

# Nonlinear Dynamics

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# 1 Flows on the line

## 1.1 Introduction

**Definition 1.1** (Fixed points). A fixed point on a phase diagram is a point in which there is no flow, i.e.  $x' = 0$ . Fixed points represent equilibrium solutions, and are denoted with an asterisk  $x^*$ .

**Definition 1.2** (Phase point). A phase point is an imaginary particle placed at a point  $x_0$  from which we can observe how it is carried along with the "flow". As time increases, the phase point moves along the  $x$ -axis according to some function  $x(t)$ .  $x(t)$  is called the trajectory based at  $x_0$ .

**Theorem 1.3.** *Consider the IVP*

$$\begin{aligned}x' &= f(x), \\x(0) &= x_0.\end{aligned}$$

*If  $f(x)$  and  $f'(x)$  are continuous on an open interval  $R$  of the  $x$ -axis, and  $x_0 \in R$ , then the initial value problem has a unique solution on some time interval  $-\tau, \tau$  about  $t = 0$ .*

**Remark.** In a first-order system, trajectories can either approach a fixed point, or diverge to infinity. Trajectories are forced to increase or decrease monotonically because  $x'$  can not hold two values for the same  $x$ . This means that phase points never 'overshoot' a fixed point to which its path converges. Therefore there are no periodic solutions to  $x' = f(x)$ .

**Definition 1.4** (Potentials). In a first-order system  $x' = f(x)$ , the potential function  $V(x)$  is defined by

$$f(x) = -\frac{dV}{dx}$$

**Remark.** Using the chain rule, we can see

$$\begin{aligned}\frac{dV}{dt} &= \frac{dV}{dx} \frac{dx}{dt} \\&= -\left(\frac{dV}{dx}\right)^2 \\&\leq 0\end{aligned}$$

Therefore potential decreases or stays constant along trajectories.

**Proposition 1.5** (Euler's method). Suppose  $x' = f(x)$  a one-dimensional dynamical system. Euler's method is a way of estimating  $x(t)$  at discrete times spaced  $\Delta t$  apart. We define  $x_n$  to be the approximate value of  $x(t)$  at  $n\Delta t$  by choosing a starting point  $x_0$ , and using the following recursive definition to find any  $x_n$ :

$$x_{n+1} = x_n + f(x_n)\Delta t.$$

## 1.2 Bifurcations

**Definition 1.6** (Bifurcation). A bifurcation is a change in the qualitative structure of the flow caused by changing a parameter in an equation. The values at which bifurcations occur are called bifurcation points.

**Remark.** Bifurcation diagrams plot fixed points on the vertical axis against parameter values  $r$  on the horizontal axis. Dashed lines represent unstable fixed points, where solid lines represent stable ones.

**Definition 1.7** (Saddle-node bifurcation). This is a bifurcation presents as fixed points colliding and annihilating as a parameter is varied. An example of this is increasing parameter  $r$  in the equation  $x' = r + x^2$ . When  $r < 0$  this equation has two zeros in the phase plane and thus two fixed points. When  $r = 0$   $x(t)$  has one phase point and when  $r > 0$  there are no phase points. Saddle-node bifurcation has normal form

$$x' = r + x^2 \text{ or } x' = r - x^2.$$

**Definition 1.8** (Normal form). The normal form of a bifurcation is the prototypical presentation of that bifurcation. For example, the partial taylor expansion of a function with a saddle-node bifurcation at  $x = x^*$  and  $r = r_c$  presents as the normal form of a saddle bifurcation

$$f(x^*, r_c) + (x - x^*) \frac{\partial f}{\partial x}(x^*, r_c) + (r - r_c) \frac{\partial f}{\partial r}(x^*, r_c) + \frac{1}{2}(x - x_0)^2 \frac{\partial^2 f}{\partial x^2}(x^*, r_c).$$

Because  $\frac{\partial f}{\partial x}(x^*, r_c) = 0$  and  $f(x^*, r_c) = 0$  at the bifurcation point, this equation can then be written in normal form

$$(r - r_c) \frac{\partial f}{\partial r}(x^*, r_c) + \frac{1}{2}(x - x^*)^2 \frac{\partial^2 f}{\partial x^2}(x^*, r_c).$$

**Definition 1.9** (Transcritical bifurcation). A fixed point that exists for all values of a parameter, but whose stability changes depending on that parameter, is said to undergo transcritical bifurcation. The normal form of a transcritical bifurcation is

$$x' = rx - x^2.$$

**Example 1.10.** Show that the first-order system  $x' = x(1 - x^2) - a(1 - e^{-bx})$  undergoes a transcritical bifurcation at  $x = 0$  when the parameters  $a, b$  satisfy a certain equation.

*Proof:*  $x = 0$  is a fixed point for all  $(a, b)$  so it is plausible that the point bifurcates transcritically. Using the second degree taylor expansion for  $x'$ , we can estimate the behavior of this function near  $x = 0$ :

$$x' = (1 - ba)x + \frac{1}{2}(b^2a)x^2 + O(x^3).$$

From this we see that transcritical bifurcation must occur when  $ba = 1$ , i.e. when the first derivative of  $x'$  with respect to  $x$  is zero. Using this equation we can estimate the location of fixed points near  $ab = 1$  by simplifying the above equation to solve for  $x^*$ :

$$x^* \approx \frac{2(ba - 1)}{b^2a}.$$

□

**Example 1.11.** Describe the behavior of the laser equation  $n' = Gn(N_0 - \alpha n) - kn$ .

*Proof:* The change in the number of photons  $n'$  is given by the rate of photon generation  $GnN$  minus the outflow rate  $kn$ . This equation has a normal form  $X' = RX - X^2$ . The rate of photon generation is proportional to the number of excited atoms  $N = (N_0 - \alpha n)$ . We can simplify the above equation into two forms which tell us about the nature of the fixed points as  $N_0$  changes.

$$\begin{aligned} n' &= -G\alpha n^2 + (GN_0 - k)n \\ n' &= -G\alpha n \left( n - \frac{GN_0 - k}{G\alpha} \right) \end{aligned}$$

Clearly transcritical bifurcation of fixed point  $n = 0$  occurs at  $GN_0 - k = 0$ . This occurs with our intuitive understanding that if the rate of photon egress is larger than photon generation for small  $n$ , then no laser action will take place. However if  $GN_0 > k$  then the fixed point at  $n = 0$  becomes unstable and a stable fixed point emerges, as given by the second equation. □

**Definition 1.12** (Supercritical pitchfork bifurcation). Pitchfork bifurcations occur in problems which have symmetry, i.e. fixed points appear and disappear in symmetrical pairs. The normal form of supercritical pitchfork bifurcation is

$$x' = rx - x^3.$$

**Definition 1.13** (Subcritical pitchfork bifurcation). Subcritical pitchfork bifurcation presents as a mirror image about the  $y$ -axis of supercritical bifurcation, but with unstable symmetrical fixed points with a stable fixed point at zero, that converge to an unstable fixed point at  $x = 0$ . The normal form of a subcritical pitchfork bifurcation is

$$x' = rx + x^3.$$