

Scrambling Encoded Information

Physics Research Project

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Introduction

Our Question

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This process is known as **Quantum Scrambling**. We aim to investigate this process in many-body quantum systems.

Quantum Scrambling

The set-up:

- ▶ System with many particles, e.g spin chain or qubit system
- ▶ Information is encoded as a string of operators e.g $\mathcal{O} = X \otimes I \otimes Z \otimes Y \otimes I$
- System evolves via unitary conjugation $\rightarrow \mathcal{O}(t) = U\mathcal{O}U^{\dagger}$

How can we tell if information has been scrambled? We should observe:

- ► An increase in the support of a quantum operator known as *operator spreading*.
 - ► Local operator: *IIIX_iIII*
 - ▶ Global operator: $X_1X_2X_3X_4X_5X_6X_7$
- ► A highly entangled system, signified by a growth in *entanglement entropy*

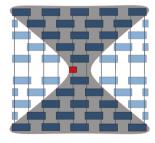


Figure: Emergent light-cone structure

Entanglement Entropy

Entanglement in state description

$$\begin{array}{c|c} |0\rangle & \hline & H \\ \hline & \\ |0\rangle & \hline \end{array} \right\} \xrightarrow[\sqrt{2}]{|00\rangle + |11\rangle}$$

Entanglement in operator space

$$T^{\dagger}XT = \frac{X - Y}{\sqrt{2}}, \quad T^{\dagger}YT = \frac{X + Y}{\sqrt{2}}$$

- Quantify entanglement using entanglement entropy.
- ▶ Initial operator string (e.g a product of Pauli Operators) becomes a product of superpositions of operators ⇒ entanglement entropy grows
- ▶ We specifically look at the (bipartite) Von Neumann Entropy,

$$S_A = -\mathsf{tr}(\rho_A \log_2(\rho_A))$$

Entanglement entropy grows linearly in time, then saturates at the Page value [1]: $S_{\text{max}} \equiv \ln|A|$.

Why?

Some motivation:

- ▶ Dates back to study of black holes and the black hole information paradox.
- Recently seen great attention in condensed matter, quantum information and quantum gravity.
- Lack of understanding of evolution in quantum chaotic systems.
- Reveals an irreversibilty of generic unitary dynamics can no longer recover information with local probes.
- Proposed to characterise dynamics in ultra-cold atom systems.

Now back to our proposition...

Simulating a Quantum System

A Proposition

For sufficiently generic unitary evolution, we expect information to become *scrambled* amongst the degrees of freedom of our system.

To prove this proposition, we simulate our system as a quantum circuit. More specifically, we utilise two circuit models:

- Clifford circuits.
- ► Non-interacting Fermion circuits.

Both models are classically simulable, such that they can be simulated with polynomial effort

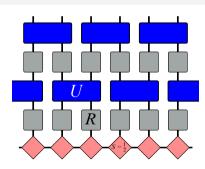
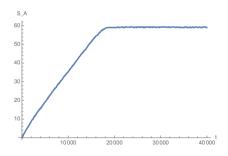


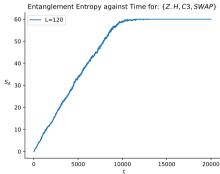
Figure: 'Brickwork' quantum circuit [2]

Clifford Circuits

Circuits comprised of gates from the Clifford Group: $\{H, CNOT, S\}$

- ▶ Not sufficient to generate operator complexity: $HXH^{\dagger} = Z$
- ► Can recover operator complexity with a non-local description of Pauli strings [3].
- ► This sacrifices any notion of operator spreading. Initial operator has support over the entire system.





Fermionic Systems

We can describe quantum circuits using operators from second quantization, namely the creation, a_i^{\dagger} and annihilation, a_i operators.

- ▶ System of N qubits: $|n_1\rangle \otimes |n_2\rangle \otimes \cdots \otimes |n_L\rangle$
- System of N Local Fermionic Modes: $|n_1, n_2, \dots, n_L\rangle \equiv (a_1^{\dagger})^{n_1} (a_2^{\dagger})^{n_2} \dots (a_r^{\dagger})^{n_L} |\mathbf{0}\rangle.$

Obey crucial anti-commutation relations:

$$\{a_i, a_j\} = \{a_i^{\dagger}, a_j^{\dagger}\} \equiv 0$$
 $\{a_i, a_j^{\dagger}\} = \delta_{ij}I$

To enforce anti-commutation relations, fermionic operators are defined using global operators. Implies inherent non-locality to fermionic description.

Non-interacting Fermion Circuits

► Use fermionic operators to construct unitary gates:

$$U = \exp(i\frac{\pi}{4}(a_0^{\dagger}a_1 + a_1^{\dagger}a_0))$$

- Build a circuit using unitary gates. For non-interacting Fermions, no terms higher than quadratic order in fermionic operators.
- ► Shown to be classically simulable by Terhal and DiVincenzo [4].
- ► Also has been shown to exhibit Operator Spreading and a growth in entanglement entropy. Good candidate for studying quantum scrambling.

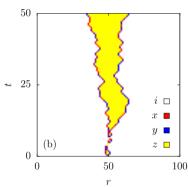


Figure: Operator spreading in Free Fermion Circuit [5]

Non-Interacting Fermion Circuits - Results

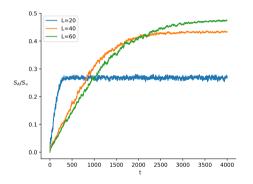


Figure: Numerical results for a free fermion circuit with system size, L=20, 40, and 60.

- ► Free fermion circuits exhibit quantum scrambling.
- ► Free Fermion circuits weakly scramble.
- ► Smaller system sizes saturate earlier.
- Entropy saturation value is dependent on system size.

Comparison

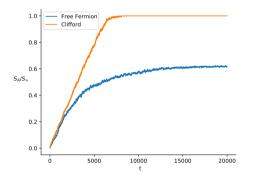


Figure: Numerical results for Clifford circuit and free fermion circuit. System size: L=100

- ► Free Fermion weakly entangle compared to Clifford circuits.
- Entropy saturation occurs later in free fermion circuit
- ► Clifford circuits: Exhibit entanglement entropy growth but no operator spreading.
- ► Free Fermion circuits: Exhibit both entanglement entropy growth and operator spreading, but weaker.

Further Directions

- ▶ Changing geometry of free-fermion circuits to boost entanglement growth.
- ► Finding new, less constrained sets of gates in the Clifford group to scramble operators.
- ► Recent work on dual-unitary circuits has shown them to be good candidates for the study of scrambling using tensor network methods.
- ▶ Including higher-order terms in fermion circuits (interaction terms).

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