

LEVERAGING MULTI-MESSENGER ASTROPHYSICS FOR DARK MATTER SEARCHES

By

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A DISSERTATION

Submitted to  
Michigan State University  
in partial fulfillment of the requirements  
for the degree of

Physics—Doctor of Philosophy  
Computational Mathematics in Science and Engineering—Dual Major

Today

## **ABSTRACT**

I did Dark Matter with HAWC and IceCube. I also used Graph Neural Networks

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Today

## **ACKNOWLEDGMENTS**

I love my friends. Thanks to everyone that helped me figure this out. Amazing thanks to the people at LANL who supported me. Eames, etc Dinner Parties Jenny and her child Kaydince Kirsten, Pat, Andrea Family. You're so far but so critical to my formation. Unconditional love. Roommate

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## LIST OF ABBREVIATIONS

<b>MSU</b>	Michigan State University
<b>LANL</b>	Los Alamos National Laboratory
<b>DM</b>	Dark Matter
<b>SM</b>	Standard Model
<b>HAWC</b>	High Altitude Water Cherenkov Observatory

## **CHAPTER 1**

### **INTRODUCTION**

Is the text not rendering right? Ah ok it knows im basically drafting the doc still

## CHAPTER 2

### DARK MATTER IN THE COSMOS

#### 2.1 Introduction

The dark matter problem can be summarized in part by following thought experiment.

Let's say you're the teacher for an elementary school classroom. You take them on a field trip to your local science museum and among exhibits is one for mass and weight. The exhibit has a gigantic scale, and you come up with a fun problem for your class.

You ask your class, "What is the total weight of the classroom? Give your best estimation to me in 30 minutes, and then we'll check your guess on the scale. If your guess is within 10% of the right answer, we will stop for ice cream on the way back"

The students are ecstatic to hear this, and they get to work. The solution is some variation of the following strategy. The students should give each other their weight or best guess if they do not know. Then, all they must do is add each student's weight and get a grand total for the class. The measurement on the giant scale should show the true weight of the class. When comparing the measured weight to your estimation, multiply the measurement by  $1.0 \pm 0.1$  in order to get the  $\pm 10\%$  tolerances for your estimation.

Two of your students, Sandra and Mario, return to you with a solution.

They say, "We weren't sure of everyone's weight. We used 65 lbs. for the people we didn't know and added everyone who does know. There are 30 of us, and we got 2,000 lbs.! That's a ton!"

You estimated 1,900 lbs. assuming the average weight of a student in your class was 60 lbs. So, you are pleased with Sandra's and Mario's answer. You instruct your students to all gather on the giant scale and read off the weight together. To all your surprise, the scale reads *10,000 lbs!* 10,000 is significantly more than a 10% error from 2,000. In fact, it is approximately 5 times more massive than either your or your students' estimates. You think to yourself and conclude there must be something wrong with the scale. You ask an employee to check the scale and verify it is calibrated well. They confirm that the scale is in working order. You weigh a couple of students individually to test that the scale is well calibrated. Sandra weighs 59 lbs., and Mario weighs 62

lbs., typical weights for their age. You then weigh each student individually and see that their weights individually do not deviate greatly from 60 lbs. So, where does all the extra weight come from?

This thought experiment serves as an analogy to the Dark Matter problem. The important substitution to make however is to replace the students with stars and the classroom with a galaxy, say the Milky Way. Individually the mass of stars is well measured and defined with the Sun as our nearest test case. However, when we set out to measure the mass of a collection of stars as large as galaxies, our well-motivated estimation is wildly incorrect. There simply is no way to account for this discrepancy except without some unseen, or dark, contribution to mass and matter in galaxies. I set out in my thesis to narrow the possibilities of what this Dark Matter could be.

This chapter is organized like the following... **TODO: Text should look like ... Chatper x has blah blah blah.**

## 2.2 Dark Matter Basics

Presently, the most compelling of cosmology that includes Dark Matter (DM) in order to explain a variety of observations is  $\Lambda$  Cold Dark Matter, or  $\Lambda$ CDM. I present the evidence supporting  $\Lambda$ CDM in 2.3, yet discuss the conclusions of the  $\Lambda$ CDM model here. According to  $\Lambda$ CDM fit to observations on the Cosmic Microwave Background (CMB), DM is 26.8% of the universe's current energy budget. Baryonic matter, stuff like atoms, gas, and stars, contributes to 4.9% of the universe's current energy budget [1, 2, 3].

DM is dark; it does not interact readily with light at any wavelength. DM also does not interact noticeably with the other standard model forces (Strong and Weak) at a rate that is readily observed [3]. DM is cold, which is to say that the average velocity of DM is below relativistic speeds [1]. 'Hot' DM would not likely manifest the dense structures we observe like galaxies, and instead would produce much more diffuse galaxies than what is observed [3, 1]. DM is old; it played a critical role in the formation of the universe and the structures within it [1, 2].

Observations of DM have so far been only gravitational. The parameter space available to what DM could be therefore is extremely broad. The broad DM parameter space is iteratively tested in

DM searches by supposing a hypothesis that has not yet been ruled out and performing observations to test them. When the observations yield a null result, the parameter space is further constrained. I present some approaches for DM searches in Section 2.4.

### 2.3 Evidence for Dark Matter

Dark Matter (DM) has been a looming problem in physics for almost 100 years. Anomalies have been observed in galactic dynamics as early as 1933 when Fritz Zwicky noticed unusually large velocity dispersion in the Coma cluster. Zwicky's measurement was the first recorded to use the Virial theorem to measure the mass fraction of visible and invisible matter in celestial bodies [4]. From Zwicky in [5], "*If this would be confirmed, we would get the surprising result that dark matter is present in much greater amount than luminous matter.*" Zwicky's and others' observation did not instigate a crisis in astrophysics because the measurements did not entirely conflict with their understanding of galaxies [4]. In 1978, Rubin, Ford, and Norbert measured rotation curves for ten spiral galaxies [6]. Rubin et al.'s 1978 publication presented a major challenge to the conventional understanding of galaxies that could no longer be accredited to measurement uncertainties. Evidence has been mounting ever since for this exotic form of matter. The following subsections sample some compelling evidence supporting DM.

#### 2.3.1 First Clues: Stellar Velocities

Zwicky, and later Rubin, measured the stellar velocities of various galaxies to estimate their virial mass. The Virial Theorem upon which these observations are interpreted is written as

$$2T + V = 0. \quad (2.1)$$

Where  $T$  is the kinetic energy and  $V$  is the potential energy in a self-gravitating system. The potential was defined as the classical Newton's law of gravity from stars and gas contained in the observed galaxies

$$V = -\frac{1}{2} \sum_i \sum_{j \neq i} \frac{m_i m_j}{r_{ij}}. \quad (2.2)$$

Zwicky et al. measured just the velocities of stars apparent in optical wavelengths [5]. Rubin et al. added by measuring the velocity of the hydrogen gas via the 21 cm emission line of Hydrogen [6].

The velocities of the stars and gas are used to infer the total mass of galaxies and galaxy clusters via Equation (2.1). An inferred mass is also obtained from the luminosity of the selected sources. The two inferences are compared to each other as a luminosity to mass ratio and typically yields [1]

$$\frac{M}{L} \sim 400 \frac{M_{\odot}}{L_{\odot}} \quad (2.3)$$

$M_{\odot}$  and  $L_{\odot}$  referring to stellar mass and stellar luminosity, respectively. These ratios clearly indicate a discrepancy in apparent light and mass from stars and gas and their velocities.

Rubin et al. [6] demonstrated that the discrepancy was unlikely to be an under-estimation of the mass of the stars and gas. The inferred "dark" mass was up to 5 times more than the luminous mass. This dark mass also needed to extend well beyond the extent of the luminous matter.

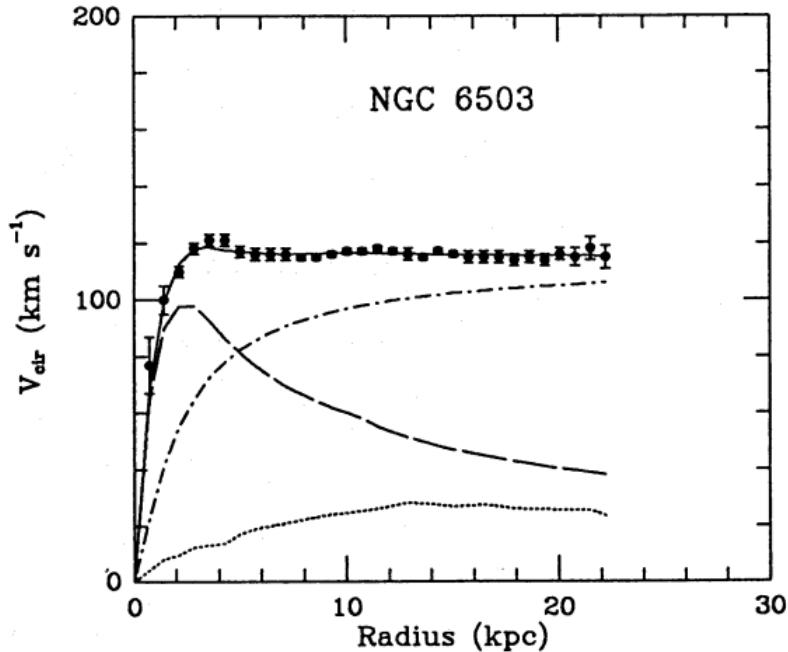


Figure 2.1 Rotation curve fit to NGC 6503 from [7]. Dashed line is the contribution from visible matter. Dotted curves are from gas. Dash-dot curves are from dark matter (DM). Solid line is the composite contribution from all matter and DM sources. Data are indicated with bold dots with error bars. Data agree strongly with matter + DM composite prediction

Figure 2.1 features one of many observations made on the stellar velocities within galaxies. The measured rotation curves mostly feature a flattening of velocities at higher radius which is not expected if the gravity was only coming from gas and luminous matter. The extension of

the flat velocity region also indicates that the DM is distributed far from the center of the galaxy. Modern velocity measurements include significantly larger objects, galactic clusters, and smaller objects, dwarf galaxies. Yet, measurements along this regime are leveraging the Virial theorem with Newtonian potential energies. We know Newtonian gravity is not a comprehensive description of gravity. New observational techniques have been developed since 1978, and those are discussed in the following sections.

### 2.3.2 Evidence for Dark Matter: Gravitational Lensing

Modern evidence for dark matter comes from new avenues beyond stellar velocities. Gravitational lensing from DM is a new channel from general relativity. General relativity predicts aberrations in light caused by massive objects. In recent decades we have been able to measure the lensing effects from compact objects and DM halos. Figure 2.2 shows how different massive objects change the final image of a faraway galaxy resulting from gravitational lensing. Gravitational lensing developed our understanding of dark matter in two important ways.

Gravitational lensing provides additional compelling evidence for DM. The observation of two merging galactic clusters in 2006, shown in Figure 2.3, provided a compelling argument for DM outside the Standard Model. These clusters merged recently in astrophysical time scales. Galaxies and star cluster have mostly passed by each other as the likelihood of their collision is low. Therefore, most of these massive objects will track the highest mass, dark and/or baryonic, density. Yet, the intergalactic gas is responsible for most of the baryonic mass in the systems [4]. These gas bodies will not phase through and will heat up as they collide together. The hot gas is observed as x-ray emission from the cluster. Two observations of the clusters were made independently of each other.

The first was the lensing of light around the galaxies due to their gravitational influences. When celestial bodies are large enough, the gravity they exert bends space and time itself. The warped space-time lenses light and will deflect in an analogous way to how glass lenses will bend light, see Figure 2.2. With a sufficient understanding of light sources behind a massive object, we can reconstruct the contours of the gravitational lenses. The gradient of the contours shown in

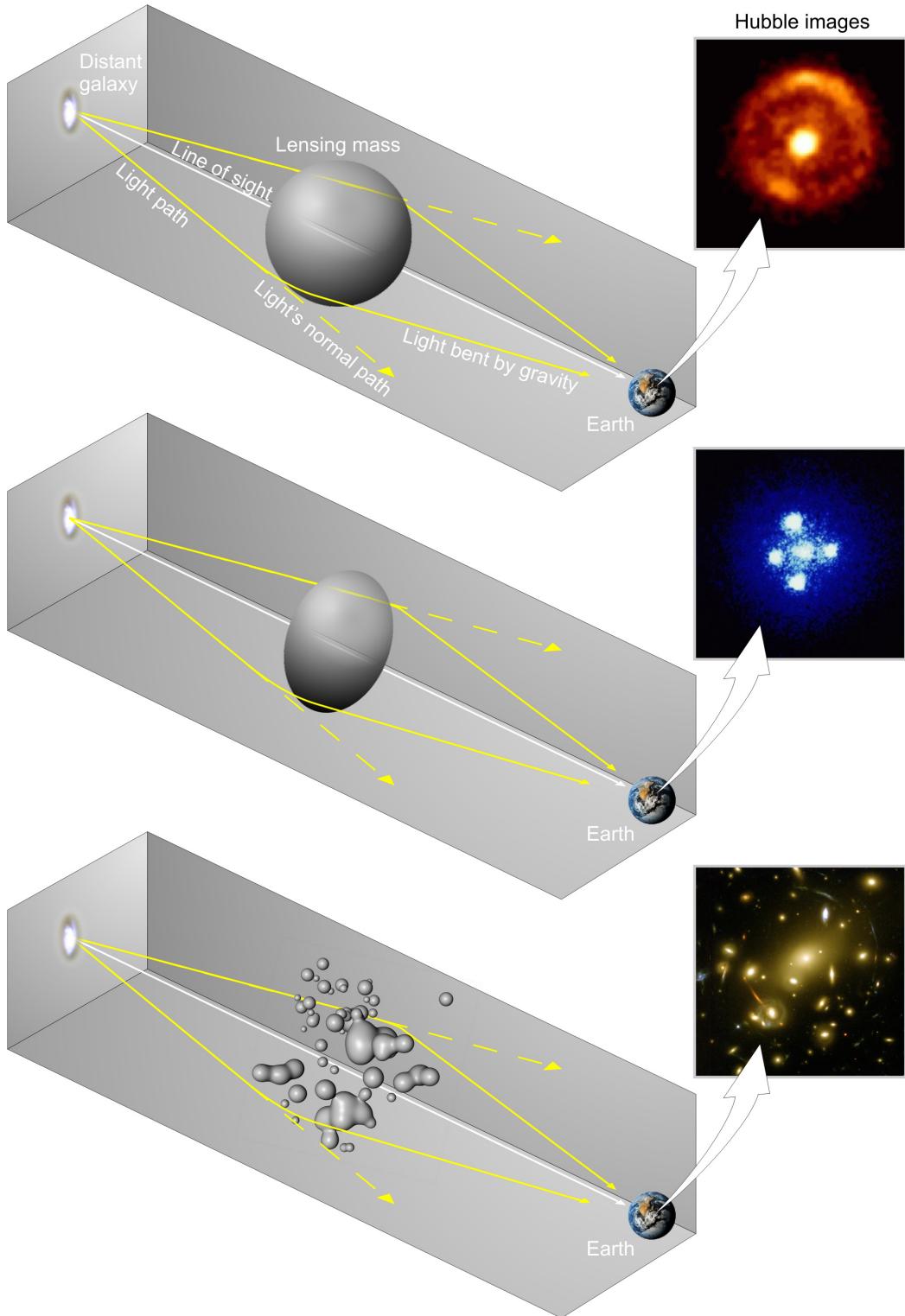


Figure 2.2 Light from distant galaxy is bent in unique ways depending on the distribution of mass between the galaxy and Earth. Yellow dashed lines indicate where the light would have gone if the matter were not present [8].

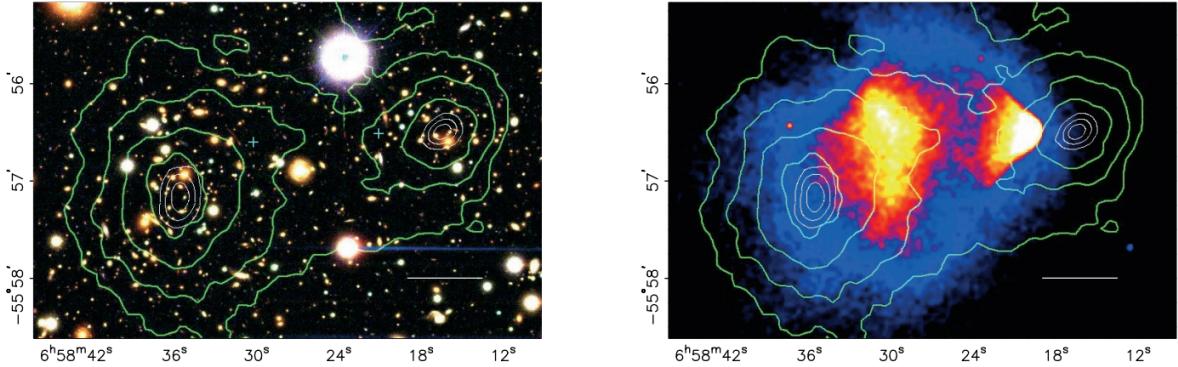


Figure 2.3 (left) Optical image of galactic cluster 1E0657-558. (right) X-ray image of the cluster with redder meaning hotter and higher baryon density. (both) Green contours are reconstruction of gravity contours from weak lensing. White rings are the best fit mass maxima at 68.3%, 95.5%, and 99.7% confidence. The maxima of the clusters are clearly separated from x-ray maxima. [9]

Figure 2.3 then indicates how dense the matter is and where it is.

The x-ray emission can then be observed from the clusters. Since these galaxies are mostly gas and are merging, then the gas should be getting hotter. If they are merging, the x-ray emissions should be the strongest where the gas is mostly moving through each other. Hence, X-ray emission maps out where the gas is in the merging galaxy cluster.

The lensing and x-ray observations were done on the Bullet cluster featured on Figure 2.3. The x-ray emissions do not align with the gravitational contours from lensing. The incongruence in mass density and baryon density suggests that there is a lot of matter somewhere that does not interact with light. Moreover, this dark matter can not be baryonic [9]. The Bullet Cluster measurement did not really tell us what DM is exactly, but it did give the clue that DM also does not interact with itself very strongly. If DM did interact strongly with itself, then it would have been more aligned with the x-ray emission [9]. There have been follow-up studies of galaxy clusters with similar results. The Bullet Cluster and others like it provide a strong case against something possibly amiss in our gravitational theories.

### 2.3.3 Evidence for Dark Matter: Cosmic Microwave Background

The Cosmic Microwave Background (CMB) is the primordial light from the early universe when Hydrogen atoms formed from the free electron and proton soup in the early universe. The

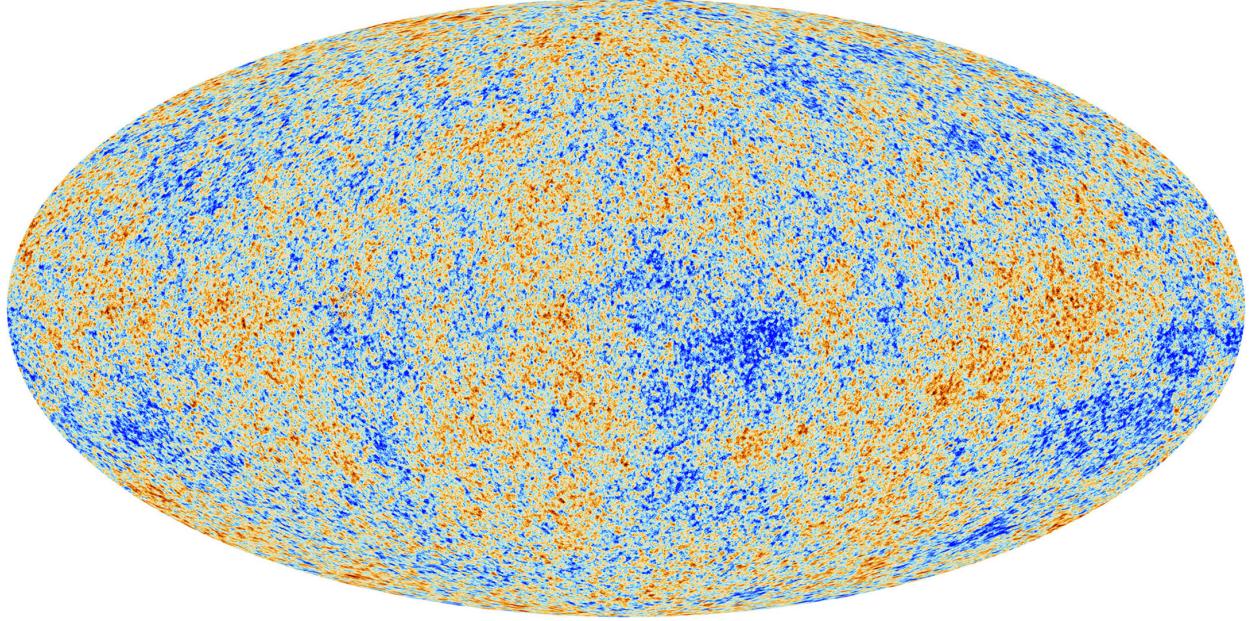


Figure 2.4 Plank CMB sky. Sky map features small variations in temperature in primordial light. These anisotropies can be used to make inferences about the universe's energy budget. [10]

CMB is the earliest light we can observe; released when the universe was about 380,000 years old. Then we look at how the simulated universes look like compared to what we see. Figure 2.4 is the most recent CMB image from the Plank satellite after subtracting the average value and masking the galactic plane [10]. Redder regions indicate a slightly hotter region in the CMB and blue indicates colder. The intensity variations are on the order of 1 in 1000 with respect to the average value.

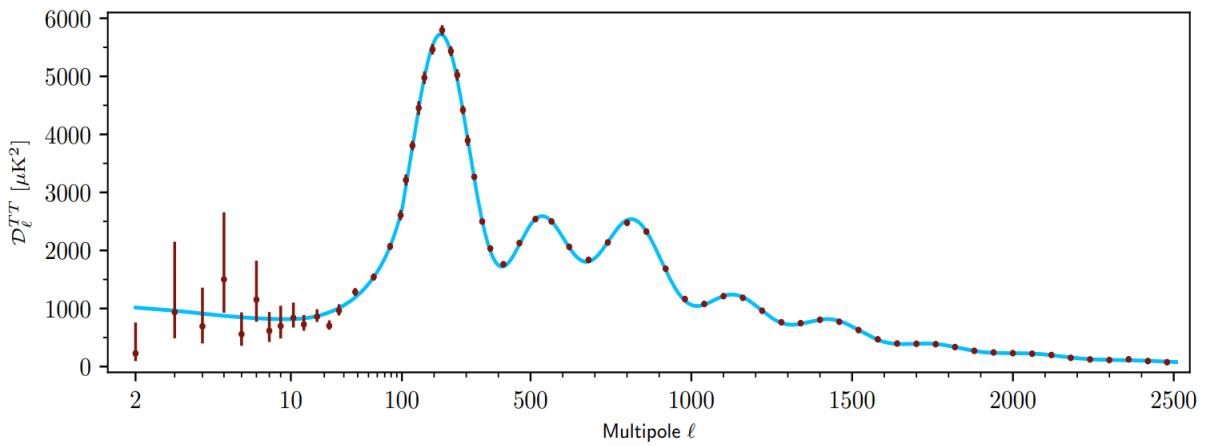


Figure 2.5 Observed Cosmic Microwave Background power spectrum as a function of multipole moment from Plank [10]. Blue line is best fit model from  $\Lambda$ CDM. Red points and lines are data and error respectively.

The Cosmic Microwave Background shows that the universe had DM in it from an incredibly early stage. To measure the DM, Dark Energy, and matter fractions of the universe from the CMB, the image is deconstructed into a power spectrum, which shows the amplitude of the fluctuations as a function of spherical multipole moments.  $\Lambda\text{CDM}$  provides the best fit to the power spectra of the CMB as shown in Figure 2.5. The CMB power spectrum is very sensitive to the fraction of each energy contribution in the early universe. Low  $l$  modes are dominated by variations in gravitational potential. Intermediate  $l$  emerge from oscillations in photon-baryon fluid from competing baryon pressures and gravity. High  $l$  is a damped region from the diffusion of photons during electron-proton recombination. [1]

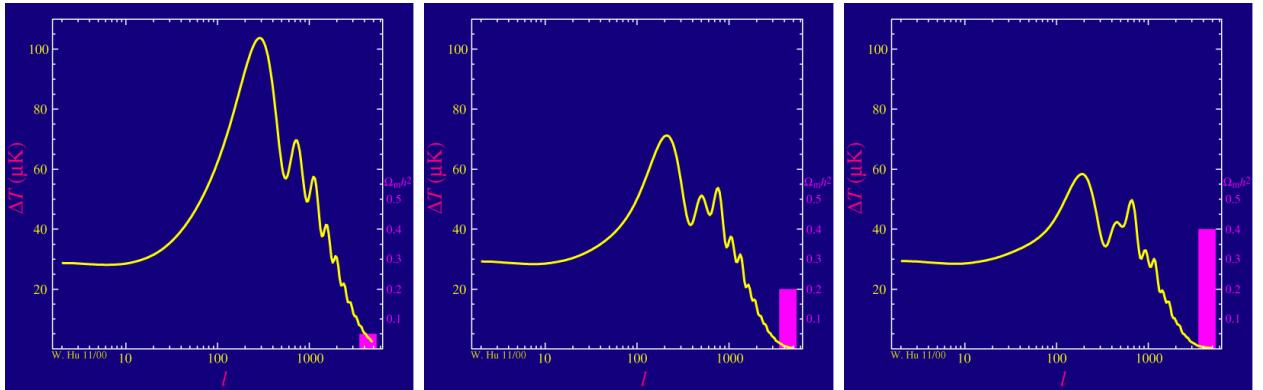


Figure 2.6 Predicted power spectra of CMB for different  $\Omega_m h^2$  values for fixed baryon density from [11]. (left) Low  $\Omega_m h^2$  increases the prominence of first and second peaks. (middle)  $\Omega_m h^2$  is most similar to the observed power spectrum. The second and third peaks are similar in height. (right)  $\Omega_m h^2$  is large which suppresses the first peak and raises the prominence of the third peak.

The harmonics would look very different for a universe with less DM. Figure 2.6 demonstrates the effect  $\Omega_m h^2$  has on the expected power spectrum for fixed baryon matter density. [11] Sweeping  $\Omega_m h^2$  in this way clearly shows the effect dark matter has on the CMB power spectrum. The observations fit well with the  $\Lambda\text{CDM}$  model and the derived fractions are as follows. The matter fraction:  $\Omega_m = 0.3153$ ; and the baryon fraction:  $\Omega_b = 0.04936$  [10]. Plank's observations also provide a measure of the Hubble constant,  $H_0$ .  $H_0$  especially has seen a growing tension in the past decade that continues to deepened with observations from instruments like the James Webb Telescope [12, 13]. As Hubble tensions deepen, we may find that perhaps  $\Lambda\text{CDM}$ , despite it's

successes, is missing some critical physics.

Overall the Newtonian motion of stars in galaxies, weak lensing from galactic clusters, and power spectra from primordial light form a compelling body of research in favor of dark matter. However, these observations really only indicate that DM is there. It takes another leap of theory and experimentation to make observations of DM that are non-gravitational in nature. In Section 2.3, the evidence for DM implies strongly that the DM is matter and not a lost parameter in the gravitational fields between massive object. Finally, if we take one axiom: that this matter has some quantum behavior, such as being described by some Bohr wavelength and abiding by some fermion or boson statistics; then we arrive at particle dark matter. One particle DM hypothesis is the Weakly Interacting Massive Particle (WIMP). This DM candidate theory is discussed further in the next section and is the focus to this thesis.

## 2.4 Searching for Dark Matter: Particle DM

Section 2.4 shows the Standard Model of particle physics and is currently the most accurate model for the dynamics of fundamental particles like electrons and photons. The current status of the SM does not have a viable DM candidate. When looking at the standard model, we can immediately exclude any charged particle because charged particles interact strongly with light. Specifically this will rule out the following charged, fundamental particles:  $e, \mu, \tau, W, u, d, s, c, t, b$  and their corresponding antiparticles. Recalling from Section 2.2 that DM must be long-lived and stable over the age of the universe which excludes all SM particles with decay half-lives at or shorter than the age of the universe. The lifetime constraint additionally eliminates the  $Z$  and  $H$  bosons. Finally, the candidate DM needs to be somewhat massive. Recall from Section 2.2 that DM is cold or not relativistic through the universe. This eliminates the remaining SM particles:  $\nu_{e,\mu,\tau}, g, \gamma$  as DM candidates. Because there are no DM candidates within the SM, the DM problem strongly hints to physics beyond the SM (BSM).

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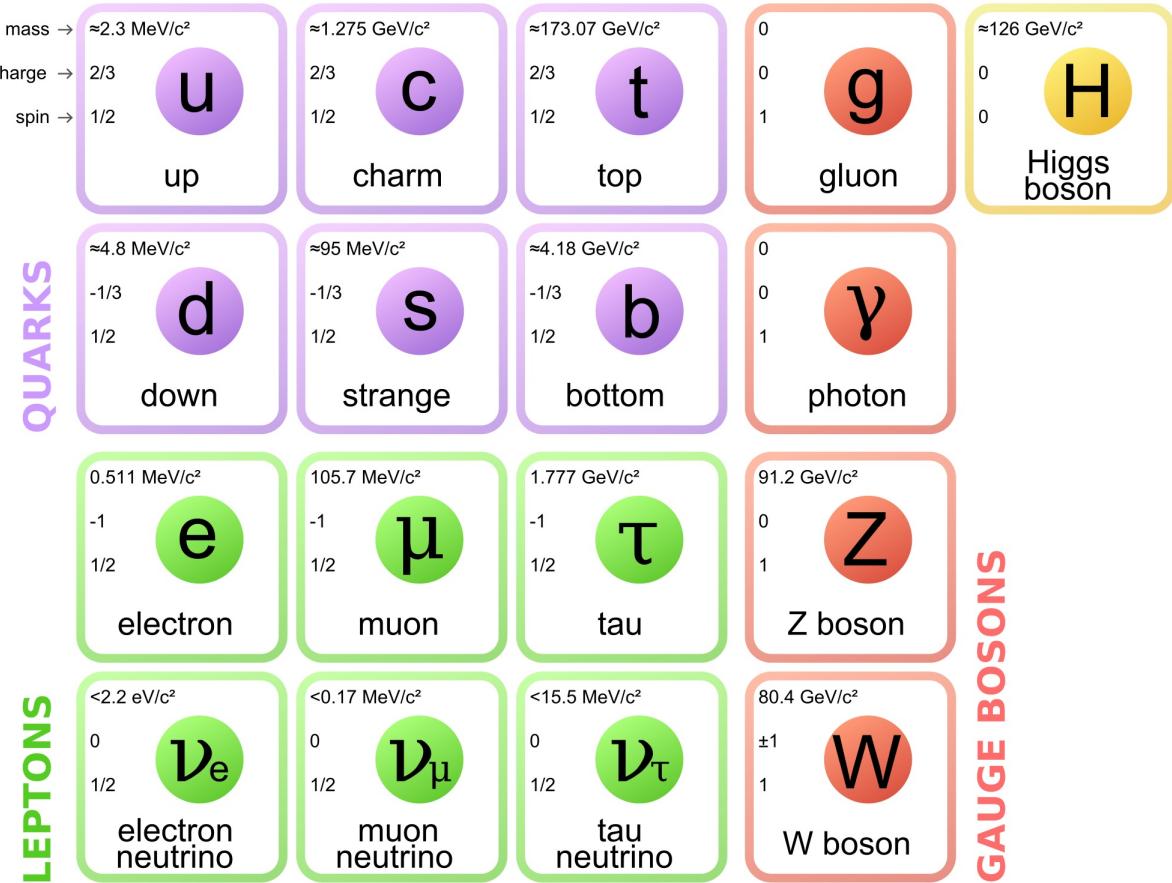


Figure 2.7 The Standard Model (SM) of particle physics. Figure taken from <http://www.quantumdiaries.org/2014/03/14/the-standard-model-a-beautiful-but-flawed-theory/>

### 2.4.1 Shake it, Break it, Make it

When considering DM that couples in some way with the SM, the interactions are roughly demonstrated by interaction demonstrated in Figure 2.8. The figure is a simplified Feynman diagram where the arrow of time represents the interaction modes of: **Shake it, Break it, Make it.**

**Shake it** refers to the direct detection of dark matter. Direct detection interactions start with a free DM particle and some SM particle. The DM and SM interact under some elastic or inelastic collision and recoil away from each other. The DM remains in the dark sector and imparts some momentum onto the SM particle. The hope is that the momentum imparted onto the SM particle is sufficiently high enough to pick up with highly sensitive instruments. Because we cannot create the DM in the lab, a direct detection experiment must wait until DM is incident on the detector. Most direct detection experiments are therefore placed in low-background environments with inert

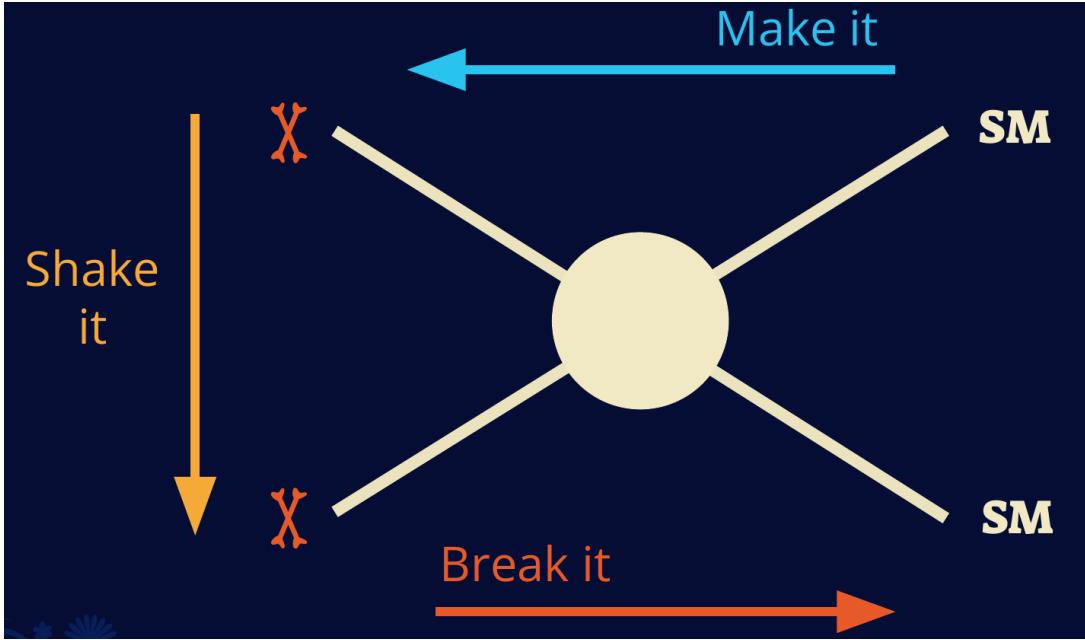


Figure 2.8 Simplified Feynman diagram demonstrating different ways DM can interact with SM particles. The 'X's refer to the DM particles whereas the SM refer to fundamental particles in the SM. The large circle in the center indicates the vertex of interaction and is purposely left vague. The colored arrows refer to different directions of time as well as their respective labels. The arrows indicate the initial and final state of the DM -SM interaction in time.

detection media like the noble gas Xenon. [14]

**Make it** refers to the production of DM from SM initial states. The experiment starts with particles in the SM. These SM particles are accelerated to incredibly high energies and then collided with each other. In the confluence of energy, DM hopefully emerges as a byproduct of the SM annihilation. Often it is the collider experiments that are able to generate energies high enough to probe DM production. These experiments include the world-wide collaborations ATLAS and CMS at CERN where protons are collided together at extreme energies. The DM searches however are complex. DM likely does not interact with the detectors and lives long enough to escape the detection apparatus of CERN's colliders. This means any DM production experiment searches for an excess of events with missing momentum or energy in the events. An example event with missing transverse momentum is shown in Figure 2.9. The missing momentum with no particle tracks implies a neutral particle carried the energy out of the detector. However, there are other neutral particles in the SM, like neutrons or neutrinos, so any analysis have to account for SM

signatures of missing momentum. [15]

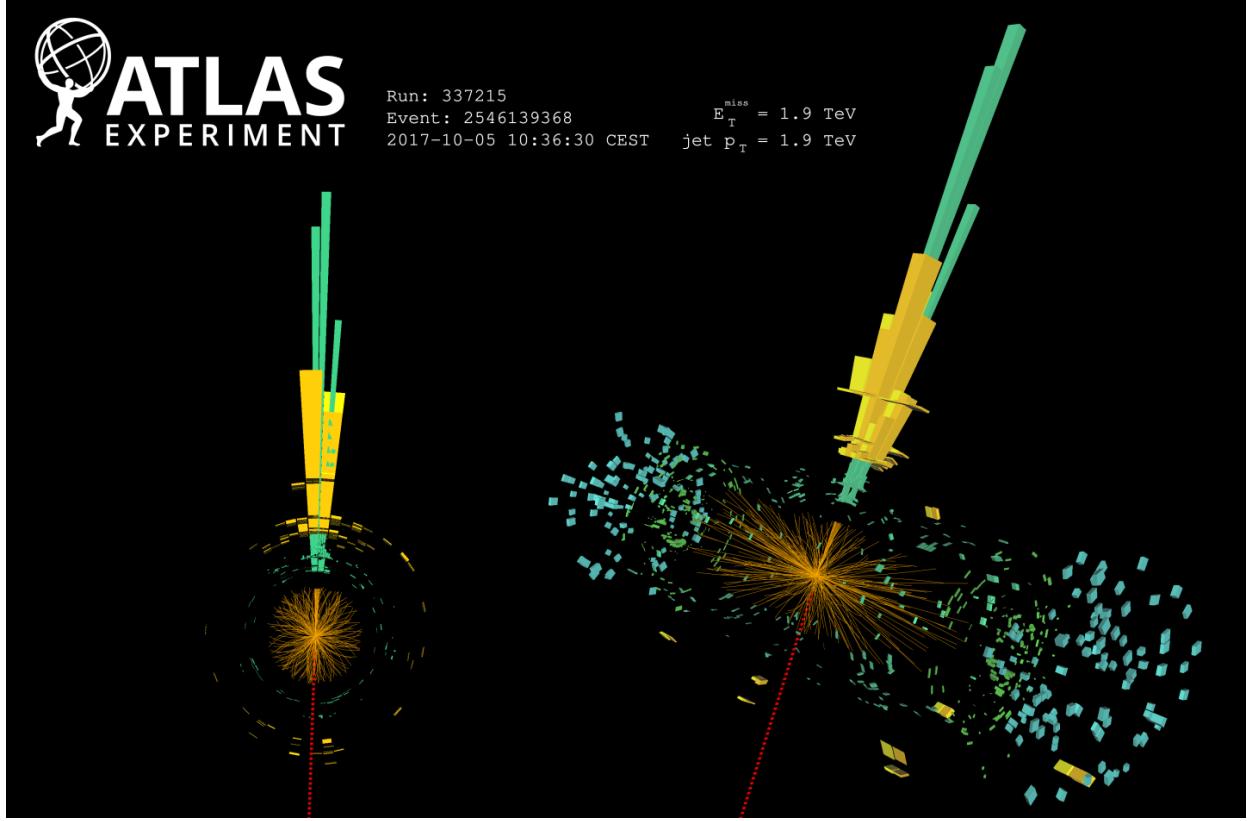


Figure 2.9 A single jet event in ATLAS detector from 2017 [16]. Jet momentum observed to be 1.9 TeV. Missing transverse momentum observed to be 1.9 TeV as the initial momentum of the event was 0. Implied MET is shown as a red dashed line in event display.

#### 2.4.2 Break it: Standard Model Signatures of Dark Matter through Indirect Searches

**Break it** refers to the creation of SM particles from the dark sector, and it is the primary focus of this thesis. The interaction begins with DM or in the dark sector. The hypothesis is that this DM will either annihilate with itself or decay and produce an SM byproduct. This method is often referred to the Indirect Detection of DM because we have no lab to directly control or manipulate the DM. Therefore, most indirect DM searches are performed using observations of known DM densities among the astrophysical sources. The strength is that we have the whole of the universe and its 13.6 billion year lifespan to use as the detector or particle accelerator. Additionally, locations of dark matter are also well understood since it was astrophysical observations that presented the problem of DM in the first place.

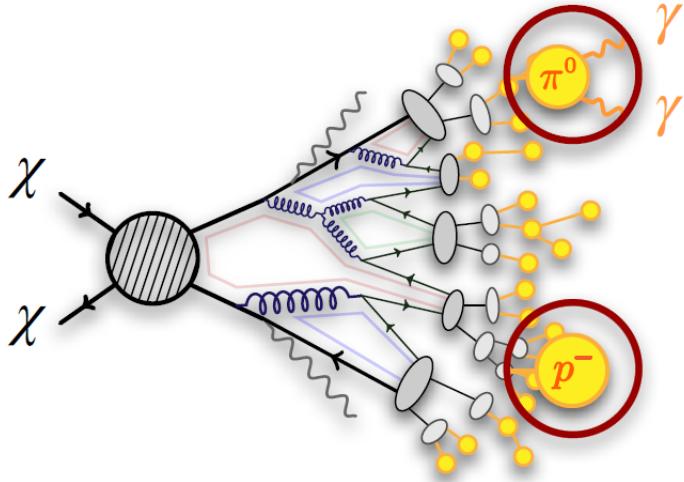


Figure 2.10 More detailed pseudo-Feynman diagram of particle cascade from dark matter annihilation into 2 quarks. The quarks hadronize and down to stable particles like  $\gamma$  or the anti-proton ( $p^-$ ). Diagram pulled from ICRC 2021 presentation on DM annihilation search [17].

However, anything can happen in the universe. There are many difficult to deconvolve backgrounds when searching for DM. One prominent example is the galactic center. We know the galactic center has a large DM content because of stellar kinematics in our Milky Way and DM halo simulations. Yet, any signal from the galactic center is challenging to parse apart from the extreme environment of our supermassive black hole, unresolved sources, and diffuse emission from the interstellar medium [18]. Despite the challenges, any DM model that yields evidence in the other two observation methods, **Shake it or Make it** must be corroborated with indirect observations of the known DM sources. Without corroborating evidence, DM observation in the lab is hard-pressed to demonstrate that it is the model contributing to the DM seen at the universal scale.

In the case of WIMP DM, signals are typically described in terms of primary SM particles produced from a DM decay or annihilation. The SM initial state particles are then simulated to stable final states such as the  $\gamma$ ,  $\nu$ ,  $p$ , or  $e$  which can traverse galactic lengths to reach Earth.

Figure 2.10 shows the quagmire of SM particles that emerges from SM initial states that are not stable [17]. There are many particles with varying energies that can be produced in such an interaction. For any arbitrary DM source and stable SM particle, the SM flux from DM annihilating

to some neutral particle in the SM,  $\phi$ , from a region in the sky is described by the following

$$\frac{d\Phi_\phi}{dE_\phi} = \frac{\langle\sigma v\rangle}{8\pi m_\chi^2} \frac{dN_\phi}{dE_\phi} \times \int_{\text{source}} d\Omega \int_{l.o.s} \rho_\chi^2 dl(r, \theta') \quad (2.4)$$

In Equation (2.4),  $\langle\sigma v\rangle$  is the velocity-weighted annihilation cross-section of DM to the SM.  $m_\chi$  refers to the mass of DM, noted with Greek letter  $\chi$ .  $\frac{dN_\phi}{dE_\phi}$  is the N particle flux weighted by the particle energy. An example is provided in Figure 2.11 for the  $\gamma$  final state. The integrated terms are performed over the solid angle,  $d\Omega$ , and line of sight, l.o.s.  $\rho$  is the density of DM for a location  $(r, \theta')$  in the sky. The terms left of the 'x' are often referred to as the particle physics component. The terms on the right are referred to as the astrophysical component. For decaying DM, the equation changes to...

$$\frac{d\Phi_\phi}{dE_\phi} = \frac{1}{4\pi\tau m_\chi} \frac{dN_\phi}{dE_\phi} \times \int_{\text{source}} d\Omega \int_{l.o.s} \rho_\chi dl(r, \theta') \quad (2.5)$$

In Equation (2.5),  $\tau$  is the decay lifetime of the DM. Just as in Equation (2.4), the left and right terms are the particle physics and the astrophysical components respectively. The integrated astrophysical component of Equation (2.4) is often called the J-Factor. Whereas the integrated astrophysical component of Equation (2.5) is often called the D-Factor.

Exact DM  $\text{DM} \rightarrow \text{SM}$  branching ratios are not known, so it is usually assumed to go 100% into an SM particle/anti-particle. Additionally, when a DM annihilation or decay produces one of the neutral, long-lived SM particles ( $\nu$  or  $\gamma$ ), the particle can be traced back to a DM source. For DM above GeV energies, there are very few SM processes that can produce particles with such a high energy. Seeing such a signal would almost certainly be an indication of the presence of dark matter. The universe fortunately provides us with the largest volume and lifetime ever for a particle physics experiment.

## 2.5 Sources for Indirect Dark Matter Searches

The first detection of DM relied on optical observations. Since then, we've developed new techniques to find DM dense regions. As described in Section 2.3.1, many DM dense regions were through observing galactic rotation curves. Our Milky Way galaxy is among DM dense regions discovered, and it is the largest nearby DM dense region to look at. Additionally, the DM halo

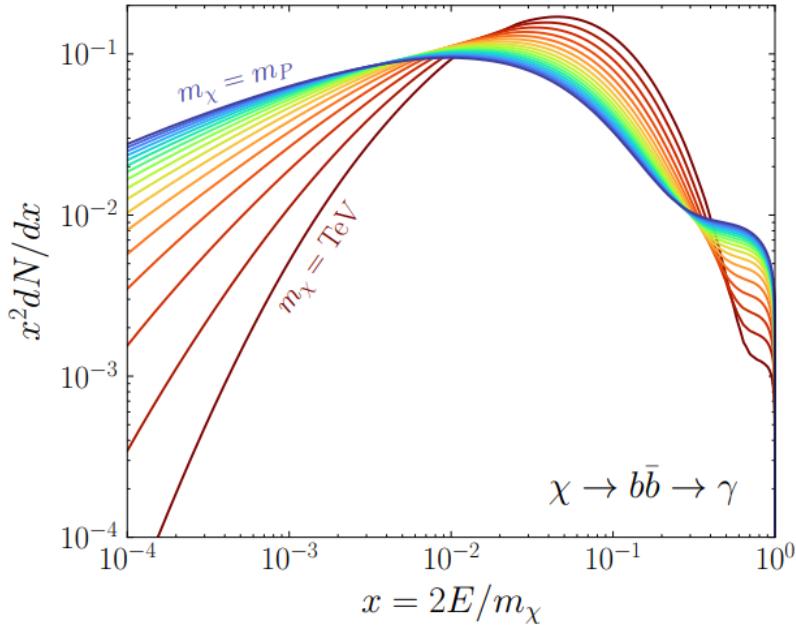


Figure 2.11 Dark Matter (DM) decay spectrum for  $b\bar{b}$  initial state and  $\gamma$  final state. Redder spectra are for larger DM masses. Bluer spectra are light DM masses.  $x$  is a unitless factor defined as the ratio of the mass of DM,  $m_\chi$ , and the final state particle energy  $E_\gamma$ . Figure from [19].

surrounding the Milky Way is somewhat clumpy [18]. There are regions in the DM halo of the Milky Way that have more DM than others that have captured gas over time. In some cases these sub-halos were dense enough to host stars. These apparent sub galaxies are known was dwarf spheroidal galaxies and are the main sources studied in this thesis. Each source type comes with different trade-offs. Galactic Center studies will be very sensitive to the assumed distribution of DM. The central DM density can vary substantially as demonstrated in Figure 2.12. At distances close to the center of the galaxy, or small  $r$ , the differences in DM densities can be 3-4 orders of magnitude. Searches toward the galactic center will therefore be quite sensitive to the assumed DM distribution.

Searches toward Dwarf Spheroidal Galaxies (dSph's) suffer from uncertainties in the DM density less than the galactic center studies. This is mostly from their diminutive size being smaller than the angular resolution of most  $\gamma$ -ray observatories [18]. The DM content dSph's are typically determined with the Virial theorem, Equation (2.1), and are usually majority DM [18] in mass.

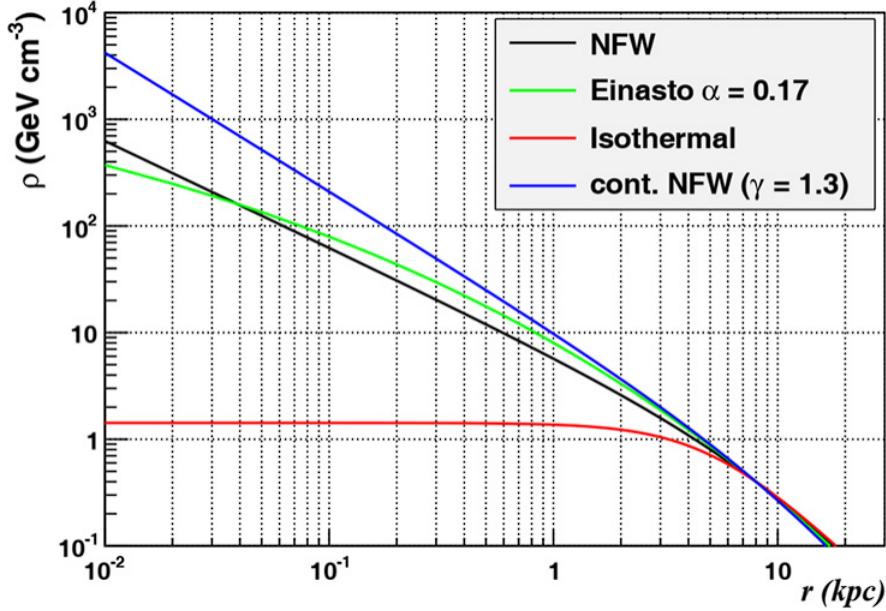


Figure 2.12 Different dark matter density profiles compared. Some models produce very large densities at small  $r$  [20].

dSph's tend to be ideal sources to look at for DM searches. Their environments are fairly quiet with little astrophysical backgrounds. Unlike the galactic center, the most active components of dSph's are the stars within them versus a violent accretion disc around a black hole. All this together means that dSph's are among the best sources to look at for indirect DM searches. dSph's are the targets of focus for this thesis.

## 2.6 Multi-Messenger Dark Matter

Astrophysics entered a new phase in the past few decades that leverages our increasing sensitivity to SM channels and general relativity (GR). Up until the 21st century, astrophysical observations were done with photons ( $\gamma$ ) only. Astrophysics with this 'messenger' is fairly mature now. Novel observations of the universe have since only adjusted the sensitivity of the wavelength of light that's observed. Gems like the CMB [10], and more have ultimately been observations of different wavelengths of light. Multi-messenger astrophysics proposes using other SM particles such the  $p^{+/-}$ , or  $\nu$  or gravitation waves predicted by general relativity.

The experiment LIGO had a revolutionary discovery in 2016 with the first detection of a binary

black hole merger [21]. This opened the collective imagination to observing the universe through gravitational waves. There's also been a surge of interest in the neutrino ( $\nu$ ) sector. IceCube demonstrated that we are sensitive to neutrinos in regions that correlate with significant photon emission like the galactic plane [22]. Neutrinos, like gravitational waves and light, travel mostly unimpeded from their source to our observatories. This makes pointing to the originating source of these messengers much easier than it is for cosmic rays that are almost always deflected from their source by magnetic fields.

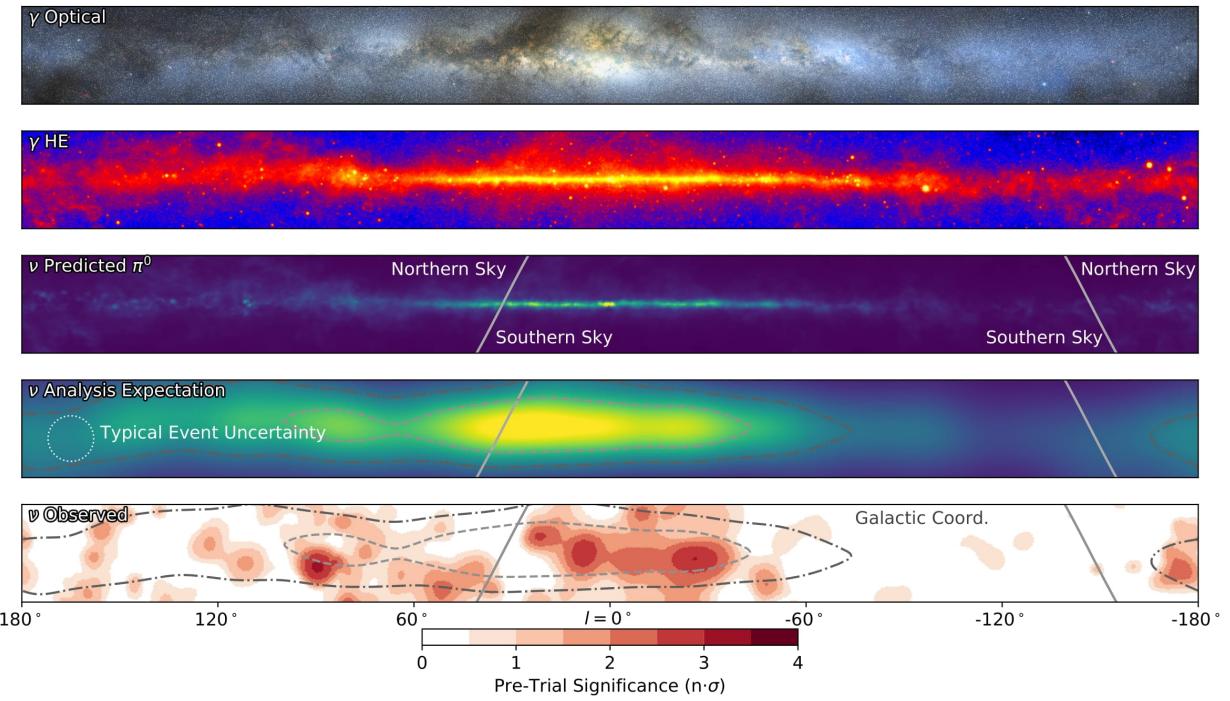


Figure 2.13 The Milky Way Galaxy in photons ( $\gamma$ ) and neutrinos ( $\nu$ ) [22]. Galactic center is at  $l=0^\circ$  and is the brightest region in all panels. (top) An Optical color image of the Milky Way galaxy seen from Earth. Clouds of gas and dust obscure some light from stars. (2nd down) Integrated flux of  $\gamma$ -rays observed by the Fermi-LAT telescope [23]. (middle) Expected neutrino emission that corresponds with Fermi-LAT observations. (2nd up) Expected neutrino emission profile after considering detector systematics of IceCube. (bottom) Observed neutrino emission from region of the galactic plane. Substantial neutrino emission is detected.

The IceCube collaboration recently published a groundbreaking result of the Milky Way in neutrinos. The recent result from IceCube, shown in Figure 2.13, proves that we can make observations under different messenger regimes. The top two panels are the appearance of the

galactic plane to different wavelengths of light. Some sources are more apparent in some panels, while others are not. This new channel is potential very powerful because neutrinos are readily able to penetrate through gas and dust in the Milky Way. This new image also refines our understanding of how high energy particles are produced since the fit to IceCube data corresponds with the decay of  $\pi^0$  over exclusively leptonic models [22].

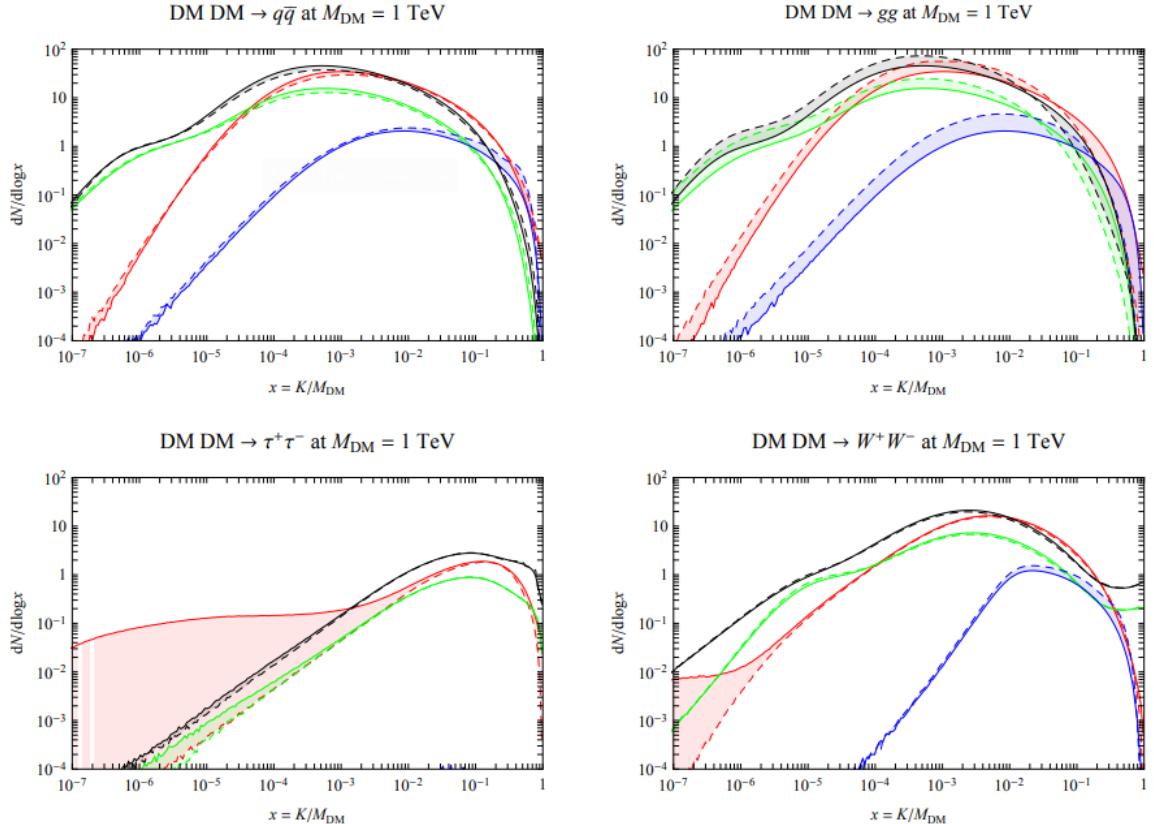


Figure 2.14 Dark Matter annihilation spectra for different final state particle and standard model annihilation channels [24]. Photons (red),  $e^\pm$  (green),  $\bar{p}$  (blue),  $\nu$  (black).

Exposing our observations to more cosmic messengers greatly increases our sensitivity to rare processes. In the case of DM, Figure 2.14, there are many SM particles produced in a dark matter annihilation. Among the final state fluxes are gammas and neutrinos. Charged particles are also produced however they would not likely make it to Earth since they will be deflected by magnetic fields between the source and Earth. This means observatories that can see the neutral messengers

are especially good for DM searches and for combining data for a multi-messenger DM search.

## CHAPTER 3

### MULTIMESSENGER ASTROPHYSICS: DETECTING HIGH ENERGY NEUTRAL MESSENGERS

#### 3.1 Introduction

Before the 20th century, all astrophysics observations were optical in nature. We literally only saw things with highly magnified optical observations. Then we discovered cosmic rays. cosmic rays are charged particles, typically naked protons or H+. This was seen by Victor Hess in 19???. Around the same time we discovered neutrinos from beta decay. Sometime around 1950 we started to build neutrino detectors which were mostly sensitive to neutrinos from the sun. Finally, it was theorized that compact objects like black holes and neutron stars would create waves in space-time when they experience mergers or collisions.

In the 21st century, we have developed new observation techniques and detectors that are no only sensitive to these four messengers - photons ([TODO: photon](#)), neutrinos ([TODO: nu](#)), Cosmic Rays (CR), and Gravitational Wave (WV) - we're collect high energy versions of these events. For the standad model particles, we're now sensitive to all messengers above the MeV eneryg range. Additionally, the GW's were sensitive to are in the stellar mass black hole region and above within our galactic neighborhood. This means were becoming sensitive to the fundamental physics occuring within the universe and we can rely on the universe as a TeV+ particle accelerator. We also have the abaility to correlate high energy events across messengers and gain new insights on the processes that occur in our universe.

This thesis focuses on very high energy (VHE) gamma rays and neutrinos. These can both be observed through the water cherenkov detection technique altho not exclusively. Methods on how to detect and observe these neutral messengers are discussed Section 3.3 and Section 3.4

#### 3.2 Charged Particles in a Medium

For high enery gamma-rays and neutrinos, we can exploit the same effect that charged particles have with water. This effect is known as Cherenkov radiation. Cherenkov Radiation occurs when a charged particle, usually electrons ( $e$ ) or muons ( $\mu$ ), traverse a medium, like water, faster than the

speed of light in that medium. This is similar to sonic boom where an object moves through air faster than the speed of sound in air. Cherenkov radiation can therefore be thought of as an 'optic boom'. Many astro-particle physics experiments will use water as the medium as because water has a unique set of properties ideal for charged particle tracking.



Figure 3.1 TODO: Show a nuclear reactor with cherenkov radiation[NEEDS A SOURCE][FACT CHECK THIS]

The frequency of light emitted due to cherenkov radiation follows the equation:

$$INSERTCherenkovwavelengthcalcHERE. \quad (3.1)$$

The absorption spectra is shown in the following figure:

### 3.3 Photons ( $\gamma$ )

### 3.4 Neutrinos ( $\nu$ )

### 3.5 Opportunities to Combine for Dark Matter



Figure 3.2 TODO: absorption spectrum of liquid and solid water[NEEDS A SOURCE][FACT CHECK THIS]

## **CHAPTER 4**

### **HIGH ALTITUDE WATER CHERENKOV (HAWC) OBSERVATORY**

#### **4.1 The Detector**

#### **4.2 Events Reconstruction and Data Acquisition**

##### **4.2.1 G/H Discrimination**

##### **4.2.2 Angle**

##### **4.2.3 Energy**

#### **4.3 Remote Monitoring**

##### **4.3.1 ATHENA Database**

##### **4.3.2 HOMER**

## **CHAPTER 5**

### **ICECUBE NEUTRINO OBSERVATORY**

#### **5.1 The Detector**

#### **5.2 Events Reconstruction and Data Acquisition**

##### **5.2.1 Angle**

##### **5.2.2 Energy**

#### **5.3 Northern Test Site**

##### **5.3.1 PIgeon remote dark rate testing**

##### **5.3.2 Bulkhead Construction**

## **CHAPTER 6**

### **COMPUTATIONAL TECHNIQUES IN PARTICLE ASTROPHYSICS**

**6.1 Neural Networks for Gamma/Hadron Separation**

**6.2 Parallel Computing for Dark Matter Analyses**

**CHAPTER 7**

**GLORY DUCK**

## **CHAPTER 8**

### **NU DUCK**

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