

Quantum Decoherence Mostly Results in White Noise

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Abstract

An overwhelming majority of quantum (pure and mixed) states, when undertaking decoherence, will result in a classical probability with no algorithmic information. Thus most quantum states decohere into white noise. This can be seen as a consequence of the vastness of Hilbert spaces.

Information non-growth laws say information about a target source cannot be increased with randomized processing. In classical information theory, we have [CT91]

$$I(g(X):Y) \leq I(X:Y).$$

where g is a randomized function, X and Y are random variables, and I is the mutual information function. Thus processing a channel at its output will not increase its capacity. Information conservation carries over into the algorithmic domain, with the inequalities [Lev84, Eps22]

$$\mathbf{I}(f(x):y) <^+ \mathbf{I}(x:y); \quad \mathbf{I}(f(a);\mathcal{H}) <^+ \mathbf{I}(a;\mathcal{H}).$$

The information function is $\mathbf{I}(x:y) = \mathbf{K}(x) + \mathbf{K}(y) - \mathbf{K}(x,y)$, where \mathbf{K} is Kolmogorov complexity. The other term is $\mathbf{I}(a;\mathcal{H}) = \mathbf{K}(a) - \mathbf{K}(a|\mathcal{H})$, where $\mathcal{H} \in \{0,1\}^\infty$ is the halting sequence. These inequalities ensure target information cannot be obtained by processing. If for example the second inequality was not true, then one can potentially obtain information about \mathcal{H} with simple functions. Obtaining information about \mathcal{H} violates the Independence Postulate, (see [Lev13]). Information non growth laws can be extended to signals [Eps23b] which can be modeled as probabilities over \mathbb{N} or Euclidean space¹. The “signal strength” of a probability p over \mathbb{N} is measured by its self information.

$$\mathbf{I}_{\text{Prob}}(p:p) = \log \sum_{i,j} 2^{\mathbf{I}(i:j)} p(i)p(j).$$

A signal, when undergoing randomized processing f , will lose its cohesion². Thus any signal going through a classical channel will become less coherent [Eps23b].

$$\mathbf{I}_{\text{Prob}}(f(p):f(p)) <^+ \mathbf{I}_{\text{Prob}}(p:p).$$

In Euclidean space, probabilities that undergo convolutions with probability kernels will lose self information. For example a signal spike at a random position will spread out when convoluted with the Gaussian function, and lose self information. The above inequalities deal with classical

¹In [Eps23b] probabilities over $\{0,1\}^\infty$ and T_0 second countable topologies were also studied.

²A probability p , when processed by a channel $f : \{0,1\}^* \times \{0,1\}^* \rightarrow \mathbb{R}_{\geq 0}$ is a new probability $fp(x) = \sum_z f(x|z)p(z)$.

transformations. One can ask, is whether, quantum information processing can add new surprises to how information signals occur and evolve.

One can start with the prepare-and-measure channel, also known as a Holevo-form channel. Alice starts with a random variable X that can take values $\{1, \dots, n\}$ with corresponding probabilities $\{p_1, \dots, p_n\}$. Alice prepares a quantum state, corresponding to density matrix ρ_X , chosen from $\{\rho_1, \dots, \rho_n\}$ according to X . Bob performs a measurement on the state ρ_X , getting a classical outcome, denoted by Y . Though it uses quantum mechanics, this is a classical channel $X \rightarrow Y$. So using the above inequality, cohesion will deteriorate regardless of X 's probability, with Th

$$\mathbf{I}_{\text{Prob}}(Y : Y) <^+ \mathbf{I}_{\text{Prob}}(X : X).$$

There remains a second option, constructing a signal directly from a mixed state. This involves constructing a mixed state, i.e. density matrix σ , and then performing a POVM measurement³ E on the state, inducing the probability $E\sigma(\cdot)$. However from [Eps23b], for elementary (even enumerable) probabilities $E\sigma$,

$$\mathbf{I}_{\text{Prob}}(E\sigma : E\sigma) <^+ \mathbf{K}(\sigma, E).$$

Thus for simply defined density matrices and measurements, no signal can appear. So experiments that are simple will result in simple measurements, or white noise. However it could be that a larger number of uncomputable pure or mixed states produce coherent signals. However, theorems in [Eps23a] say otherwise, in that given a POVM measurement E , a vast majority of pure and mixed states will have negligible self-information. Thus for uniform distributions Λ and μ over pure and mixed states⁴⁵,

$$\int 2^{\mathbf{I}_{\text{Prob}}(E|\psi\rangle : E|\psi\rangle)} d\Lambda = O(1); \quad \int 2^{\mathbf{I}_{\text{Prob}}(E\sigma : E\sigma)} d\mu(\sigma) = O(1).$$

This can be seen as a consequence of the vastness of Hilbert spaces as opposed to the limited discriminatory power of quantum measurements. In addition, there could be non-uniform distributions of pure or mixed states that could be of research interest. In quantum decoherence, a quantum state becomes entangled with the environment, losing decoherence. The off diagonal elements of the mixed state become dampened, as the state becomes more like a classical mixture of states. Let p_σ be the idealized classical probability that σ decoheres to, with $p_\sigma(i) = \sigma_{ii}$. The following theorem from [Eps23a] states that for an overwhelming majority of pure or mixed states σ , p_σ is noise, that is, has negligible self-information.

$$\int 2^{\mathbf{I}_{\text{Prob}}(p_{|\psi\rangle} : p_{|\psi\rangle})} d\Lambda = O(1); \quad \int 2^{\mathbf{I}_{\text{Prob}}(p_\sigma : p_\sigma)} d\mu(\sigma) = O(1).$$

These results mark a continuing investigation of algorithmic information and physics. Certainly a measurement can potentially produce a signal if the (randomized) preparation of quantum states σ is a complex process (i.e. high $\mathbf{K}(\sigma)$). If Λ or μ is the apriori distribution of quantum states, then virtually all measurements are signalless. Does this necessitate the assumption of a new apriori distribution for quantum states?

³A POVM measurement E is a collection of positive-semi definite Hermitian matrices $\{E_k\}$ such that $\sum_k E_k = 1$. Given a state σ , E induces a probability over the measurements of the form $E\sigma(k) = \text{Tr} E_k \sigma$.

⁴The mixed state integral is $\int f(\sigma) d\mu(\sigma) = \int_{\Delta_M} \int_{\Lambda_1} \dots \int_{\Lambda_M} f\left(\sum_{i=1}^M p_i |\psi_i\rangle \langle \psi_i|\right) d\Lambda_1 \dots d\Lambda_M d\eta(p_1, \dots, p_M)$, where η is any distribution over the M -simplex Δ_M .

⁵The proof to these inequalities is in the running for the strangest in AIT, relying on a lower computable combination of *upper* computable tests.

References

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