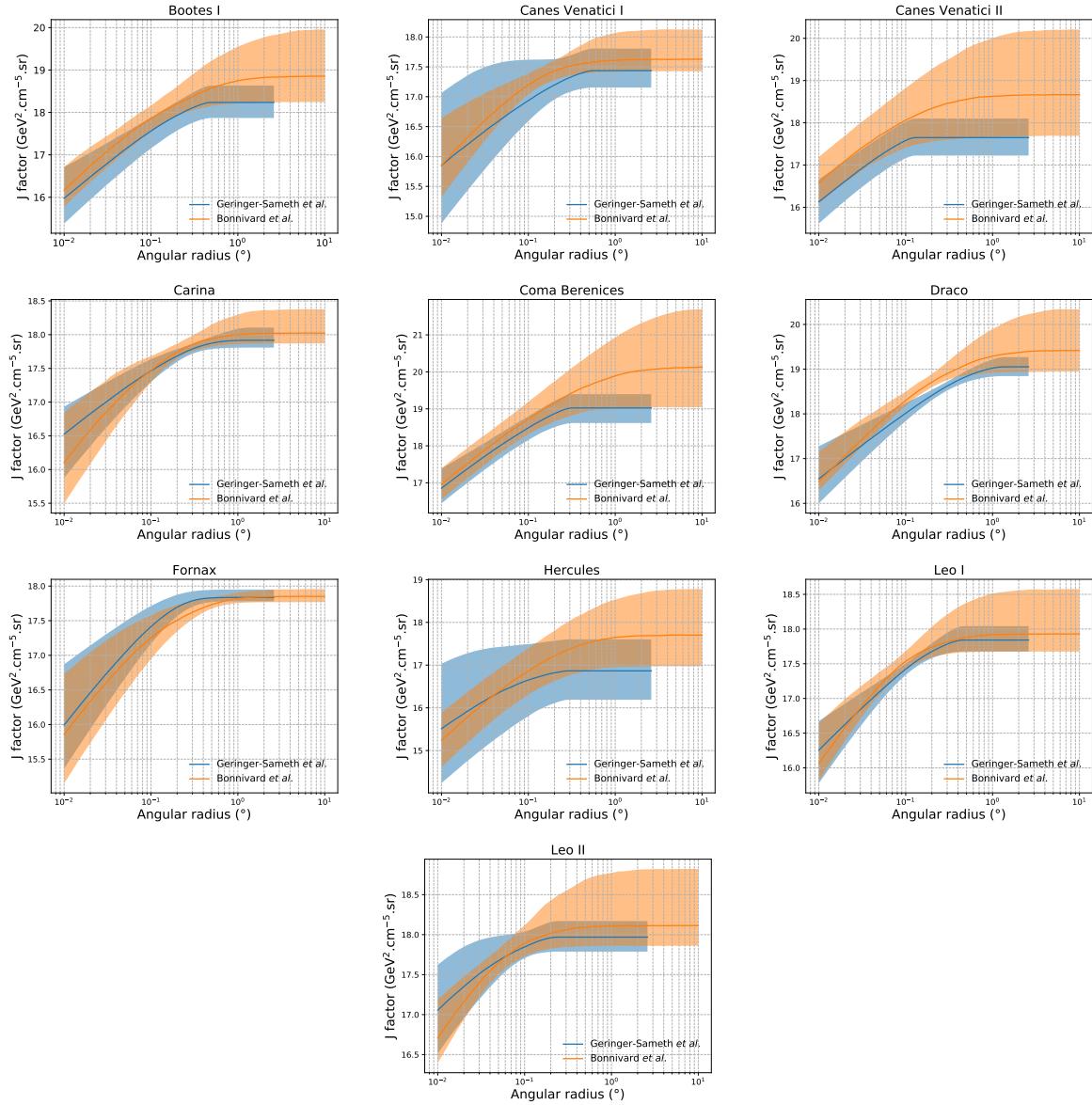


Figure 5.16 Comparisons between the  $J$ -factors versus the angular radius for the computation of  $J$  factors from Ref. [55] ( $\mathcal{GS}$  set in Table 5.1) in blue and for the computation from Ref. [49, 57] ( $\mathcal{B}$  set in Tab. 5.1). The solid lines represent the central value of the  $J$ -factors while the shaded regions correspond to the  $1\sigma$  standard deviation.

## 1127 5.9 Discussion and Conclusions

1128 In this multi-instrument analysis, we have used observations of 20 dSphs from the gamma-ray  
 1129 telescopes Fermi-LAT, H.E.S.S., MAGIC, VERITAS, and HAWC to perform a collective DM  
 1130 search annihilation signals. The data were combined across sources and detectors to significantly

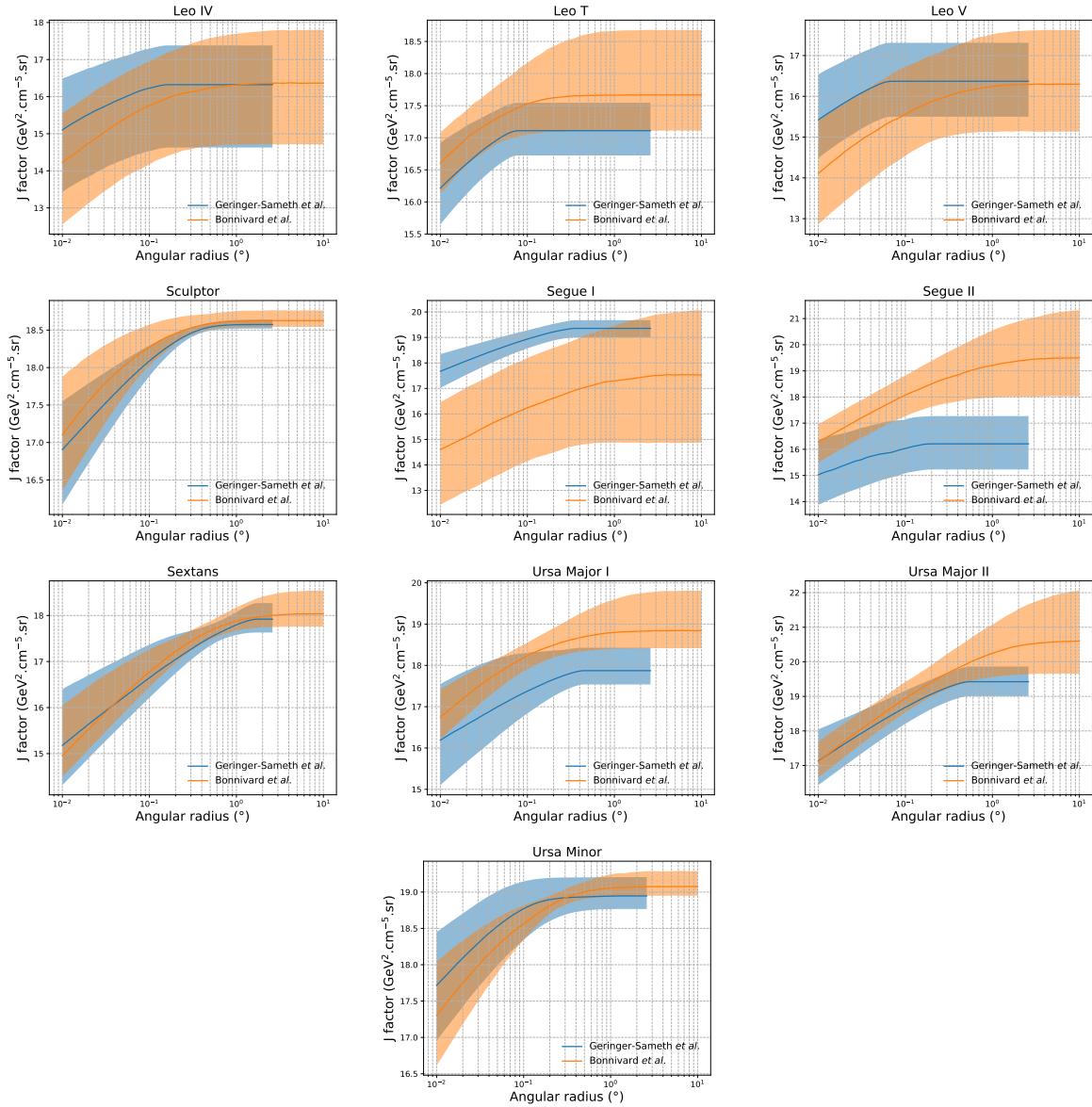


Figure 5.17 Comparisons between the  $J$ -factors versus the angular radius for the computation of  $J$  factors from Ref. [55] ( $\mathcal{GS}$  set in Tab. 5.1) in blue and for the computation from Ref. [49, 57] ( $\mathcal{B}$  set in Tab. 5.1) in orange. The solid lines represent the central value of the  $J$ -factors while the shaded regions correspond to the  $1\sigma$  standard deviation.

increase the sensitivity of the search. We have observed no significant deviation from the null, no DM hypothesis, and so present our results in terms of upper limits on the annihilation cross-section for seven potential DM annihilation channels.

Fermi-LAT brings the most stringent constraints for continuum channels below approximately 1 TeV. The remaining detectors dominate at higher energies. Overall, for multi-TeV DM mass,

1136 the combined DM constraints from all five telescopes are 2-3 times stronger than any individual  
1137 telescope for multi-TeV DM.

1138     Derived from observations of many dSphs, our results produce robust limits given the DM  
1139 content of the dSphs is relatively well constrained. The obtained limits span the largest mass  
1140 range of any WIMP DM search. Our combined analysis improves the sensitivity over previously  
1141 published results from each detector which produces the most stringent limits on DM annihilation  
1142 from dSphs. These results are based on deep exposures of the most promising known dSphs with  
1143 the currently most sensitive gamma-ray instruments. Therefore, our results constitute a legacy of  
1144 a generation of gamma-ray instruments on WIMP DM searches towards dSphs. Our results will  
1145 remain the reference in the field until a new generation of more sensitive gamma-ray instruments  
1146 begin operations, or until new dSphs with higher  $J$ -factors are discovered.

1147     This analysis serves as a proof of concept for future multi-instrument and multi-messenger  
1148 combination analyses. With this collaborative effort, we have managed to sample over four orders  
1149 in magnitude in gamma-ray energies with distinct observational techniques. Determining the nature  
1150 of DM continues to be an elusive and difficult problem. Larger datasets with diverse measurement  
1151 techniques could be essential to tackling the DM problem. A future collaboration using similar  
1152 techniques as the ones described in this paper could grow even beyond gamma rays. The models we  
1153 used for this study include annihilation channels with neutrinos in the final state. Advanced studies  
1154 could aim to merge our results with those from neutrino observatories with large data sets. Efforts  
1155 are already underway to add data from the IceCube, ANTARES, and KM3NeT observatories to  
1156 these gamma-ray results.

1157     From this work, a selection of the best candidates for observations, according to the latest  
1158 knowledge on stellar dynamics and modelling techniques for the derivation of the  $J$ -factors on  
1159 the potential dSphs targets, is highly desirable at the time that new experiments are starting their  
1160 dark matter programs using dSphs. Given the systematic uncertainty inherent to the derivation of  
1161 the  $J$ -factors, an informed observational strategy would be to select both objects with the highest  
1162  $J$ -factors that could lead to DM signal detection, and objects with robust  $J$ -factor predictions, i.e.

1163 with kinematic measurements on many bright stars, which would strengthen the DM interpretation  
1164 reliability of the observation outcome.

1165 This analysis combines data from multiple telescopes to produce strong constraints on astro-  
1166 physical objects. From this perspective, these methods can be applied beyond just DM searches.  
1167 Almost every astrophysical study can benefit from multi-telescope, multi-wavelength gamma-ray  
1168 studies. We have enabled these telescopes to study the cosmos with greater precision and detail.  
1169 Many astrophysical searches can benefit from multi-instrument gamma-ray studies, for which our  
1170 analysis lays the foundation.

## CHAPTER 6

### MULTITHREADING HAWC ANALYSES FOR DARK MATTER SEARCHES

#### 6.1 Introduction

HAWC's current software suite, plugins to 3ML and HAL [54, 42], do not fully utilize computational advancements of recent decades. Said advancements include the proliferation of Graphical Processing Units (GPUs), and multithreading on multicore processors. The analysis described in chapter 5 took up to 3 months of wall time waiting for the full gambit of data analysis and simulation of background to compute. Additionally, with the updated 2D energy binning scheme,  $f_{\text{hit}}$  and Neural Network (NN), the time needed to compute expected to grow. Although excessive computing time was, in part, from an intense use of a shared computing cluster, it was evident that there was room for improvement. In HAWC's next generation dSph DM search, I decided to develop codes that would utilize the multicore processors on modern high performance computing clusters. The results of this work are featured in this chapter and brought a human timing improvement to computation that scales approximately as  $1/N$  where  $N$  is the number of threads.

#### 6.2 Dataset and Background

This section enumerates the data and background methods used for HAWC's multithreaded study of dSphs. Section 6.2.1 and Section 6.2.2 are most useful for fellow HAWC collaborators looking to replicate a multithreaded dSph DM search.

##### 6.2.1 Itemized HAWC files

These files are only available within HAWC's internal documentation and collaborators. They are not meant for public access, and are presented here so that HAWC collaborators can reproduce results accurately.

- Detector Resolution: `refit-Pass5-Final-NN-detRes-zenith-dependent.root`
- Data Map: `Pass5-Final-NN-maptree-ch103-ch1349-zenith-dependent.root`
- Spectral Dictionary: `HDMspectra_dict_gamma.npy`

1195 **6.2.2 Software Tools and Development**

1196 This analysis was performed using HAL and 3ML [42, 43] in Python3. I built software  
1197 in collaboration with Michael Martin and Letrell Harris to implement the *Dark Matter Spectra*  
1198 *from the Electroweak to the Planck Scale* (HDM) [65] and dSphs spatial model from [66] for  
1199 HAWC analysis. A NumPy dictionary of HDM, `HDMspectra_dict_gamma.npy`, was made for  
1200 portability within the collaboration. These dictionaries were generated from the [git repository](#) [65].  
1201 The analysis was performed using the Neural Network energy estimator for Pass 5.F. A description  
1202 of this estimator was provided in chapter 3. [TODO: Define a subsection when it's written](#), and its  
1203 key, relevant improvements are an improved energy estimation and improved sensitivities at higher  
1204 zenith angles. All other software used for data analysis, DM profile generation, and job submission  
1205 to SLURM are also kept in my sandbox in the [Dark Matter HAWC](#) project. The above repository  
1206 also incorporates the model inputs used previously in Glory Duck, described in chapter 5, so Glory  
1207 Duck remains compatible with modern software.

1208 **6.2.3 Data Set and Background Description**

1209 The HAWC data maps used for this analysis contain 2565 days of data between runs 2104 and  
1210 7476. They were generated from pass 5.f reconstruction. The analysis is performed using the NN  
1211 energy estimator with bin list:

1212     B1C0Ea, B1C0Eb, B1C0Ec, B1C0Ed, B1C0Ee, B2C0Ea, B2C0Eb, B2C0Ec, B2C0Ed, B2C0Ee,  
1213     B3C0Ea, B3C0Eb, B3C0Ec, B3C0Ed, B3C0Ee, B3C0Ef, B4C0Eb, B4C0Ec, B4C0Ed, B4C0Ee,  
1214     B4C0Ef, B5C0Ec, B5C0Ed, B5C0Ee, B5C0Ef, B5C0Eg, B6C0Ed, B6C0Ee, B6C0Ef, B6C0Eg,  
1215     B6C0Eh, B7C0Ee, B7C0Ef, B7C0Eg, B7C0Eh, B7C0Ei, B8C0Ee, B8C0Ef, B8C0Eg, B8C0Eh,  
1216     B8C0Ei, B8C0Ej, B9C0Ef, B9C0Eg, B9C0Eh, B9C0Ei, B9C0Ej, B10C0Eg, B10C0Eh,  
1217     B10C0Ei, B10C0Ej, B10C0Ek, B10C0El

1218 Bin 0 was excluded as it has substantial hadronic contamination and poor angular resolution.

1219 Background considerations and source selection was identical to Section 5.2.3, and no additional  
1220 arguments are provided here. Many of the HAWC systematics explored in Section 5.7 also apply

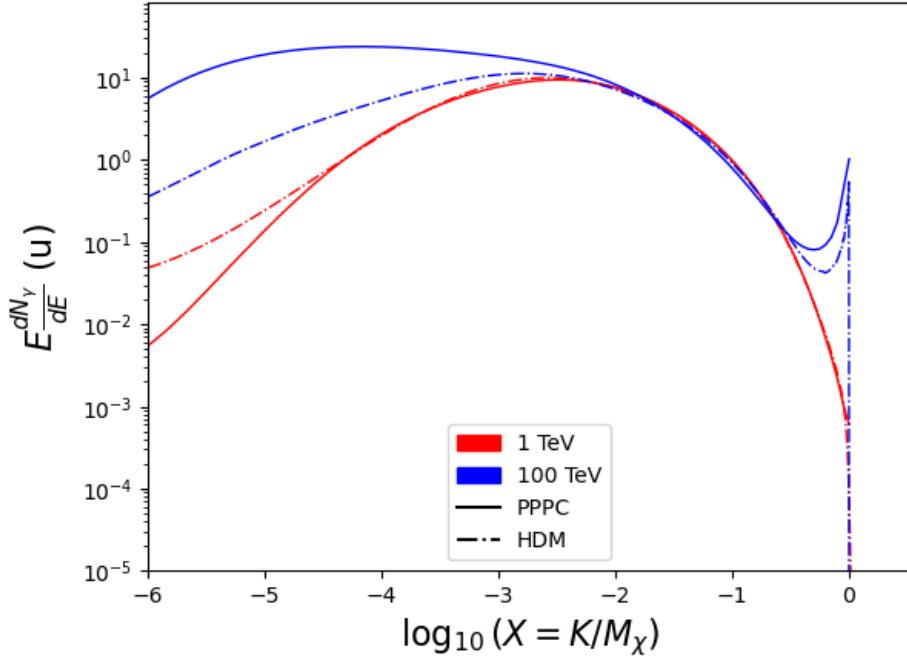


Figure 6.1 Spectral hypotheses from PPPC [44] and HDM [65] for DM annihilation:  $\chi\chi \rightarrow W^-W^+$ . Solid lines are spectral models with EW corrections from the PPPC. Dash-dot lines are spectral models from HDM. Red lines are models for  $M_\chi = 1$  TeV. Blue lines represent models for  $M_\chi = 100$  TeV.

1221 for this DM search and are not added upon here.

### 1222 6.3 Analysis

1223 The analysis and its systematics are almost identical to Section 5.3. Importantly, we use the  
 1224 same **TODO: fix this ref** Equation (7.1) and Equation (5.2) for estimating the gamma-ray flux at  
 1225 HAWC from our sources.

#### 1226 6.3.1 $\frac{dN_\gamma}{dE_\gamma}$ - Particle Physics Component

1227 For these spectra, we import HDM with Electroweak (EW) corrections and additional correc-  
 1228 tions for neutrinos above the EW scale [65]. The spectra are implemented as a model script in  
 1229 astromodels for 3ML. A comprehensive description of EW corrections and neutrino considerations  
 1230 are provided later in Sec. 8.

1231 Figure 6.1 demonstrates the impact of changes implemented in HDM on DM annihilation to W  
 1232 bosons. A class in astromodels was developed to include HDM and is aptly named **HDMspectra**  
 1233 within `DM_models.py`. The SM DM annihilation channels studied here are  $\chi\chi \rightarrow$ :

1234       $e^+e^-$ ,  $\mu^+\mu^-$ ,  $\tau^+\tau^-$ ,  $b\bar{b}$ ,  $t\bar{t}$ ,  $gg$ ,  $W^+W^-$ ,  $ZZ$ ,  $c\bar{c}$ ,  $u\bar{u}$ ,  $d\bar{d}$ ,  $s\bar{s}$ ,  $\nu_e\bar{\nu}_e$ ,  $\nu_\mu\bar{\nu}_\mu$ ,  $\nu_\tau\bar{\nu}_\tau$ ,  $\gamma\gamma$ ,  $hh$ .

1235    For  $\gamma\gamma$  and  $ZZ$ , a substantial fraction of the signal photons are expected to have  $E_\gamma = m_\chi$  [65].  
 1236    This introduces  $\delta$ -function that is much narrower than the energy resolution of the HAWC detector.  
 1237    To ensure that this feature is not lost in the likelihood fits, the 'line' feature is convolved with a  
 1238    Gaussian kernel with a  $1\sigma$  width of  $0.05 \cdot m_\chi$  and total kernel window of  $\pm 4\sigma$ . This differs from  
 1239    HAWC's previous line study where 30% of HAWC's energy resolution was used for the kernel [67].  
 1240    The NN energy estimator's strength compared to  $f_{\text{hit}}$  at low gamma-ray energy enables narrower  
 1241    kernels [65].  $\chi\chi \rightarrow \gamma\gamma$  and  $ZZ$  spectral hypotheses are shown in Figure 6.2. We did not explore  
 1242    how well we reconstruct injected signal events for various kernels widths. This is a systematic  
 1243    that should be tested before publication to journal. Spectral models for the remaining annihilation  
 1244    channels are plotted for each  $m_\chi$  in Figure B.1.

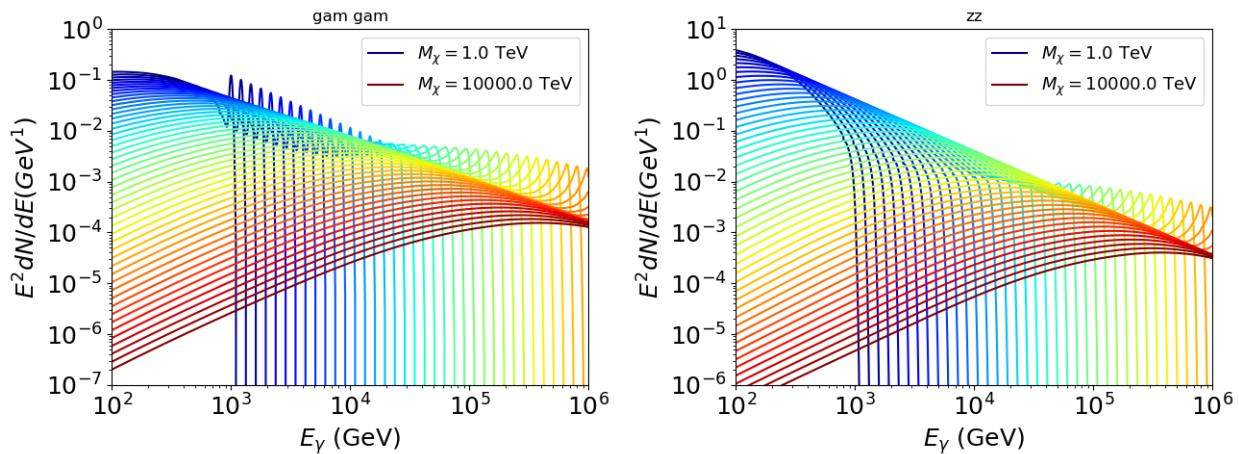


Figure 6.2 Photon spectra for  $\chi\chi \rightarrow \gamma\gamma$  (left) and  $\chi\chi \rightarrow ZZ$  (right) after Gaussian convolution of line features. Both spectra have  $\delta$ -features at photon energies equal to the DM mass. Bluer lines are annihilation spectra with lower DM mass. Redder lines are spectra from larger DM mass. All spectral models are sourced from the Heavy Dark Matter models [65]. Axes are drawn roughly according to the energy sensitivity of HAWC.

1245    **6.3.2 J Astrophysical Components**

1246    The J-factor profiles for each source are imported from Louis Strigari et al. (referred to with  
 1247     $\mathcal{LS}$ ) [66]. The  $\mathcal{LS}$  catalog fits a Navarro–Frenk–White (NFW) [48] spatial DM distributions to

1248 the dSphs which has a DM density of

$$\rho(r) = \frac{\rho_0}{\frac{r}{R_s} \left(1 + \frac{r}{R_s}\right)^2}. \quad (6.1)$$

1249  $\rho_0$  and the scale radius,  $R_s$  are free parameters fit for each dSph.  $r$  is the distance from the center  
1250 of the dSph.

1251 Profiles in  $\frac{dJ}{d\Omega}(\theta)$  up to an angular separation  $\theta = 0.5^\circ$  were provided directly from the authors.

1252 Map generation from these profiles were almost identical to Section 5.3.2 except that a higher order  
1253 trapezoidal integral was used for the normalization of the square, uniformly-spaced map:

$$p^2 \cdot \sum_{i=0}^N \sum_{j=0}^M w_{i,j} \frac{dJ}{d\Omega}(\theta_{i,j}, \phi_{i,j}) \quad (6.2)$$

1254  $p$  is the angular side of one pixel in the map.  $w_{i,j}$  is a weight assigned the following ways:

1255  $w_{i,j} = 1$  if  $(\theta_{i,j}, \phi_{i,j})$  is fully within the region of integration

1256  $w_{i,j} = 1/2$  if  $(\theta_{i,j}, \phi_{i,j})$  is on an edge of the region of integration

1257  $w_{i,j} = 1/4$  if  $(\theta_{i,j}, \phi_{i,j})$  is on a corner of the region of integration

1258 Figure 6.3 shows the median and  $\pm 1\sigma$  maps used as input for this DM annihilation study.

### 1259 6.3.3 Source Selection and Annihilation Channels

1260 HAWC's sources for this multithreaded analysis include Coma Berenices, Draco, Segue 1, and  
1261 Sextans.  $\mathcal{LS}$  observed up to 43 sources in its publication, however only 4 of the best fit profiles  
1262 were provided at the time this thesis was written. A full description of each source used in this  
1263 analysis is found in Table 6.1.

1264 This analysis improves on chapter 5 in the following ways. Previously, the particle physics  
1265 model used for gamma-ray spectra from DM annihilation was from the PPPC [44] which missed  
1266 important considerations relevant for the neutrino sector. HDM is used to account for this shortfall  
1267 [65]. HDM also models DM to the Planck scale which permits HAWC to probe PeV scale DM. For  
1268 this study, we sample DM masses: 1 TeV - 10 PeV with 6 mass bins per decade in DM mass. In

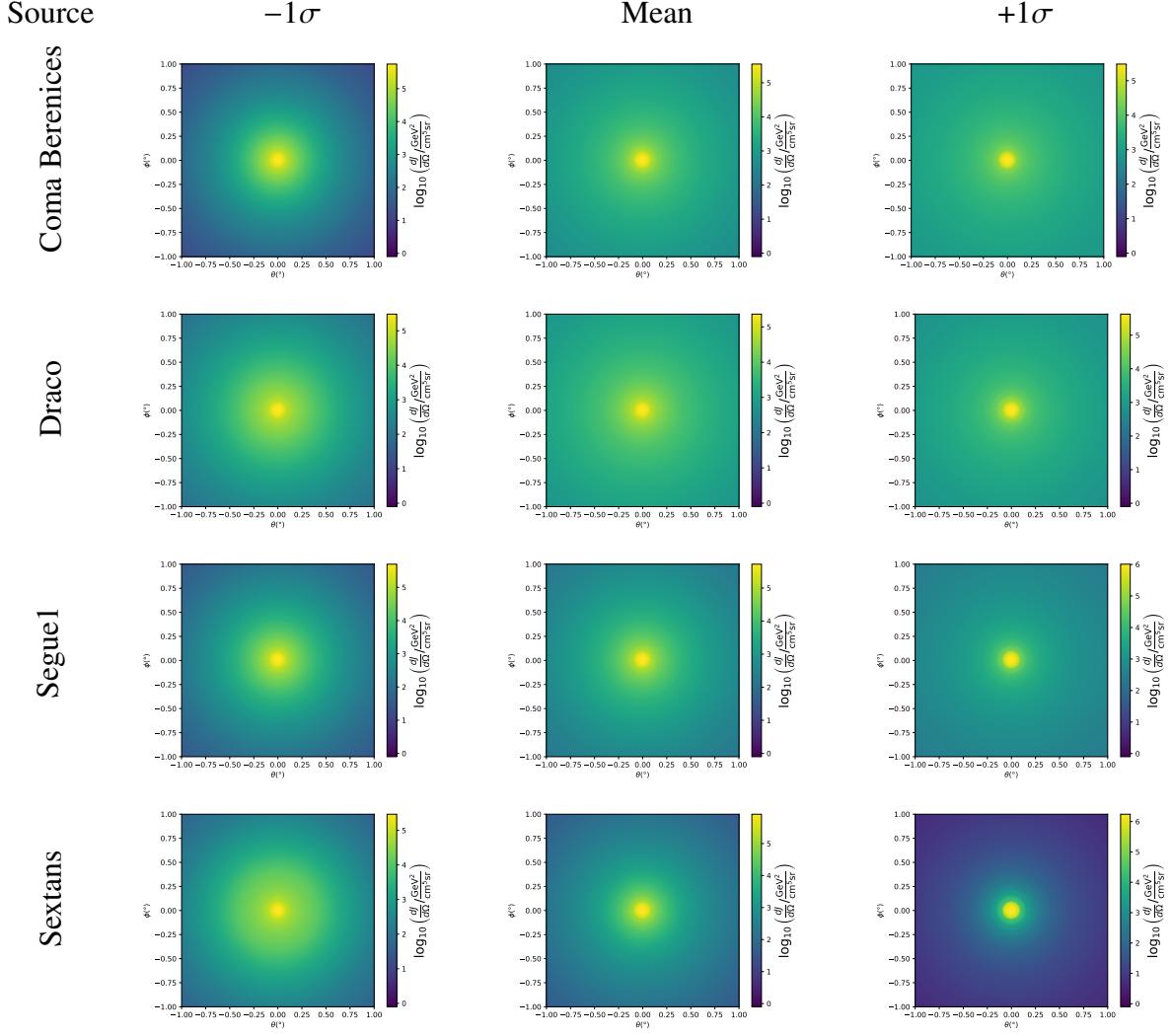


Figure 6.3  $\frac{dJ}{d\Omega}$  maps for Coma Berenices, Draco, Segue1, and Sextans. Columns are divided for the  $\pm 1\sigma$  uncertainties in  $dJ/d\Omega$  around the mean value from  $\mathcal{LS}$  [66]. Origin is centered on the specific dwarf spheroidal galaxies (dSph).  $\theta$  and  $\phi$  axes are the angular separation from the center of the dwarf. Profiles are truncated at  $1^\circ$  and flattened beyond.

1269 the case of line spectra ( $\chi\chi \rightarrow \gamma\gamma$ , or  $ZZ$ ), we double the mass binning to 12 DM mass bins per  
 1270 decade in DM mass.

1271  $\mathcal{LS}$  provides 25 sources within HAWC's field of view. Additionally, NFW [48] DM distributions  
 1272 have fewer parameters than Zhao [47], so  $\mathcal{LS}$  fits ultra-faint dwarves which expands the number of  
 1273 sources. However, all sources were not provided by the authors in time for the completion of this  
 1274 dissertation. Finally, the gamma-ray ray dataset is much larger. The study performed here analyzes  
 1275 2565 days of data compared to 1017 days analyzed in chapter 5.

1276 **6.4 Likelihood Methods**

1277 These are identical to Section 5.4.1 and no additional changes are made to the likelihood. Bins  
 1278 in this analysis are expanded to include HAWC’s NN energy estimator.

1279 **6.5 Computational Methods: Multithreading**

1280 Previously, as in Section 5.3, the likelihood was minimized for one model at a time. One model  
 1281 in this case representing a DM annihilation channel (CHAN), DM mass ( $m_\chi$ ), and dSph ((SOURCE)).  
 1282 In an effort to conserve human and CPU time, jobs submitted for high performance computing  
 1283 contained a list of  $m_\chi$  to iterate over for likelihood fitting. Jobs were then trivially parallelized  
 1284 for each permutation of the two lists: CHANS and SOURCES. The lists for CHANS and SOURCES are  
 1285 found in Section 6.3.1 and Table 6.1, respectively. Initially, 11  $m_\chi$  were serially sampled for one  
 1286 job defined by a [CHAN, SOURCE] tuple. Computing the likelihoods would take between 1.5 to 2 hrs,  
 1287 stochastically, for a job. We expect to compute likelihoods for data and 300 Poisson background  
 1288 trials. The estimated CPU time based on the above for all CHAN (N = 17) and SOURCE (M = 25)  
 1289 was estimated to 127,925 jobs. In total, 1,407,175 likelihood fits and profiles would be computed  
 1290 for the 11 mass bins we wished to study. The estimated CPU time ranged between 8k CPU days  
 1291 to 10k CPU days. Human time is more challenging to estimate as job allocation is stochastic and  
 1292 highly dependent on what other users are submitting. Yet, it is unlikely that all jobs would run  
 1293 simultaneously. Therefore, we can expect human time to be about as long as was seen in chapter 5

Name	Distance (kpc)	$l, b$ ( $^\circ$ )	$\log_{10} J$ ( $\mathcal{LS}$ set) $\log_{10}(\text{GeV}^2 \text{cm}^{-5} \text{sr})$
Coma Berenices	44	241.89, 83.61	$19.00^{+0.36}_{-0.35}$
Draco	76	86.37, 34.72	$18.83^{+0.12}_{-0.12}$
Segue I	23	220.48, 50.43	$19.12^{+0.49}_{-0.58}$
Sextans	86	243.50, 42.27	$17.73^{+0.13}_{-0.12}$

Table 6.1 Summary of the relevant properties of the dSphs used in the present work. Column 1 lists the dSphs. Columns 2 and 3 present their heliocentric distance and galactic coordinates, respectively. Column 4 reports the  $J$ -factors of each source given from the  $\mathcal{LS}$  studies and estimated  $\pm 1\sigma$  uncertainties. The values  $\log_{10} J$  ( $\mathcal{LS}$  set) [66] correspond to the mean  $J$ -factor values for a source extension truncated at  $0.5^\circ$ .

1294 which was on the order of months to fully compute on a smaller analysis. A visual aid to describe  
 1295 how jobs were organized is provided in Figure 6.4.

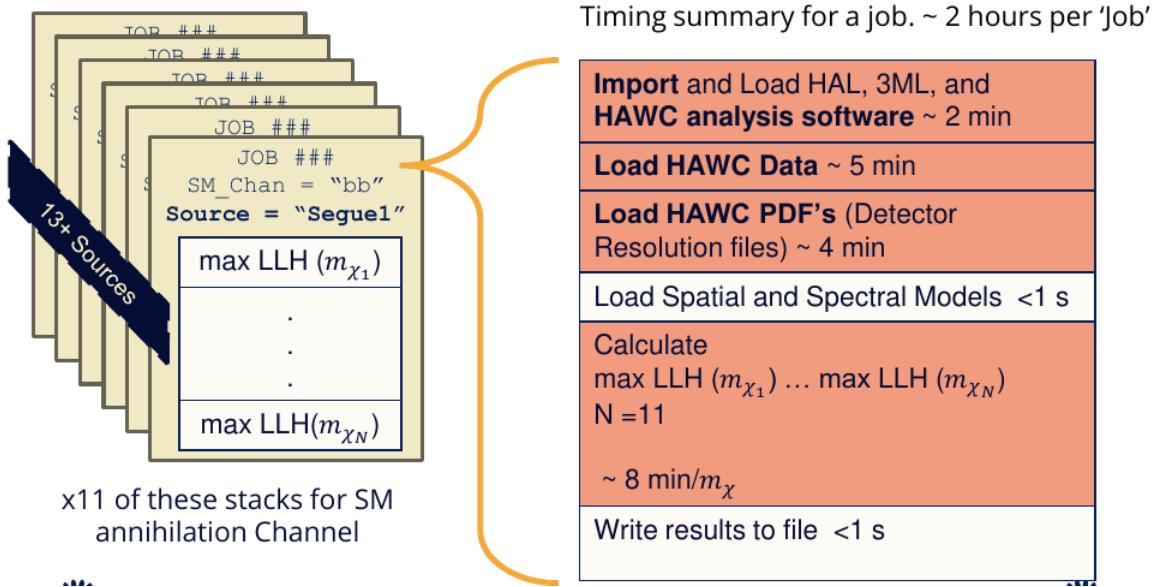


Figure 6.4 Infographic on how jobs and DM computation was organized in Section 5.3. Jobs were built for each permutation of CHANS and SOURCES shown by the left block in the figure. Each job, which took on the order 2 hrs to compute, had the following work flow: 1. Import HAWC analysis software, 2 min to run. 2. Load HAWC count maps, 5 min to run. 3. Load HAWC energy and spatial resolutions, 4 min. 5. Load DM spatial source templates and spectral models, less than 1 s. 6. Perform likelihood fit on data and model, about 8 min per DM mass. 7. Write results to file, less than 1s.

1296 The computational needs for this next generation DM analysis are extreme and is unlike other  
 1297 analyses performed on HAWC. It became clear that there was a lot to gain from optimizing how  
 1298 the likelihoods are computed. This section discusses how multi-threading was applied to solve and  
 1299 reduce HAWC's computing of likelihoods for large parameter spaces like in DM searches.

### 1300 6.5.1 Relevant Foundational Information

1301 The profiling of the likelihood for HAWC is done via gradient descent where the normalization  
 1302 of Equation (7.1) (linearly correlated with  $\langle \sigma v \rangle$ ) is rescaled in the descent. Additionally, we sample  
 1303 the likelihood space for a defined list of  $\langle \sigma v \rangle$ 's described in Section 5.4.2. The time to compute  
 1304 these values is not predictable or consistent because many variables can change across the full  
 1305 model-space. Comprehensively, these variables are:

1306       •  $m_\chi$  : DM rest mass  
 1307       • CHAN : DM annihilation channel in SM.  
 1308       • SOURCE : dSph. Involves a spatial template AND coordinate in HAWC data.  
 1309       •  $\langle\sigma v\rangle$ : Effectively the flux normalization and free parameter in the likelihood fit.  
 1310      Therefore, we develop an asynchronous, functional-parallel coding pattern. Asynchronous meaning  
 1311      the instructions within a function are independent and permitted to be out of sync with sibling  
 1312      computations. Functional-parallel meaning that instructions are the subject of parallelization  
 1313      rather than threading the likelihood computation. This is close to trivial parallelization seen in  
 1314      Figure 6.4 except that we seek to consolidate the loading stages (software, data, and detector  
 1315      resolution loading). Multiple asynchronous threads are expected to reduce total serial processing  
 1316      time and total overhead across the entire project in addition to saving human time.

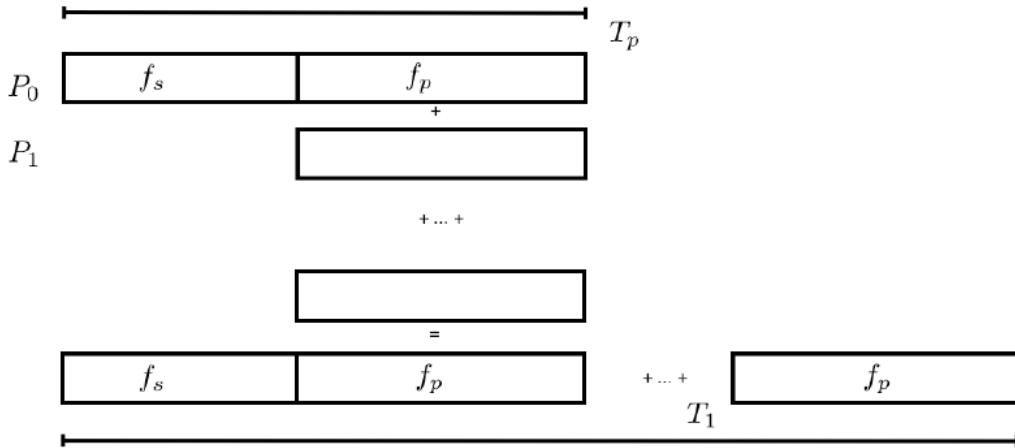


Figure 6.5 Graphic of Gustafson parallel coding pattern.  $f_s$  is the fraction of a program, in time, spent on serial computation.  $f_p$  is the fraction of computing time that is parallelizable.  $T_p$  is the total time for a parallel program to run.  $T_1$  is the total time for a parallel program to run if only 1 processor is allocated.  $P_N$  is the  $N$ -th processor where it's row is the computation the processor performs. The Gustafson pattern is most similar to what is implemented for this analysis. Figure is pulled from [68].

1317      We need a way to measure and compare the expected speedup and efficiency gain for this  
 1318      asynchronous coding pattern. I pull inspiration for timing measurement from [68] and use *Amdahl's*

1319 law with hybrid programming. Hybrid programming meaning that the computation is a mix of  
 1320 distributed and shared memory programming. If we assume the code is fully parallelizable over  $p$   
 1321 processors and  $c$  threads, the ideal speedup is simply  $pc$ , and ideal run-time is  $T_1/(pc)$ .  $T_1$  is the  
 1322 total time for a parallel program to run if only 1 processor is allocated. However, the coding pattern  
 1323 contains some amount of unavoidable serial computation, as shown in Figure 6.5. In our case, the  
 1324 run time,  $T_{p,c}$ , is estimated to be

$$T_{p,c} = \frac{T_1}{pc} (1 + F_s(c - 1)). \quad (6.3)$$

1325  $F_s$  is the fraction of CPU time dedicated to serial computation. The expected speedup,  $S_{p,c}$ , is

$$S_{p,c} \equiv \frac{T_1}{T_{p,c}} = \frac{pc}{1 + F_s(c - 1)}. \quad (6.4)$$

1326 From Equation (6.4), we can see that the speed-up scales with  $p/F_s$ . We are free to minimize  $F_s$   
 1327 asymptotically by enlarging the total models that are submitted to the thread pool, thereby shrinking  
 1328 the CPU fraction dedicated to serial operation. We are also free to define exactly how many threads  
 1329 and processors we utilize, yet eventually hit a hard cap at the hardware available on our computing  
 1330 cluster. HAWC uses Intel Xeon™processors with 48 cores and 96 threads. We see that a successful  
 1331 code will scale well as the expected speedup is inversely correlated with  $F_s$ . As the total number  
 1332 of models sampled grows, the speedup will also.

### 1333 6.5.2 Implementation

1334 The multithreaded code was written in Python3 and is documented in the `dark_matter_hawc`  
 1335 [repository](#) within the script named `mpu_analysis.py`. A version of the script as of April 25  
 1336 **TODO: make sure to update on this date** is also provided in Section B.2. It has many dependencies  
 1337 including the HAWC analysis software. Figure 6.6 displays the workflow of a job with 3 threads.  
 1338 Within a job, SOURCE is kept fixed and CHANS remains 17 elements long. More  $m_\chi$  are sampled  
 1339 from 11 bins up to 49 (for  $\gamma\gamma$  and ZZ) and 25 (for remaining CHANS) which amounts to 12 or 6  
 1340 mass bins per decade.  $m_\chi$  and CHANS are permuted into a 473 element list which is split evenly  
 1341 across N threads where N is [2, 8, 16]. For each  $m_\chi$ -CHAN tuple, 1001  $\langle\sigma v\rangle$  values are sampled in

Timing summary for a multi-threaded job.

Import and Load HAL, 3ML, and <b>HAWC analysis software</b> ~ 2 min		
Load HAWC Data ~ 5 min		
Load HAWC PDF's (Detector Resolution files) ~ 4 min		
Load Spatial Model < 1s		
Load Spectra Models <1 s	Load Spectra Models <1 s	Load Spectra Models <1 s
Calculate max LLH (Chan_0, $m_{\chi_1}$ ) ... max LLH (Chan_?, $m_{\chi_N}$ ) N = TOTAL/N_THREADS  ~ 8 min/Spec_model	Calculate max LLH (Chan_?, $m_{\chi_1}$ ) ... max LLH (Chan_?, $m_{\chi_N}$ ) N = TOTAL/N_THREADS  ~ 8 min/Spec_model	Calculate max LLH (Chan_?, $m_{\chi_1}$ ) ... max LLH (Chan_?, $m_{\chi_N}$ ) N = TOTAL/N_THREADS  ~ 8 min/Spec_model
Write results to file <1 s	Write results to file <1 s	Write results to file <1 s
Join Threads and terminate < 1s		

Figure 6.6 Task chart for one multithreaded job developed for this project. Green blocks indicate a shared resource across the threads AND computation performed serially. Red blocks indicate functional parallel processing within each thread. 3 threads are represented here, yet many more can be employed during the full analysis. Jobs are defined by the SOURCE as these require unique maps to be loaded into the likelihood estimator. The  $m_{\chi}$ , CHAN, and  $\langle \sigma v \rangle$  variables are entered into the thread pool and allocated as evenly as possible across the threads.

- 1342 the likelihood, and the value of  $\langle \sigma v \rangle$  that maximizes the likelihood is found. Although rare, fits  
 1343 that failed are handled on a case by case basis.

### 1344 6.5.3 Performance

M Tasks	$T_{p,c}$ (hr:min:s)			
	$T_{1,1}$	$T_{1,2}$	$T_{1,8}$	$T_{1,16}$
50	1:40:37.5	0:52:43.7	0:19:13.8	0:13:44.0
74	2:22:30.0	1:15:00.6	0:25:21.3	0:15:49.8
100	(3:07:51.9)	1:40:10.5	0:30:44.4	0:20:01.4
200	(6:02:20.6)	-	1:00:32.0	0:30:35.0
473	(13:58:40.3)	-	2:01:41.4	1:07:53.2

Table 6.2 Timing summaries for analyses for serial and multithreaded processes. M tasks is the number of functional-parallel tasks ran for the computation.  $T_{p,c}$  is a single run time in hours:minutes:seconds for runs utilizing  $p$  nodes and  $c$  threads. Runs are run interactively on the same computer to maximize consistency. Empty entries are indicated with '-'. (·) entries are estimated entries extrapolated from data earlier in the column.

- 1345 We see a significant reduction to wall time needed for our dSph analyses to run. Table 6.2

1346 shows the timing summaries for analyses of different sizes and thread counts. Additionally, the  
 1347 efficiency gained when consolidating the serial loading of data is also apparent in our ability to  
 1348 study many more tasks in about the same amount of wall time as a smaller serial computation. Trials  
 1349 represented in the table were run on an AMD Opteron® processor 6344 with 48 cores, 2 threads per  
 1350 core; 2.6 GHz clock. This is not the same architecture used for analysis on the HAWC computing  
 1351 cluster however they are similar enough that results shown here are reasonably representative of  
 1352 computing on the HAWC computing cluster. I use Tab. 6.2 for the inferences and conclusions in  
 1353 the following paragraphs.

1354 First, we want to find  $T_s$ , the time of serial computation. From Fig. 6.5, the timing for our  
 1355 coding pattern can be written as

$$T_s + Mt_p = T_{1,1}^M. \quad (6.5)$$

1356  $M$  is the number of functional-parallel tasks (represented as column 1 of Tab. 6.2), and  $t_p$  is the  
 1357 average time to complete a single parallel task.  $T_{1,1}^M$  is the total time for a parallel program to run if  
 1358 only 1 processor is allocated for  $M$  parallel task. With two runs of different  $M$  ( $M_1$  and  $M_2$ ), we  
 1359 can use a system of equations to compute

$$T_s = 803.1\text{s} \text{ and } t_p = 104.6\text{s} \quad (6.6)$$

1360 Now, we have specific estimation for the fraction of serial computing time,  $F_s$ :

$$F_s = \frac{803.1}{803.1 + 104.6 \cdot M}. \quad (6.7)$$

1361 The maximum  $M$  for this study is 473 which evaluates to:  $F_s = 0.016$  or 1.6% of computing time.  
 1362 Table 6.3 shows the resulting speedups.

1363 We see a speedup that generally exceeds expectations from Eq. (6.4) for real trail runs. We also  
 1364 see that there are diminishing returns as the number of threads increases. For small jobs with large  $c$ ,  
 1365 both the expected and observed speedup are significantly smaller than  $c$ . One thing not considered  
 1366 in Eq. (6.4) is the time incurred via communication latency. Communication latency increases  
 1367 with the number of threads and contributes to diminishing returns. Additionally, these values are

M Tasks	$F_s$	$S_{1,2}$	$S_{1,8}$	$S_{1,16}$
50	1.33 E-1	1.90 [1.76]	5.23 [4.14]	6.35 [5.34]
74	9.40 E-2	1.90 [1.83]	5.62 [4.82]	9.00 [6.64]
100	7.13 E-2	1.88 [1.87]	6.11 [5.34]	9.38 [7.73]
200	3.70 E-2	- [1.93]	5.98 [6.36]	11.85 [10.29]
473	1.60 E-2	- [1.97]	6.89 [7.20]	12.35 [12.91]

Table 6.3 Speed up summaries for analyses for serial and multithreaded processes. M tasks is the number of functional-parallel tasks ran for the computation.  $S_{p,c}$  is a single speedup comparison for runs utilizing  $p$  nodes and  $c$  threads. [·] are the estimated speedups calculated from Tab. 6.2, Eq. (6.7), and Eq. (6.4). Empty entries are indicated with '-'.

1368 for single runs and do not consider the stochastic variation expected in a shared high performance  
 1369 computing resource. Therefor, these results are not strictly conclusive, yet demonstrate the merits  
 1370 of multi-threading. We see very clearly that there is a lot to gain, and this new coding pattern will  
 1371 expand HAWC's analysis capabilities.

## 1372 6.6 Analysis Results

1373 3 of the 43  $\mathcal{L}\mathcal{S}$  dSphs considered for the multithreaded analysis. These dSph are analyzed  
 1374 for emission from DM annihilation according to the likelihood method described in Section 5.4.  
 1375 The three likelihood profiles are then stacked to synthesize a combined limit on the dark matter  
 1376 annihilation cross-section,  $\langle\sigma v\rangle$ . This combination is done each of the 17 SM annihilation channels.  
 1377 Figure 6.7 and Fig. 6.8 show the combined limits for all annihilation channels with HAWC's  
 1378 observations. Test statistics of the best fit  $\langle\sigma v\rangle$  values for each  $m_\chi$  and CHAN are shown in Fig. 6.9  
 1379 and Fig. 6.10. We also compare these limits to HAWC's Glory Duck limits shown in Section 5.5.  
 1380 The comparison to Glory Duck are featured in Fig. 6.11 for all the DM annihilation channels studied  
 1381 for Glory Duck. A full comparison is provided in the appendix in Fig. B.2, Fig. B.3, and Fig. B.4.  
 1382 Here, we show updated limits for  $\chi\chi \rightarrow b\bar{b}, e\bar{e}, \mu\bar{\mu}, \tau\bar{\tau}, t\bar{t}, W^+W^-$ ,  $\gamma\gamma$  and  $ZZ$ . For the first time  
 1383 ever, we show limits for  $\chi\chi \rightarrow c\bar{c}, s\bar{s}, u\bar{u}, d\bar{d}, \nu_e\bar{\nu}_e, \nu_\mu\bar{\nu}_\mu, \nu_\tau\bar{\nu}_\tau, gg$ , and  $hh$ .

1384 No DM was found in HAWC observations. The largest excess found in HAWC data was for DM  
 1385 annihilating to  $W$ -bosons or  $\nu_e\bar{\nu}_e$  for  $m_\chi = 10$  TeV at significance  $2.11\sigma$  and  $2.14\sigma$  respectively.  
 1386 HAWC's limits and excesses are dominated by Segue1. Coma Berenices shows excesses at higher  
 1387 DM mass, yet no similar excesses were observed in Segue1 or Sextans. Sextans did not contribute

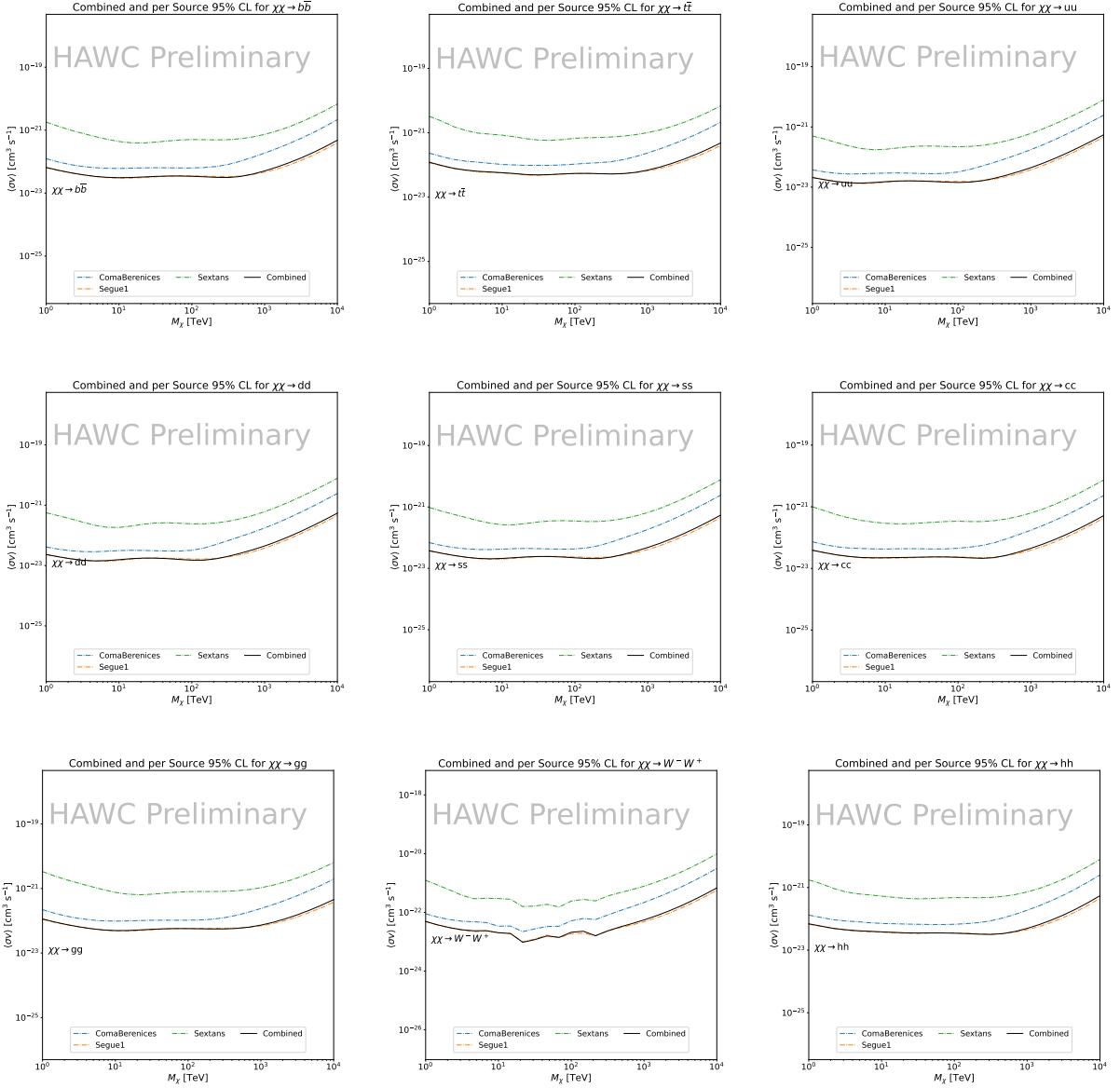


Figure 6.7 HAWC upper limits at 95% confidence level on  $\langle\sigma v\rangle$  versus  $m_\chi$  for  $\chi\chi \rightarrow b\bar{b}$ ,  $t\bar{t}$ ,  $u\bar{u}$ ,  $d\bar{d}$ ,  $s\bar{s}$ ,  $c\bar{c}$ ,  $gg$ ,  $W^+W^-$ , and  $hh$ . Limits are with  $\mathcal{L}\mathcal{S}$   $J$ -factors [66]. The solid line represents the observed combined limit. Dashed lines represent limits from individual dSphs.

1388 significantly to signal excesses or the combined limit as it is at high zenith. Draco was not included  
 1389 as the PDF of some of our analysis bins were wider than what is reasonable for a point source  
 1390 analysis. Draco is at a high zenith for HAWC, so the effort required to include it was not justified  
 1391 by the benefits.

1392 We did not generate background trials in time of writing this thesis. These are not shown and

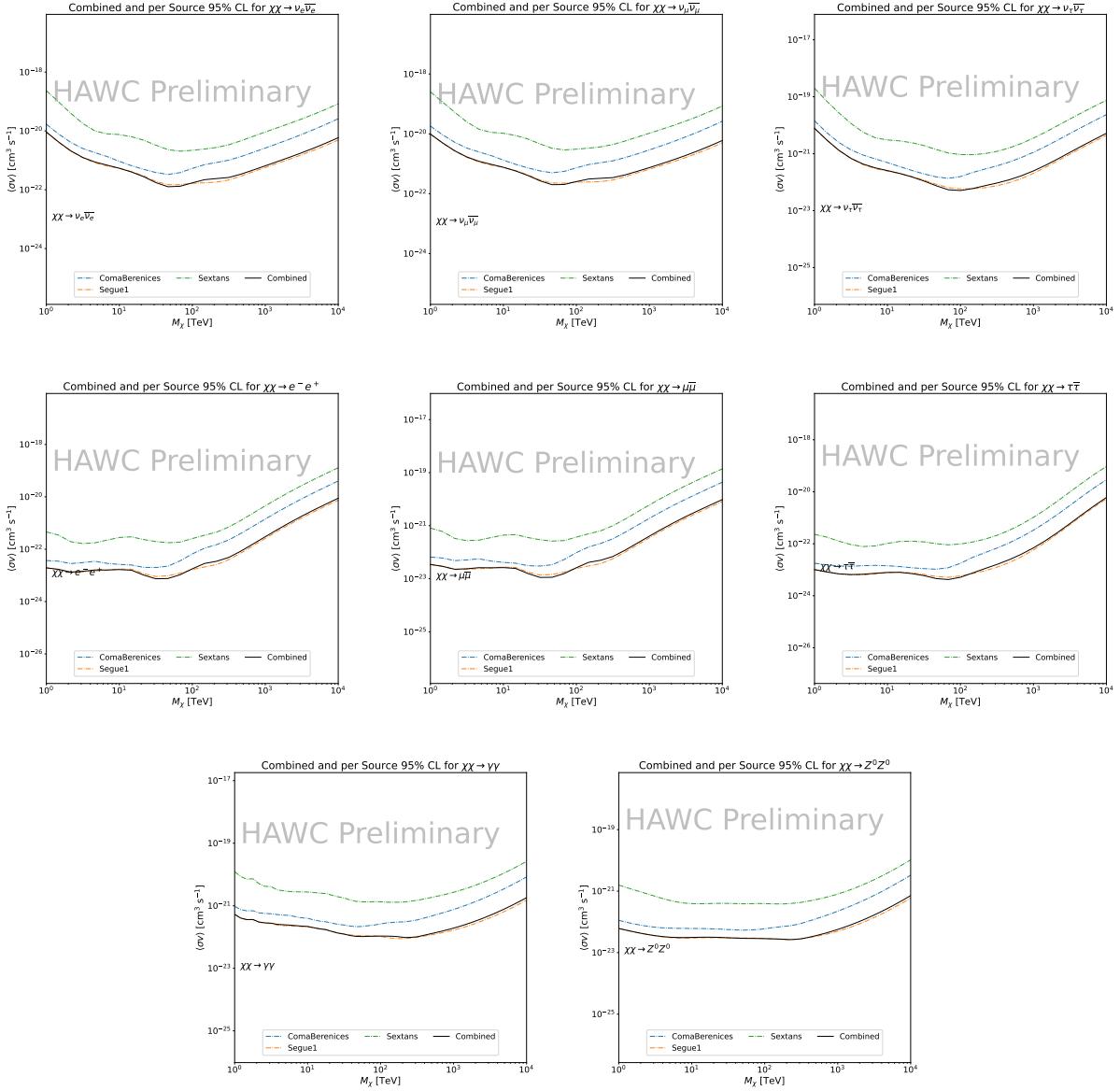


Figure 6.8 HAWC upper limits at 95% confidence level on  $\langle\sigma v\rangle$  versus  $m_\chi$  for  $\chi\chi \rightarrow \nu_e \bar{\nu}_e$ ,  $\nu_\mu \bar{\nu}_\mu$ ,  $\nu_\tau \bar{\nu}_\tau$ ,  $e \bar{e}$ ,  $\mu \bar{\mu}$ ,  $\tau \bar{\tau}$ ,  $\gamma\gamma$  and  $ZZ$ . Limits use  $\mathcal{L}S$   $J$ -factors [66]. The solid line represents the observed combined limit. Dashed lines represent limits from individual dSphs.

1393 are an immediate next step for this analysis before publication.

1394 When comparing these results to Section 5.5, we see an overall decrease to the confidence limit  
 1395 therefore improvement to HAWC's expected sensitivity. This improvement is generally stronger  
 1396 than a doubling of data, or a factor  $\sqrt{2}$  decrease. The comparison is somewhat complex and  
 1397 dependent on the dSph and SM annihilation channel. Figure 6.11 shows the comparisons of limits

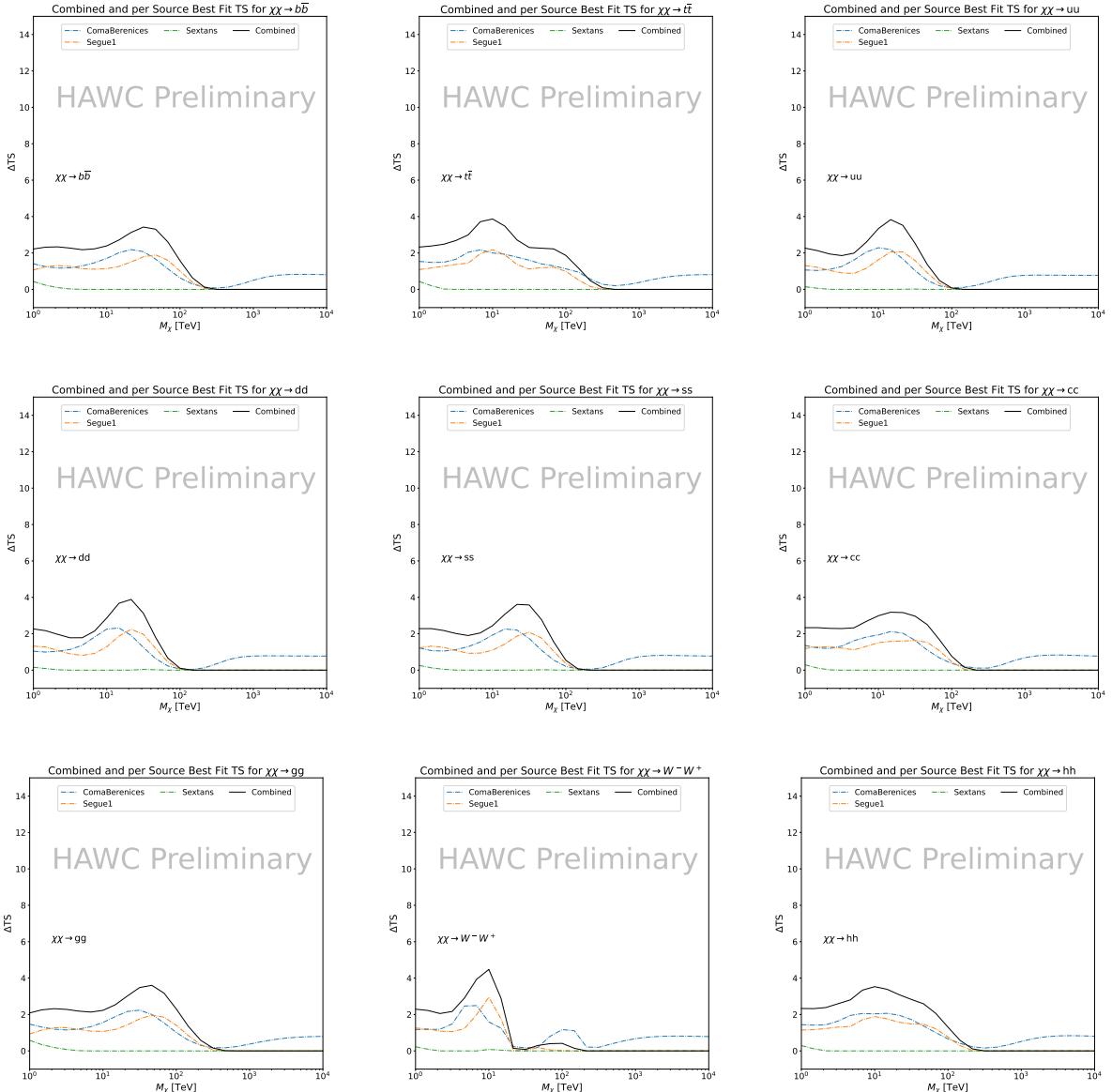


Figure 6.9 HAWC TS values for best fit  $\langle\sigma v\rangle$  versus  $m_\chi$  for SM annihilation channels:  $\chi\chi \rightarrow b\bar{b}$ ,  $t\bar{t}$ ,  $u\bar{u}$ ,  $d\bar{d}$ ,  $s\bar{s}$ ,  $c\bar{c}$ ,  $gg$ ,  $W^-W^+$ , and  $hh$ . Limits use  $\mathcal{LS}$   $J$ -factors. The solid black line shows the combined best fit TS values. The colored, dashed lines are the TS values from each dSph.

1398 calculated for this analysis and Glory Duck (Section 5.5). Segue 1 and Coma Berenices are low  
 1399 zenith where improvements to HAWC's analysis come only from energy estimation. Differences  
 1400 between these two are dominantly from their differences in  $J$ -factor, half-light radii of the dSphs,  
 1401 and the particle physics inputs. Substantial gains in HAWC's analysis methods (pass 5.F) were  
 1402 made at high zenith which is important for sources like Sextans. The HDM particle physics model

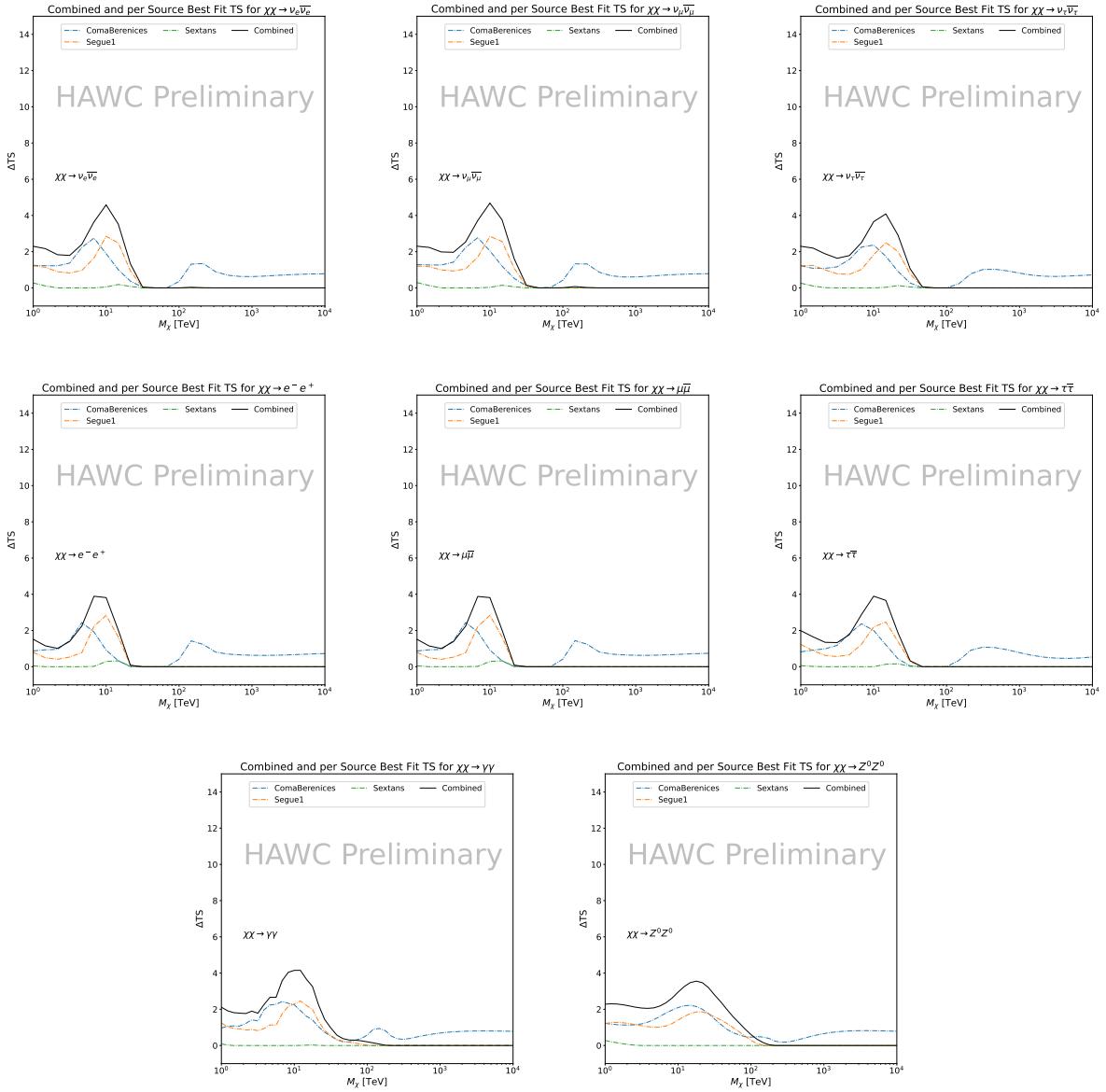


Figure 6.10 HAWC TS values for best fit  $\langle\sigma v\rangle$  versus  $m_\chi$  for SM annihilation channels:  $\chi\chi \rightarrow \nu_e \bar{\nu}_e$ ,  $\nu_\mu \bar{\nu}_\mu$ ,  $\nu_\tau \bar{\nu}_\tau$ ,  $e^- e^+$ ,  $\mu \bar{\mu}$ ,  $\tau \bar{\tau}$ ,  $\gamma\gamma$  and  $ZZ$ . Limits use  $\mathcal{LS}$   $J$ -factors. The solid black line shows the combined best fit TS values. The colored, dashed lines are the TS values from each dSph.

1403 produces almost identical spectra to the PPPC for  $\chi\chi \rightarrow e^- e^+$ . This channel can be used to  
 1404 compare limits between dSph spatial models. Overhead sources see minimal improvement to the  
 1405 limits, while high zenith sources see an order of magnitude improvement for all DM masses. Softer  
 1406 SM annihilation channels see broad improvements to the limit compared to harder channels.

1407 **6.7 Systematics**

1408     Systematics to this analysis are identical to what was performed earlier in Glory Duck, Sec-  
1409     tion 5.7. We are also sensitive to the choice in spatial template, and this was explored in Section 5.7.2  
1410     and Section 5.8.2.

1411 **6.8 Conclusion and Discussion**

1412     In this multithreaded analysis, we have used observations of 3 dSphs from HAWC to perform  
1413     a collective DM annihilation search towards dSphs. The data were combined across sources  
1414     to significantly increase the sensitivity of the search. Advanced computational techniques were  
1415     deployed to accelerate wall-time spent analyzing by an order of magnitude. We have observed  
1416     no significant deviation from the null, no DM hypothesis, and so present our results in terms of  
1417     upper limits on the velocity-weighted cross-section,  $\langle\sigma v\rangle$ , for seventeen potential DM annihilation  
1418     channels across four decades of DM mass.

1419     This analysis serves as a proof of concept for multithreaded HAWC analyses with large parameter  
1420     spaces. Larger datasets with fast techniques will be essential to tackling the DM problem. The  
1421     models we used for this study include annihilation channels with neutrinos in the final state.  
1422     Advanced studies could aim to merge our results with those from neutrino observatories with large  
1423     data sets.

1424     A full HAWC analysis will include systematic studies of the  $J$ -factor distributions. Additionally,  
1425     because of the timing reduction, the study can be doubled in size to include DM decay. We have not  
1426     yet received the remaining spatial profiles to the  $\mathcal{LS}$  catalog, and limits can be quickly computed  
1427     once these are received. Finally, statistical studies with Poisson variation of HAWC's background  
1428     are essential to a comprehensive understanding of our observed excesses.

Segue1

Coma Berenices

Sextans

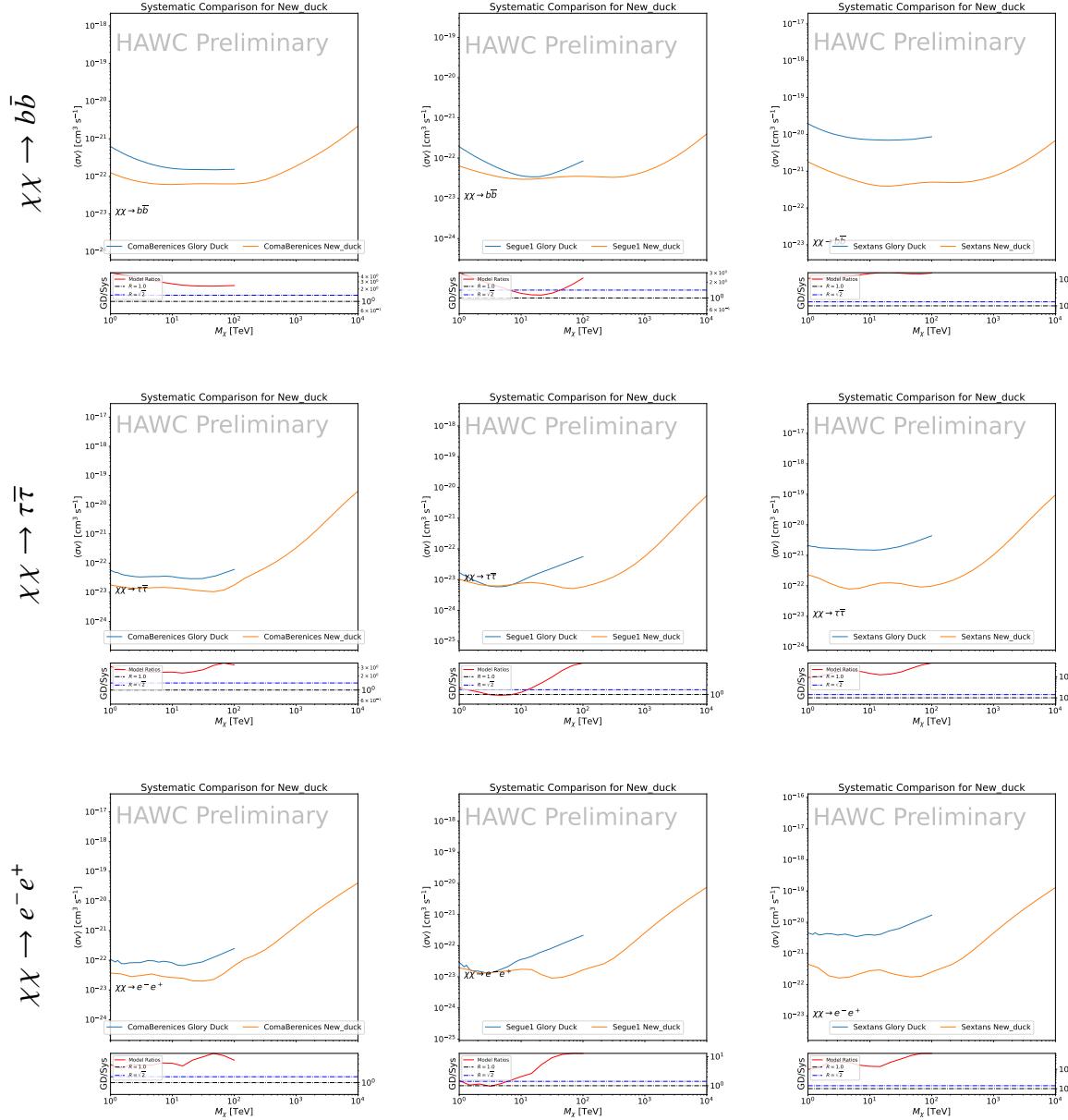


Figure 6.11 Comparison of HAWC limits from this analysis to Glory Duck (Fig. 5.5) for 3 dSphs and 3 DM annihilation channels:  $b\bar{b}$ ,  $\tau\bar{\tau}$ , and  $e^-e^+$ . Each sector shows the 95% confidence limit from Glory Duck (blue line) and this analysis (orange line) in the top plot. The lower plot features the ratio in log scale of Glory Duck to this analysis in a red solid line. Horizontal dashed lines are for ratios of 1.0 (black) and  $\sqrt{2}$  (blue). Ratios larger than 1.0 are for limits smaller, or stricter, than Glory Duck. Ratios larger than  $\sqrt{2}$  indicates limits are stricter than a simple doubling of the Glory Duck data.

## CHAPTER 7

### 1429 HEAVY DARK MATTER ANNIHILATION SEARCH WITH ICECUBE'S NORTH SKY 1430 TRACK DATA

#### 1431 7.1 Introduction

1432 Neutrinos are another astrophysical messenger than can travel long distances without interaction.  
1433 Uniquely, they interact less readily than photons especially above PeV energies. Neutrinos thereofre  
1434 provide another window through which we can perform dark matter searches. Neutrinos come in  
1435 three flabors and so this triples the multiplicity of the particles we are searching for.

1436 Icecube has not done a DM annihilation analysis towards dwarf galaxies for a while. **TODO:**  
1437 [cite 2013 paper](#). This is in spite of the potentially crucial sensitivity afforded from neutrino spectral  
1438 lines [TODO: cite dan hooper and neutrino lines](#). A lot has changed in IC3 since that last analysis (we  
1439 have more strings, we have much more sophisticated analysis methods, and the theory modeling  
1440 has made significant leaps.) Therefore it is time to finally do a DM search toward dSphs. The hope  
1441 is that by laying down the important statistical foundation as well, that this work can be meshed  
1442 with gamma-ray data. IceCube is sensitive to annihilating DM to the DM ranges above 1 TeV  
1443 and can produce competitive results relative to gamma ray observatories in spectral models that  
1444 produce sharp neutrino features. The goal of this analysis is to perform a DM annihilation search  
1445 using the new datasets NST. The search will only be towards dwarf spheroidal galaxies (dSph).  
1446 These sources are known for their low backgrounds and high DM contents. Since the dataset is  
1447 sensitive to the north and south, as many dSph as possible will be included. Additionally, with  
1448 annihilation, these sources can be treated as point sources with little loss to sensitivity or model  
1449 dependence on how the DM is distributed. DM masses from 500 GeV to 100 PeV are considered  
1450 for this analysis. All standard model annihilation channels available from the HDMspectra are  
1451 studied in this analysis.

1452 Additional work is done to extract the Likelihood profiles for each DM, source hypothesis so  
1453 that these data can be combined with gamma-ray observatories. This work is considered a separate  
1454 project as the statistical treatment is unique from many IceCube analyses. The wiki for [ the

1455 combined analysis] **TODO: instead point to chapter**This chapter presents the analysis work for  
1456 IC3 for DM searches toward dSphs. This section describes the various steps and features of the  
1457 analysis. It is structure first introduces the data and how it is treated, then systematic studies of the  
1458 dwarves individually. Finally, the stacked analysis and results are presented.

1459 **7.2 Dataset and Background**

1460 This section enumerates the data and background methods used for IceCube's study of dSphs.  
1461 Section 7.2.1 and Section 7.2.2 are most useful for fellow IceCube collaborators looking to replicate  
1462 this analysis.

1463 **7.2.1 Itemized IceCube files**

- 1464 • Software Environment: CVMFS Py3-v4.1.1
- 1465 • Data Sample: Northern Tracks NY86v5p1
- 1466 • Analysis Software: csky ([nu\\_dark\\_matter](#))
- 1467 • Analysis wiki: [https://wiki.icecube.wisc.edu/index.php/Dark\\_Matter\\_Annihilation\\_Search\\_towards\\_dwarf\\_spheroidals\\_with\\_NST\\_and\\_DNN\\_Cascades](https://wiki.icecube.wisc.edu/index.php/Dark_Matter_Annihilation_Search_towards_dwarf_spheroidals_with_NST_and_DNN_Cascades)
- 1469 • Project repository

1470 **7.2.2 Software Tools and Development**

1471 This analysis was performed inside IceCube's CVMFS (3.4.1.1) software environment using  
1472 csky for likelihood calculations. Csky did not come with dark matter spectral models nor could  
1473 accomodate custom flux models. We developed these capacities for single source and stacked  
1474 source studies for this analysis. The analysis code is held in a separate repository from csky. The  
1475 [nu\\_dark\\_matter](#) branch of csky manages the input of custom dark matter spectra and accompanied  
1476 DM astrophysical source then calculates likelihoods with a selected data sample. The [IceCube Dark](#)  
1477 [Matter dSph repository](#) manages the generation of spectral models for neutrinos, physics parameter  
1478 extraction from  $n_{\text{sig}}$ ,  $J$ -factor per source inputs, and bookkeeping for the large parameter space.  
1479 The project repository required a secondary software environment for neutrino oscillations. How

1480 to launch and run those calculations are documented in the project repository and the Docker image  
1481 is additionally saved in Section C.1

### 1482 **7.2.3 Data Set and Background Description**

1483 For this analysis, I use the Northern Sky Tracks (NST) Sample (Version V005-P01). The sample  
1484 contains up-going track-like events, usually from  $\nu_\mu$  and  $\nu_\tau$  and has a superior angular resolution  
1485 compared to the cascade dataset. This sample covers 10.4 years of data (IC86\_2011-2021). The  
1486 accepted neutrino energy range used for the analysis is unique from most other IceCube searches  
1487 because DM spectra are very hard. The sampled energy range is  $1 < \log(E_\nu/\text{GeV}) < 9.51$  with  
1488 step size 0.125.

1489 The strength of a dwarf analysis is that there is no additional background consideration beyond  
1490 nominal, baseline background estimations. For NST, the nominal contribution comes from atmo-  
1491 spheric neutrinos and isotropic astrophysical neutrinos. We estimate the background by scrambling  
1492 NST data along Right Ascension.

## 1493 **7.3 Analysis**

1494 The expected differential neutrino flux from DM-DM annihilation to standard model particles,  
1495  $d\Phi_\nu/dE_\nu$ , over solid angle,  $\Omega$  is described by the familiar equation.

$$\frac{d\Phi_\nu}{dE_\nu} = \frac{\langle\sigma v\rangle}{8\pi m_\chi^2} \frac{dN_\nu}{dE_\nu} \times \int_{\text{source}} d\Omega \int_{l.o.s} \rho_\chi^2 dl(r, \theta') \quad (7.1)$$

1496 This is identical to past examples except that there are 3 neutrino flavors, so there are a corresponding  
1497 3 flux equations. Section 5.3 has a complete description of all the terms. Additionally, neutrinos  
1498 oscillate between flavors which needs to be considered for the expected neutrino flux at Earth.  
1499 Section 7.3.1 presents the particle physics model for DM annihilation. Section 7.3.2 presents the  
1500 spatial distributions built for each dSph.

### 1501 **7.3.1 $\frac{dN_\nu}{dE_\nu}$ - Particle Physics Component**

1502 Neutrino spectra from heavy dark matter annihilation were generated using HDMSSpectra [65]  
1503 and  $\chi$ arrov [69]. HDMSSpectra simulates the decay and annihilation of heavy dark matter, for  
1504 different dark matter masses and SM primary annihilation channels. The simulation includes

1505 electroweak radiative corrections and higher order loop corrections with quarks. This publication  
1506 also pushes the simulated DM mass to the Plank scale (1 EeV), however this study will not explore  
1507 that high.

1508 An important novel feature in the spectra is that neutrino line channels will be accompanied  
1509 with a low energy tail. Thus the earth will not fully attenuate a neutrino SM channel signal from  
1510 high declination sources where the neutrino flux must first traverse through the Earth. The SM  
1511 annihilation channels that feature lines include all leptonic channels. ( $\nu_{e,\mu,\tau}$ ,  $e$ ,  $\mu$ , and  $\tau$ ) We use  
1512 [χarov](#) to propagate and oscillate the neutrinos from the source to Earth. Because these sources are  
1513 quite large in absolute terms, and also far (order 10 kpc or more), the resulting flavor spectra are  
1514 the averages of the transition probabilities [69]:

$$P(\nu_\alpha \rightarrow \nu_\beta) \approx \begin{bmatrix} 0.55 & 0.18 & 0.27 \\ 0.18 & 0.44 & 0.38 \\ 0.27 & 0.38 & 0.35 \end{bmatrix} \quad (7.2)$$

1515 When calculating the expected contribution to  $n_s$ , only  $\nu_\mu$ ,  $\nu_\tau$  are considered as NST's effective  
1516 area to  $\nu_e$  is essentially 0 [70]. With these consideration, the expected composite neutrino spectrum  
1517 is sum of the two flavors:  $\nu_\mu + \nu_\tau$ . The spectral tables are then converted to splines to condense  
1518 information, enable random sampling of the spectra, and enable faster computation times. The  
1519 spectral splines are finally implemented as a DM class in csky. Examples of the spectra before and  
1520 after propagation are shown in Fig. 7.1.

$M_\chi$ 

1 TeV

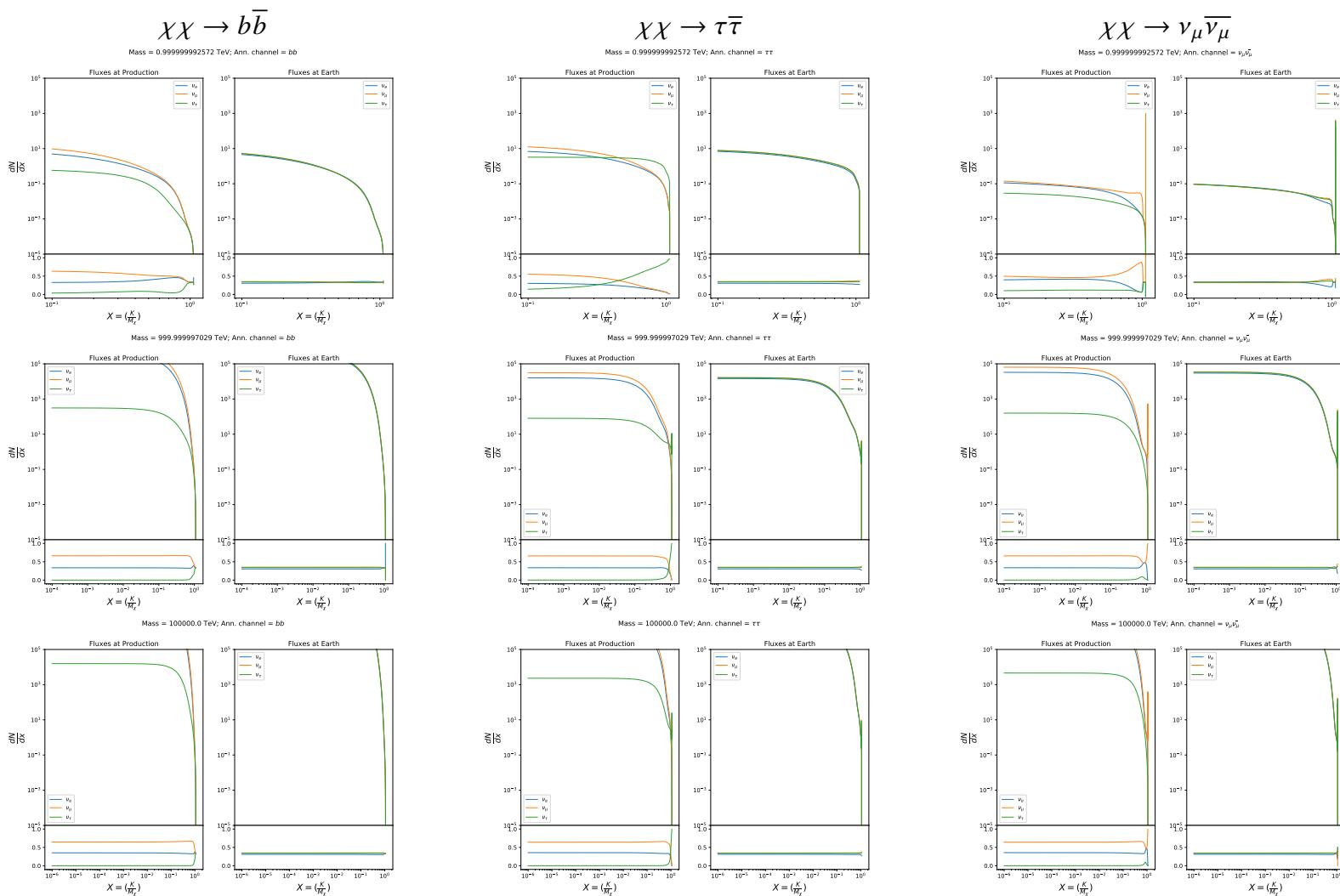


Figure 7.1 Neutrino spectra at production (left panels) and after oscillation at Earth (right panels). Blue, orange, and green lines are the  $\nu_e$ ,  $\nu_\mu$ , and  $\nu_\tau$  spectra respectively. Top panels show the spectra in  $\frac{dN}{dE}$ . Lower panels plot the flavor ratio to  $\nu_e + \nu_\mu + \nu_\tau$ . SM annihilation channels  $b\bar{b}$ ,  $\tau\bar{\tau}$ , and  $\nu_\mu\bar{\nu}_\mu$  are shown for  $M_\chi = 1 \text{ PeV}$ ,  $1 \text{ TeV}$ , and  $1 \text{ EeV}$ .

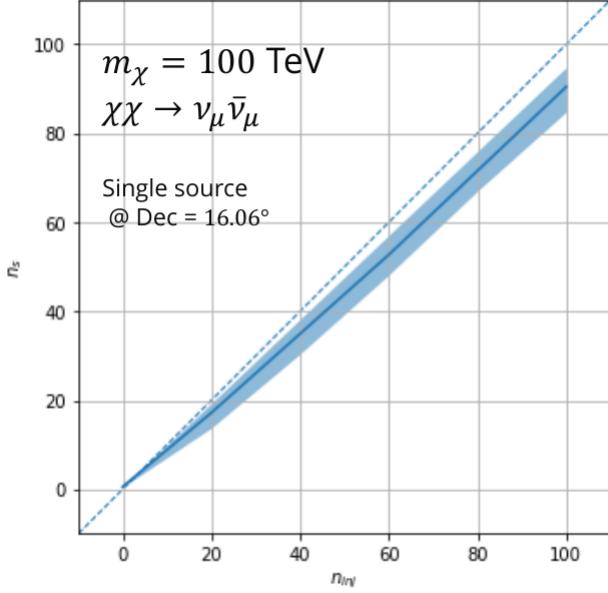


Figure 7.2 Signal recovery for 100 TeV DM annihilation into  $\nu_\mu \bar{\nu}_\mu$  for a source at Dec = 16.06°.  $n_{\text{inj}}$  is the number of injected signal events in simulation.  $n_s$  is the number of reconstructed signal events from the simulation. Although the uncertainties are small and tight, the reconstructed  $n_s$  are systematically underestimated.

### 7.3.1.1 Treatment of Neutrino Line Features

All leptonic DM annihilation channels  $\chi\chi \rightarrow [\nu_{e,\mu,\tau}, e, \mu, \tau]$  develop a prominent and narrow

spectral line feature. For all neutrino flavors, this line is visible and prominent in all mass models

studied for this analysis. For charged leptons, the feature only really shows up at the larger DM

mass models. Examples for lines in both neutrinos and charged leptons annihilation are provided

in Fig. 7.1. This line feature is so narrow relative the sampled energy range that the MC rarely

samples within the neutrino line. As a result, often the best fit to simulation of background will

always floor to TS = 0 and the signal recovery tends to be conservative.

To remedy this, a similar approach to the IceCube’s decay analysis [TODO: refer to Minjin’s](#)

[page](#). 2 kernels were tested (Gaussian, uniform (flat)) to smooth out the line feature. The widths

were tuned such that the signal recovery approached unity for DM mass 100 TeV to 1 PeV.

Additionally, the tuning was performed only for a source at declination 16.06 (Segue 1). This is

to avoid confusion loss in signal recovery from too narrow a line and from Earth’s attenuation of

high energy neutrinos. The convolution also needed to as close as possible preserve the integrated

counts of neutrinos. The optimized kernel window for all lines is summarized as:

- 1536 • Gaussian kernel w/ $2\sigma$  width =  $3.5E-3 \cdot m_\chi$
- 1537 • Minimum energy included in convolution =  $\text{MIN}[0.995 \cdot m_\chi, En(\nu_{line}) - 4\sigma]$
- 1538 • Maximum energy included in convolution =  $\text{MAX}[1.005 \cdot m_\chi, En(\nu_{line}) + 4\sigma]$

1539 where  $En(\nu_{line})$  is the neutrino energy where the neutrino line is at the maximum.

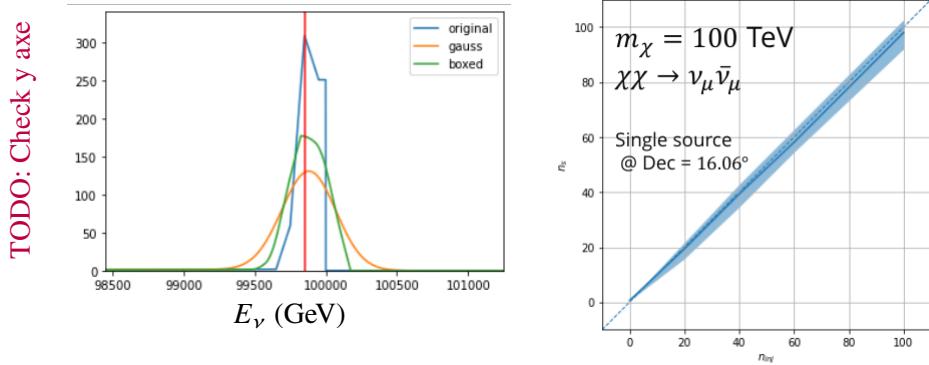


Figure 7.3 Top left panel shows the two kernels overlaid the original spectrum from Charon. delta I is the difference in the integral of the peaks with respect to the original spectrum. The vertical red line indicated where the original neutrino line is maximized. Lower right shows the signal recoveries of the DM model using the Gaussian kernel with parameters enumerated above.

1540 These parameters broadly improved the signal recovery of the line spectra. An example is  
1541 provided below. Signal recovery plots of the full analysis are provided much further down.

### 1542 7.3.1.2 Spline Fitting

1543 In an effort to reduce computational work, memory burden, and align with point source methods  
1544 used for NGC1068 and Seyfert analyses, spectral splines were created and adopted for estimating  
1545 the neutrino flux for the different annihilation. Software was written to generate, handle, and  
1546 calculate values on the splines. When using splines, one has to be careful of the goodness to fit.  
1547 There are critical caveats when testing the goodness to fit to MC generated above for all channels.

- 1548 • The splines must be Log10(\*) in Energy and dN/dE to account for the exponential nature of  
1549 the flux
- 1550 • The fidelity of the fit matters more at  $E_\nu \approx m_\chi$  where the model uncertainties are minimal  
1551 and physical considerations (like the cut-off) are most apparent.

- 1552 • The fidelity of the fit matters less at low  $E_\nu$  as the model uncertainties are large AND  
 1553 IceCube's sensitivity diminishes significantly below 500 GeV

- 1554 • Total integrated counts should be well preserved, however, the resolution of the MC is much  
 1555 higher than IceCube's energy resolution.

1556 – Meaning over several steps in E, the integral is preserved

1557 – the step size enters the cost function

1558 – Oscillating residuals, so long as they are very small and well centered, are not penalized  
 1559 as this gets averaged out.

1560 The resulting cost function to evaluate the goodness of fit was used to account for the above  
 1561 considerations.

$$e_i = x_i \cdot \left( \frac{dN_i}{dE_i} - 10^{\hat{e}_i} \right) \quad (7.3)$$

1562 Where  $\hat{e}_i$  is the spline estimator's value for  $x_i$ .  $x_i = E_{\nu_i}/m_\chi$ .  $\frac{dN_i}{dE_i}$  is the flux value from MC.

$$\text{err} = \sqrt{\frac{1}{x_{\max} - x_{\min}} \int_{x_{\min}}^{x_{\max}} \mathbf{e}^2 dx} \quad (7.4)$$

1563 I then take the RMS of the error distribution and the resulting value (err) is used to evaluate  
 1564 the fidelity of the spectral spline. Each SM channel had different tolerances for 'err'. Channels  
 1565 with very hard cut-offs had looser tolerance for err because a lot of error would be generated from  
 1566 the cut-off being estimated to occur slightly early or late. Soft channels don't have this issue and  
 1567 therefore the tolerance is very strict. The table blow summarizes the tolerances for the SM channels.

1568  
 1569 The errors are then plotted in two ways. First, FAIL and OK are directly plotted with  $e_i$  as a  
 1570 function of x, and the full spline and MC. Second, a summary plot of all the splines is plotted and  
 1571 colors coded.

1572 Figure C.1 are the spline summaries and represent the current, up-to-date status of the splines.  
 1573 The goal broadly is to eliminate all red and inspect yellow.  $\nu_e$  is not considered in this analysis  
 1574 among the neutrino final states and so no work was done to converge the spline fits for this flavor.

$\chi\chi \rightarrow$	GOOD	OK	FAIL	Limits of err calc [ $X_{min}, X_{max}$ ]
$Z^0 Z^0, W^+ W^-$	1.0E-3	1.0E-3, 1.0E-2	1.0E-2	MAX[100GeV/ $m_\chi$ , $10^{-6}$ ], 1.0
$t\bar{t}, hh$	1.0E-5	1.0E-5, 1.0E-4	1.0E-4	MAX[100GeV/ $m_\chi$ , $10^{-6}$ ], 1.0
$b\bar{b}, d\bar{d}, u\bar{u}$	9.0E-7	9.0E-7, 9.0E-6	9.0E-6	MAX[100GeV/ $m_\chi$ , $10^{-6}$ ], 1.0
$\nu\bar{\nu}_{e,\mu,\tau}$	1.0E-3	1.0E-3, 1.0E-2	1.0E-2	MAX[100GeV/ $m_\chi$ , $10^{-6}$ ], MIN[0.995, ( $E_n(\nu_{line}) - 4\sigma$ )/ $M_\chi$ ]
$e\bar{e}, \mu\bar{\mu}, \tau\bar{\tau}$	1.0E-3	1.0E-3, 1.0E-2	1.0E-2	MAX[100GeV/ $m_\chi$ , $10^{-6}$ ], MIN[0.995, ( $E_n(\nu_{line}) - 4\sigma$ )/ $M_\chi$ ]

Table 7.1 TODO: fill me daddy

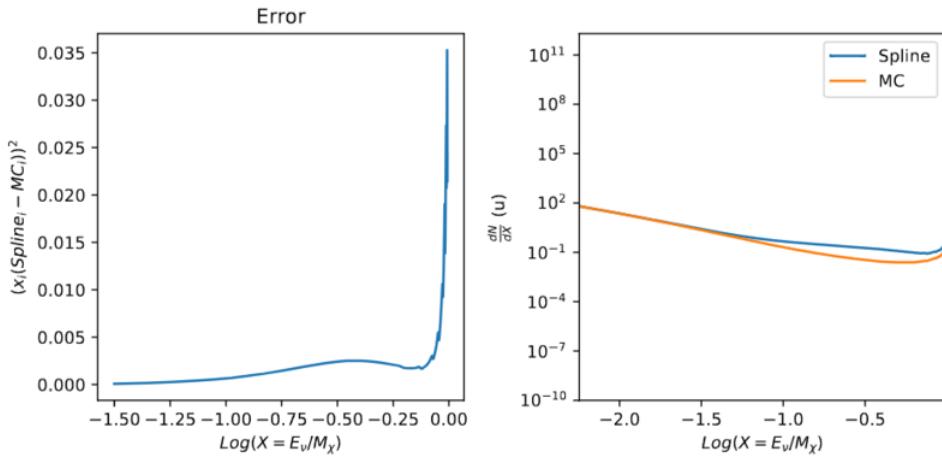


Figure 7.4 Example spline that failed the fit. Failed splined are corrected on a case by case basis unless the SM channel has a systematic problem fitting the splines. In this case, I made a bookkeeping error and loaded the incorrected neutrino flavor

1575     A Final inspection of the splines by eye was done to verify that the spline fitting did not introduce  
 1576     spurious features into the distribution that would corrupt the LLH fitting.

### 1577 7.3.1.3 Composite Neutrino Spectra

1578     With all of the previously mentioned pieces, we are ready to fully assemble a comprehensive  
 1579     description of the particle physics term  $dN/dE$  in Eq. (7.1).

$$\frac{dN_\nu}{dE_{\nu \oplus}} = \left( \frac{dN_{\nu_e}}{dE_{\nu_e}} + \frac{dN_{\nu_\mu}}{dE_{\nu_\mu}} + \frac{dN_{\nu_\tau}}{dE_{\nu_\tau}} \right)_{\text{src}} \cdot \mathbf{P}(\nu_\alpha \rightarrow \nu_\beta) \quad (7.5)$$

1580     Figure 7.5 shows the spectral models that required Gaussian smoothing, the leptonic annihilation  
 1581     channels. The remaining models where the only processing was the spline fitting are documented  
 1582     in the TODO: refer to apdxNotice that the different neutrino flavors are unique, especially in their  
 1583     low energy tails. Therefore, this analysis will be sensitive to DM annihilating to the distinct neutrino

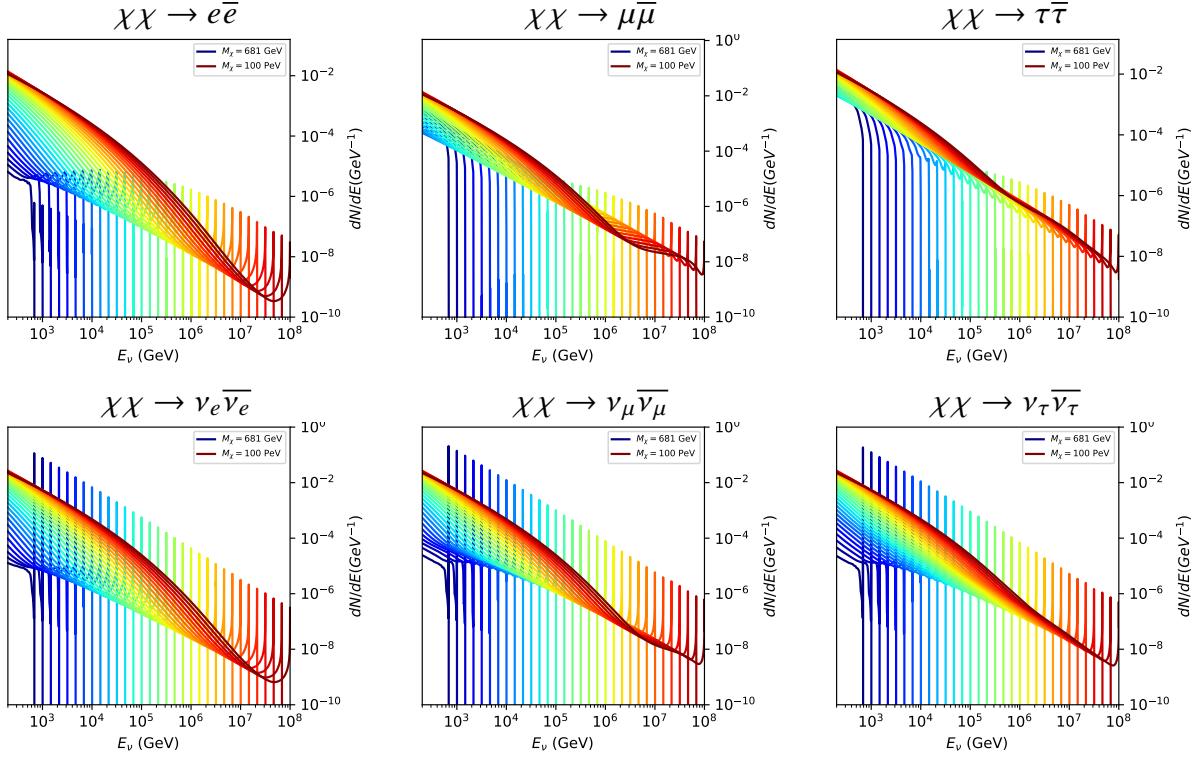


Figure 7.5 Summary of input spectral models that were smoothed with Gaussian kernel. Spectral models are for  $\chi\chi \rightarrow e\bar{e}$ ,  $\mu\bar{\mu}\tau\bar{\tau}$ ,  $\nu_e\bar{\nu}_e$ ,  $\nu_\mu\bar{\nu}_\mu$ , and  $\nu_\tau\bar{\nu}_\tau$ . These spectra are the composite ( $\nu_\mu + \nu_\tau$ ) of neutrino flavors. Every spectral model used for this analysis is featured as a colored solid line. Bluer lines are for lower DM mass spectral models. DM masses range from 681 GeV to 100 PeV. HDM [65],  $\chi$ arov [69], and Photospline [71] are used to generate these spectra. Energy (x-axis) was chosen to roughly represent the energy sensitivity of NST.

1584 flavors. All leptonic channels show prominent, hard features around  $E_\nu = m_\chi$ .

### 1585 7.3.2 J- Astrophysical Component

1586 The expected neutrino counts from a dwarf spheroidal galaxy depends also on the the 'astro-  
 1587 physical factor'. The value for this (in our specific case) J-factor for a target depends on its dark  
 1588 matter density distribution,  $\rho_\chi$  and how far it is  $l$ . For this analysis, we adopt the  $\mathcal{GS}$  model used  
 1589 in Sec. 5 for dSph from [45]. These models are based on a modified Navarro-Frenk-White (NFW)  
 1590 profile where the indices of the NFW (traditionally 1,3,1) are allowed to float. More specifically,  
 1591 these DM distributions are described using the Zhao profile. The Zhao profile is written as:

1592 where  $\theta$  is the angular distance from the center of the source. For the case annihilation, the  
 1593 source diameter, [<https://iopscience.iop.org/article/10.1088/0004-637X/801/2/74> here] defined as

1594 the  $2\theta_{\max}$ , of these dwarves is typically under  $1^\circ$  with the largest in the catalog, Fornax, extending  
 1595 to  $2.61^\circ$ . Fornax is not in the northern sky and the remaining sources are notably below this angular  
 1596 size. Therefore, the sources are treated as point sources because the typical source diameter is under  
 1597 1 degree. The J-factor used for the point source assumption is the total J emitted from  $\theta_{\max}$ . These  
 1598 values are enumerated in Geringer-Sameth 2015 and again in the table below with their coordinates.  
 1599 Coordinates are given in J200.0 equatorial coordinates. IceCube uses identical sources to Tab. 5.1  
 1600 except we analyze source with declinations above  $0.0$  degrees.

### 1601 7.3.3 Source Selection and Annihilation Channels

1602 We use all of the dSphs presented in IceCube’s previous dSph DM search [70]. IceCube’s  
 1603 sources for these simulation studies include Bootes I, Canes VenaticiI, Canes Venatici II, Coma  
 1604 Berenices, Draco, Hercules, Leo I, Leo II, Leo V, Leo T, Segue 1, Segue 2, Ursa Major I, Ursa Major  
 1605 II, and Ursa Minor. A full description of all sources used in Table 5.1. Sources with declinations  
 1606 less than  $0.0$  are excluded from this analysis.

1607 This analysis improves on the previous IceCube dSph paper [70] in the following ways. Pre-  
 1608 viously, the IceCube detector was not yet completed to the 86 string configuration. Many more  
 1609 dSphs will be observed, from 4 to 15. Previously, the particle physics model used for neutrino-ray  
 1610 spectra from DM annihilation did not have EW corrections where they are now included [65]. The  
 1611 spectral models also predict substantial differences between the neutrino flavors, so this analysis  
 1612 will be the first DM dwarf analysis to discriminate between primary neutrino flavors. The study  
 1613 performed here studies 10.4 years of data.

1614 The SM annihilation channels probed for this study include  $b\bar{b}$ ,  $t\bar{t}$ ,  $u\bar{u}$ ,  $d\bar{d}$ ,  $e\bar{e}$ ,  $\mu\bar{\mu}$ ,  $\tau\bar{\tau}$ ,  $ZZ$ ,  
 1615  $W^+W^-$ ,  $\nu_e\bar{\nu}_e$ ,  $\nu_\mu\bar{\nu}_\mu$ , and  $\nu_\tau\bar{\nu}_\tau$ .

### 1616 7.4 Likelihood Methods

1617 I use the Point-Source search likelihood which is widely used in IceCube analyses. The  
 1618 likelihood function is defined as the following:

$$L(n_s) = \prod_{i=1}^N \left[ \frac{n_s}{N} S_i + \left(1 - \frac{n_s}{N}\right) B_i \right] \quad (7.6)$$

1619 where  $i$  is an event index,  $S$  and  $B$  are the signal PDF and background PDF respectively. For a joint  
1620 analysis where the sources are stacked the likelihood is expanded in the simplified way:

$$L(n_s) = \prod_{i=1}^{N_{\text{sources}}} L_i(n_s) \quad (7.7)$$

1621 Where  $L_i$  is the likelihood from the  $i$ -th source in the stacked analysis. The test statistic definition  
1622 remains the same as Eq. (5.7)

## 1623 7.5 Background Simulation

1624 Before we look at data, we must first analyze background and signal injection to validate our  
1625 analysis. The following sections show the results of the likelihood fitting for a suite of background  
1626 trials for the DM models we set out to study in [TODO: refer to the section](#). We study the TS  
1627 distributions first for each source, then for the stacked analysis.

1628 The TS distributions are not expected to behave according to a chi-squared distribution with 1  
1629 degree of freedom. This is in large part due to the distinct spectral shapes demonstrated earlier.  
1630 These can vary significantly between DM mass and annihilation models. Therefore, Wilks' theorem  
1631 may not be applicable to the analysis. Instead, a critical value is defined from a large number of  
1632 background trials.

1633 I assume that TS values are physical:  $\text{TS} \geq 0$ .  $\eta$  denotes the fraction of positive TS values  
1634 above the threshold and written in the legend.  $\epsilon[x]$  indicate the fraction of events where  $\text{TS} < x$ .  
1635 For TS plots shown here, the decimal values of  $x$  are 1.0e-2 and 1.0e-3. The following plots show  
1636 the background TS distributions obtained from Segue1, a source with little Earth attenuation and  
1637 large J-factor, assuming that dark matter annihilates into  $b\bar{b}$ . I also show the 15 source stack TS  
1638 distributions with identical DM models.

### 1639 7.5.1 TS per Source

1640 Below I present the TS distributions for Segue1 and  $\chi\chi \rightarrow b\bar{b}$ . All remaining channels and  
1641 source TS panels are hosted on [TODO: Change this text, it will all be here](#).

1642 Although it was not expected, almost every distribution produced follows a chi2 distribution  
1643 with 1 degree of freedom. This is important for future assumptions made (in multi-messenger) and

1644 may justify statistical calculations assuming Wilk's theorem is valid.

1645 **TODO:** add text saying that you show: bb, numu, and tau??? specs for Seg1 and UMa2?

## 1646 7.5.2 Stacked TS

1647 The presentation of these plots are identical to the previous 'per Source' section. I use csky  
1648 source software to calculate the TS distributions. Bugs were found when implementing, however  
1649 were rectified. Warning to future users performing a stacked analysis with custom spectra. In  
1650 using the above, I am making the implicit assumption that the primary/only cause to a difference in  
1651 neutrino counts from the sources is accounted for through the J-factors. The J-factors are therefor  
1652 used as weights for the stacking where an individual source's weight is defined as:

1653 Each subplot, except the final, is the TS distribution for a specific DM mass listed in the subplot.  
1654 The final subplot plots the all DM spectral models used as input for the TS distribution calculations  
1655 with bluer lines indicating lower DM mass and redder indicating higher DM mass. Below is an  
1656 image of bb. The full resolution pdfs were provided in links above.

## 1657 7.6 Signal Recovery

### 1658 7.6.1 Sensitivities

1659 In IceCube, we usually define the 90% confidence level (CL), as the minimum number of signal  
1660 events ( $n_s$ ) required to have a Type I error rate smaller than 0.5 and Type II error rate of 0.1. Csky  
1661 performs the sweep to find  $n_s$  that satisfies the previous condition, and from  $n_s$  I use the following  
1662 equation

$$n_s = T_{live} \int_0^{\Delta\Omega} d\Omega \int_{E_{min}}^{E_{max}} dE_\nu A_{eff}(\hat{n}, E_\nu) \frac{d\Phi_\nu}{d\Omega dE_\nu}(\hat{n}, E_\nu), \quad (7.8)$$

1663 to extract the sensitivity on the dark matter annihilation cross-section.  $T_{live}$  is the detector  
1664 livetime,  $A_{eff}$  is the effective area of the detector, and  $E_{min}$ ,  $E_{max}$  are the minimum, maximum  
1665 energies of the expected neutrinos, respectively.

1666 Sensitivities are calculated for each source individually as if they were the only source and as  
1667 a stack. Example plots of these plots are shown below and organized by the single source/stacked  
1668 studies. Finally, I generated a plot with all hypotheses which is presented at the very end.

1669 **7.7 Systematics**

1670 Lol What Systematics. Beside signal recovery we don't have many additional studies for here.  
1671 The current analysis plan is to compare these sensitivities to another  $J$ -factor catalog such as  $\mathcal{LS}$   
1672 [66]. Additionally, we set out to perform a standard suite of IceCube systematic studies which  
1673 include: **TODO: THE BIG 4: ICE MODEL ETC**

1674 **7.8 Conclusions**

1675 We built many things for this analysis. We utilized advanced computing techniques like  
1676 parallel programming and spline fitting of particle physics Monte Carlo to greatly expand and  
1677 refine IceCube's sensitivity to DM annihilation from dSphs. We imported updated astrophysical  
1678 and particle physics models that better represent what we believe neutrino signals from DM  
1679 annihilation should look like. We, for the first time, build an analysis that is sensitivity to PeV DM  
1680 annihilation.

1681 When we compare to previous IceCube publications of dSphs [70], we see an order of magnitude  
1682 improvement to our sensitivity. This analysis has been working group approved within IceCube and  
1683 has begun the unblinding process. This processes did not complete in time for this dissertation.  
1684 Therefor we do not show data for this thesis and is the clear next step.

1685 The test statistic distributions in this analysis also demonstrate more characteristic behaviour  
1686 compared to previous DM analyses. With a 10 year dataset, we finally have enough statistics to  
1687 almost trivially combine with other photon observatories, such as HAWC. The first ground work for  
1688 a multi-messenger DM search is provided with concluding remarks in Sec. 8.

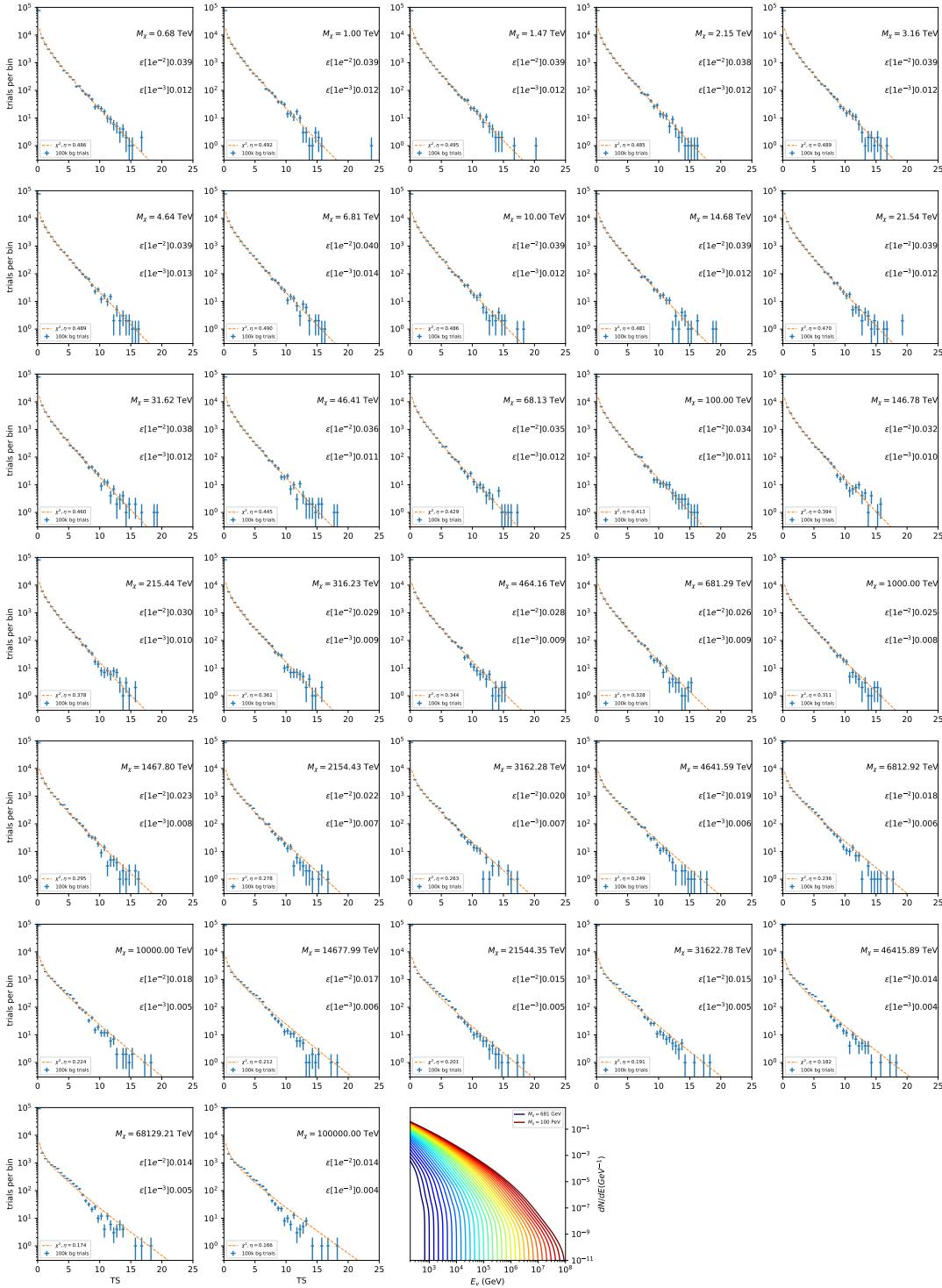


Figure 7.6 Test statistic (TS) distributions for Segue 1 and  $\chi\chi \rightarrow b\bar{b}$ . Each subplot, except the final, is the TS distribution for a specific DM mass listed in the subplot. Orange dashed lines are the traces for a  $\chi^2$  distribution with 1 degree of freedom.  $\epsilon[\cdot]$  is the fraction of trials smaller than the bracketed value. The final subplot plots the all DM spectral models, similar to Fig. 7.5, used as input for the TS distributions.

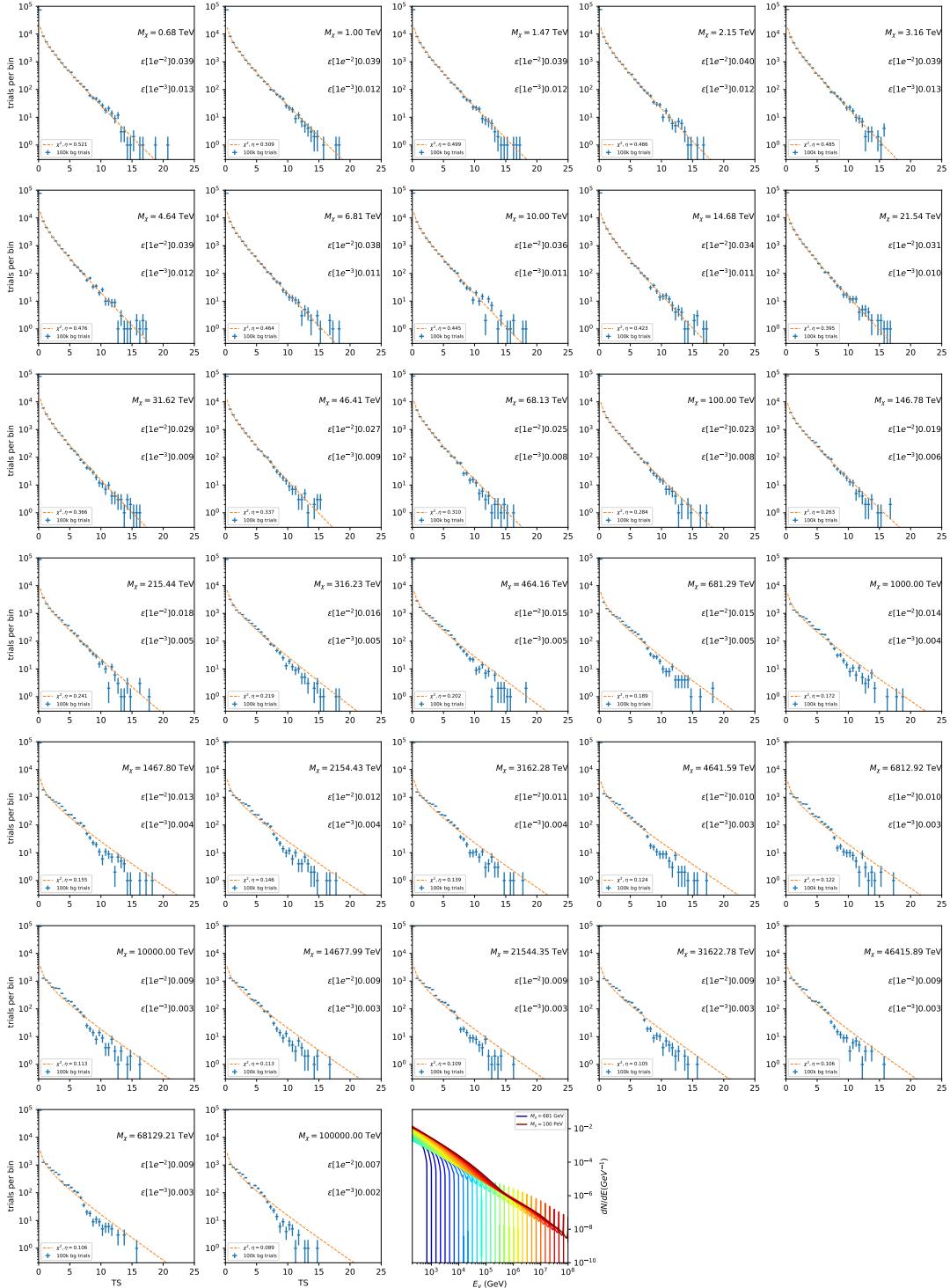


Figure 7.7 Same as Fig. 7.6 for Segue 1  $\chi\chi \rightarrow \tau\bar{\tau}$ .

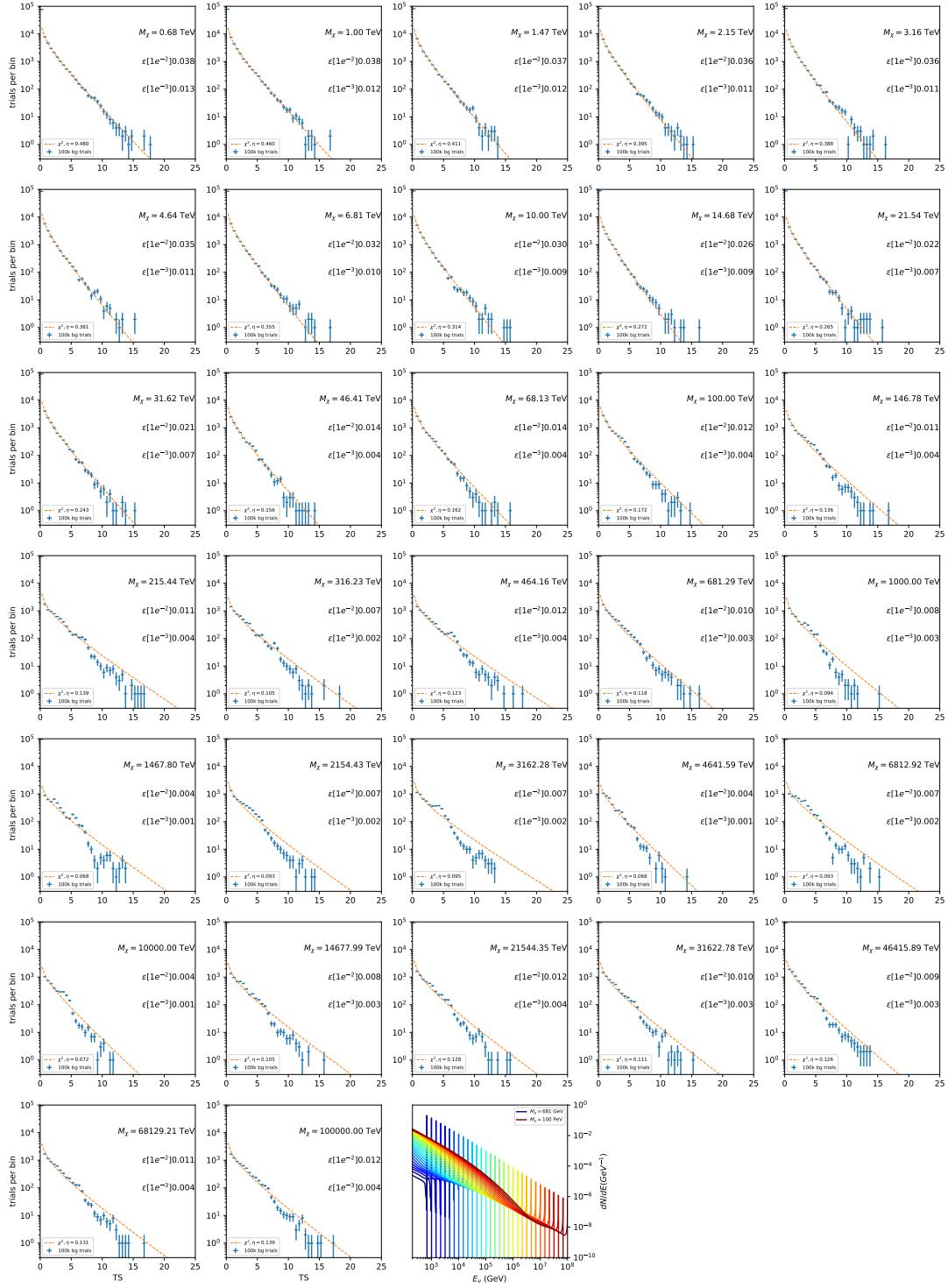


Figure 7.8 Same as Fig. 7.6 for Segue 1  $\chi\chi \rightarrow \nu_\mu\bar{\nu}_\mu$ .

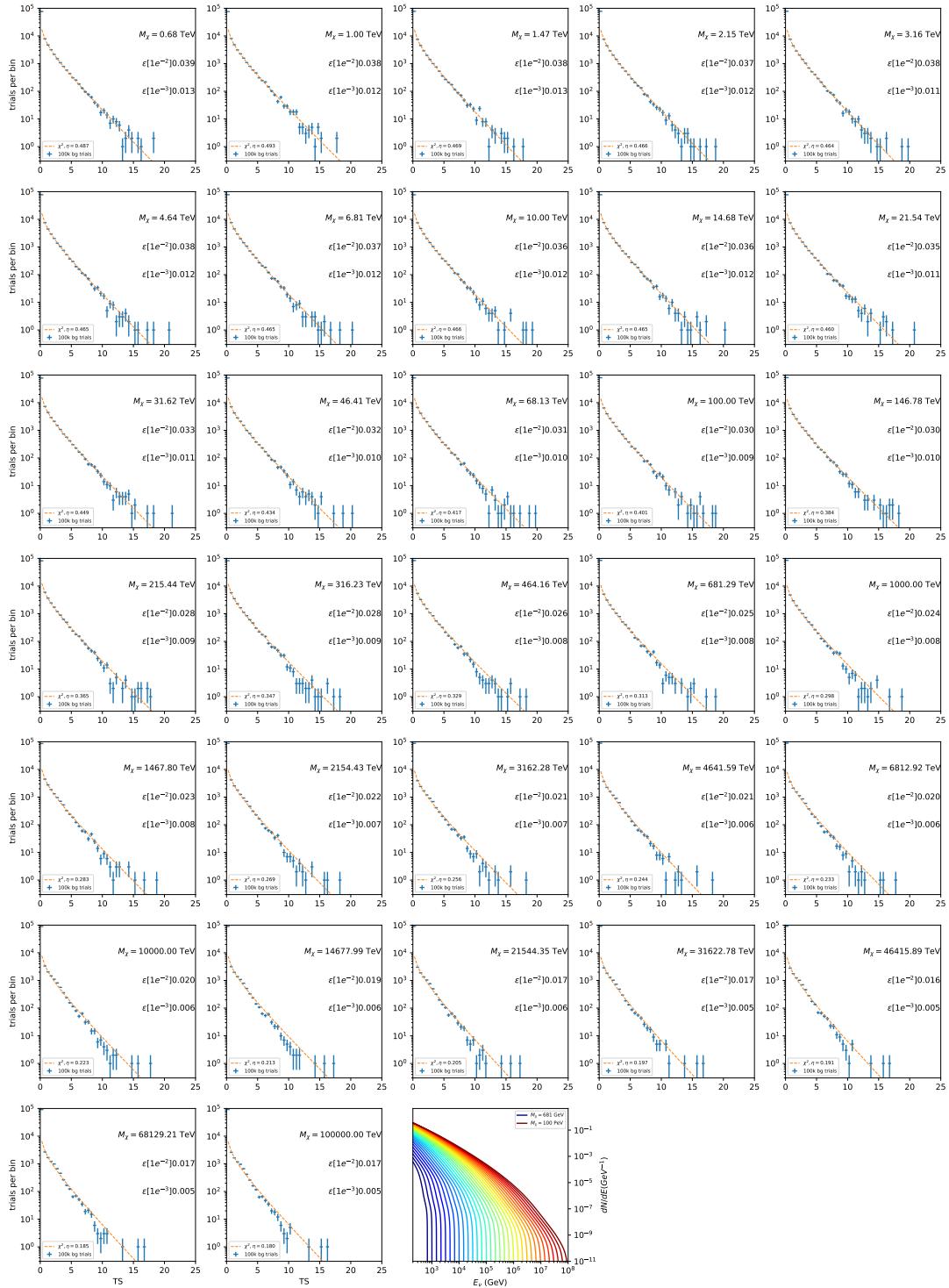


Figure 7.9 Same as Fig. 7.6 for Ursa Major II 1  $\chi\chi \rightarrow b\bar{b}$ .

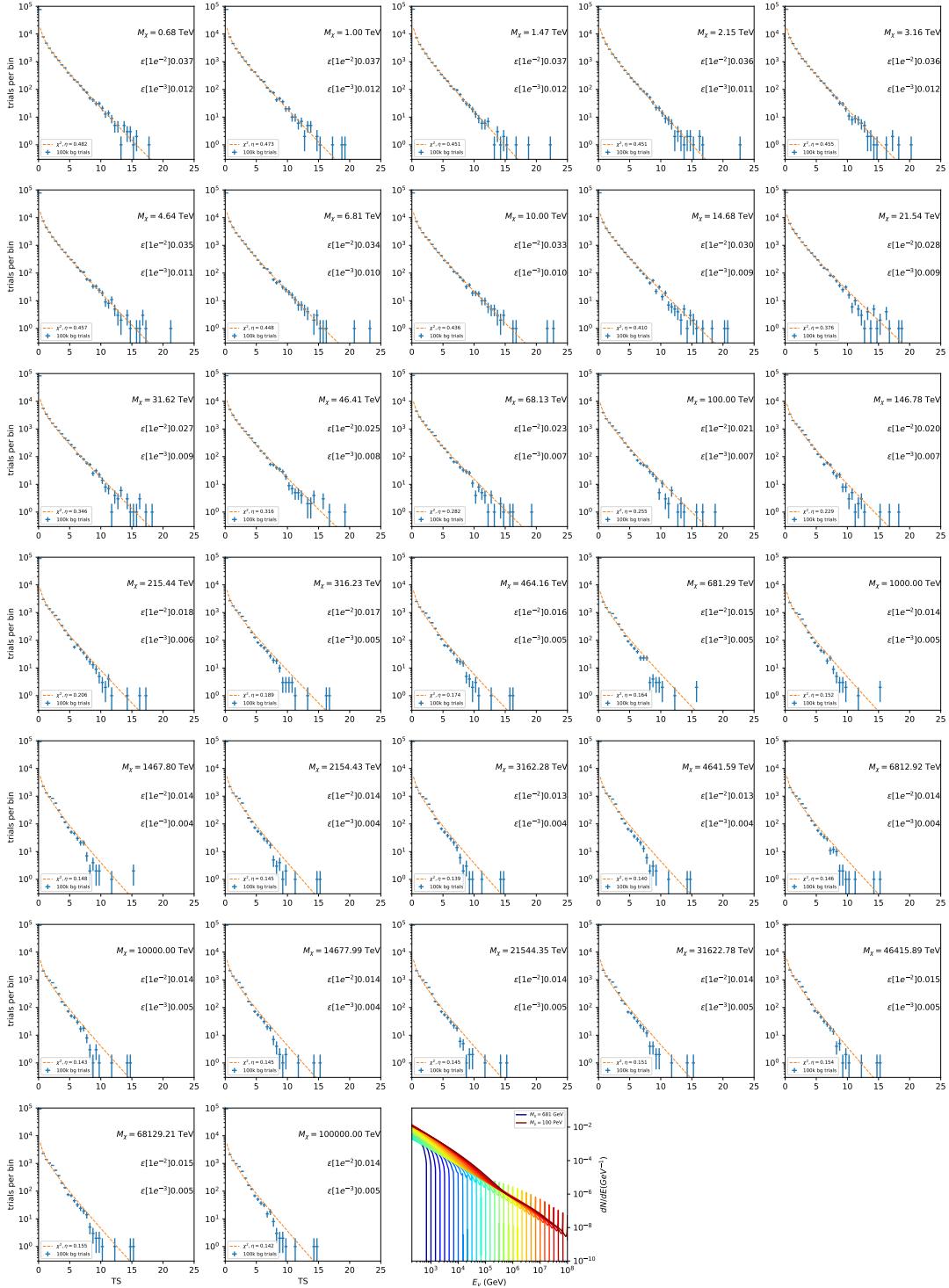


Figure 7.10 Same as Fig. 7.6 for Ursa Major II 1  $\chi\chi \rightarrow \tau\bar{\tau}$ .

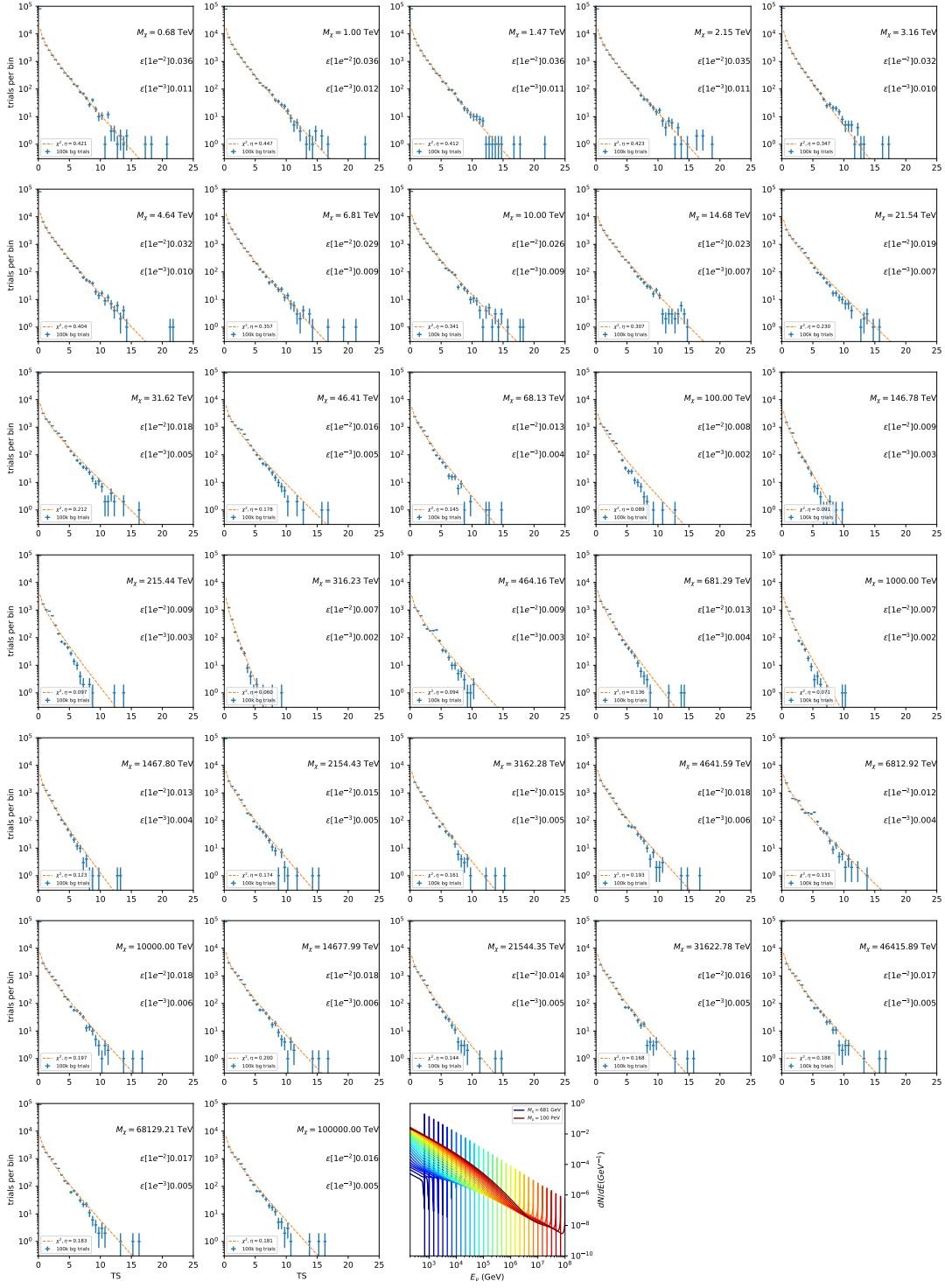


Figure 7.11 Same as Fig. 7.6 for Ursus Major II 1  $\chi\chi \rightarrow \nu_\mu\bar{\nu}_\mu$ .

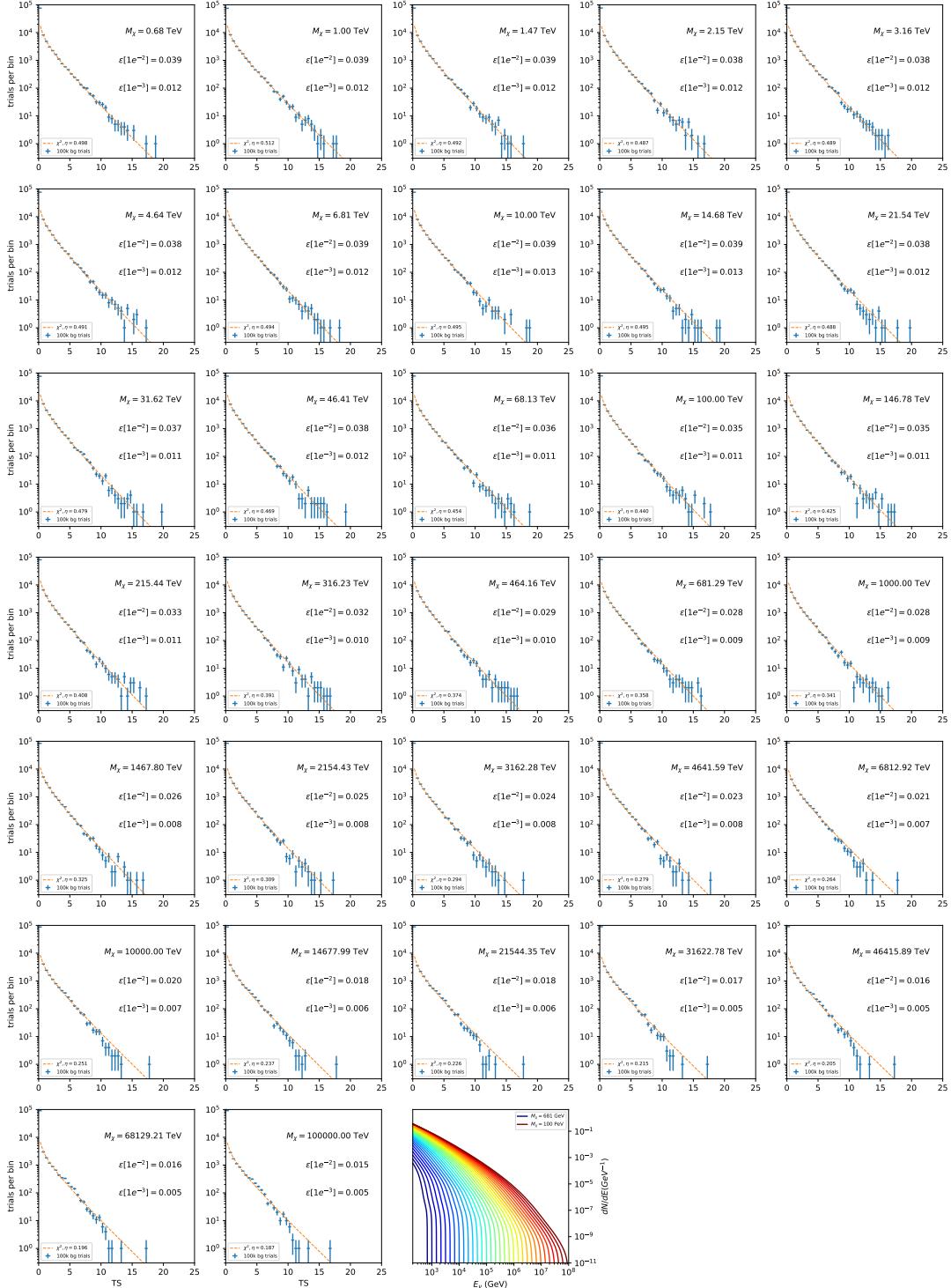


Figure 7.12 Same as Fig. 7.6 for 15, GS J-factor, stacked sources and  $\chi\chi \rightarrow b\bar{b}$ .

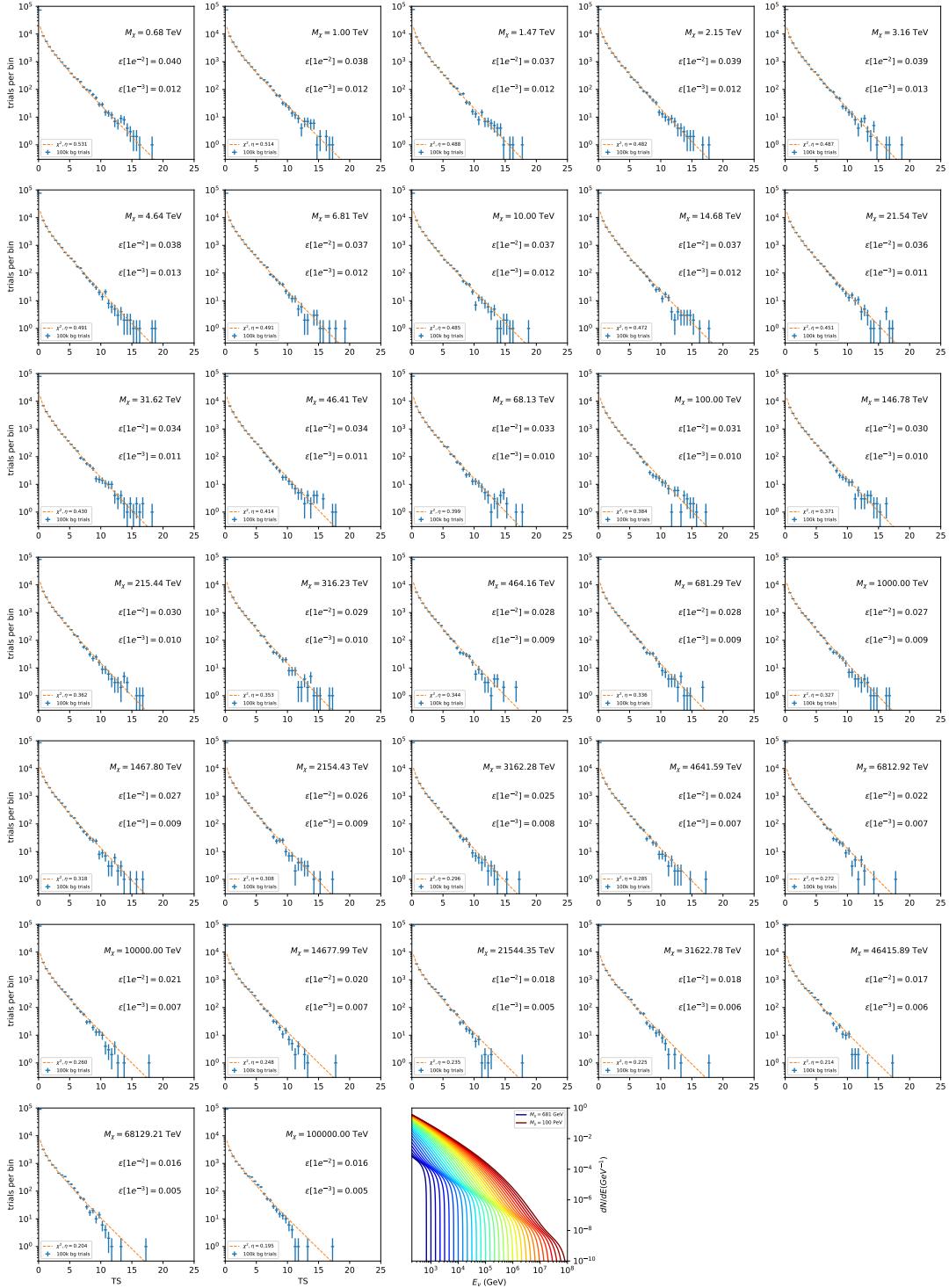


Figure 7.13 Same as Fig. 7.6 for 15,  $\mathcal{GS}$   $J$ -factor, stacked sources and  $\chi\chi \rightarrow t\bar{t}$ .

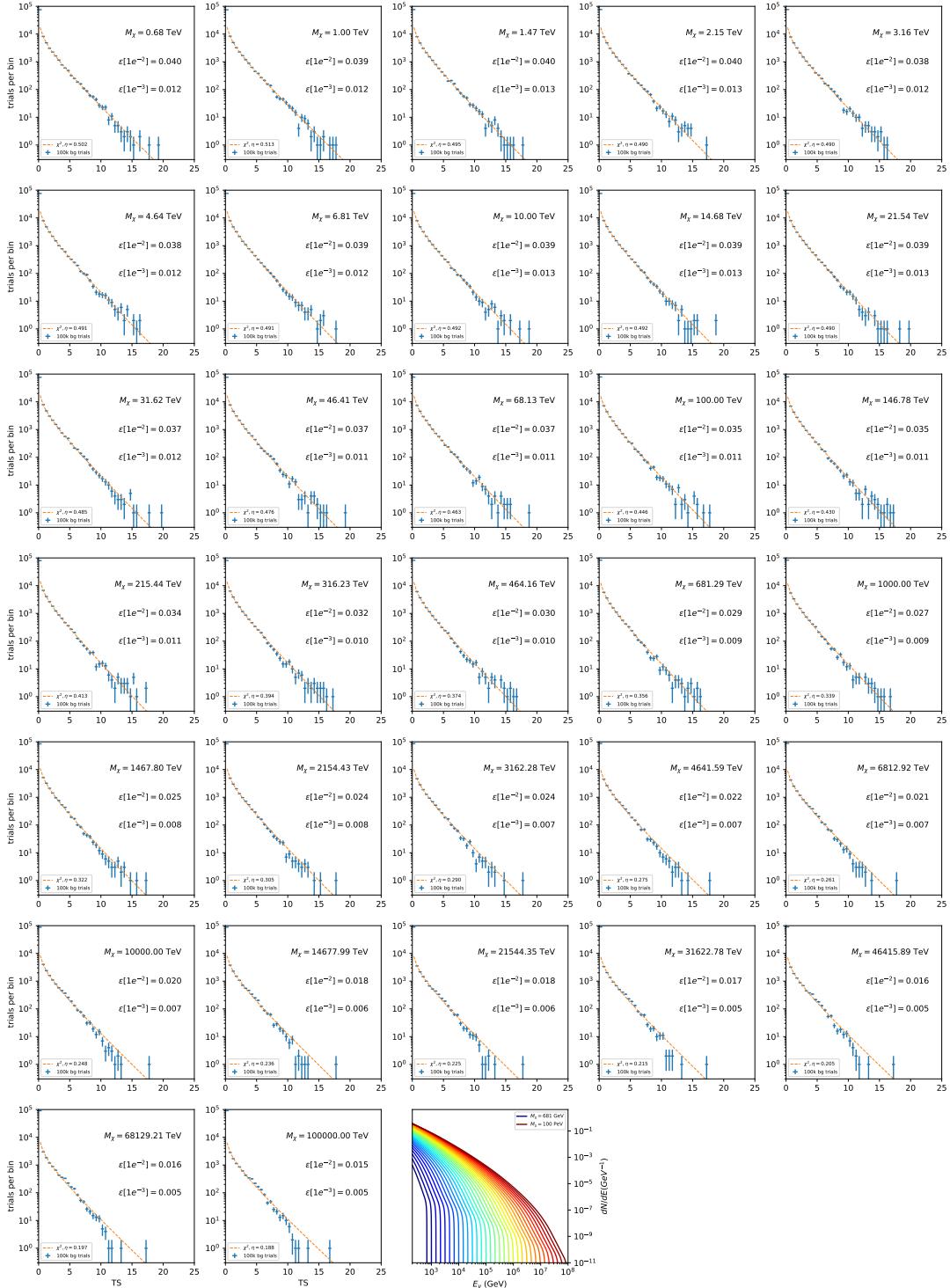


Figure 7.14 Same as Fig. 7.6 for 15,  $\mathcal{GS}$  J-factor, stacked sources and  $\chi\chi \rightarrow u\bar{u}$ .

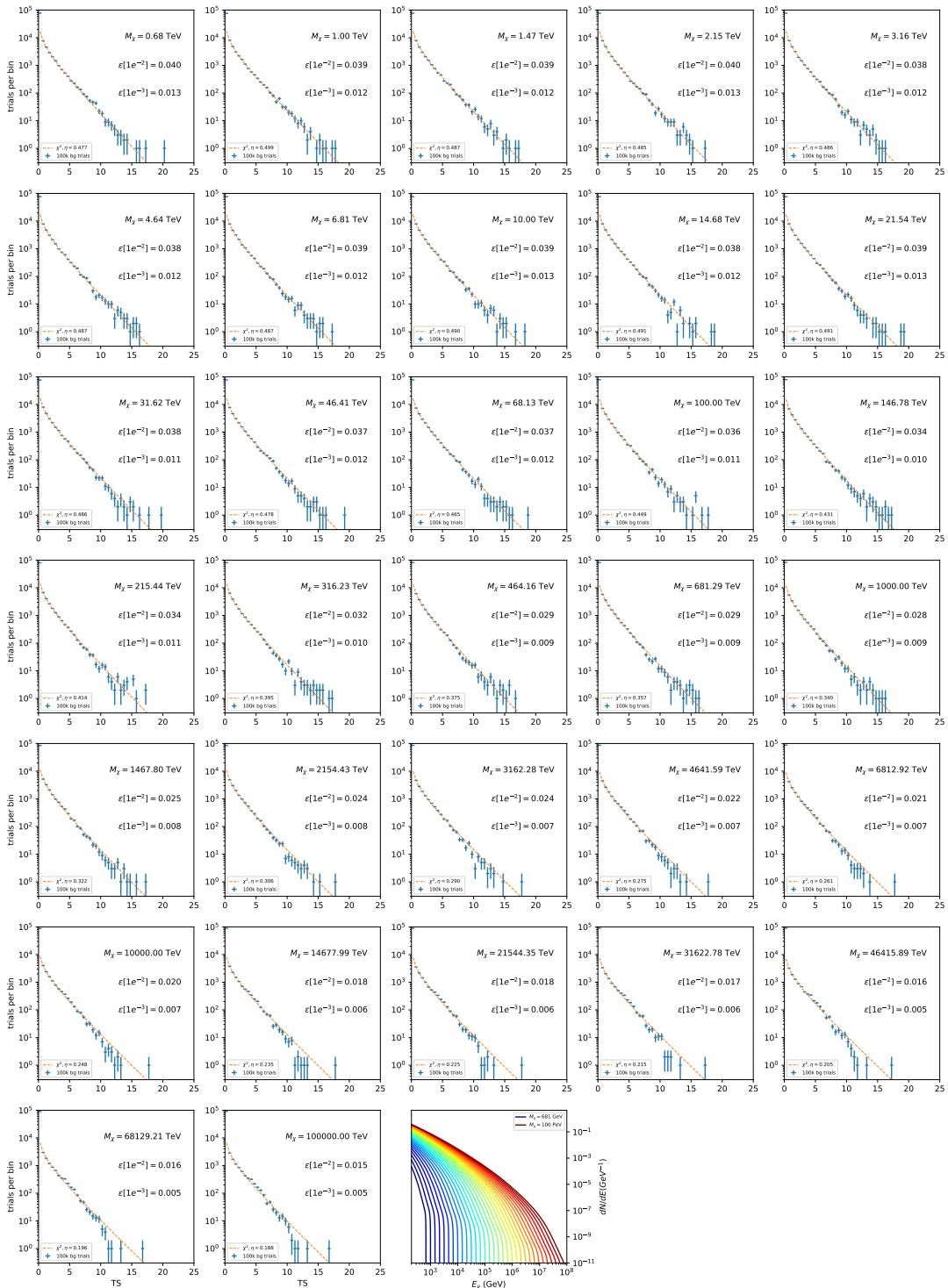


Figure 7.15 Same as Fig. 7.6 for 15,  $\mathcal{GS}$  J-factor, stacked sources and  $\chi\chi \rightarrow d\bar{d}$ .

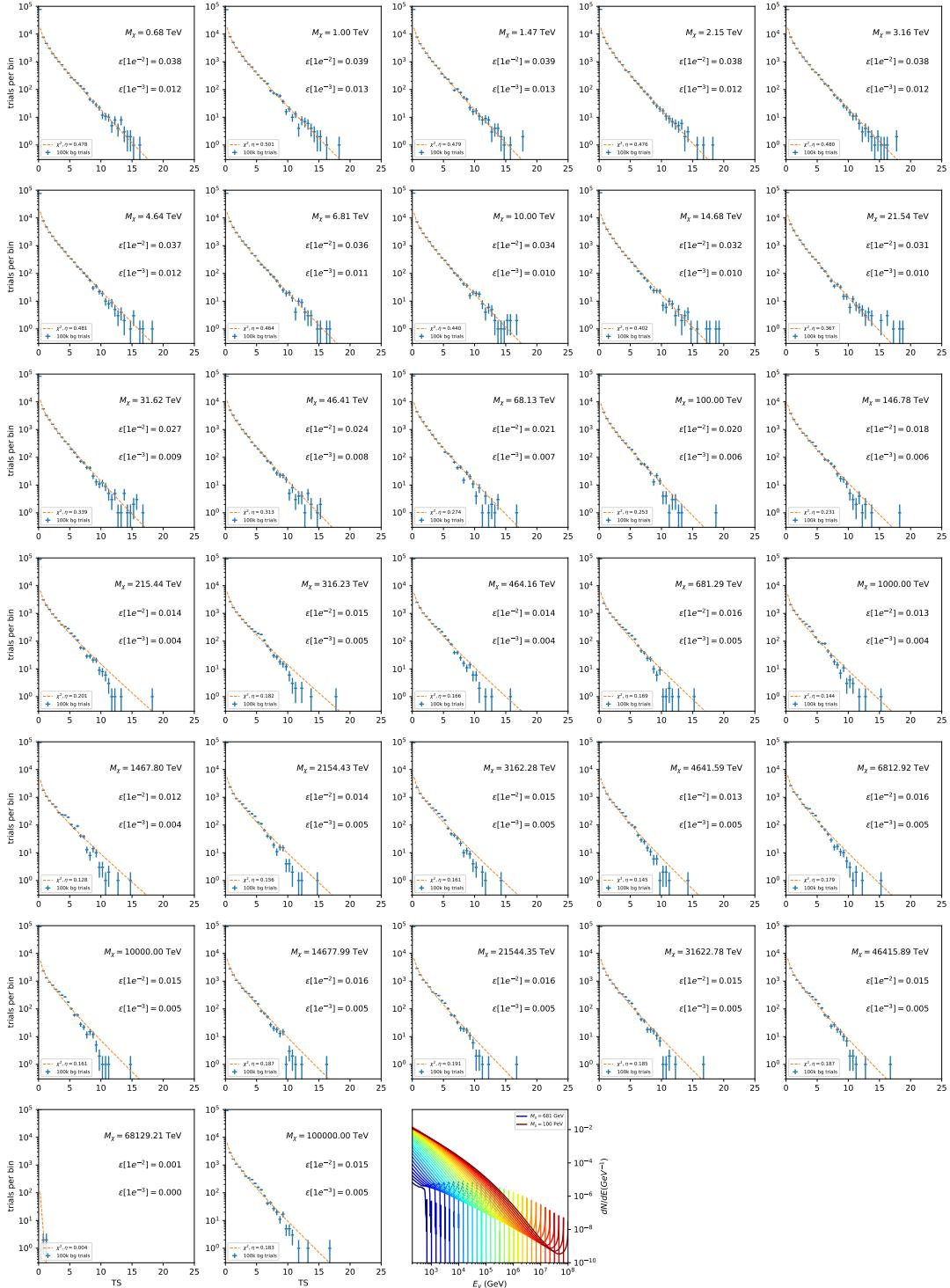


Figure 7.16 Same as Fig. 7.6 for 15,  $\mathcal{GS}$  J-factor, stacked sources and  $\chi\chi \rightarrow e\bar{e}$ .

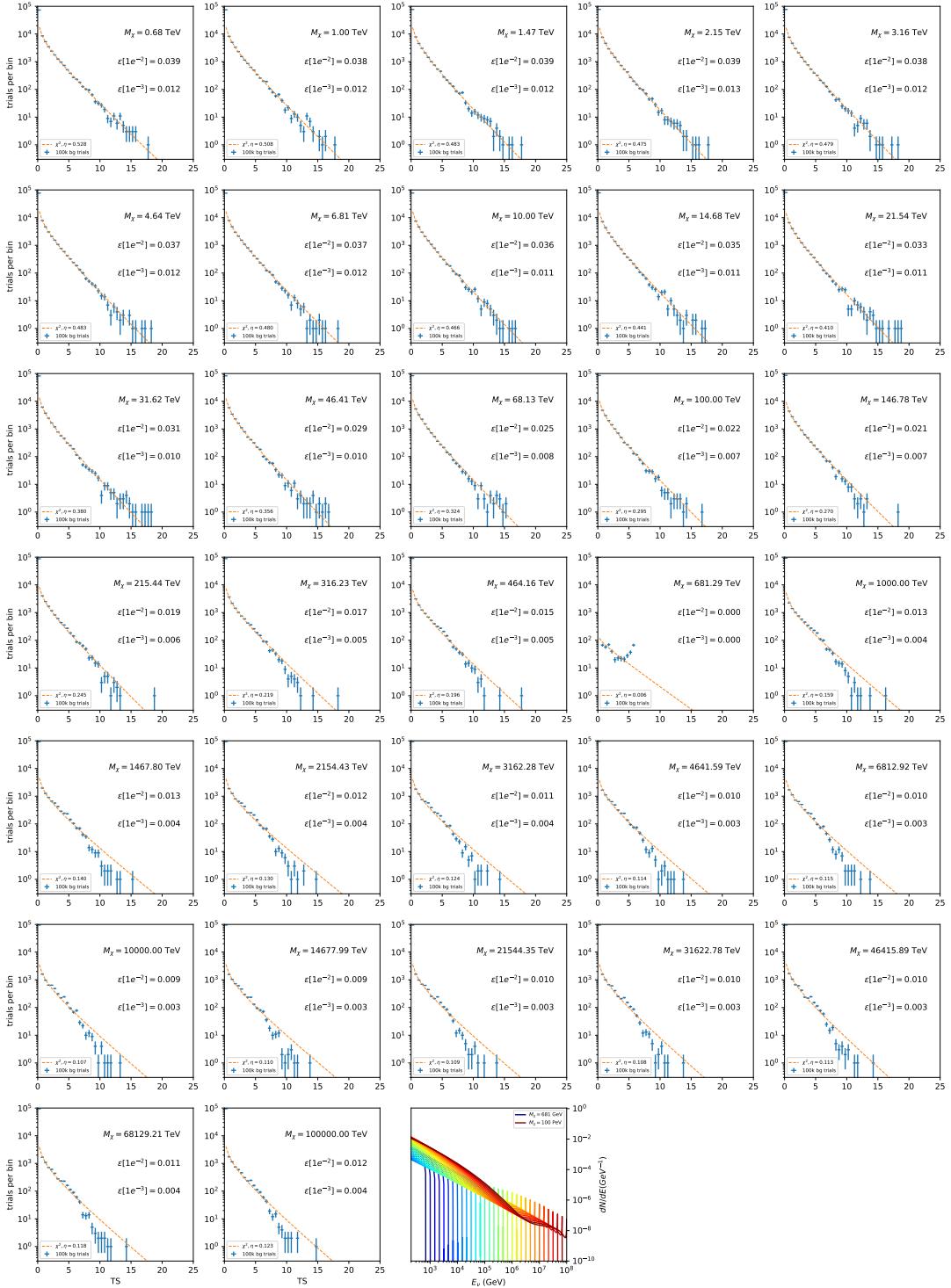


Figure 7.17 Same as Fig. 7.6 for 15,  $\mathcal{GS}$   $J$ -factor, stacked sources and  $\chi\chi \rightarrow \mu\bar{\mu}$ .

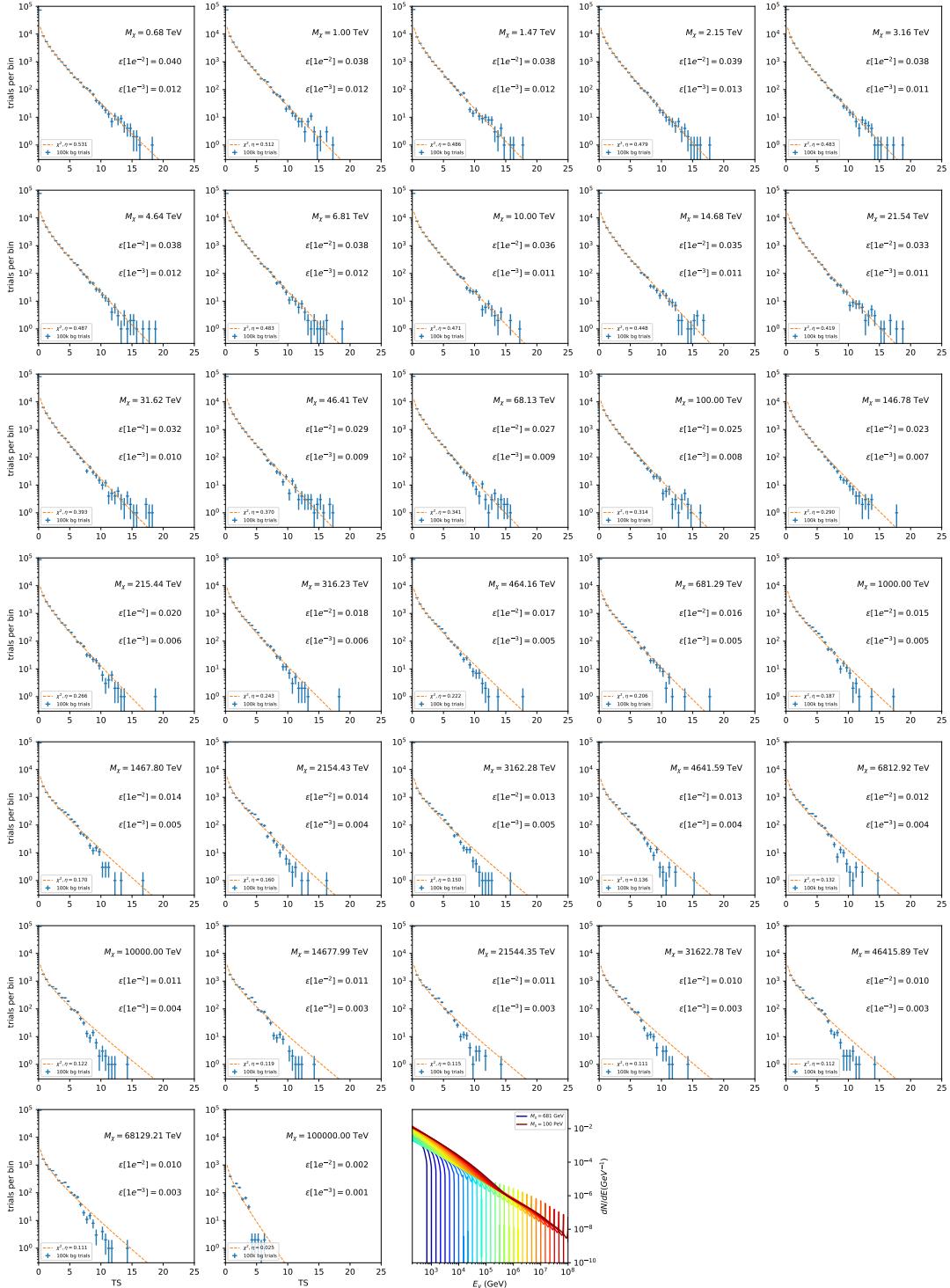


Figure 7.18 Same as Fig. 7.6 for 15,  $\mathcal{GS}$   $J$ -factor, stacked sources and  $\chi\chi \rightarrow \tau\bar{\tau}$ .

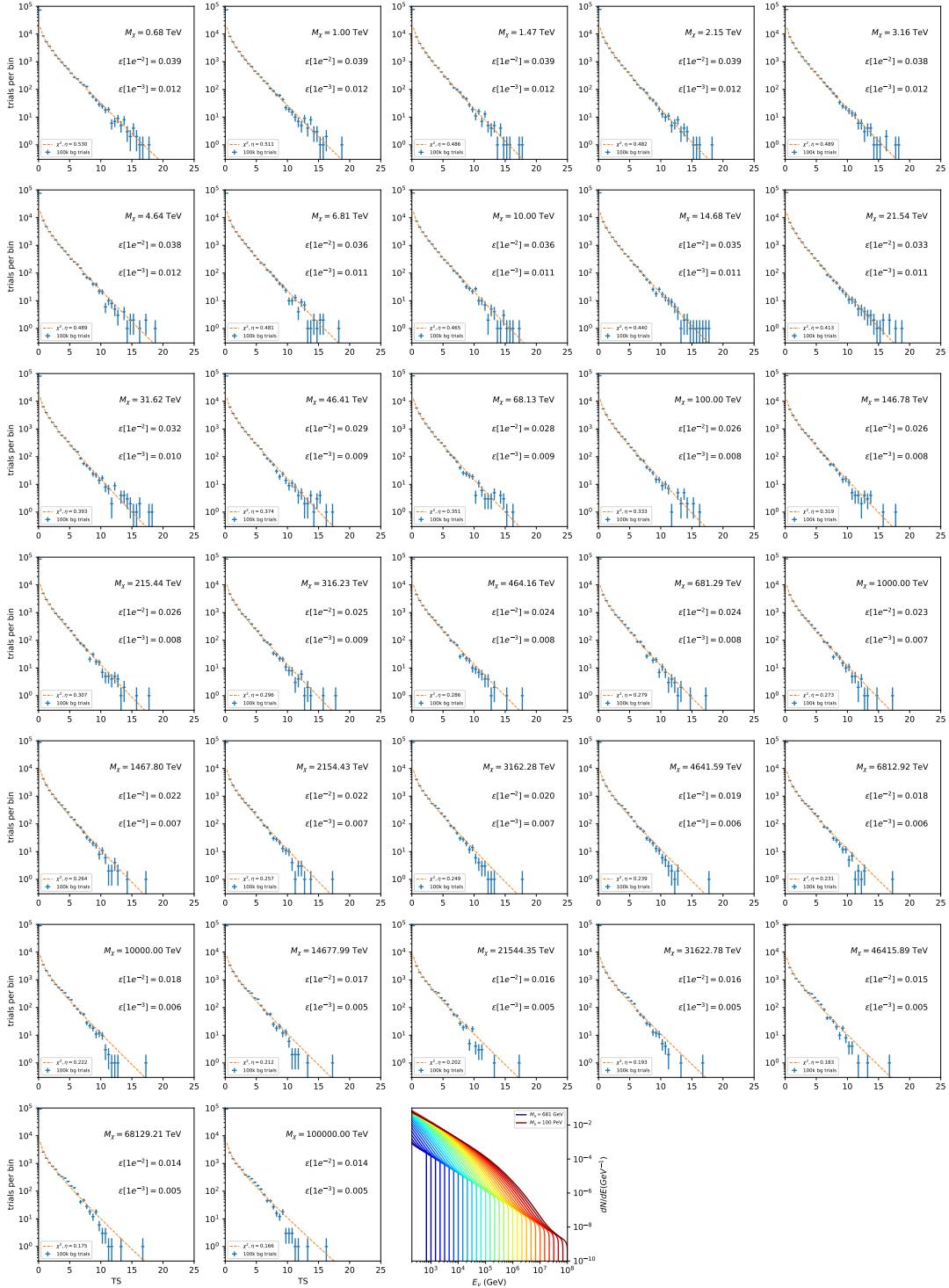


Figure 7.19 Same as Fig. 7.6 for 15,  $\mathcal{GS}$  J-factor, stacked sources and  $\chi\chi \rightarrow W^+W^-$ .

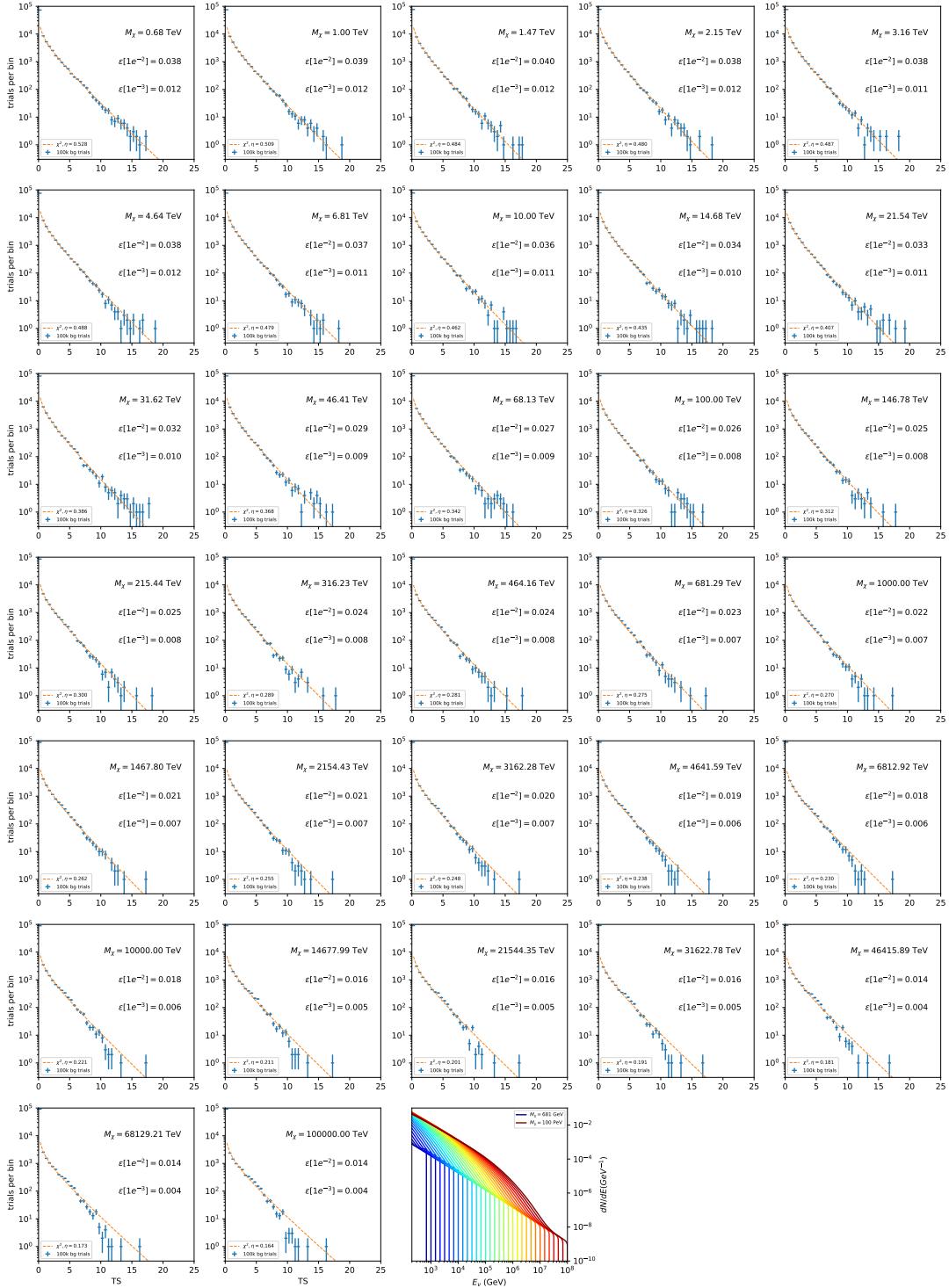


Figure 7.20 Same as Fig. 7.6 for 15,  $\mathcal{GS}$  J-factor, stacked sources and  $\chi\chi \rightarrow ZZ$ .

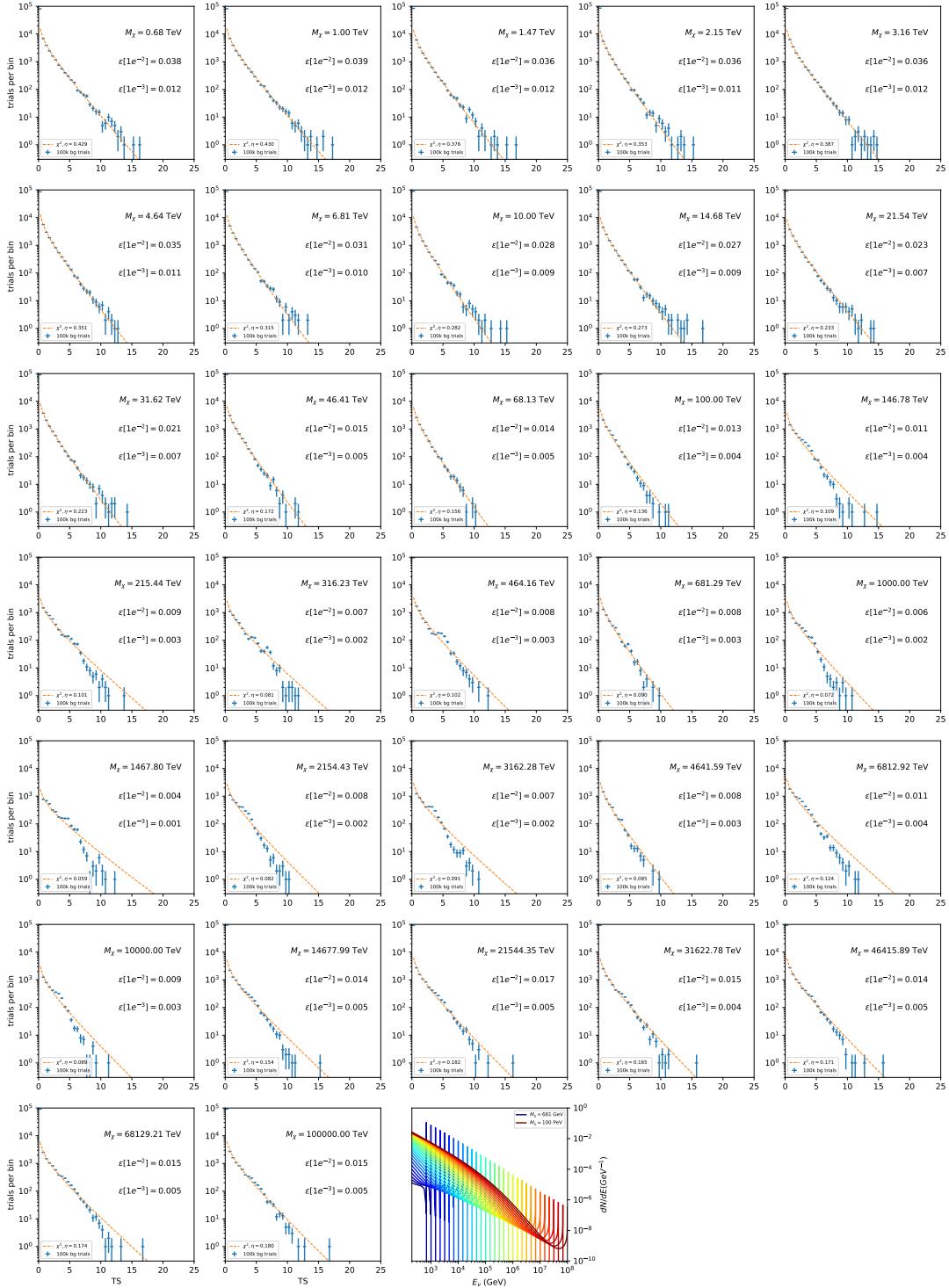


Figure 7.21 Same as Fig. 7.6 for 15,  $\mathcal{GS}$   $J$ -factor, stacked sources and  $\chi\chi \rightarrow \nu_e \bar{\nu}_e$ .

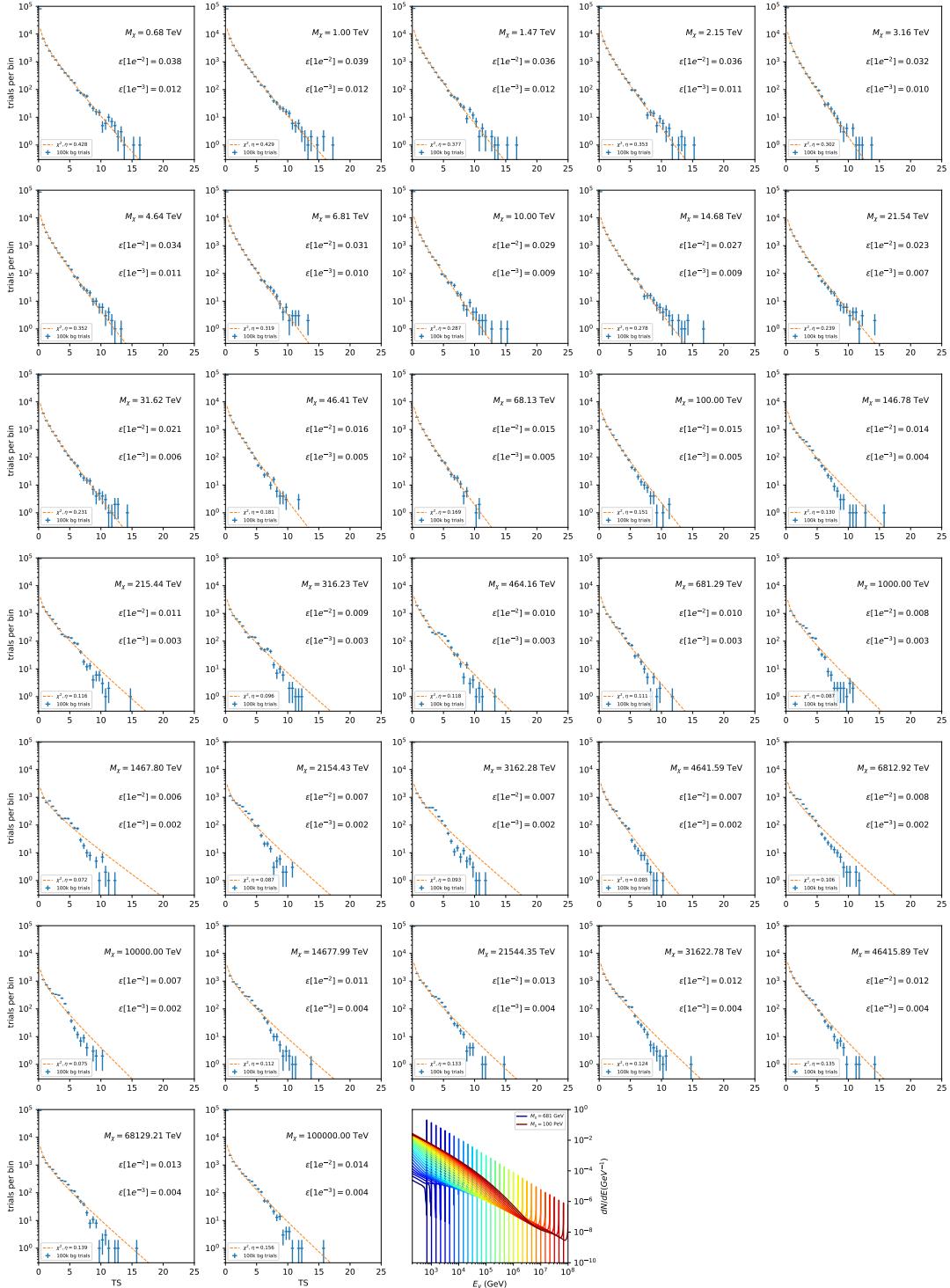


Figure 7.22 Same as Fig. 7.6 for 15, GS J-factor, stacked sources and  $\chi\chi \rightarrow \nu_\mu \bar{\nu}_\mu$ .

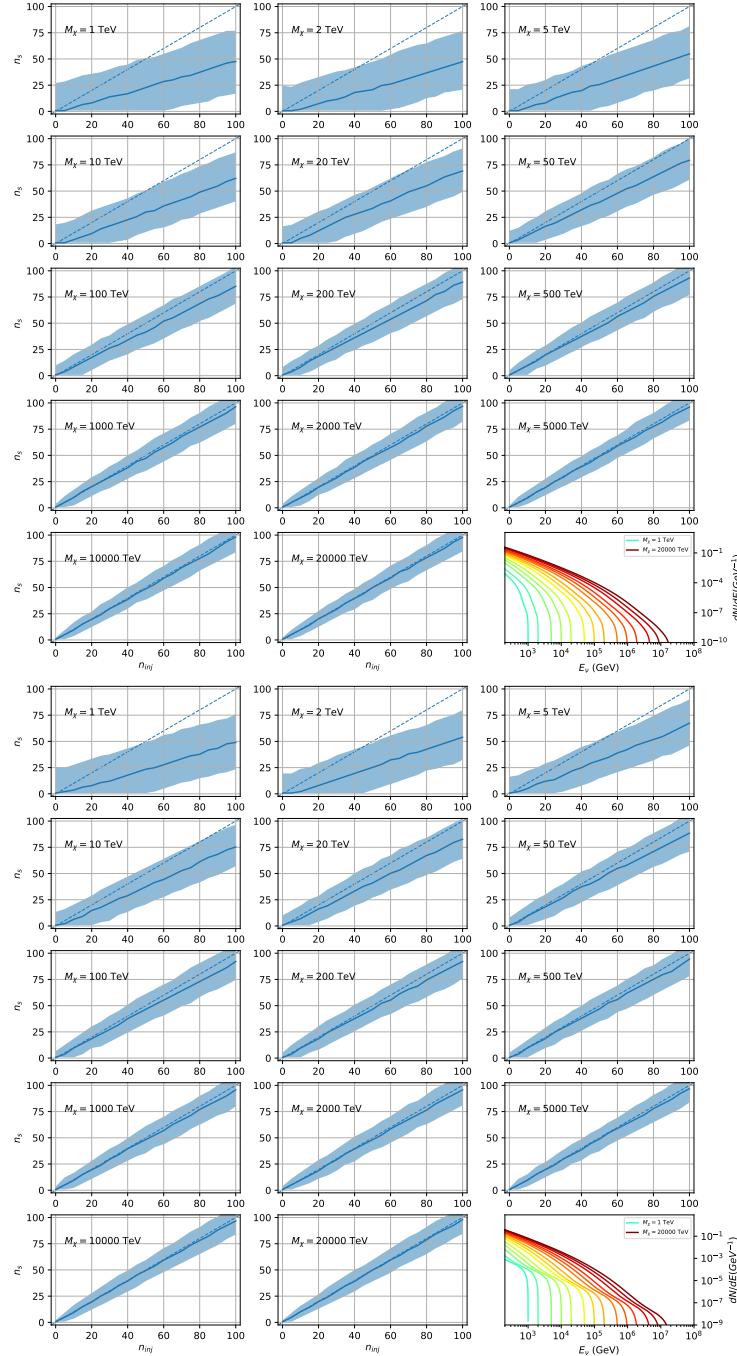


Figure 7.23 Signal Recovery study for an analysis with 15 stacked sources using the  $\mathcal{GS}$   $J$ -factors [45]. Each panel block represents 14 studies for DM mass ranging between 1 TeV and 20 PeV and one annihilation channel. Top panel block is for  $\chi\chi \rightarrow b\bar{b}$ . Bottom panel block is for  $\chi\chi \rightarrow t\bar{t}$ . Each panel block features every spectral model used as input in the bottom-right subpanel. The remaining panels show  $N_{inj}$  as the number of signal events injected into background simulation. Whereas,  $N_s$  is the number of signal events recovered from analyzing the injected simulation. Blue line represents the median values of 100 simulations. Light blue bands show the  $1\sigma$  statistical uncertainty around the median.

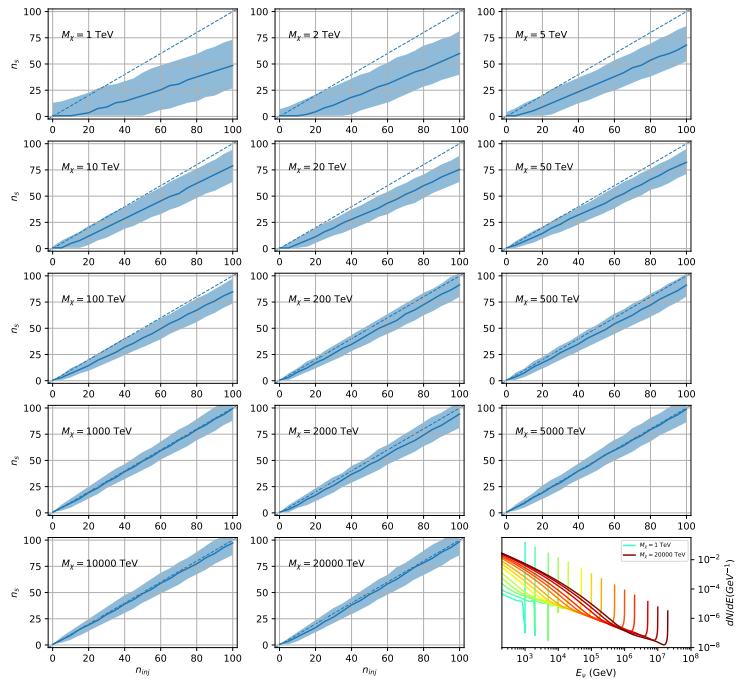


Figure 7.24 Same as Fig. 7.23 but for  $\chi\chi \rightarrow \nu_\mu\bar{\nu}_\mu$ .

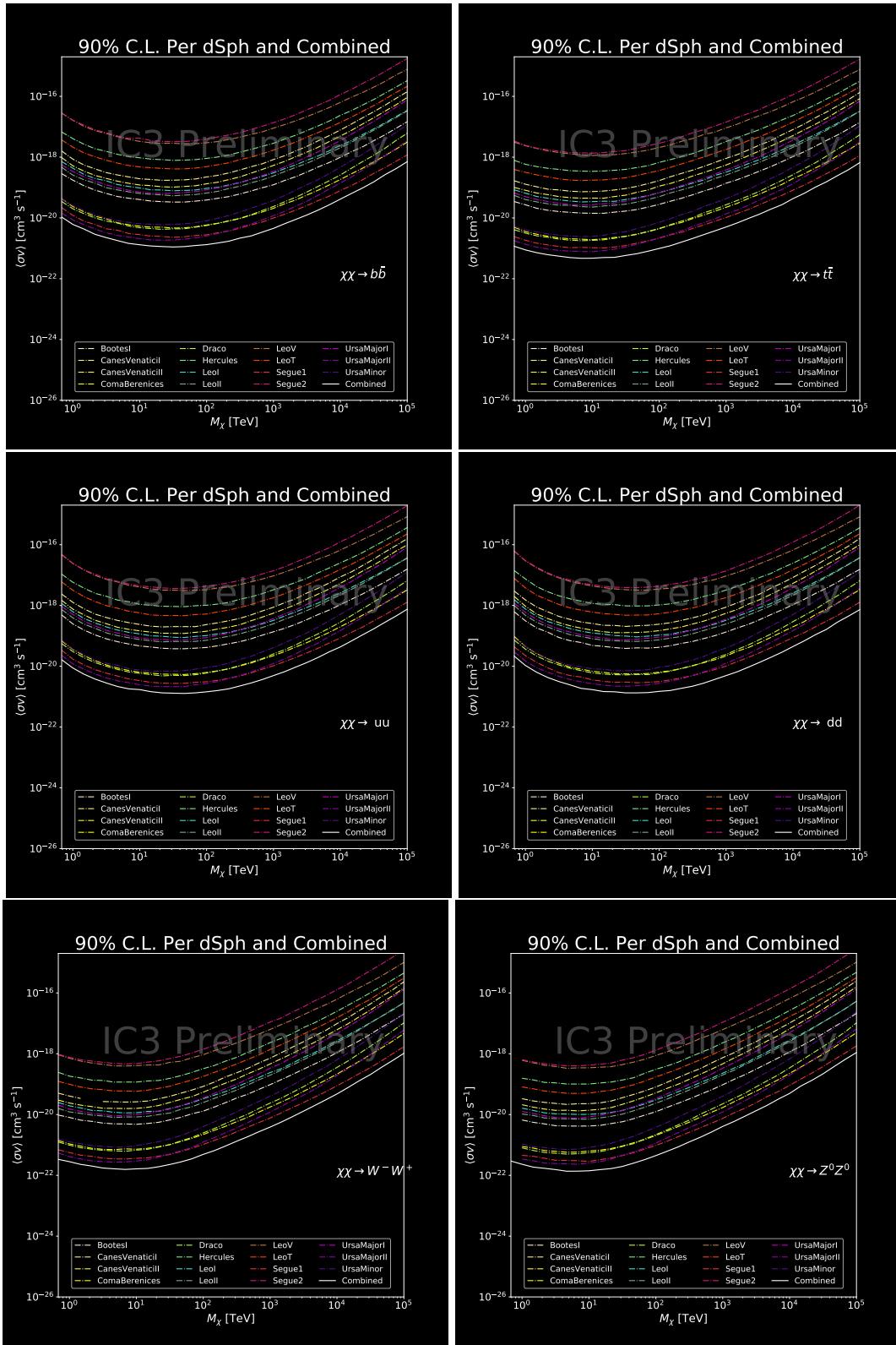


Figure 7.25 Words. I prent Icecibe Sensitivities weeee

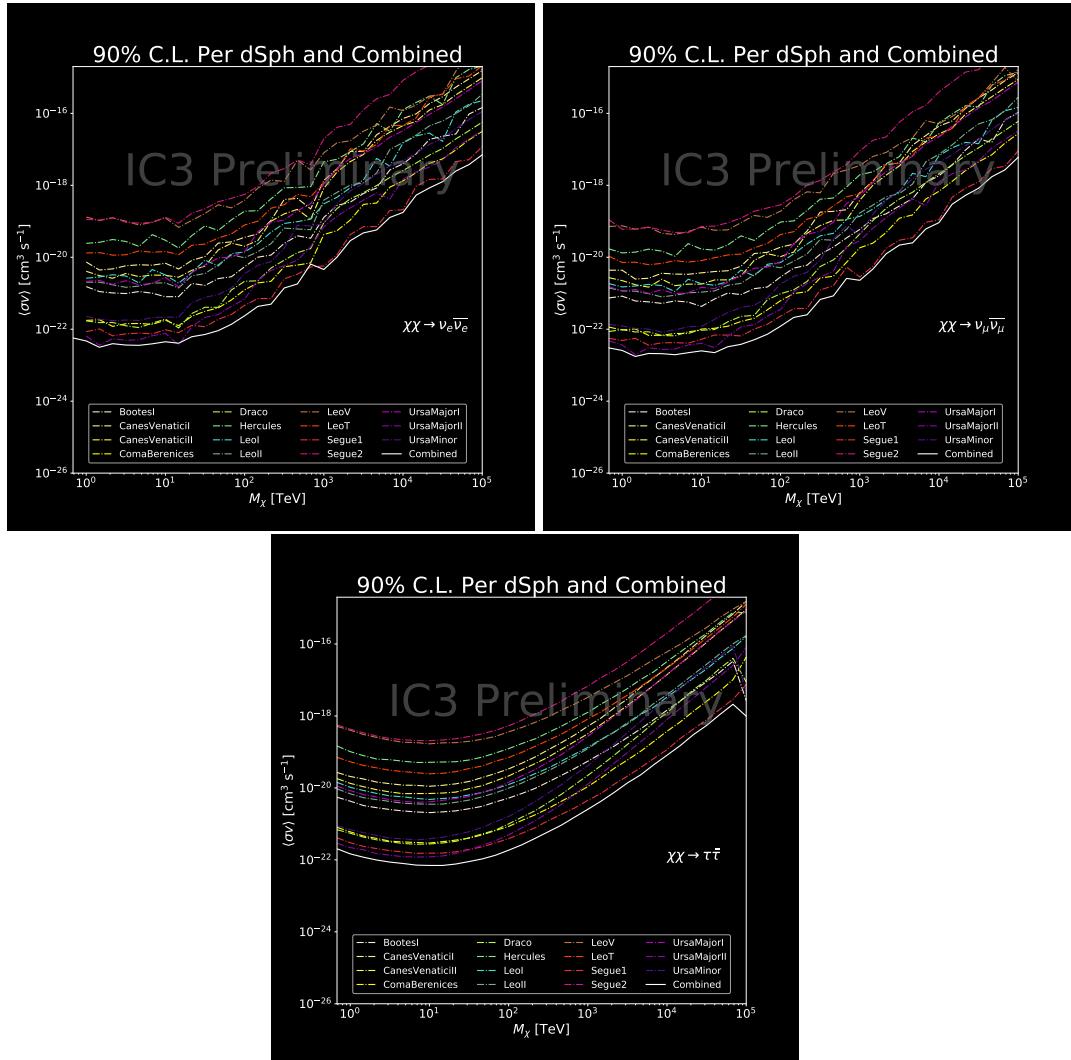


Figure 7.26 Words. I prent Icecibe Sensitivities weeee

## **CHAPTER 8**

### **NU DUCK**

1689

## MULTI-EXPERIMENT SUPPLEMENTARY FIGURES

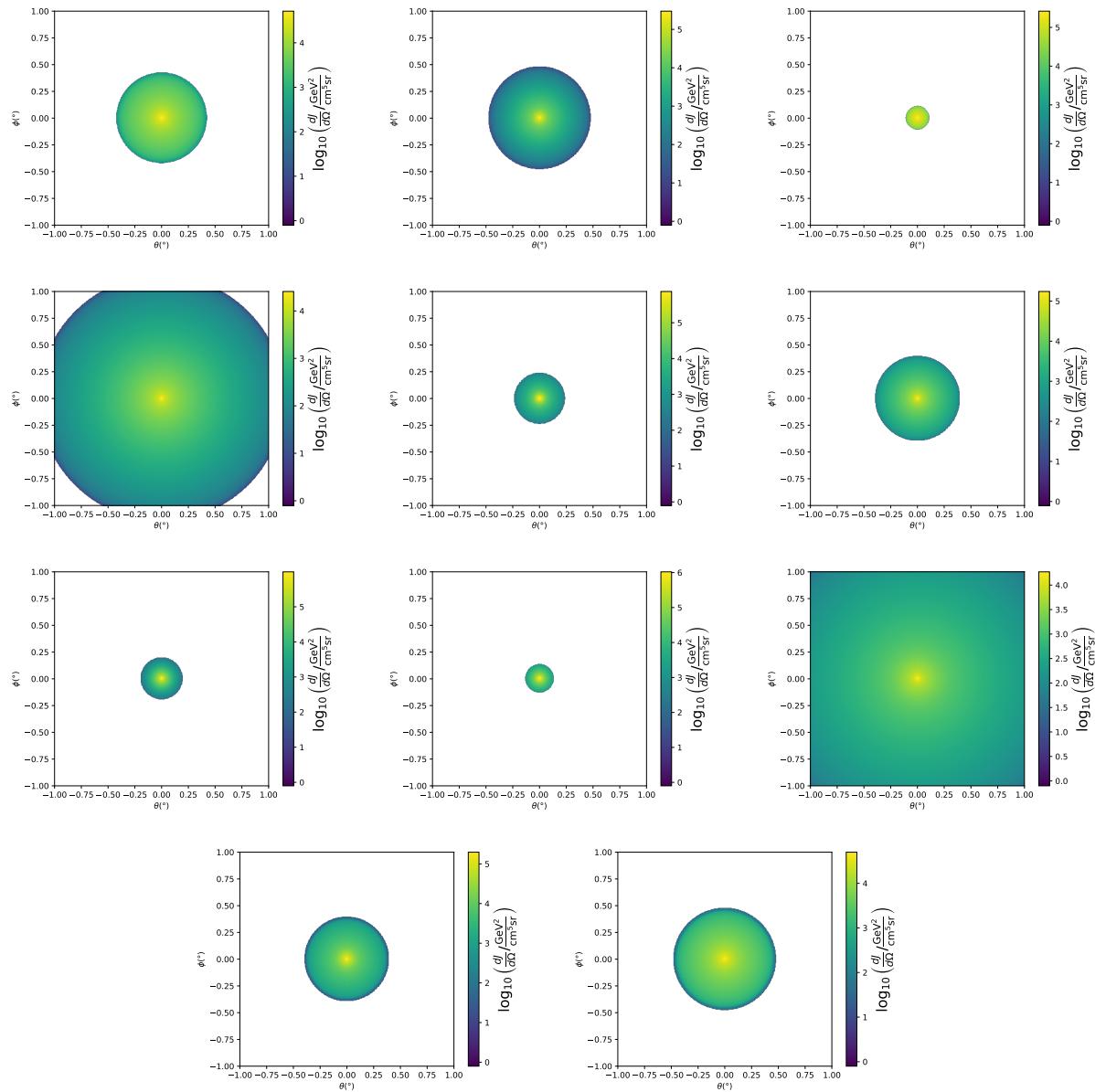


Figure A.1 Sister figure to Figure 5.3. Sources in the first row from left to right: Bootes I, Canes Venatici I, II. In second row: Draco, Hercules, Leo I. In the first row: Leo II, Leo IV, Sextans. In the final row: Ursa Major I, Ursa Major II.

## APPENDIX B

### 1691 MULTITHREADING DARK MATTER ANALYSES SUPPLEMENTARY MATERIAL

#### 1692 B.1 Remaining Spectral Models

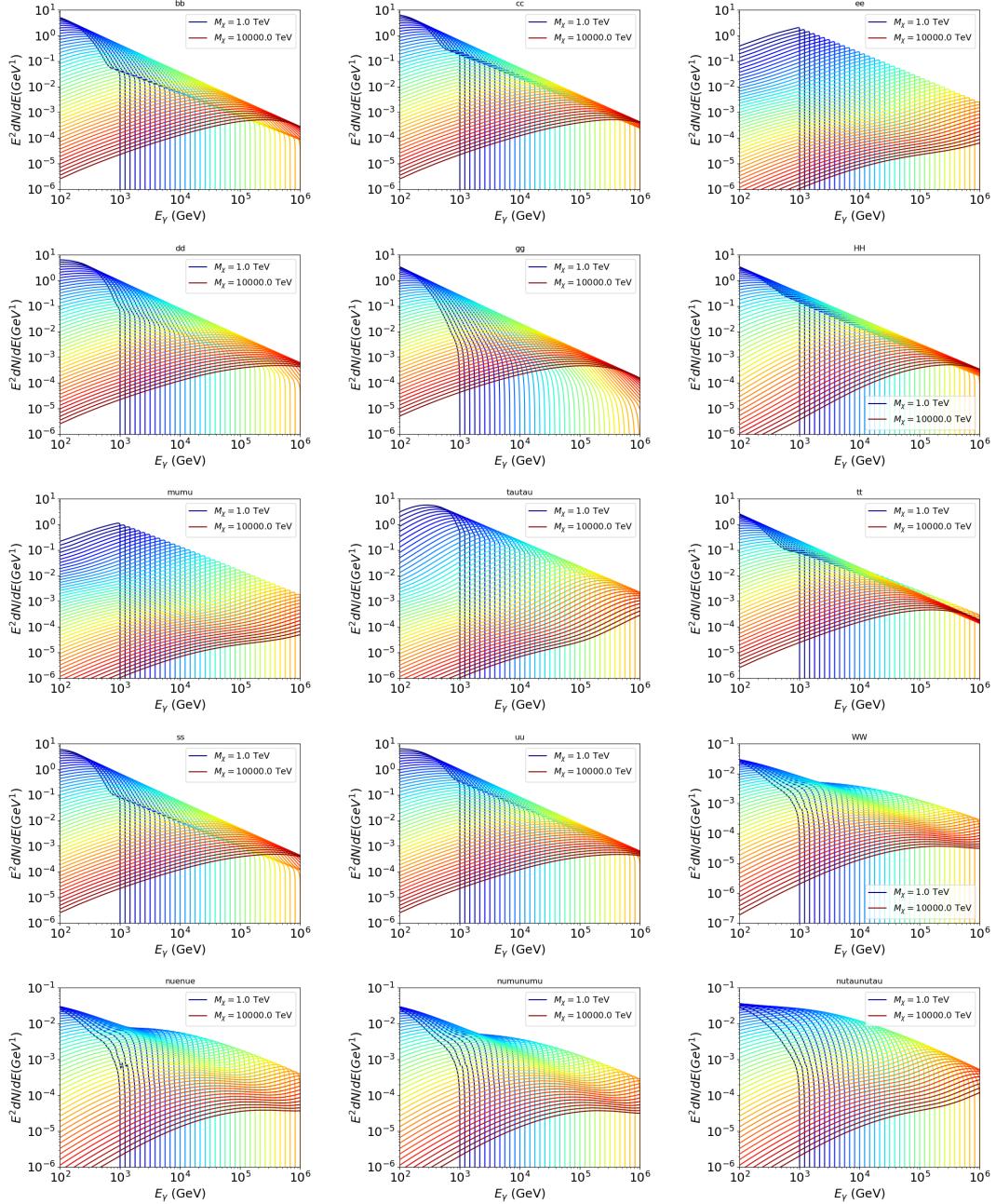


Figure B.1 Sister figure to Figure 6.2 for remaining SM primary annihilation channels studied for this thesis. These did not require any post generation smoothing and so are directly pulled from [65] with a binning scheme most helpful for a HAWC analysis.

## 1693 B.2 mpu\_analysis.py

```
16941 import warnings
16952 with warnings.catch_warnings():
16963     warnings.simplefilter("ignore")
16974 # Python base libraries
16985 import os
16996 import sys
17007 import time
17018 # Import general libraries with namespace
17029 import matplotlib
17030 # Necessary for computing on cluster
17041 matplotlib.use("agg")
17052 import numpy as np
17063 import multiprocessing as mp
17074 # Import HAWC software
17085 sys.path.insert(1, os.path.join(os.environ['HOME'], 'hawc_software', 'threeML-
1709     analysis-scripts', 'fitModel'))
17106 from analysis_modules import *
17117 from threeML import *
17128 from hawc_hal import HAL, HealpixConeROI
17139 from threeML.minimizer.minimization import FitFailed
17140 # Import Dark Matter HAWC Libraries
17151 import analysis_utils as au
17162 import spectra as spec
17173 import sources as srcs
17184
17195 #* READ ONLY PATHS This block will change eventually
17206 MASS_LIST = './plotting/studies/nd/masses.txt'
17217 CHAN_LIST = './plotting/studies/nd/chans.txt'
17228
17239 #* WRITE PATHS, default location is to scratch
17240 OUT_PATH = os.path.join(os.environ["SCRATCH"], os.environ["USER"], 'New_duck')
```

```

17251 print('Our out path is going to be {}'.format(OUT_PATH))
17262
17273 # Define parallel Function. Can also be run serially
17284 def some_mass_fit(sigVs, datalist, shape, dSph, jfactor, mass_chan,
17295                 progress=None, log_file='', queue=None, i_job=0):
17306
17317     if progress is None:
17328         progress = [0]
17339     else: # Create log files for each thread
17340         log_file = log_file.replace('.log', '_ThreadNo_')
17351         log_file = log_file + str(i_job) + ".log"
17362         sys.stdout = open(log_file, "w")
17373
17384     fits = []
17395
17406     try:
17417         for m_c in mass_chan:
17428             print(f'Mass chan tuple: {m_c}')
17439             mass = int(m_c[0])
17440             ch = m_c[1]
17451             # Build path to output files
17462             outPath = os.path.join(OUT_PATH, ch, dSph)
17473             au.ut.ensure_dir(outPath)
17484
17495             if progress[i_job] < 0:
17506                 # If the master gets a Keyboard interrupt, commit suicide.
17517                     break
17528
17539                 ### Start Model Building for DM mass and SM channel #####
17540                 spectrum = spec.DM_models.HDMSpectra()
17551                 spectrum.set_channel(ch)
17562
17573                 myDwarf = ExtendedSource(dSph, spatial_shape=shape,

```

```

17584                     spectral_shape=spectrum)
17595
17606             spectrum.J = jfactor * u.GeV**2 / u.cm**5
17617             spectrum.sigmav = 1e-24 * u.cm**3 / u.s
17628             spectrum.set_dm_mass(mass * u.GeV)
17639
17640             spectrum.sigmav.bounds = (1e-30, 1e-12)
17651             model = Model(myDwarf)
17662             ##### End model Building #####
17673
17684             jl = JointLikelihood(model, datalist, verbose=False)
17695
17706             try:
17717                 result, lhdf = jl.fit(compute_covariance=False)
17728                 ts = -2.* (jl.minus_log_like_profile(1e-30) - jl.
1773 _current_minimum)
17749                 # Also profile the LLH vs sv
17750                 ll = jl.get_contours(spectrum.sigmav, sigVs[0],
17761                                     sigVs[-1], len(sigVs),
17772                                     progress=False, log=['False'])
17783
17794                 sigv_ll = au.ut.calc_95_from_profile(ll[0], ll[2])
17805                 # Write results to file
17816                 outFileLL = outPath+f"/LL_{dSph}_{ch}_mass{int(mass)}_GD.txt"
17827                 np.savetxt(outFileLL, (sigVs, ll[2]),
17838                               delimiter='\t', header='sigV\tLL\n')
17849
17850                 with open(outPath+f"/results_{dSph}_{ch}_mass{int(mass)}_GD.
1786 txt", "w") as results_file:
17871                     results_file.write("mDM [GeV]\tsigV_95\tsigV_B.F.\n")
17882
17893                     results_file.write("%i\t%.5E\t%.5E\t%.5E"%(mass, sigv_ll,
17904                                         ts, result.value[0]))

```

```

17915         # End write to file
17926     except FitFailed: # Don't kill all threads if a fit fails
17937         print("Fit failed. Go back and calculate this spectral model
1794     later")
17958         fits.append((ch, mass, -1, -1))
17969         with open(log_file+'.fail', 'w') as f_file:
17970             f_file.write(f'{ch}, {mass}\n')
17981
17992         progress[i_job] += 1
18003         matplotlib.pyplot.close() # Prevent leaky memory
18014
18025         fits.append((ch, mass, result.value[0], ts))
18036         progress[i_job] += 1
18047         matplotlib.pyplot.close()
18058     except KeyboardInterrupt:
18069         progress[i_job] = -1
18070
18081     fits = np.array(fits)
18092     if queue is None:
18103         return fits
18114     else:
18125         queue.put((i_job, fits))
18136
18147 def main(args):
18158     masses = np.loadtxt(MASS_LIST, dtype=int)
18169     chans = np.loadtxt(CHAN_LIST, delimiter=',', dtype=str)
18170     mass_chan = au.ut.permute_lists(chans, masses)
18181
18192     print(f"DM masses for this study are: {masses}")
18203     print(f"SM Channels for this study are XX -> {chans}")
18214     print(mass_chan)
18225
18236 # extract information from input argument

```

```

18247 dSph = args.dSph
18258 data_mngr = au.ut.Data_Selector('P5_NN_2D')
18269 dm_profile = srcts.Spatial_DM(args.catalog, dSph, args.sigma, args.decay)
18270
18281     ### Extract Source Information ####
18292 if dm_profile.get_dec() < -22.0 or dm_profile.get_dec() > 62.0:
18303     raise ValueError("HAWC can't see this source D: Exitting now...")
18314
18325 print(f'{dSph} information')
18336 print(f'jfac: {dm_profile.get_factor()}\tRA: {dm_profile.get_ra()}\tDec: {dm_profile.get_dec()}\n')
18347
18358 shape = SpatialTemplate_2D(fits_file=dm_profile.get_src_fits())
18379     ### Finish Extract Source Information ####
18380
18391     ### LOAD HAWC DATA ####
18402 roi = HealpixConeROI(data_radius=2.0, model_radius=8.0,
18413                         ra=dm_profile.get_ra(), dec=dm_profile.get_dec())
18424 bins = choose_bins.analysis_bins(args, dec=dm_profile.get_dec())
18435
18446 hawc = HAL(dSph, data_mngr.get_datmap(), data_mngr.get_detres(), roi)
18457 hawc.set_active_measurements(bin_list=bins)
18468 datalist = DataList(hawc)
18479     ### FINISH LOAD HAWC DATA ####
18480
18491 # set up SigV sampling. This sample is somewhat standardized
18502 sigVs = np.logspace(-28,-18, 1001, endpoint=True) # NOTE This will change
1851 with HDM
18523
18534 if args.n_threads == 1:
18545     # No need to start || programming just iterate over the masses
18556     kw_arg = dict(sigVs=sigVs, datalist=datalist, shape=shape, dSph=dSph,
18567                     jfactor=dm_profile.get_factor(), mass_chan=mass_chan,

```

```

18578                 log_file=args.log)
18589             some_mass_fit(**kw_arg)
18590         else:
18601             # I Really want to suppress TQMD output
18612             from tqdm import tqdm
18623             from functools import partialmethod
18634             tqdm.__init__ = partialmethod(tqdm.__init__, disable=True)
18645
18656             x = np.array_split(mass_chan, args.n_threads)
18667             n_jobs = len(x)
18678
18689             print("Thread jobs summary by mass and SM channel")
18690             for xi in x:
18701                 print(f'{xi}')
18712
18723             queue = mp.Queue()
18734             progress = mp.Array('i', np.zeros(n_jobs, dtype=int))
18745
18756             # Define task pool that will be split amongsts threads
18767             kw_args = [dict(sigVs=sigVs, datalist=datalist, shape=shape,
18778                             dSph=dSph, jfactor=dm_profile.get_factor(),
18789                             mass_chan=mass_chan, progress=progress,
18790                             queue=queue, i_job=i, log_file=args.log)
18801                     for i, mass_chan in enumerate(x)]
18812
18823             # Define each process
18834             procs = [mp.Process(target=some_mass_fit, kwargs=kw_args[i]) \
18845                     for i in range(n_jobs)]
18856
18867             ### Start MASTER Thread only code block ###
18878             # Begin running all child threads
18889             for proc in procs: proc.start()
18890

```

```

18901     try:
18912         # In this case, the master does nothing except monitor progress of
1892         the threads
18933         # In an ideal world, the master thread also does some computation.
18944         n_complete = np.sum(progress)
18955         while_count = 0
18966
18977             while n_complete < len(mass_chan):
18988
18999                 if np.any(np.asarray(progress) < 0):
19000                     # This was no threads are stranded when killing the script
19011                     raise KeyboardInterrupt()
19022                     if while_count%1000 == 0:
19033                         print(f"{np.sum(progress)} of {len(mass_chan)} finished")
19044
19055                     n_complete = np.sum(progress)
19066                     time.sleep(.25)
19077                     while_count += 1
19088
19099             except KeyboardInterrupt:
19100                 # signal to jobs that it's time to stop
19111                     for i in range(n_jobs):
19122                         progress[i] = -2
19133                         print('\nKeyboardInterrupt: terminating early.')
19144                     ### End MASTER Thread only code block ###
19155
19166                     fitss = [queue.get() for proc in procs]
19177                     print(fitss)
19188                     print(f'Thread statuses: {progress[:]}')
19199
19200                     # putting results in a file
19211
19222                     print("QUACK! All Done!")

```

```

19233
19244
19255 if __name__ == '__main__':
19266     import argparse
19277
19288     p = argparse.ArgumentParser(description="Run a DM annihilation analysis on
1929      a dwarf spheroidal for a specific SM channel for DM masses [1 TeV to 10
1930      PeV]")
19319
19320     # Dwarf spatial modeling arguements
19331     p.add_argument("-ds", "--dSph", type=str,
19342             help="dwarf spheroidal galaxy to be studied", required=
1935     True)
19363     p.add_argument("-cat", "--catalog", type=str, choices=['GS15', 'LS20'],
19374             default='LS20', help="source catalog used")
19385     p.add_argument("-sig", "--sigma", type=int, choices=[-1, 0, 1], default=0,
19396             help="Spatial model uncertainty. 0 corresponds to the
1940 median. +/-1 correspond to the +/-1sigma uncertainty respectively.")
19417
19428     # Arguements for the energy estimators
19439     p.add_argument("-e", "--estimator", type=str,
19440             choices=['P5_NHIT', 'P5_NN_2D'],
19451             default="P5_NN_2D", required=False,
19462             help="The energy estimator choice. Options are: P5_NHIT,
1947 P5_NN_2D. GP not supported (yet).")
19483     p.add_argument("--use-bins", default=None, nargs="*",
19494             help="Bins to use for the analysis", dest="use_bins")
19505     p.add_argument('--select_bins_by_energy', default=None, nargs="*",
19516             help="Does nothing. May fill in later once better
1952 understood")
19537     p.add_argument('--select_bins_by_fhit', default=None, nargs="*",
19548             help="Also does nothing see above")
19559     p.add_argument( '-ex', '--exclude', default=None, nargs="*",

```

```

19560         help="Exclude Bins", dest="exclude")

19571

19582 # Computing and logging arguements.

19593 p.add_argument('-nt', '--n_threads', type=int, default=1,
19604                         help='Maximum number of threads spawned by script. Default
1961      is 4')

19625 p.add_argument('-log', '--log', type=str, required=True,
19636                         help='Name for log files. Especially needed for threads')

19647

19658 p.add_argument('--decay', action="store_true",
19669                         help='Set spectral DM hypothesis to decay')

19670

19681 args = p.parse_args()

19692 print(args.decay)

19703 if args.exclude is None: # default exclude bins 0 and 1
19714     args.exclude = ['B0C0Ea', 'B0C0Eb', 'B0C0Ec', 'B0C0Ed', 'B0C0Ee']

19725

19736 if args.decay: OUT_PATH += '_dec'
19747 else: OUT_PATH += '_ann'

19758

19769 OUT_PATH = OUT_PATH + '_' + args.catalog
19770 if args.sigma != 0: OUT_PATH += f'_{args.sigma}sig'

19781

19792 main(args)

```

1980 **B.3 Comparison with Glory Duck**

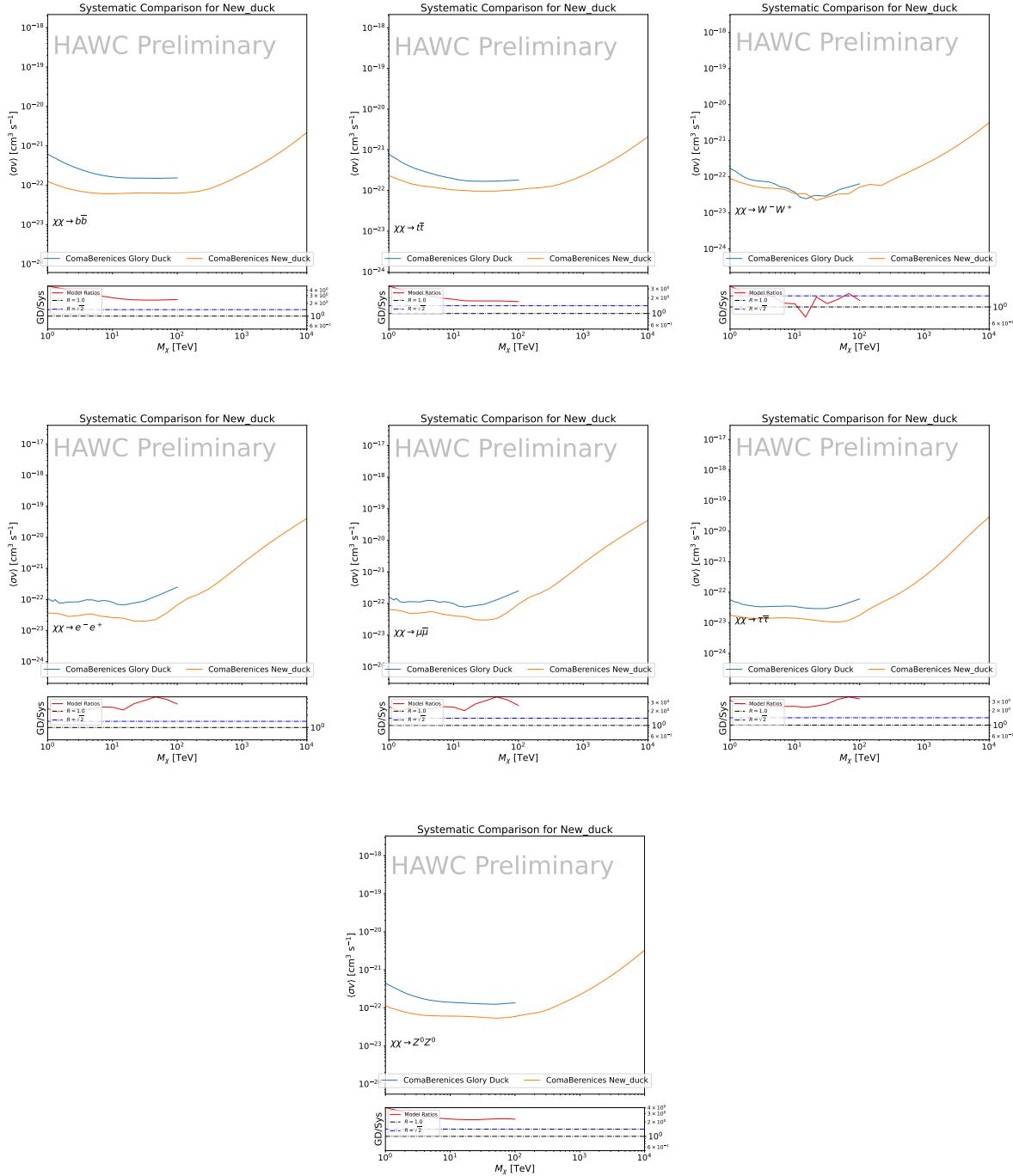


Figure B.2 TODO: fill this out

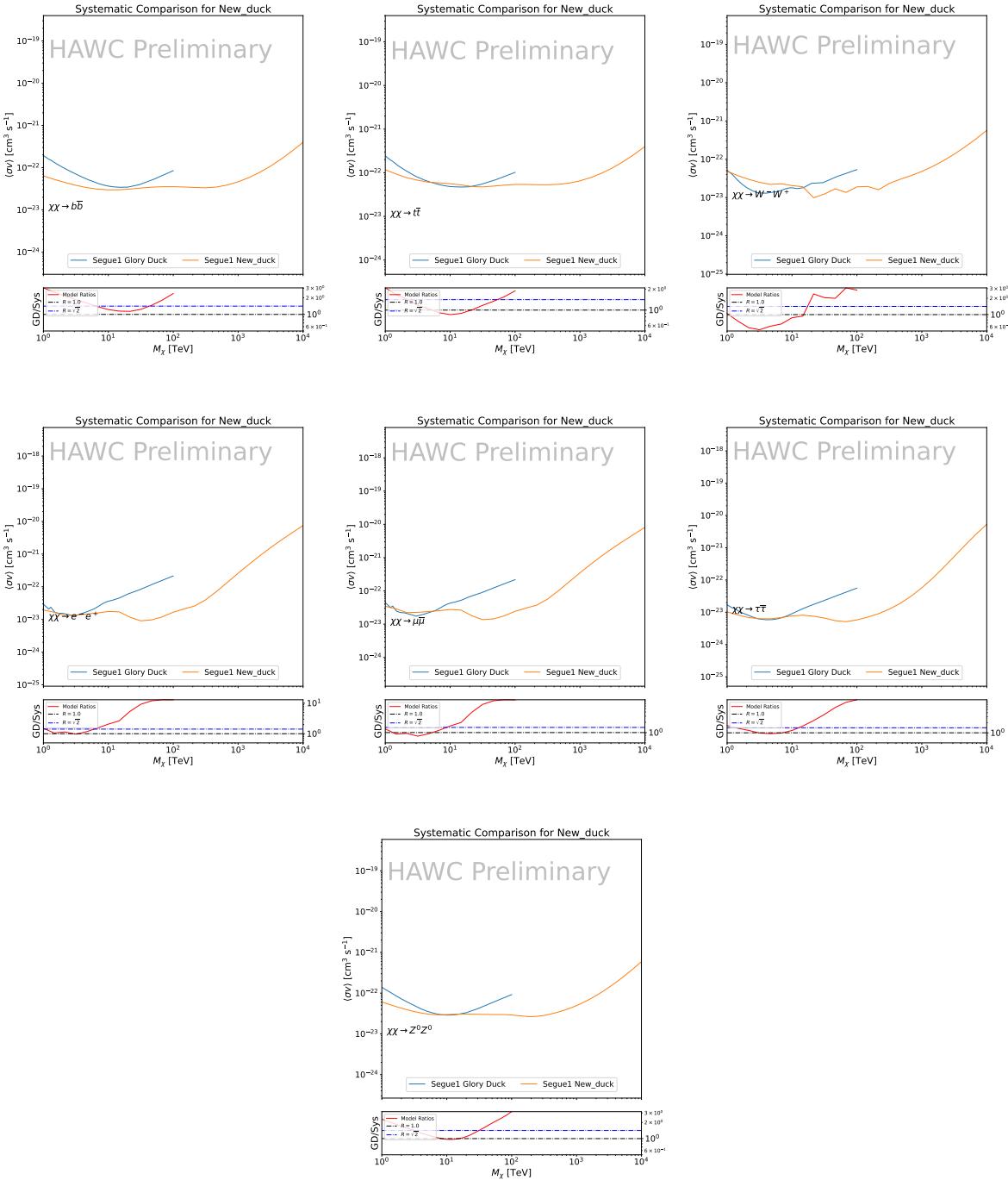


Figure B.3 TODO: fill this out

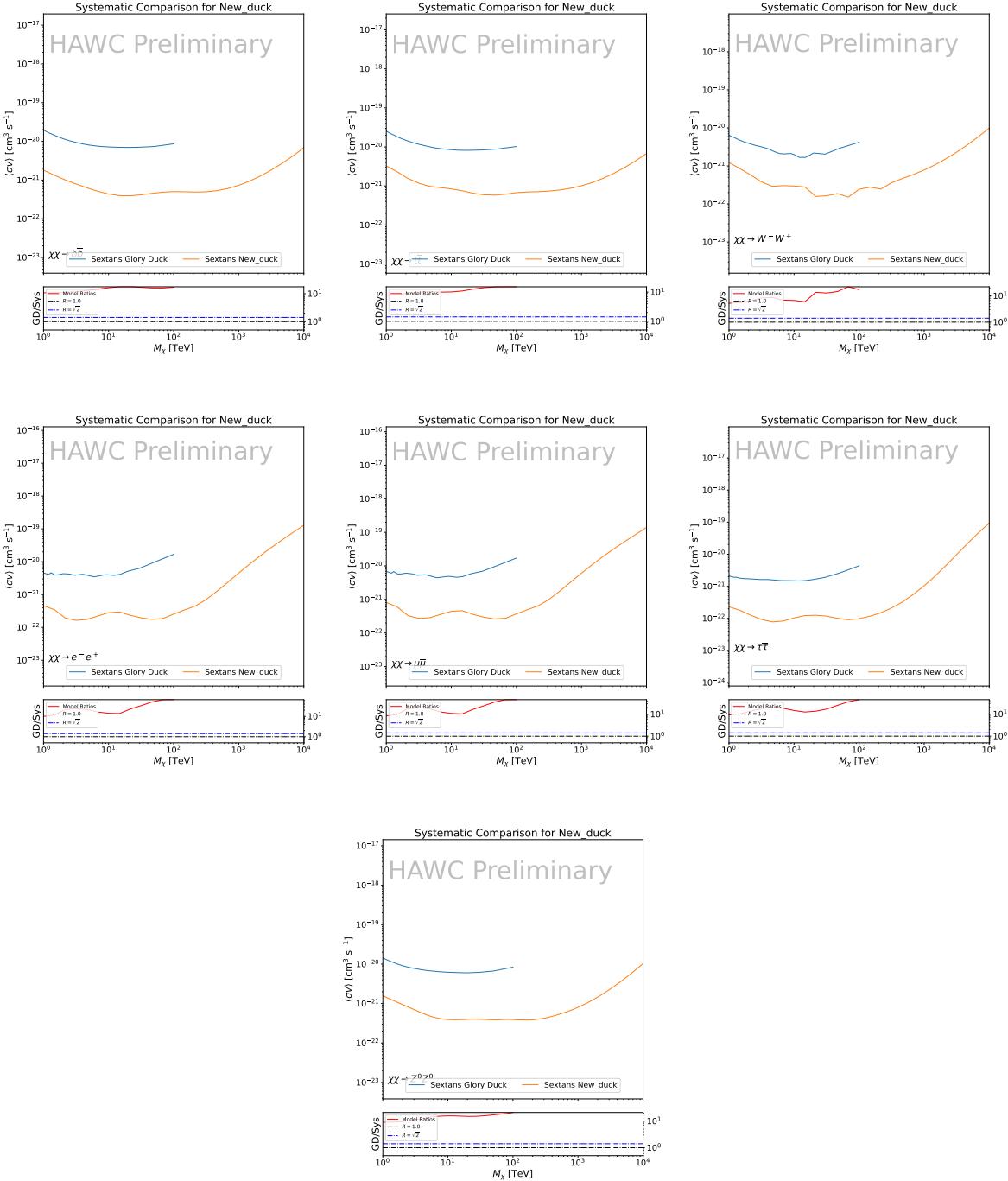


Figure B.4 TODO: fill this out

## APPENDIX C

### 1981 ICECUBE HEAVY DARK MATTER ANALYSIS SUPPLEMENTARY MATERIAL

#### 1982 C.1 Docker Image for Oscillating Neutrino Spectra

```
1983 1 FROM ubuntu:18.04
1984 2
1985 3 # Execute commands to install software packages
1986 4 RUN apt -y update
1987 5
1988 6     # Install utility programs
1989 7 RUN apt -y install vim wget git cmake
1990 8
1991 9 ARG DEBIAN_FRONTEND=noninteractive
1992 0
1993 1     # Install python
1994 2 RUN apt -y install python python-dev python-pip libjpeg-dev zlib1g-dev
1995 3
1996 4     # We need Python2 for installing Charon.
1997 5 RUN apt -y install python-numpy python-sympy python-matplotlib \
1998 6             python-sympy python-h5py python-astropy python-ipython
1999 7
2000 8     # Install dependencies of Charon : SQuIDS, NuSQuIDS
2001 9 RUN apt -y install libgsl-dev libgslcblas0 libgsl23 gsl-bin pkg-config
2002 0     # Install SQuIDS
2003 1 RUN mkdir /home/SQuIDS /home/SQuIDS_install
2004 2 WORKDIR /home/SQuIDS
2005 3 RUN git clone https://github.com/jsalvado/SQuIDS.git
2006 4 WORKDIR /home/SQuIDS/SQuIDS
2007 5 RUN git checkout 7ad9ba7c6ad06d1f0fa8418f937ebf1a403fef90
2008 6     # Before executing "make install" an environmental variable has to be set.
2009 7 ENV PKG_CONFIG_PATH=/home/SQuIDS/SQuIDS/lib
2010 8 RUN ./configure --prefix=../SQuIDS_install \
```

```

20119  && make
20120 RUN make install
20131
20142 # Set up an environmental variable that is required to install nuSQuIDS..
20153 ENV SQuIDS=/home/SQuIDS/SQuIDS
20164 ENV LD_LIBRARY_PATH=$SQuIDS/lib:$LD_LIBRARY_PATH
20175
20186 # Install NuSQuIDS
20197 RUN mkdir /home/nuSQuIDS
20208 WORKDIR /home/nuSQuIDS
20219 RUN git clone https://github.com/qrliu/nuSQuIDS.git
20220 WORKDIR /home/nuSQuIDS/nuSQuIDS
20231 RUN git checkout 072d8ef740e2fc7330f1fabaea94f0f4540c46f9
20242 RUN apt -y install libhdf5-dev hdf5-tools
20253 RUN apt -y install libboost1.65-all-dev
20264 RUN ./configure --with-squids=$SQuIDS --with-python-bindings --prefix=../
2027     nuSQuIDS_install \
20285     && make \
20296     && make install
20307
20318 # Set up an environmental variable for nuSQuIDS.
20329 ENV nuSQuIDS=/home/nuSQuIDS/nuSQuIDS
20330 ENV LD_LIBRARY_PATH=$nuSQuIDS/lib:$LD_LIBRARY_PATH
20341
20352 # Build the python bindings
20363 WORKDIR /home/nuSQuIDS/nuSQuIDS/resources/python/src
20374 RUN make
20385
20396 # Set up an environmental variable for the python bindings.
20407 ENV PYTHONPATH=$nuSQuIDS/resources/python/bindings/:$PYTHONPATH
20418
20429 # Install Charon in the /home/Charon/charon directory.
20430 RUN mkdir /home/Charon

```

```
20441 WORKDIR /home/Charon
20452 RUN git clone https://github.com/icecube/charon.git \
20463     && apt -y install unzip python-scipy
20474 WORKDIR charon
20485 RUN git checkout c531efe4e01dc364a60d1c83f950f04526ccd771
20496 RUN unzip ./charon/xsec/xsec.zip -d charon/xsec/ \
20507
20518 # Download neutrino spectra tables in the /home/Charon/charon/data directory
2052 .
20539 && mkdir ./charon/data
20540 WORKDIR ./charon/data
20551 RUN wget --no-check-certificate https://icecube.wisc.edu/~qliu/charon/
2056     SpectraEW.hdf5 \
20572 && wget --no-check-certificate https://icecube.wisc.edu/~qliu/charon/
2058     Spectra_PYTHIA.hdf5 \
20593 && wget --no-check-certificate https://icecube.wisc.edu/~qliu/charon/
2060     Spectra_noEW.hdf5
20614
20625 WORKDIR ../../
20636 RUN python setup.py install
20647 WORKDIR /home
```

2065 C.2 Spline Fitting Statuses

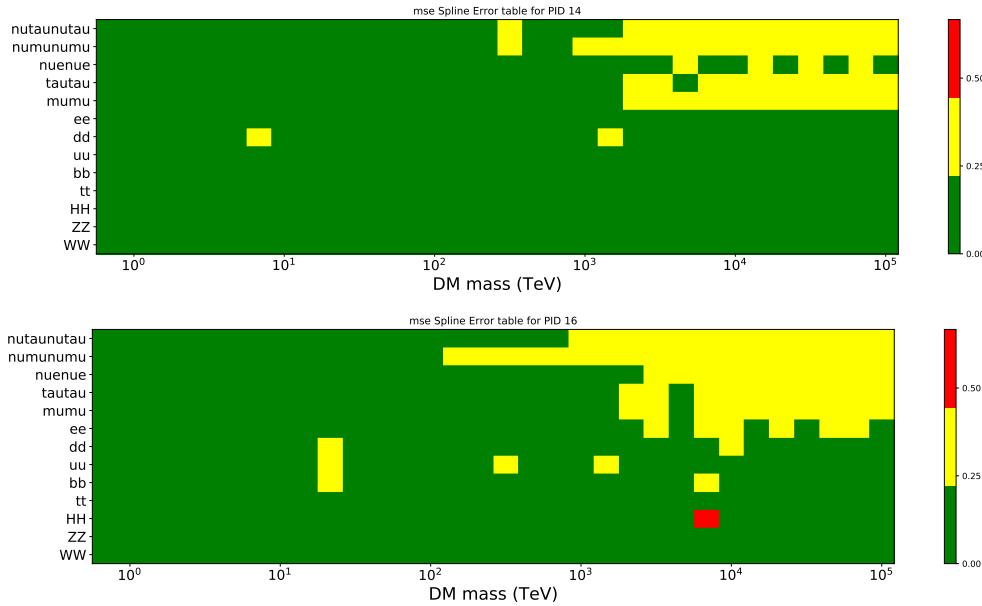


Figure C.1 TODO: fill me daddy

2066 C.3 Segue 1 And Ursa Major II Signal Recovery



Figure C.2 TODO: Fill this out eventually. I think I want all the plots generated first[NEEDS A SOURCE][FACT CHECK THIS]

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