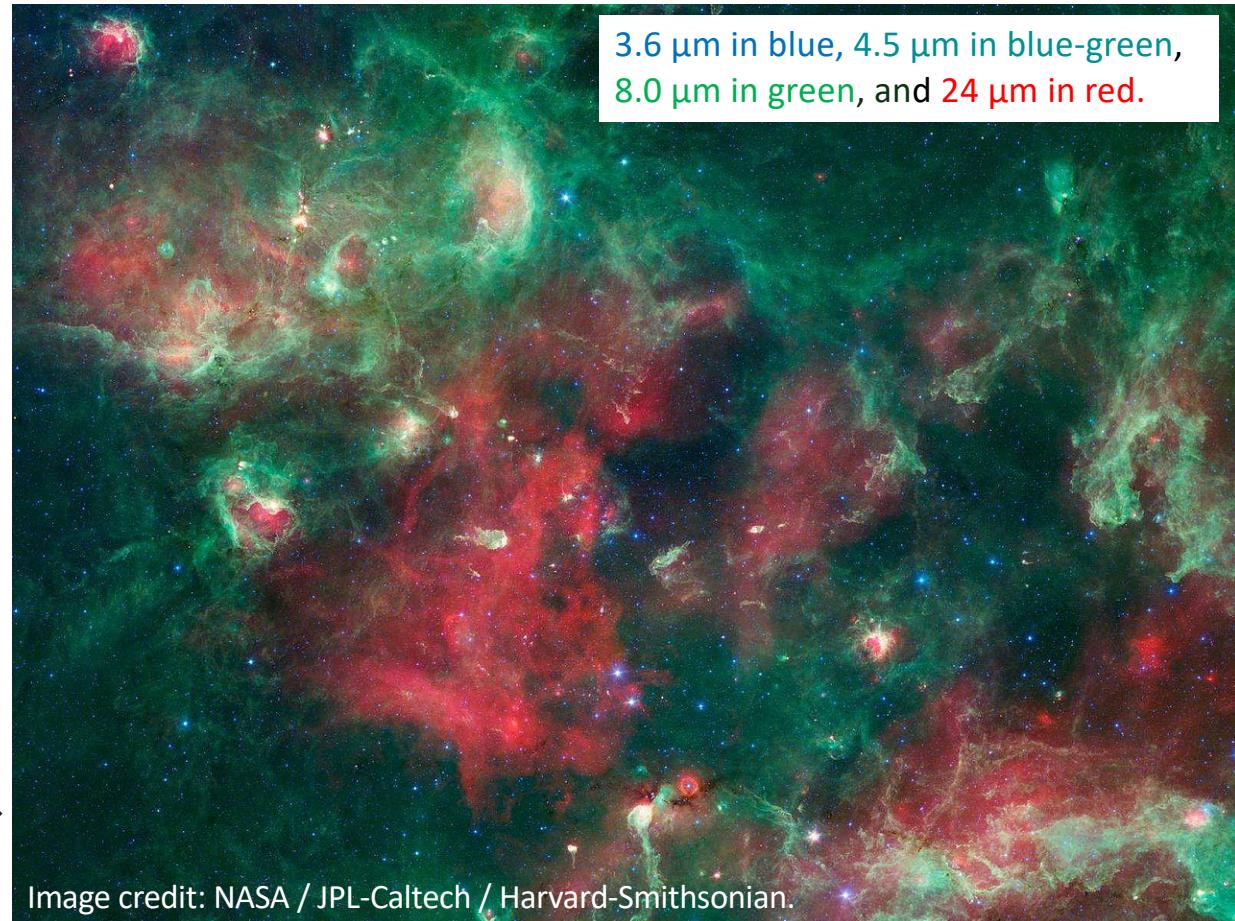


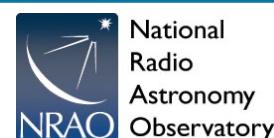
Ionized Gas

HII Regions
Recombination Lines
Free-Free Emission

Cygnus X Star Forming
region as imaged by the
Spitzer Space Telescope →



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Radio Recombination Lines (ERA 7.2, 7.6)

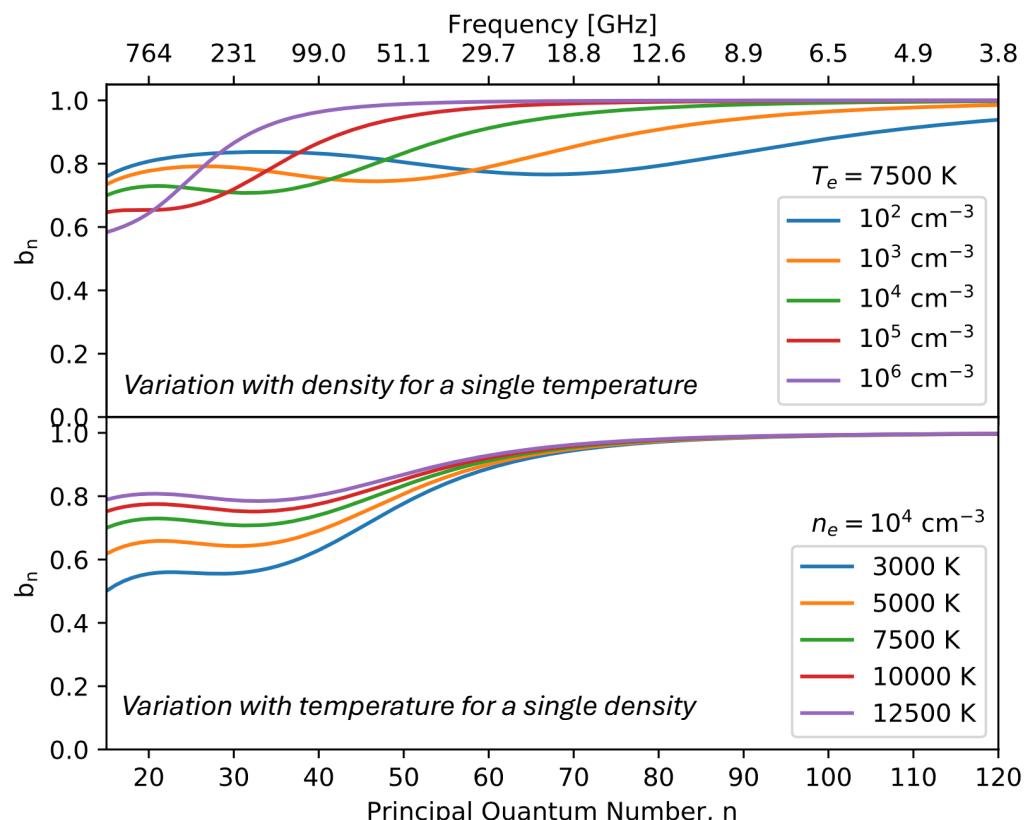
Radiative Transfer! (7.6.1)

LTE does not apply to all recombination lines

The departure (from LTE) coefficient of spontaneous emission, b_n , shown as a function of principal quantum number (frequency) →

Deviation from LTE at low density, low temperature and low quantum number

Emig Thesis (2021)



Radio Recombination Lines (ERA 7.2, 7.6)

Radiative Transfer! (7.6.1)

Remember, the line opacity coefficient in LTE is:

$$\kappa(\nu) = \frac{c^2}{8\pi\nu_0^2} \frac{g_U}{g_L} n_L A_{UL} \left[1 - \exp\left(-\frac{h\nu_0}{kT}\right) \right] \phi(\nu) \quad (7.67)$$

The book finds for the $n \rightarrow n + 1$ electronic transition of hydrogen (see book for derivation w/ Saha Equation)

Valid for all radio recombination lines with $n \gg 1$. Dependence on 'n' gone!

If we are in LTE limit:

$$\kappa(\nu_0) \approx \left(\frac{n_e^2}{T_e^{5/2} \Delta\nu} \right) \left(\frac{4\pi e^6 h}{3m_e^{3/2} k^{5/2} c} \right) \left(\frac{\ln 2}{2} \right)^{1/2}. \quad (7.94)$$

Radio Recombination Lines (ERA 7.2, 7.6)

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The **optical depth** becomes defined by an **Emission Measure**:

$$\frac{\text{EM}}{\text{pc cm}^{-6}} \equiv \int_{\text{los}} \left(\frac{n_e}{\text{cm}^{-3}} \right)^2 d\left(\frac{s}{\text{pc}} \right). \quad (7.95)$$

Where the line opacity is,

$$\tau_L \approx 1.92 \times 10^3 \left(\frac{T_e}{\text{K}} \right)^{-5/2} \left(\frac{\text{EM}}{\text{pc cm}^{-6}} \right) \left(\frac{\Delta\nu}{\text{kHz}} \right)^{-1}. \quad (7.96)$$

Radio Recombination Lines (ERA 7.2, 7.6)

Radiative Transfer! (7.6.1)

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Where the line opacity is,

**Notice what this depends on!
See where this is going...?**

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Radio Recombination Lines (ERA 7.2, 7.6)

Radiative Transfer! (7.6.1)

Because the line opacity in HII region is **optically thin** $\tau \ll 1$, the **brightness temperature** contributed by a recombination emission line at its center frequency is,

$$T_L \approx T_e \tau_L \approx 1.92 \times 10^3 \left(\frac{T_e}{K} \right)^{-3/2} \left(\frac{EM}{pc \text{ cm}^{-6}} \right) \left(\frac{\Delta\nu}{kHz} \right)^{-1}. \quad (7.97)$$

A direct probe of electron temperature, T_e , and EM!

Radio Recombination Lines (ERA 7.2, 7.6)

Radiative Transfer! (7.6.1)

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A direct probe of electron temperature, T_e , and EM!

And at frequencies high enough that the line continuum (free-free; we will get to this next!), T_C is also optically thin the peak line-to-continuum ratio in LTE can be used to explicitly find T_e

$$\left(\frac{T_e}{K} \right) \approx \left[7.0 \times 10^3 \left(\frac{\nu}{\text{GHz}} \right)^{1.1} 1.08^{-1} \left(\frac{\Delta\nu}{\text{km s}^{-1}} \right)^{-1} \left(\frac{T_C}{T_L} \right) \right]^{0.87}. \quad (7.99)$$



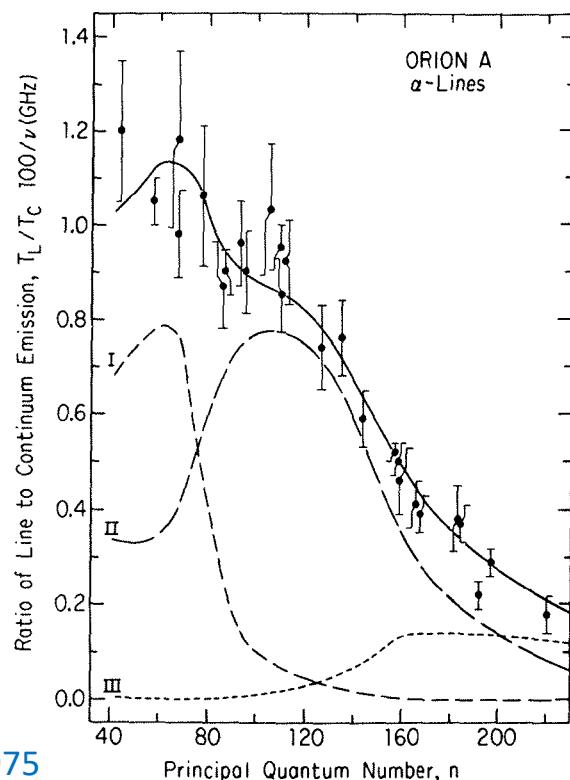
Main Takeaway:

The **line-to-continuum ratio yields an estimate of the electron temperature T_e** that is independent of the emission measure (opacity) so long as the frequency is high enough that the continuum optical depth is small.

Radio Recombination Lines (ERA 7.2, 7.6)

Radiative Transfer! (7.6.1)

Mapping the temperature distribution in the HII Region of the Orion Nebula:



Lockman and Brown 1975

Main Takeaway:

The line-to-continuum ratio yields an estimate of the electron temperature T_e that is independent of the emission measure (opacity) so long as the frequency is high enough that the continuum optical depth is small.

Fig. 7.8 (ERA)

Table 2.3. Lockman-Brown model of the Orion Nebula.^a

Region	Temperature (K)	$N_e(\text{cm}^{-3})$	Angular size (arcmin)	Linear size (pc)
I	(7,500)	$10^{4.5}$	0.33	0.043
II	(10,000)	$10^{3.5}$	4.3	0.56
III	(12,500)	$(10^{2.3})$	19.1	2.5

^a Quantities in parentheses are determined by the recombination-line observations.

Radio Recombination Lines (ERA 7.2, 7.6)

Differences between the rest and observed frequencies of radio recombination lines are attributed to Doppler shifts from nonzero radial velocities

With a simple rotational model for the disk of our Galaxy, astronomers can convert radial \rightarrow velocities to distances

They roughly outline the major spiral arms.

Georgelin & Georgelin 1976

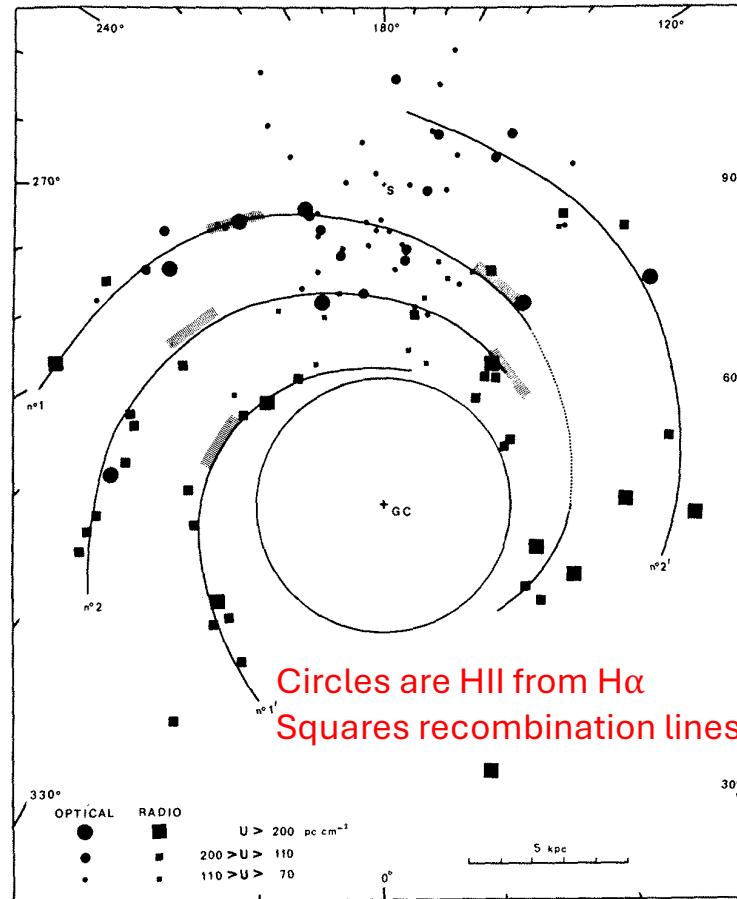


Fig. 7.9 (ERA)

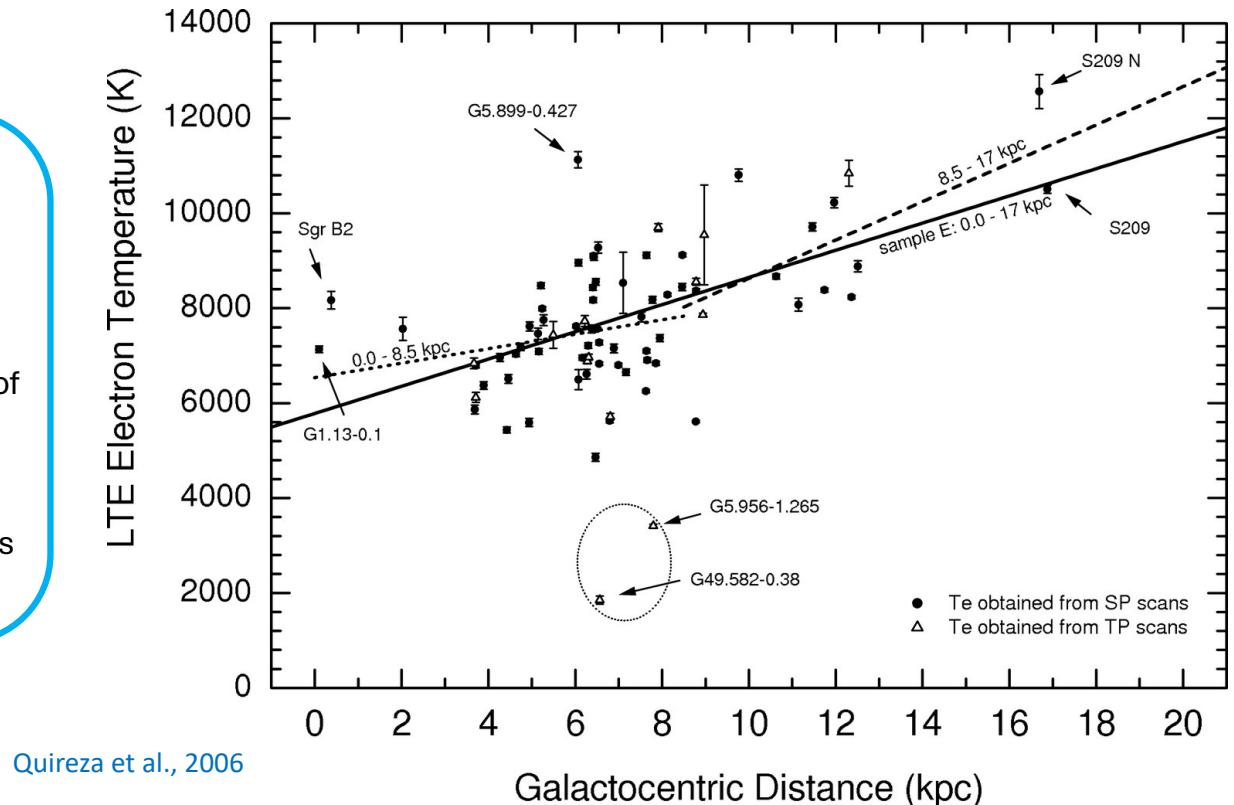
Radio Recombination Lines (ERA 7.2, 7.6)

Fig. 7.10 (ERA)

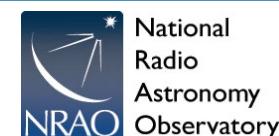
Observed electron temperatures of Galactic HII regions: **temperature increases with distance from the Galactic center** →

The explanation for this trend is the observed **decrease in metallicity** (relative abundance of elements heavier than helium) with galactocentric distance

Power radiated by emission lines of “metals” is the principal cause of HII region cooling.

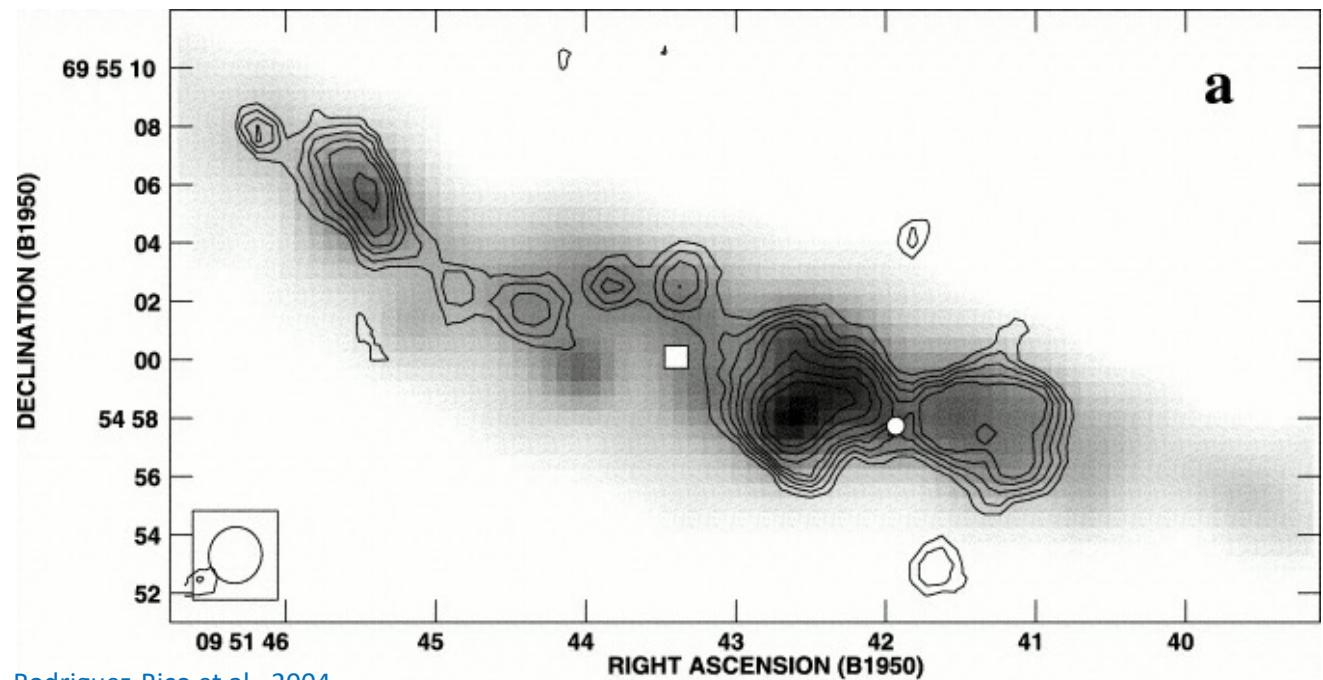


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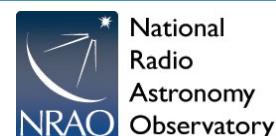
Radio Recombination Lines (ERA 7.2, 7.6)

Fig. 7.11 (ERA)

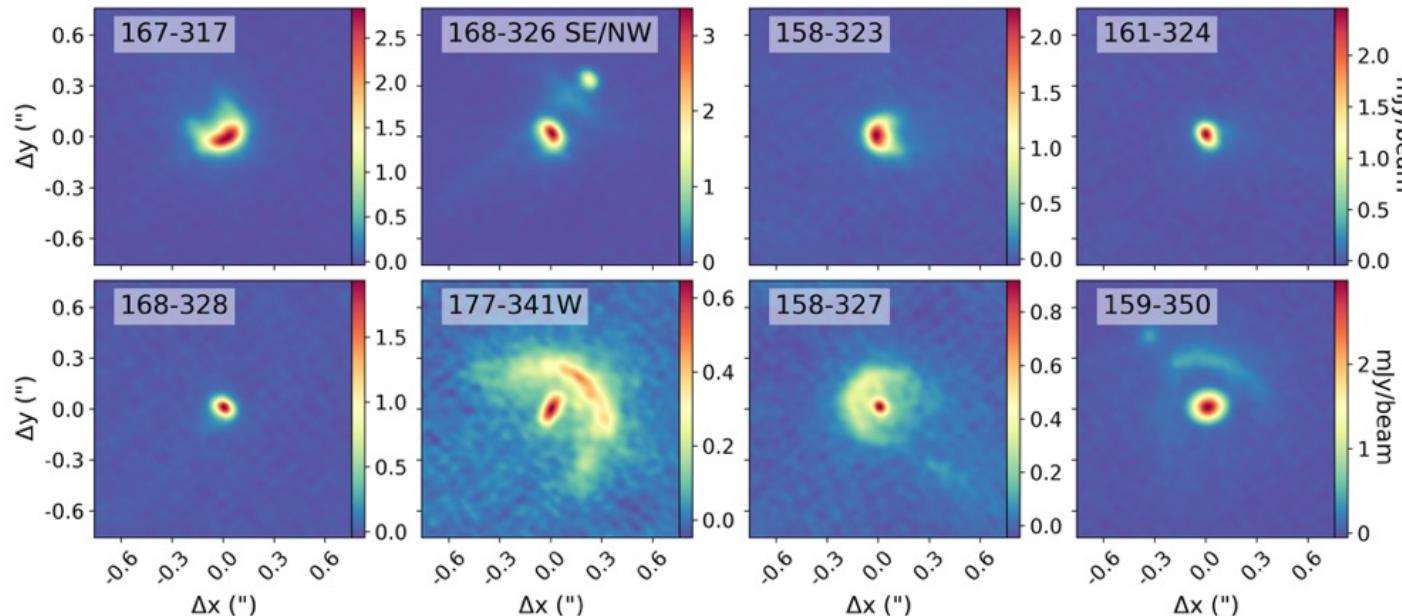


M82 imaged in H92 α (contours) and 8.3 GHz free-free continuum (gray scale)

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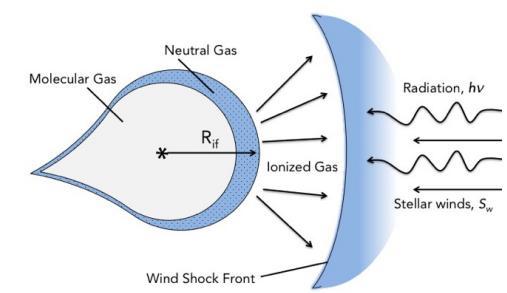


Radio Recombination Line Emission From *Proplyds*



ALMA, Band 3 Continuum
(Ballering et al. 2023)

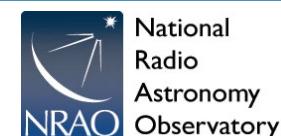
* Formed by ionization/evaporation of gas in circumstellar disks around young, low-mass stars



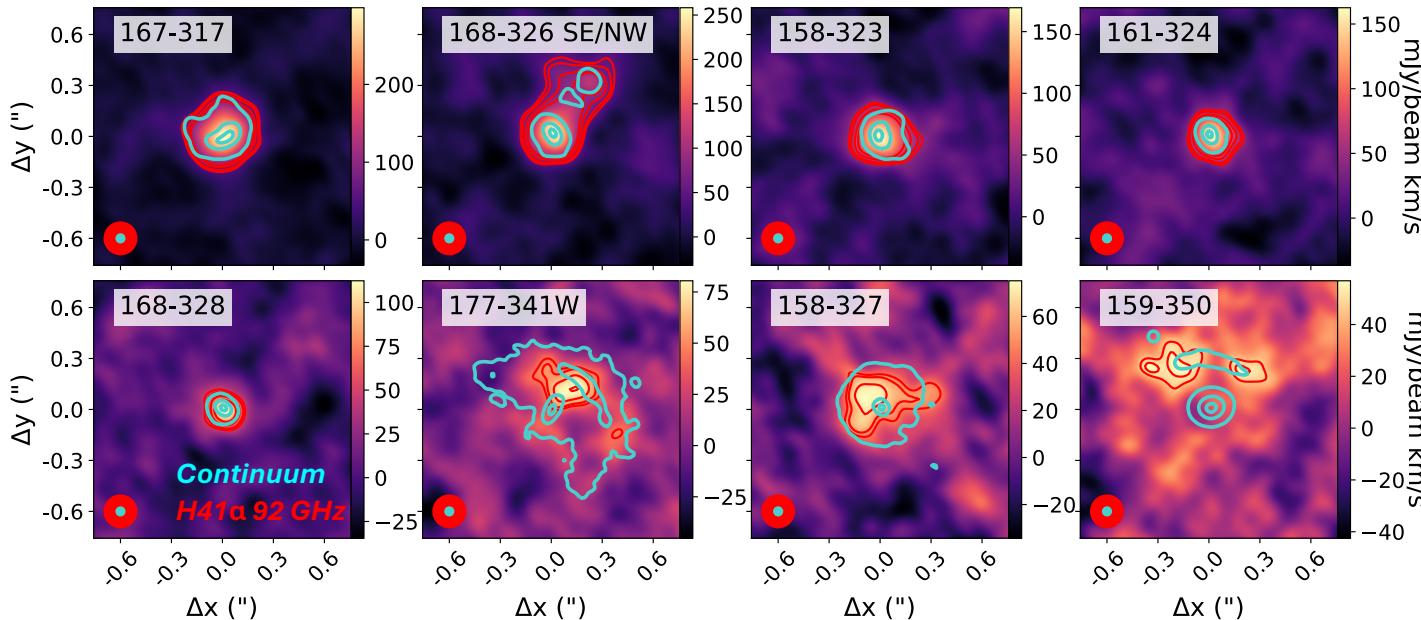
Mini ultra-compact HII regions

Slide Credit: Ryan Boyden

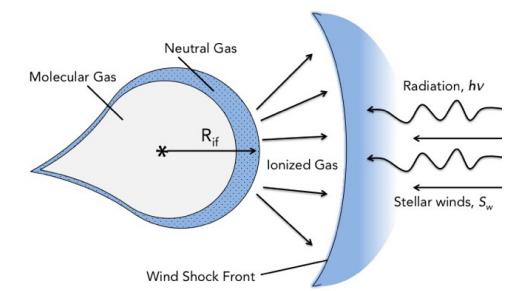
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Radio Recombination Line Emission From *Proplyds*



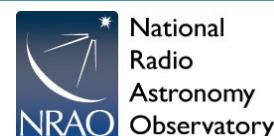
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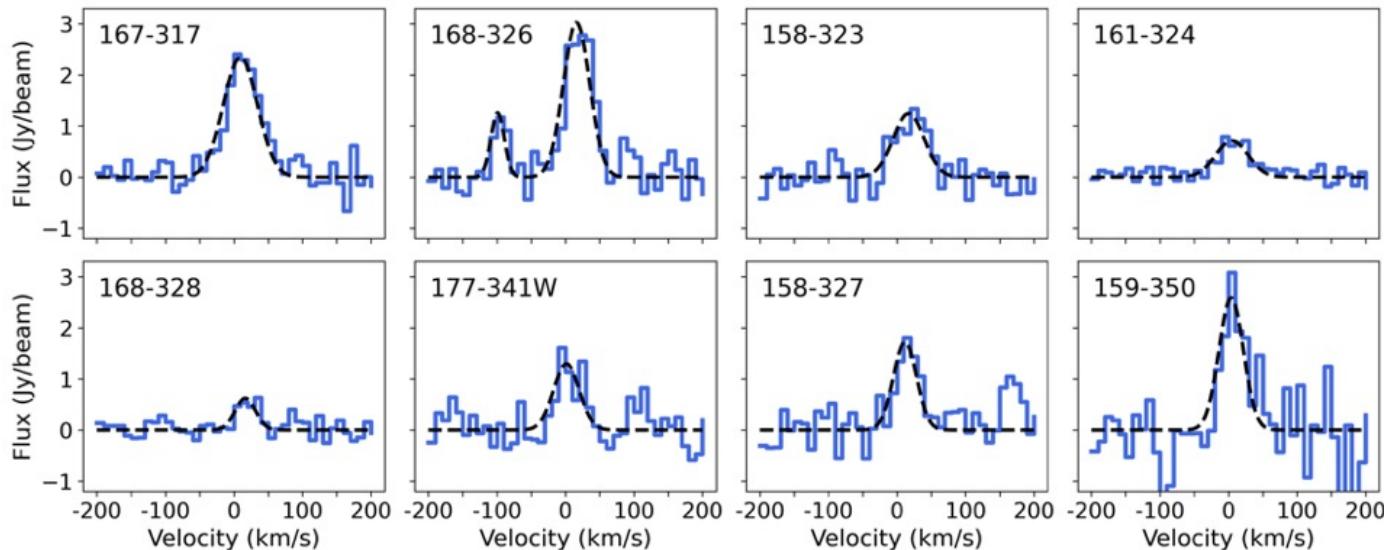
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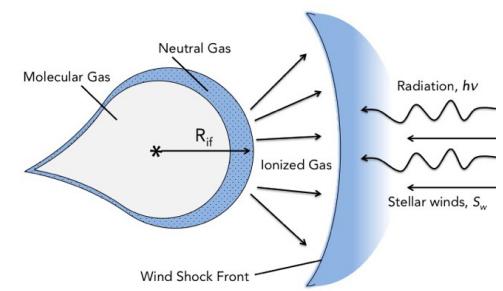
Radio Recombination Line Emission From *Proplyds*



**Line Widths: ~50 km/s
Dominated by ionized gas motions**

ALMA, H41 α
(Boyden et al. in prep)

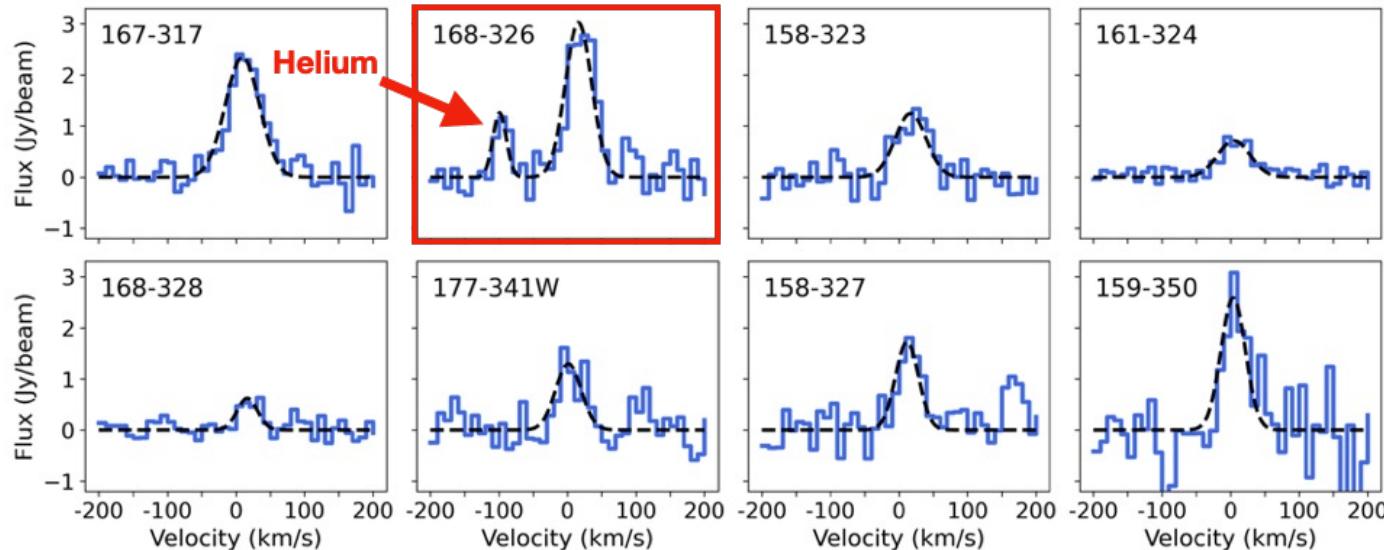
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Slide Credit: Ryan Boyden

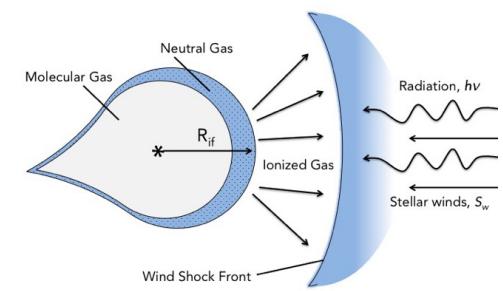
Radio Recombination Line Emission From *Proplyds*



**He/H abundance ratio: 0.2
(Helium-rich disk)**

ALMA, H41α
(Boyden et al. in prep)

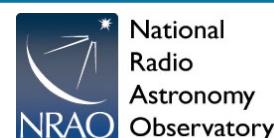
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Mini ultra-compact HII regions

Slide Credit: Ryan Boyden

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Emission Mechanisms

Spectral Lines (ERA Chap. 7)



Free-Free (ERA Chap. 4)



Synchrotron (ERA Chap. 5)



Pulsars (ERA Chap. 6)



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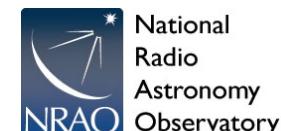
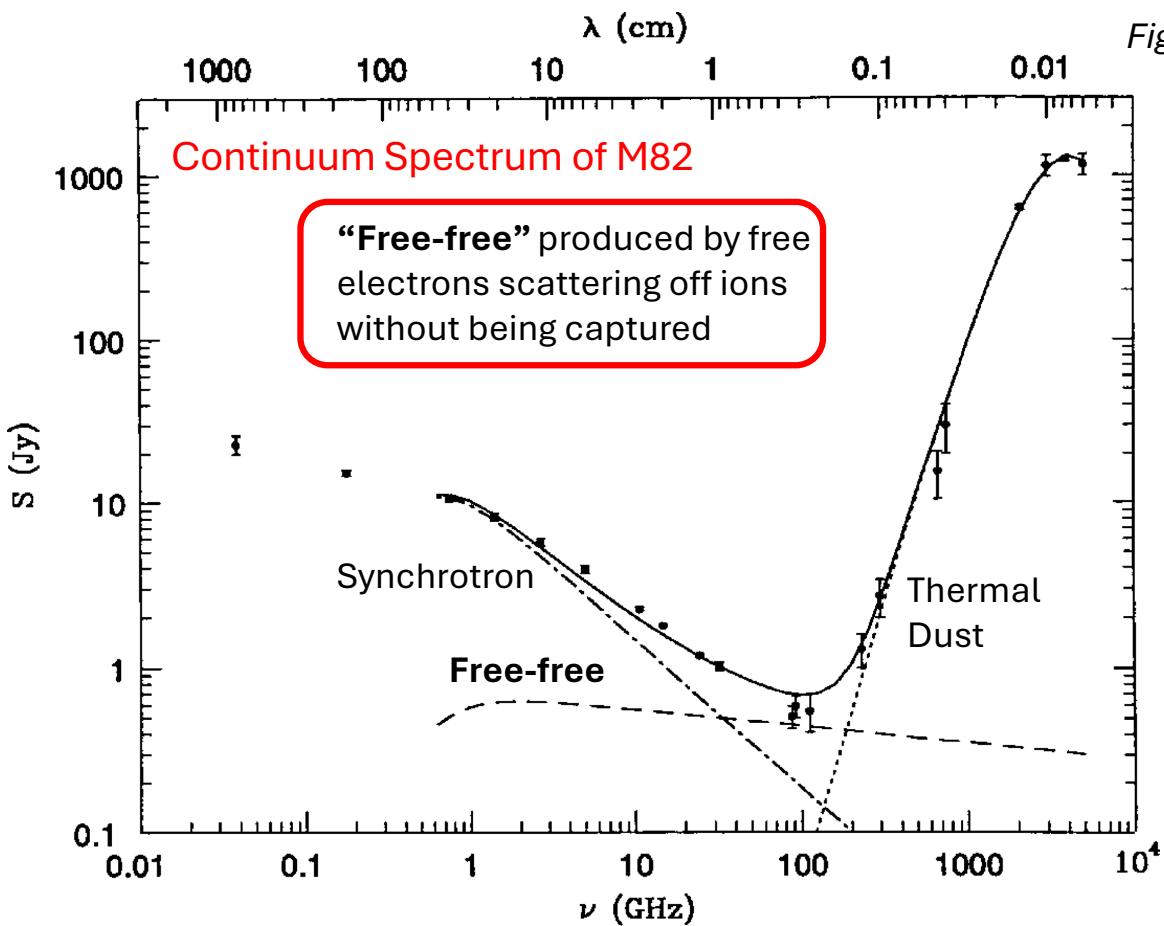
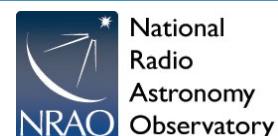


Fig. 2.24 (ERA)



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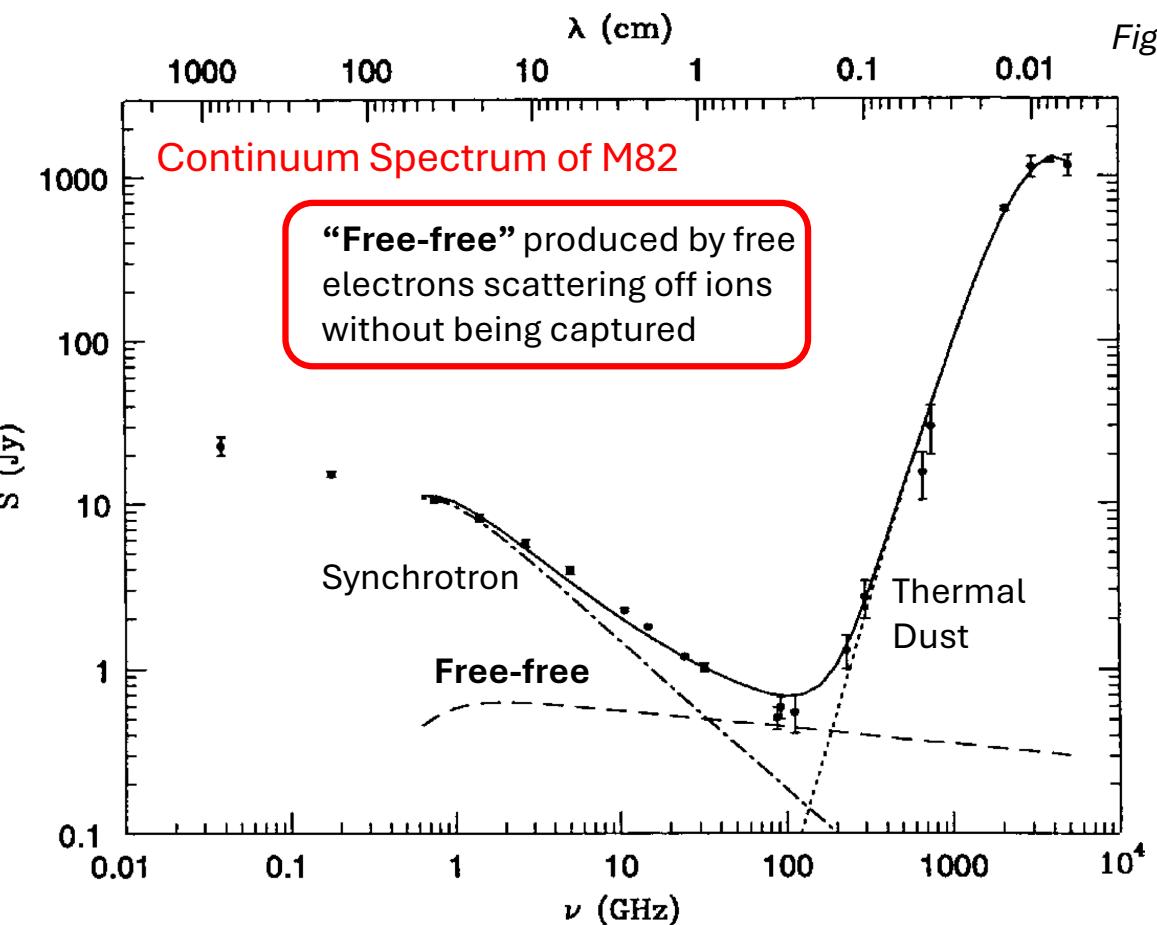
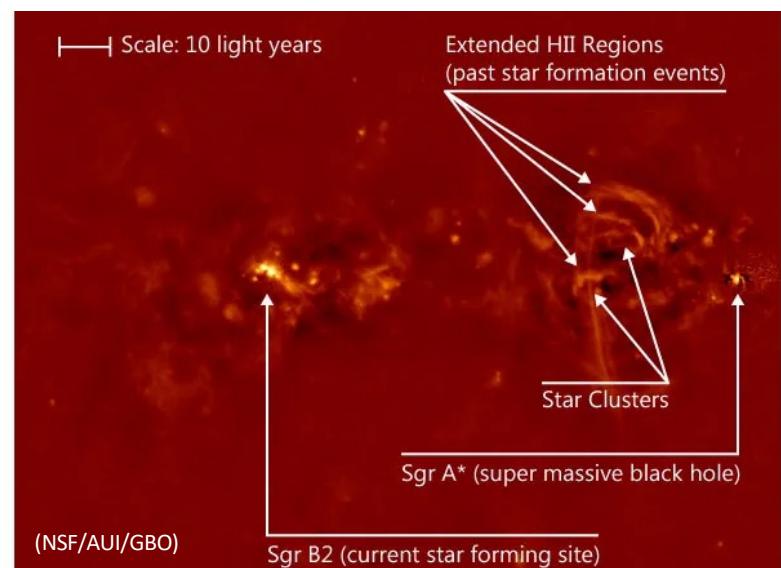


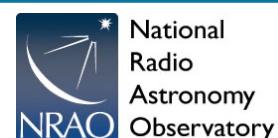
Fig. 2.24 (ERA)

Results from the **Green Bank Observatory**
Continuum Instrument, MUSTANG-2, that
observes the sky at wavelengths of 3mm (90 GHz)



Ginsburg et al., 2020

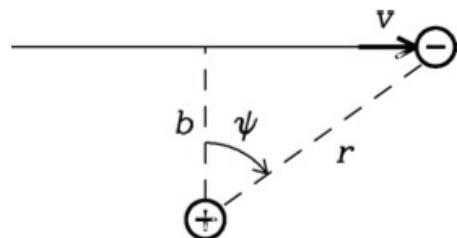
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Free-free Radiation (ERA 4.1)

Thermal and Nonthermal Emission

Free-free: the emission from a charge (e.g., electron) in the Coulomb field of another charge (ion, electron) when it experiences a small deviation in its path



The distance of closest approach, b , is called the impact parameter and the interval $\tau = b/v$ is the collision time.

Remember Larmor radiation power is:

$$P = \frac{2q^2\dot{v}^2}{3c^3} \quad (4.1)$$

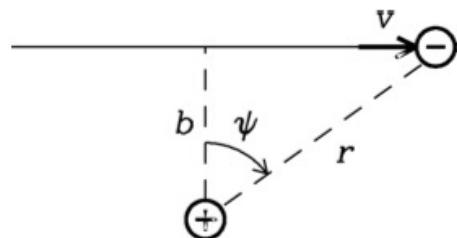
More generally, ‘**bremsstrahlung**’ radiation: **electromagnetic radiation** with power P produced by accelerating (or decelerating) an electric charge q

NOTE: the magnetic counterpart **magnetobremsstrahlung** or “**magnetic braking radiation**” (e.g., synchrotron radiation) is covered in Chapter 5!

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Typically, ‘**thermal**’ and electrons follow Maxwellian distribution

NOTE: the magnetic counterpart **magnetobremsstrahlung** or “**magnetic braking radiation**” (e.g., synchrotron radiation) is covered in Chapter 5!

Typically, ‘**nonthermal**’ relativistic electrons w/ power-law energy distribution

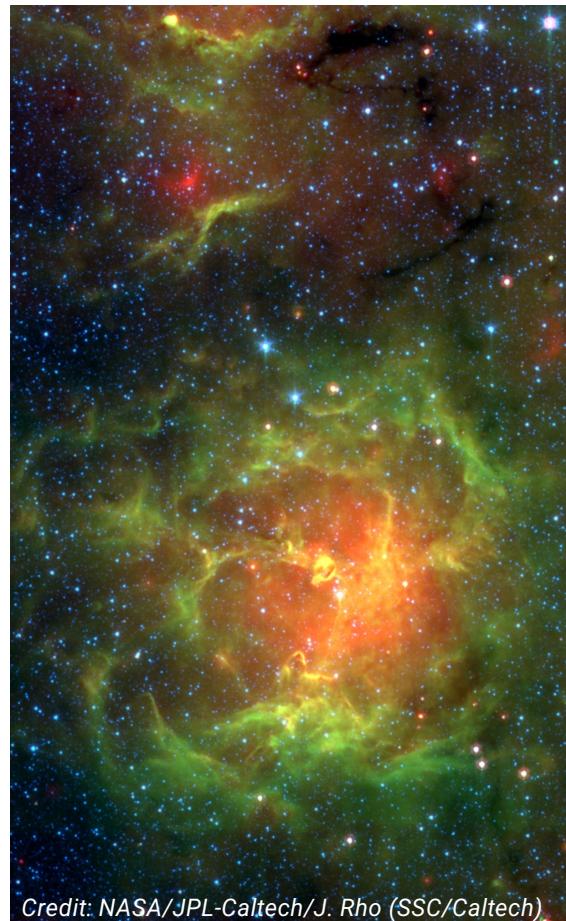
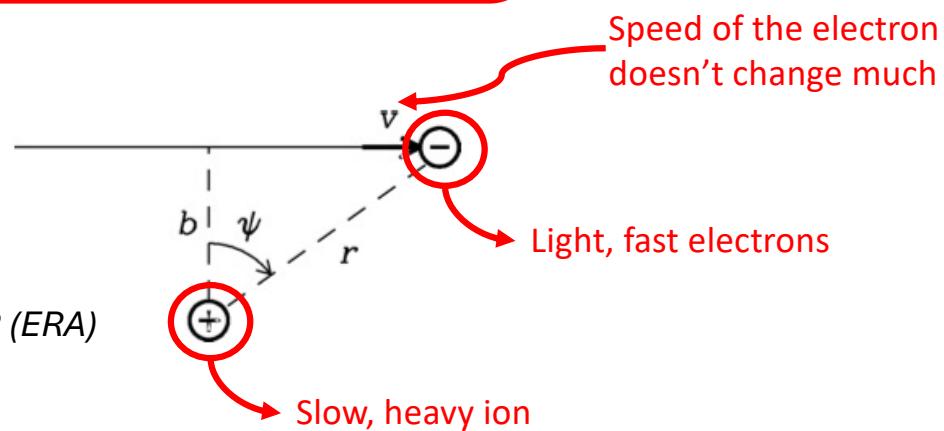
Free-free Radiation (ERA 4.3)

Free-Free in HII Regions – we need to simplify the problem

Main Takeaway

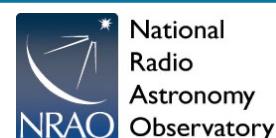
Only electron-ion collisions are important, and only the electrons radiate significantly

Fig. 4.2 (ERA)



Credit: NASA/JPL-Caltech/J. Rho (SSC/Caltech)

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The glowing Trifid Nebula HII region is revealed with near- and mid-infrared views from NASA's Spitzer Space Telescope.