

# Search for chargino and neutralino production in final states with two same-sign leptons, jets and missing transverse momentum at $\sqrt{s} = 13$ TeV with the ATLAS detector

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# Introduction

# Standard Model : Fundamental Particles

- Standard Model(SM) is the current mainstream theory to describe the electromagnetic force, weak force and strong force.

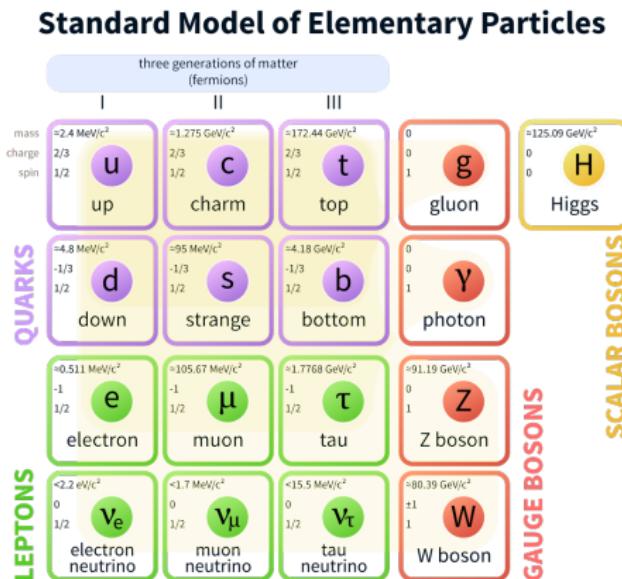
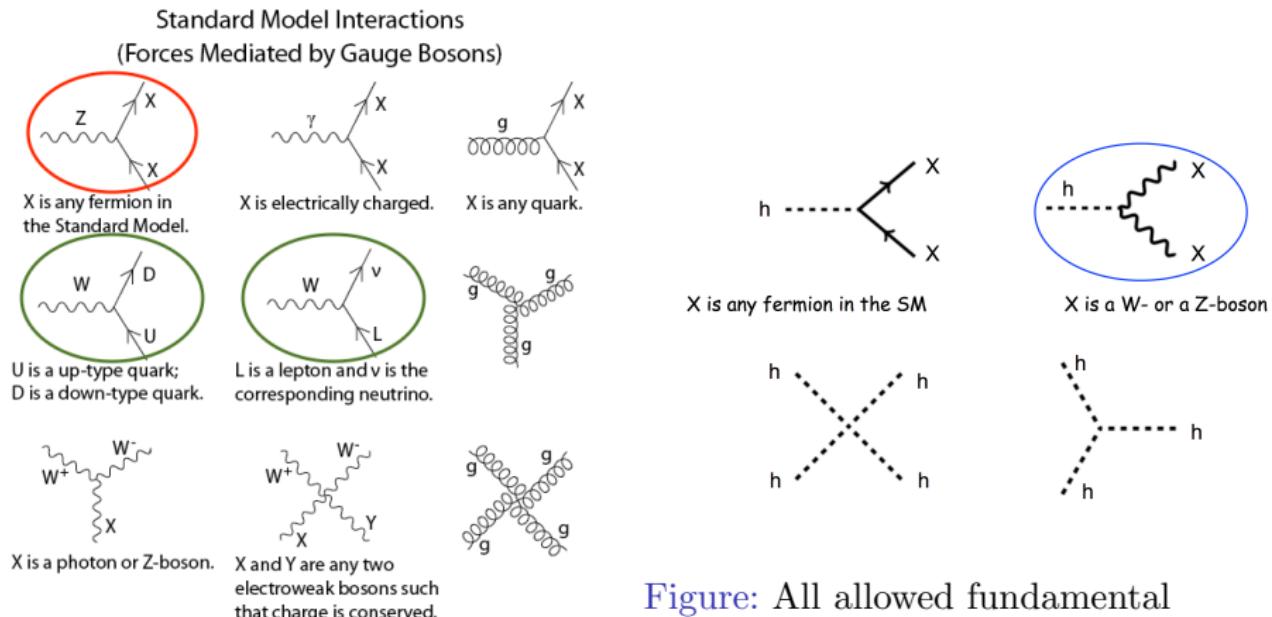


Figure: The “periodic” table for all fundamental particles in SM.

# Standard Model : Fundamental Interaction



**Figure:** All allowed fundamental higgs-related Feynman vertices in SM.

**Figure:** All allowed fundamental Feynman vertices in SM, except higgs-related vertices.

# Standard Model : Success

SM gained a huge success in predicting experimental results.

- ▶ With the W boson explaining the beta decay, the existence of the neutral Z boson was predicted and discovered later in 1980s.
- ▶ The existence of a third generation of quarks was predicted in 1970s and top quark was discovered in 1990s.
- ▶ Higgs boson was proposed in 1960s and discovered recently in 2012.

# Standard Model : Limitation

- ▶ SM cannot explain the nature of dark matter.
- ▶ The hierarchy problem
  - ▶ Why the weak force is stronger than the gravity by  $10^{24}$ .
  - ▶ The Higgs boson mass is quadratically divergent due to the quantum corrections.
- ▶ The electroweak interaction and the strong interaction cannot be unified in SM.

# Supersymmetry

- ▶ Supersymmetry is one of the most promising theory.
- ▶ Supersymmetry(SUSY) is a theoretical extension of the Standard Model.
- ▶ It can solve the hierarchy problem of Higgs mass.
- ▶ It can explain the nature of dark matter.
- ▶ It unifies the electroweak interaction and the strong interaction at  $\sim 10^{16}$  GeV. The gauge couplings change with the energy scale.

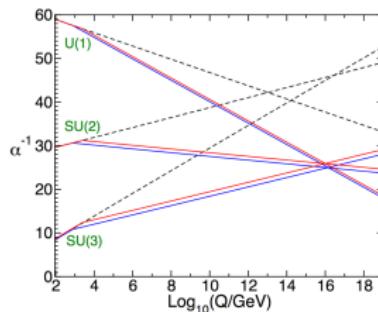
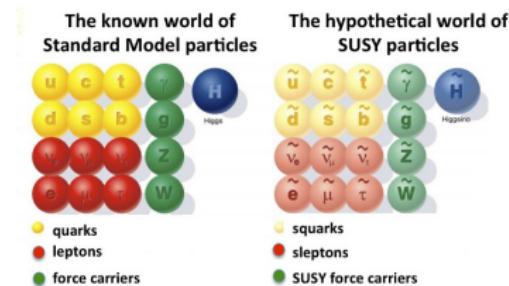


Figure: Black dashed lines are for SM. Red and blue lines are for SUSY.

# Supersymmetry

- ▶ It predicts that each particle in the Standard Model has its own partner particle, called the superpartner (SUSY particle).
- ▶ The spin of the superpartner will differ from the Standard Model particle by  $1/2$ .
- ▶ A symmetry between the fermions and bosons.



**Figure:** The particles in Standard Model and their corresponding superpartners and their names.

# Supersymmetry : Superpartners

- ▶ In the simple SUSY models, we need two neutral Higgs fields ( $H_u^0$  and  $H_d^0$ ) and two charged Higgs fields ( $H_u^+$  and  $H_d^-$ ).
- ▶ The superpartners of 4 neutral bosons  $\gamma$ ,  $Z$ ,  $H_u^0$  and  $H_d^0$  together form 4 mass eigenstates, called neutralinos:  $\tilde{\chi}_1^0$ ,  $\tilde{\chi}_2^0$ ,  $\tilde{\chi}_3^0$  and  $\tilde{\chi}_4^0$ .
- ▶ The superpartners of 4 charged bosons  $W^+$ ,  $W^-$ ,  $H_u^+$  and  $H_d^-$  together form two mass eigenstates with electric charge  $\pm 1$ , called charginos:  $\tilde{\chi}_1^\pm$  and  $\tilde{\chi}_2^\pm$ .

Type	SM particle	Symbol	Spin	R-parity	Superpartner	Symbol	Spin	R-parity
Fermions	Quark	$q$	$\frac{1}{2}$	+1	Squark	$\tilde{q}$	0	-1
	Lepton	$l$	$\frac{1}{2}$	+1	Slepton	$\tilde{l}$	0	-1
Gluon	Gluon	$g$	1	+1	Gluino	$\tilde{g}$	$\frac{1}{2}$	-1
Neutral EW Bosons	Photon	$\gamma$	1	+1	Neutralinos	$\tilde{\chi}_1^0, \tilde{\chi}_2^0, \tilde{\chi}_3^0, \tilde{\chi}_4^0$	$\frac{1}{2}$	-1
	Z Boson	$Z$	1	+1				
	Neutral Higgs	$H_u^0, H_d^0$	0	+1				
Charged EW Bosons	W Boson	$W^+, W^-$	1	+1	Charginos	$\tilde{\chi}_1^\pm, \tilde{\chi}_2^\pm$	$\frac{1}{2}$	-1
	Charged Higgs	$H_u^+, H_d^-$	0	+1				

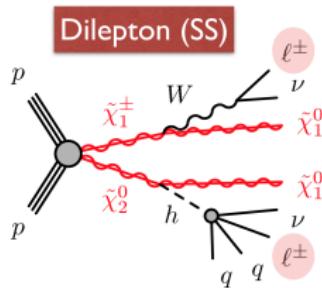
**Table:** The spin and R-parity for the Standard Model particles and their superpartners.

# Supersymmetry : R-parity

- ▶ In order to avoid the proton decay, a quantum number R-parity is introduced. All SM particles have R-parity +1, while all SUSY particles have R-parity -1.
- ▶ If the R-parity is conserved, the lightest SUSY particle (LSP) cannot decay (stable).
- ▶ If the LSP is electrically neutral and interacts with matter only by the weak interaction and gravity, it could be a candidate for dark matter, for example the lightest neutralinos  $\tilde{\chi}_1^0$ .
- ▶ In this thesis, the R-parity is assumed to be conserved, and the lightest neutralino  $\tilde{\chi}_1^0$  is assumed to be the LSP.
- ▶ Due to the conservation of R-parity, the SUSY particles can only be pair-produced, and will eventually decay into SM particles and the lightest neutralino  $\tilde{\chi}_1^0$  (i.e. LSP).

# Our Signal Scenario : Motivation

- ▶ In the recent searches for SUSY particles involved in strong interaction, their masses are suggested to be larger than 1 TeV.
- ▶ In this case, the direct pair production of SUSY electroweak boson  $\tilde{\chi}_1^\pm$  and  $\tilde{\chi}_2^0$  may be the dominant SUSY production process at the LHC, if their masses are below 1 TeV.
- ▶ In this thesis, their masses are assumed to be the same, and denoted by  $m_{\tilde{\chi}_1^\pm, \tilde{\chi}_2^0}$  (in simplified models).



# Our Signal Scenario : Decay Processes

- If all the slepton ( $\tilde{l}$ ) are heavier than  $\tilde{\chi}_1^\pm$  and  $\tilde{\chi}_2^0$  :
  1.  $\tilde{\chi}_1^\pm$  will decay to W boson and  $\tilde{\chi}_1^0$  :  
 $\tilde{\chi}_1^\pm \rightarrow W^\pm + \tilde{\chi}_1^0$
  2.  $\tilde{\chi}_2^0$  will decay to the lightest Higgs boson  $h$  and  $\tilde{\chi}_1^0$  :  
 $\tilde{\chi}_2^0 \rightarrow h + \tilde{\chi}_1^0$

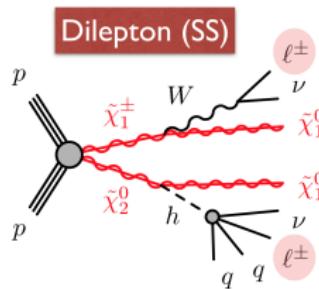
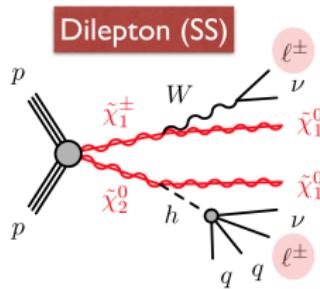


Figure: The Feynman diagram for our Wh same-sign signal scenario.

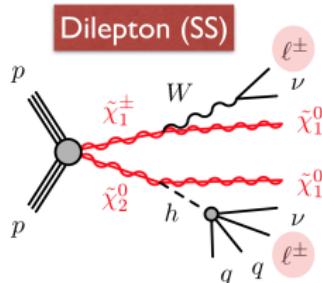
# Our Signal Scenario : Decay Processes

- ▶ The W boson will decay leptonically to one lepton (electron or muon) and one neutrino with the SM branching ratio :  $W^\pm \rightarrow \ell^\pm + \nu$
- ▶ The Higgs boson  $h$  will eventually decay to one lepton (electron or muon), quarks (i.e. jets) and neutrino(s) by various decay modes with the SM branching ratios. (For example,  $h \rightarrow W^+W^-$  and  $h \rightarrow \tau^+\tau^-$ )
- ▶ From now on, a lepton  $\ell^\pm$  only refer to an electron or muon, but not  $\tau$  lepton or neutrino.



# Our Signal Scenario : Signal Signature in the Final State

- ▶ In this thesis, we search for two same-sign(SS) leptons, in order to reduce the SM backgrounds.
- ▶ A large missing energy is expected, due to the undetected neutralinos  $\tilde{\chi}_1^0$  and neutrinos  $\nu$ .
- ▶ Each quark will eventually become a particle shower within a narrow cone, called a jet, by the process of hadronization.
- ▶ The mass difference between the two lightest neutralinos ( $m_{\tilde{\chi}_2^0} - m_{\tilde{\chi}_1^0}$ ) should be larger than the Higgs mass ( $\sim 125$  GeV).



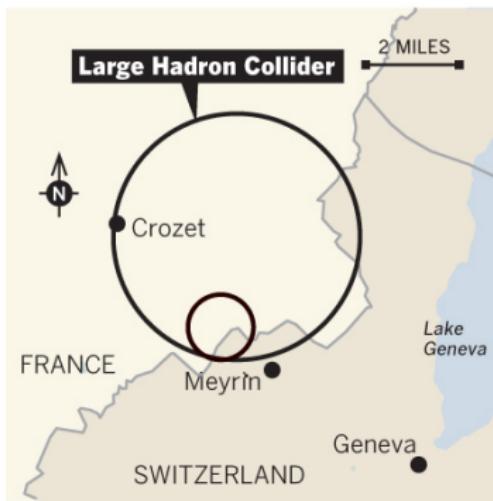
# Our Signal Scenario : Sensitive Region

- ▶ If the mass difference ( $m_{\tilde{\chi}_1^\pm, \tilde{\chi}_2^0} - m_{\tilde{\chi}_1^0}$ ) is slightly larger than the Higgs mass, it is called the compressed region.
- ▶ In the compressed region, one of the lepton may have low energy, due to the low momentum of the Higgs boson, and hence it may not be detected.
- ▶ In this case, only 2 out of 3 leptons are detected.
- ▶ Our search channel would be more sensitive than the 3-leptons channel in the compressed region.

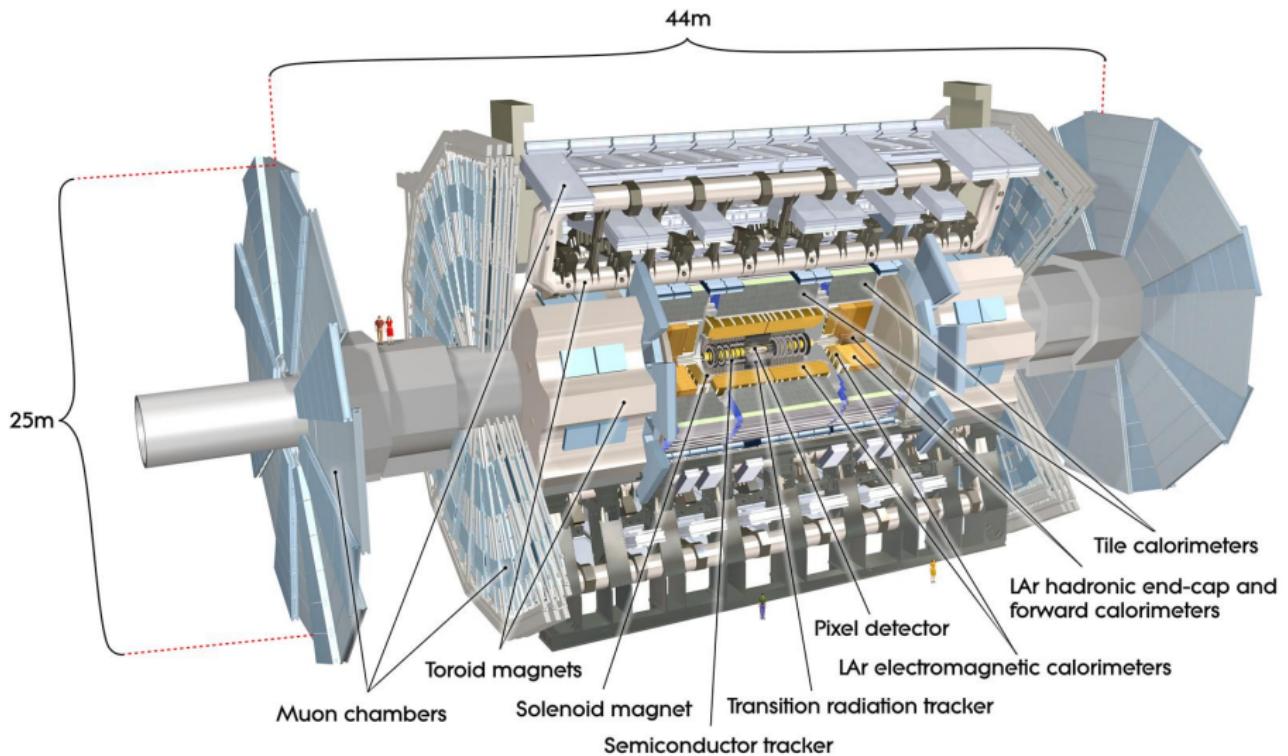
# Experimental Setup

# Experimental Setup : LHC

- ▶ The Large Hadron Collider (LHC) is the most powerful circular particle accelerator in the world.
- ▶ Its circumference is 27 km.
- ▶ Two beams of protons will be accelerated in opposite direction to centre-of-mass energy 13 TeV.
- ▶ They will be finally collided at the ATLAS detector.

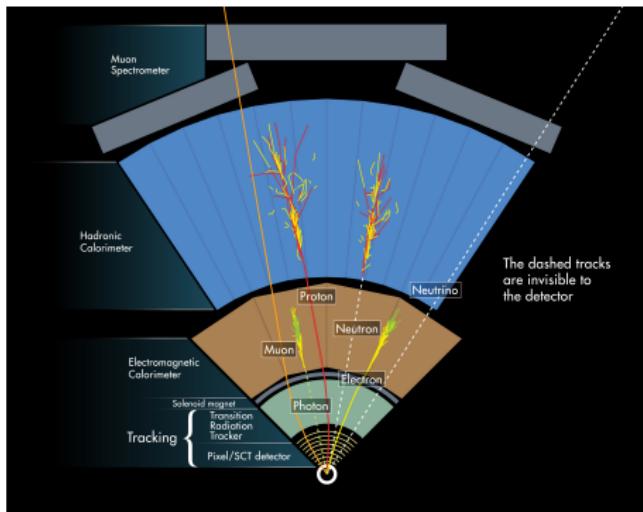


## Experimental Setup : ATLAS detector



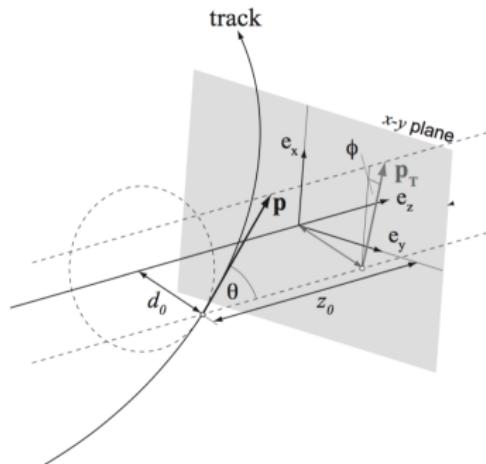
# Experimental Setup : ATLAS detector

- The ATLAS detector consists of 3 main components:
  1. Inner detector: It is a particle tracker. It measures the tracks of charged particles.
  2. Calorimeter: It measures the energy of the particle and stops the particle. It consists of electromagnetic and hadronic calorimeters.
  3. Muon spectrometer: It is a particle tracker for muons.



# Basic kinematic variables

- ▶ The z-axis is in the proton beam direction.
- ▶ The direction of the momentum  $\mathbf{p}$  of the particle can be specified by the azimuthal angle  $\phi$  and the polar angle  $\theta$ , as usual in the spherical coordinate system.

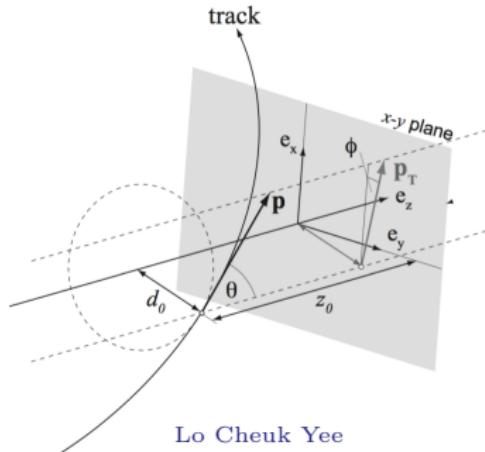


## Basic kinematic variables

- The transverse momentum  $\mathbf{p}_T$  is the projection of  $\mathbf{p}$  onto the x-y plane. The magnitude of  $\mathbf{p}_T$  is denoted by  $p_T$

$$p_T = \sqrt{p_x^2 + p_y^2}$$

- $p_T^1$  and  $p_T^2$  are  $p_T$  of the two leptons, with  $p_T^1 \geq p_T^2$ . The lepton with larger  $p_T$  is called the leading lepton, while the lepton with smaller  $p_T$  is called the sub-leading lepton.



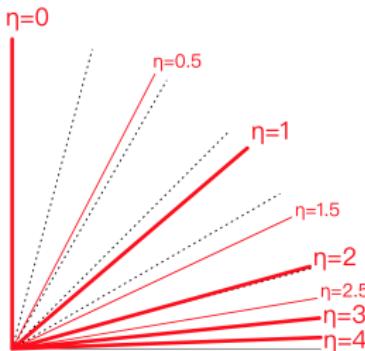
## Basic kinematic variables

- The pseudorapidity  $\eta$  is defined by

$$\eta = -\ln \left( \tan \frac{\theta}{2} \right)$$

- The angle separation  $\Delta R$  of two particles is defined as:

$$\Delta R = \sqrt{(\Delta\phi)^2 + (\Delta\eta)^2} = \sqrt{(\phi_2 - \phi_1)^2 + (\eta_2 - \eta_1)^2}$$



# Challenges of Background Estimation

In order to search for SUSY, we need to understand the SM background very well.

1. Reducible SM background: They are due to the limitation of our detector and the wrong identification of the particle. There are two dominant reducible SM backgrounds, which are estimated by the data-driven method.
  - ▶ Charge flip background
  - ▶ Fake lepton background (dominant background)
2. Irreducible SM background: They are estimated by Monte Carlo (MC)
  - ▶ VV: WZ, WW, ZZ
  - ▶ tt+V
  - ▶ Rare: VVV, multitop, Higgs

# Charge Flip Background

# Charge Flip Background : Sources

- ▶ The charge flip background is due to the mis-identification of the sign of the charge of a lepton (mainly electron) in the reconstruction.
- ▶ The sign of the charge is determined by the direction of the curvature of the track. The charge may be reversed, or flipped.

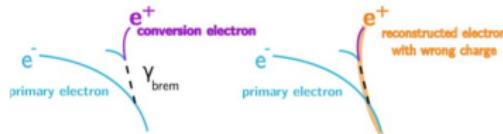


Figure: Bremsstrahlung and  $\gamma$  conversion.

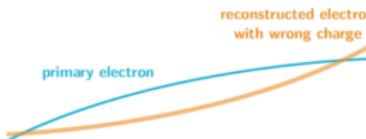


Figure: Track with high velocity  
Lo Cheuk Yee

# Charge Flip Background : Charge Flip Rate

- ▶ The probability that the charge of an electron is flipped is denoted by the charge-flip rate  $\epsilon_i$ , where the index i represents the dependency on the  $p_T$  and  $|\eta|$  of the electron.
- ▶ The value of index i is defined by the index of the following grids in the table.

Variable	Boundary of the bins
$p_T$ (GeV)	25, 60, 90, 130, 150, 1000
$ \eta $	0, 0.50, 1.00, 1.37, 1.52, 1.80, 2.00, 2.47

Table: Binning in  $p_T$  and  $|\eta|$  for the charge-flip rate  $\epsilon_i$ .

- ▶ The probability  $p_{ij}$  that an OS data event becomes a SS data event (with the leading lepton in bin  $i$  and the subleading lepton in bin  $j$ ) is

$$p_{ij} = (1 - \epsilon_i)\epsilon_j + (1 - \epsilon_j)\epsilon_i$$

- ▶ It is called the charge flip weight.

# Charge Flip Background : Charge Flip Rate

- ▶ The charge flip rates are measured by the likelihood method described in the next slide.
- ▶ After getting the charge flip rates, the charge flip background can be estimated by applying the charge flip weights on all OS events in data.

# Charge Flip Background : Likelihood Method

- ▶ The charge flip rate  $\epsilon_i$  is measured by using the data in the control region, within the Z mass window cut of 80-100 GeV.
- ▶ The events inside the Z mass window is then subtracted by the non-Zee processes, by using the sideband technique.
- ▶ After the subtraction, the total number of event is denoted by  $N^{ij}$ , and the number of events with two same-sign leptons is denoted by  $N_{SS}^{ij}$ .
- ▶ The expected value of  $N_{SS}^{ij}$  is then given by  $\lambda = N^{ij} p_{ij}$ . And, if  $N_{SS}^{ij}$  is described by a Poisson distribution, then

$$P(N_{SS}^{ij}|N^{ij}, \epsilon_i, \epsilon_j) = \frac{(N^{ij} p_{ij})^{N_{SS}^{ij}} e^{-N^{ij} p_{ij}}}{N_{SS}^{ij}!}$$

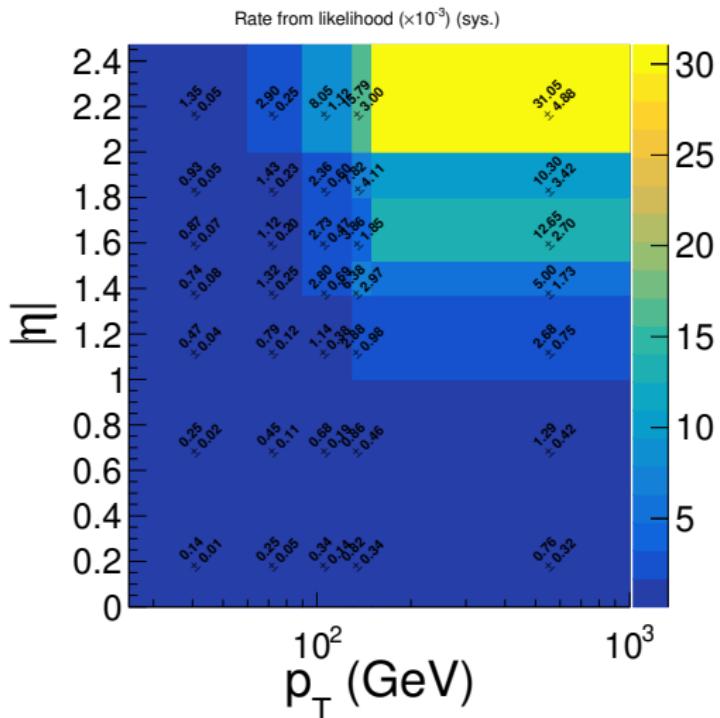
# Charge Flip Background : Likelihood Method

- ▶ Converting it to the likelihood function  $L$  and taking the negative natural log yields

$$\begin{aligned}-\ln L &= -\ln \prod_{ij} P(N_{SS}^{ij} | N^{ij}, \epsilon_i, \epsilon_j) \\&= -\sum_{ij} \left[ N_{SS}^{ij} \ln(N^{ij}[\epsilon_i(1-\epsilon_j) + (1-\epsilon_i)\epsilon_j]) \right. \\&\quad \left. - N^{ij}[\epsilon_i(1-\epsilon_j) + (1-\epsilon_i)\epsilon_j] \right] + \text{constant}\end{aligned}$$

- ▶ where the summation over  $i$  and  $j$  is taken over all  $p_T$  and  $|\eta|$  bins of both electrons.
- ▶ By minimizing this likelihood, the charge-flip rate can be estimated.

# Charge Flip Background : Result for Charge Flip Rate



# Fake Lepton Background

# Fake Lepton Background : Sources

The source of the fake lepton is that other object (like jet and photon) is mis-identified as a lepton. They can be mainly classified into 3 types:

- ▶ Heavy-flavor fakes:
  - ▶ leptons from the semi-leptonic decays of heavy-quark (b or c) hadrons.
- ▶ Light-flavor fakes:
  - ▶ leptons from the semi-leptonic decays of light-quark hadrons.  
OR
  - ▶ It is due to mis-identification of light-quark hadrons.
- ▶ photon conversion:
  - ▶ leptons from the pair production of a photon

# Fake Lepton Background : Matrix Method

- ▶ The fake leptons is estimated by the matrix method.
- ▶ The probability that a real lepton passes the tight selection criteria (i.e. tight lepton) is denoted by the real efficiency  $\epsilon$ .
- ▶ The probability that a real lepton does not pass the tight selection criteria (i.e. loose lepton) is denoted by  $\bar{\epsilon} = 1 - \epsilon$ .
- ▶ The probability that a fake lepton passes the tight selection criteria (i.e. tight lepton) is denoted by the fake efficiency  $f$ .
- ▶ The probability that a fake lepton does not pass the tight selection criteria (i.e. loose lepton) is denoted by  $\bar{f} = 1 - f$ .

$$\begin{pmatrix} N_T \\ N_L \end{pmatrix} = \begin{pmatrix} \epsilon & f \\ \bar{\epsilon} & \bar{f} \end{pmatrix} \begin{pmatrix} N_R \\ N_F \end{pmatrix}$$

# Fake Lepton Background : Matrix Method

Suppose we know  $\begin{pmatrix} N_T \\ N_L \end{pmatrix}$ , we can find  $N_F$  by inverting the matrix.

$$\begin{pmatrix} 0 \\ N_F \end{pmatrix} = \begin{pmatrix} 0 & 0 \\ 0 & 1 \end{pmatrix} \begin{pmatrix} N_R \\ N_F \end{pmatrix} = \begin{pmatrix} 0 & 0 \\ 0 & 1 \end{pmatrix} \begin{pmatrix} \epsilon & f \\ \bar{\epsilon} & \bar{f} \end{pmatrix}^{-1} \begin{pmatrix} N_T \\ N_L \end{pmatrix}$$

Then, we can find the number of tight lepton due to the fake lepton,  $N'_T$

$$\begin{aligned} \begin{pmatrix} N'_T \\ N'_L \end{pmatrix} &= \begin{pmatrix} \epsilon & f \\ \bar{\epsilon} & \bar{f} \end{pmatrix} \begin{pmatrix} 0 \\ N_F \end{pmatrix} \\ &= \begin{pmatrix} \epsilon & f \\ \bar{\epsilon} & \bar{f} \end{pmatrix} \begin{pmatrix} 0 & 0 \\ 0 & 1 \end{pmatrix} \begin{pmatrix} \epsilon & f \\ \bar{\epsilon} & \bar{f} \end{pmatrix}^{-1} \begin{pmatrix} N_T \\ N_L \end{pmatrix} \end{aligned}$$

# Fake Lepton Background : Matrix Method

We can generalize the one-lepton case to the two-leptons case.

$$\begin{pmatrix} N_{TT} \\ N_{TL} \\ N_{LT} \\ N_{LL} \end{pmatrix} = \begin{pmatrix} \epsilon_1 \epsilon_2 & \epsilon_1 f_2 & f_1 \epsilon_2 & f_1 f_2 \\ \epsilon_1 \bar{\epsilon}_2 & \epsilon_1 \bar{f}_2 & f_1 \bar{\epsilon}_2 & f_1 \bar{f}_2 \\ \bar{\epsilon}_1 \epsilon_2 & \bar{\epsilon}_1 f_2 & \bar{f}_1 \epsilon_2 & \bar{f}_1 f_2 \\ \bar{\epsilon}_1 \bar{\epsilon}_2 & \bar{\epsilon}_1 \bar{f}_2 & \bar{f}_1 \bar{\epsilon}_2 & \bar{f}_1 \bar{f}_2 \end{pmatrix} \begin{pmatrix} N_{RR} \\ N_{RF} \\ N_{FR} \\ N_{FF} \end{pmatrix}$$

where  $\epsilon_1$  is the probability that a leading real lepton passes the signal selection (i.e. tight lepton),  $\epsilon_2$  is the probability that a subleading real lepton passes the signal selection (i.e. tight lepton), etc.

# Fake Lepton Background : Matrix Method

Similarly, we can find the number of tight-tight leptons due to the fake leptons,  $N'_{TT}$

$$\begin{pmatrix} N'_{TT} \\ N'_{TL} \\ N'_{LT} \\ N'_{LL} \end{pmatrix} = \begin{pmatrix} \epsilon_1 \epsilon_2 & \epsilon_1 f_2 & f_1 \epsilon_2 & f_1 f_2 \\ \epsilon_1 \bar{\epsilon}_2 & \epsilon_1 \bar{f}_2 & f_1 \bar{\epsilon}_2 & f_1 \bar{f}_2 \\ \bar{\epsilon}_1 \epsilon_2 & \bar{\epsilon}_1 f_2 & \bar{f}_1 \epsilon_2 & \bar{f}_1 f_2 \\ \bar{\epsilon}_1 \bar{\epsilon}_2 & \bar{\epsilon}_1 \bar{f}_2 & \bar{f}_1 \bar{\epsilon}_2 & \bar{f}_1 \bar{f}_2 \end{pmatrix} \begin{pmatrix} 0 & 0 & 0 & 0 \\ 0 & 1 & 0 & 0 \\ 0 & 0 & 1 & 0 \\ 0 & 0 & 0 & 1 \end{pmatrix}$$

$$\begin{pmatrix} \epsilon_1 \epsilon_2 & \epsilon_1 f_2 & f_1 \epsilon_2 & f_1 f_2 \\ \epsilon_1 \bar{\epsilon}_2 & \epsilon_1 \bar{f}_2 & f_1 \bar{\epsilon}_2 & f_1 \bar{f}_2 \\ \bar{\epsilon}_1 \epsilon_2 & \bar{\epsilon}_1 f_2 & \bar{f}_1 \epsilon_2 & \bar{f}_1 f_2 \\ \bar{\epsilon}_1 \bar{\epsilon}_2 & \bar{\epsilon}_1 \bar{f}_2 & \bar{f}_1 \bar{\epsilon}_2 & \bar{f}_1 \bar{f}_2 \end{pmatrix}^{-1} \begin{pmatrix} N_{TT} \\ N_{TL} \\ N_{LT} \\ N_{LL} \end{pmatrix}$$

# Fake Lepton Background : Results of Fake Efficiencies

The method for measuring the fake efficiencies is described in the backup slides.

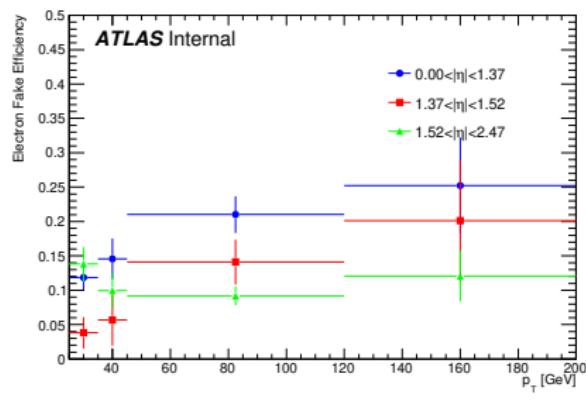


Figure: Electron fake efficiency

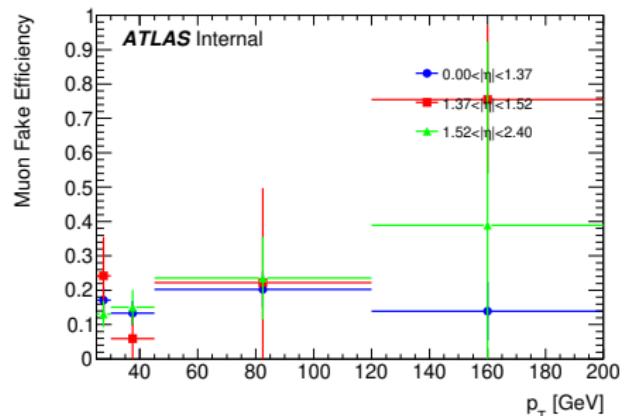


Figure: Muon fake efficiency

# Signal Region

# Signal Region: Challenges

Estimation of the signal sensitivity

- ▶ Cross section for our signal  $\sim 0.1 \text{ pb}$
- ▶ Total luminosity in one year  $\sim 10^4 \text{ pb}^{-1}$
- ▶ Expected number of signal events in one year  $\sim 10^3$
- ▶ Time interval between each collision = 25 ns
- ▶ Total number of collisions in one year  $\sim 10^{15}$
- ▶ The probability to have a signal event  $\sim 10^{-12}$
- ▶ Conclusion: the SUSY signal is very rare.

# Signal Region : Strategy

In order to extract the signal, the following strategies are used.

- ▶ A trigger system are used when taking the data (online), to only store the interested events in the disk.  
(40 MHz → 100 kHz → 1 kHz)
- ▶ Some discriminant variables (based on the kinematic variables) are defined, to help distinguish the signal and the background.
- ▶ Two dedicated signal regions are defined (SRjet1 and SRjet23) based on the discriminant variables, to maximize the signal sensitivity.

# Signal Region : discriminant variable

- ▶  $n_{\text{jets}}$ : Number of signal jets:
- ▶  $n_{b\text{-jets}}$ : Number of  $b$ -jets (jet from  $b$  quark).
- ▶  $\Delta\eta_{ll}$  is designed to estimate the angle between the two leptons.

$$\Delta\eta_{ll} = |\eta_1 - \eta_2|$$

- ▶  $m_{\text{eff}}$  is designed to estimate the total energy and momentum of the charginos  $\tilde{\chi}_1^\pm$  and neutralinos  $\tilde{\chi}_2^0$ .

$$m_{\text{eff}} = p_T^1 + p_T^2 + E_T^{\text{miss}} + \sum_{\text{signal jets}} p_T$$

- ▶  $m_{ll}$ : The invariant mass of the 4-momentum sum of the two leptons

$$(m_{ll})^2 = (p_1 + p_2)^2$$

# Signal Region : discriminant variable

- The missing momentum  $\mathbf{p}^{\text{miss}}$  due to undetected particles is defined by

$$\mathbf{p}^{\text{miss}} = - \sum_{\text{All detected particles}} \mathbf{p}$$

- $E_T^{\text{miss}}$  or MET: The transverse missing energy is designed to estimate the missing energy from the undetected particles.

$$E_T^{\text{miss}} = |\mathbf{p}_T^{\text{miss}}| = \sqrt{(p_x^{\text{miss}})^2 + (p_y^{\text{miss}})^2}$$

# Signal Region : discriminant variable

- ▶  $m_T$ : It is designed to reconstruct the mass of the W boson.

$$m_T = \sqrt{2p_T^1 E_T^{\text{miss}} (1 - \cos \Delta\phi)}$$

where  $\Delta\phi$  is the azimuthal angle between the leading lepton and the missing transverse momentum.

- ▶  $m_{T2}$ : It is designed to set a lower bound on the masses of charginos  $\tilde{\chi}_1^\pm$  and neutralinos  $\tilde{\chi}_2^0$ .

$$m_{T2} = \min_{\mathbf{q}_T} \left[ \max \left( m_T(\mathbf{p}_T^1, \mathbf{q}_T), m_T(\mathbf{p}_T^2, \mathbf{p}_T^{\text{miss}} - \mathbf{q}_T) \right) \right]$$

$$m_T(\mathbf{p}_T, \mathbf{q}_T) = \sqrt{2p_T q_T (1 - \cos \Delta\phi)}$$

# Signal Region : discriminant variable

- ▶  $m_{lj}$  or  $m_{ljj}$ : It is designed to reconstruct the mass of the Higgs boson.

$$p_{\text{jet-system}} = \begin{cases} p_{\text{jet1}} & \text{for SRjet1 (1 jet)} \\ p_{\text{jet1}} + p_{\text{jet2}} & \text{for SRjet23 (2 or 3 jets)} \end{cases}$$

$$p_{\text{closest-lepton}} = \begin{cases} p_{\ell 1} & \text{if } \Delta R(p_{\ell 1}, p_{\text{jet-system}}) \leq \Delta R(p_{\ell 2}, p_{\text{jet-system}}) \\ p_{\ell 2} & \text{if } \Delta R(p_{\ell 1}, p_{\text{jet-system}}) > \Delta R(p_{\ell 2}, p_{\text{jet-system}}) \end{cases}$$

$$(m_{lj(j)})^2 = (p_{\text{closest-lepton}} + p_{\text{jet-system}})^2$$

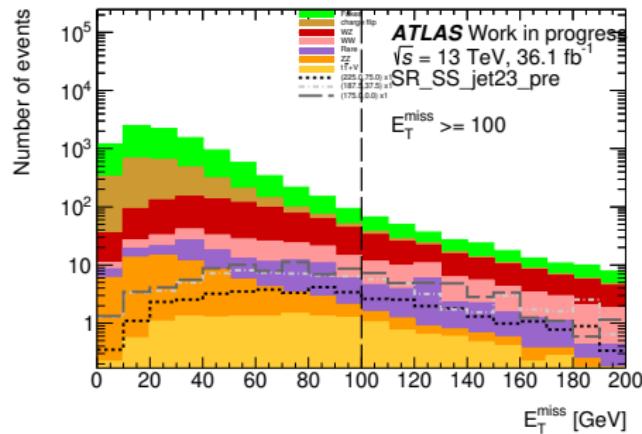
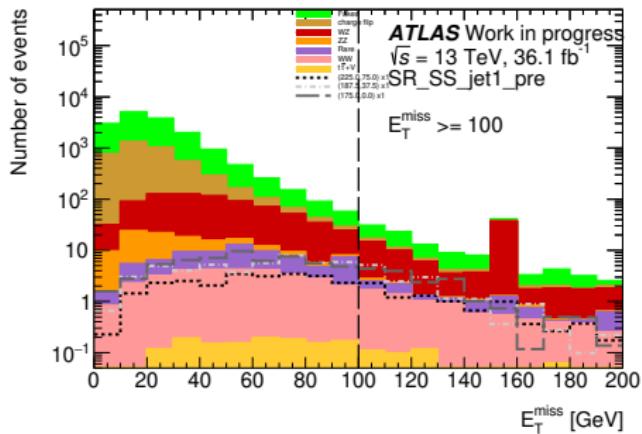
# Signal Region : Pre-selection

- ▶ Signal Region(SR) is a set of cuts of the discriminant variables, such that the signal sensitivity is high.
- ▶ By trying different sets of cuts, the signal region with the highest signal sensitivity can be found. This process is called the signal region optimization.
- ▶ In the SR optimization, the MC signal and background are used to study the expected sensitivity.

The following cuts are the pre-selection, before the signal region optimization.

- ▶ Exactly 2 leptons with  $p_T > 25$  GeV.
- ▶ Electron:  $|\eta| < 2.47$  ; Muon:  $|\eta| < 2.4$
- ▶ The two leptons are same-sign in electric charge (SS).
- ▶  $n_{b\text{-jets}} = 0$ , to reduce the SM background from top quark.
- ▶ Two signal regions: SRjet1 with  $n_{\text{jets}} = 1$ , SRjet23 with  $n_{\text{jets}} = 2$  or 3.

# Signal Region : Pre-selection plots



# Signal Region : Optimized Event Selection

	SRjet1	SRjet23
$n_{jets}$	1	2 or 3
Leading lepton $p_T$ [GeV]	$\geq 25$	$\geq 25$
Sub-leading lepton $p_T$ [GeV]	$\geq 25$	$\geq 25$
$ \Delta\eta_{ll} $	< 1.5	-
$E_T^{\text{miss}}$ [GeV]	$\geq 100$	$\geq 100$
$m_{\text{eff}}$ [GeV]	$\geq 260$	$\geq 240$
$m_T^{l1}$ [GeV]	$\geq 140$	$\geq 120$
$m_{lj}/m_{ljj}$ [GeV]	< 180	< 130
$m_{T2}$ [GeV]	$\geq 80$	$\geq 70$

# Signal Region : Expected Yields and Sensitivity

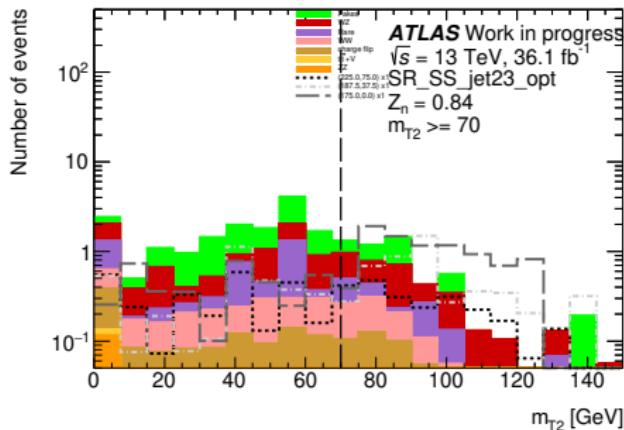
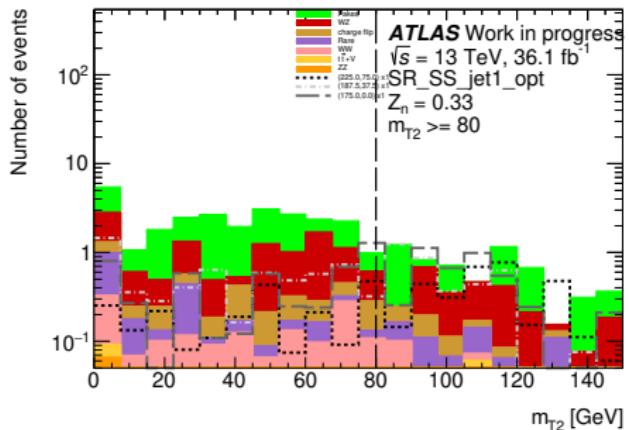
	Number of events	Significance
Fakes	$3.295 \pm 0.819$	
WZ	$2.176 \pm 0.398$	
charge flip	$0.472 \pm 0.053$	
Rare	$0.444 \pm 0.111$	
WW	$0.166 \pm 0.023$	
$t\bar{t} + V$	$0.125 \pm 0.046$	
ZZ	$0.055 \pm 0.028$	
Total BG	$6.733 \pm 0.921$	
(225.0,75.0)	$3.33 \pm 0.60$	<b>0.74</b>
(187.5,37.5)	$3.77 \pm 0.95$	0.86
(175.0,0.0)	$4.29 \pm 0.73$	0.98

Table: Yields in SRjet1

	Number of events	Significance
WZ	$1.849 \pm 0.273$	
Fakes	$1.765 \pm 0.709$	
Rare	$0.731 \pm 0.195$	
WW	$0.514 \pm 0.037$	
charge flip	$0.267 \pm 0.029$	
$t\bar{t} + V$	$0.142 \pm 0.031$	
ZZ	$0.067 \pm 0.025$	
Total BG	$5.335 \pm 0.787$	
(225.0,75.0)	$2.35 \pm 0.34$	0.57
(187.5,37.5)	$4.72 \pm 0.74$	<b>1.26</b>
(175.0,0.0)	$8.60 \pm 1.51$	2.24

Table: Yields in SRjet23

# Signal Region : N-1 plots



# Validation Region

# Validation Region: Introduction

- ▶ The validation region(VR) is designed to ensure that the estimation of background is reliable.
- ▶ Method: The backgrounds are compared with the data, to see whether they agree with each other.
- ▶ One VR is defined for each SR: VRjet1 and VRjet23.

There are three requirements for the validation regions:

- ▶ The signal contribution should be small.
- ▶ The validation region is orthogonal to the corresponding signal region.
- ▶ The background composition is similar to the corresponding signal region.

# Validation Region: Method

- ▶ In order to make the VR orthogonal to its SR,  $E_T^{\text{miss}}/m_T$  cut is inverted.
- ▶ Adjust its lower cut, to have a similar background composition as in its SR.
- ▶ Invert the  $m_{lj}/m_{ljj}$  cut, to reduce the signal.
- ▶ Remove  $m_{T2}$  and  $m_{\text{eff}}$  cuts, to increase statistics.

Cut	VRjet1	VRjet23
$n_{\text{jets}}$	1	[2, 3]
$\Delta\eta_{ll}$	< 1.5	—
$E_T^{\text{miss}}$ [GeV]	[70, 100]	> 100
$m_T$ [GeV]	> 140	[65, 120]
$m_{\text{eff}}$ [GeV]	—	> 240
$m_{lj(j)}$ [GeV]	> 130	> 130
$m_{T2}$ [GeV]	—	—

# Validation Region: Yield

Process	VRjet1	VRjet23
Rare	$0.775 \pm 0.389^{+0.661}_{-0.362}$	$2.469 \pm 0.674^{+0.998}_{-0.899}$
$t\bar{t}V$	$0.039 \pm 0.013^{+0.018}_{-0.012}$	$0.959 \pm 0.082^{+0.152}_{-0.146}$
ZZ	$0.298 \pm 0.060^{+0.089}_{-0.063}$	$0.247 \pm 0.045^{+0.113}_{-0.047}$
WZ	$4.909 \pm 0.530^{+0.960}_{-0.899}$	$19.325 \pm 0.643^{+4.393}_{-4.346}$
WW	$0.801 \pm 0.051^{+0.123}_{-0.060}$	$10.477 \pm 0.176^{+0.796}_{-0.726}$
Charge flip	$1.997 \pm 0.128^{+0.260}_{-0.260}$	$2.065 \pm 0.085^{+0.166}_{-0.166}$
Fakes	$8.021 \pm 1.390^{+5.806}_{-5.806}$	$19.990 \pm 2.013^{+13.461}_{-13.461}$
Total BG	$16.839 \pm 1.545^{+5.915}_{-5.912}$	$55.534 \pm 2.228^{+14.396}_{-14.332}$
Data	17	54

**Table:** The uncertainties include the statistical and systematic uncertainties.

The distribution plots for different variables are in the backup slides.

# Results

# Results: Observed Events

- The observed events are consistent with the SM.

	Number of events
Fakes	$3.295 \pm 2.100$
WZ	$2.176 \pm 0.398$
charge flip	$0.472 \pm 0.053$
Rare	$0.444 \pm 0.111$
WW	$0.166 \pm 0.023$
$t\bar{t} + V$	$0.125 \pm 0.046$
ZZ	$0.055 \pm 0.028$
Total BG	$6.733 \pm 2.142$
Data	2

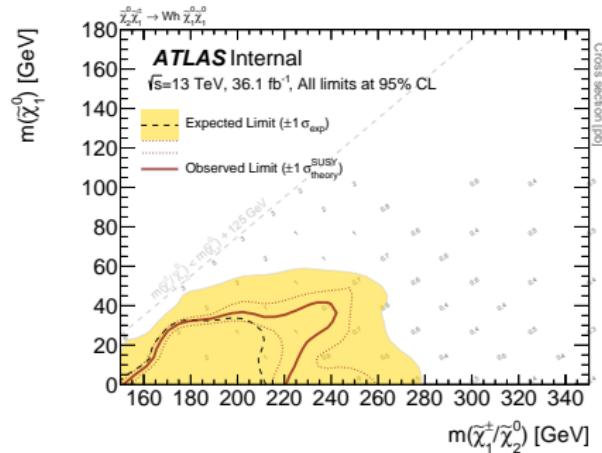
Table: Yields in SRjet1

	Number of events
WZ	$1.849 \pm 0.273$
Fakes	$1.765 \pm 1.460$
Rare	$0.731 \pm 0.195$
WW	$0.514 \pm 0.037$
charge flip	$0.267 \pm 0.029$
$t\bar{t} + V$	$0.142 \pm 0.031$
ZZ	$0.067 \pm 0.025$
Total BG	$5.335 \pm 1.499$
Data	8

Table: Yields in SRjet23

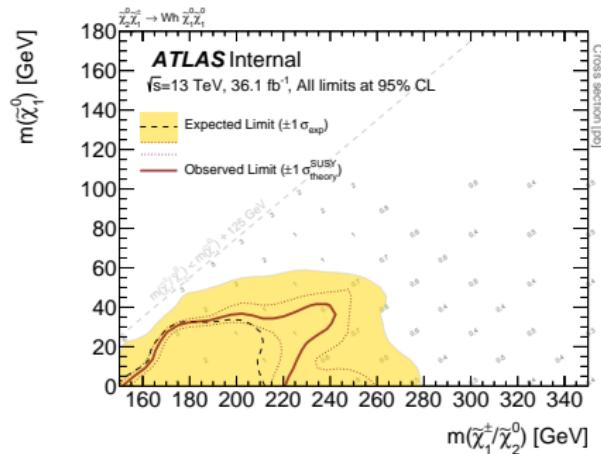
# Results: Exclusion Limits

- ▶ New exclusion limits were set, by combining the two SR.
- ▶ The contour line is in 95% confidence level.



# Results: Exclusion Limits

- ▶ The exclusion limits for the masses of  $\tilde{\chi}_1^\pm$  and  $\tilde{\chi}_2^0$  are extended up to 245 GeV, and the exclusion limits for the mass of  $\tilde{\chi}_1^0$  are extended up to 40 GeV.
- ▶ This means that the mass points inside the contour line are rejected with 95% confidence level.



# Conclusion

- ▶ By using the ATLAS data with total integrated luminosity 36.1  $\text{fb}^{-1}$ , we searched for the electroweak pair production of a chargino and a neutralino ( $p + p \rightarrow \tilde{\chi}_1^\pm + \tilde{\chi}_2^0$ ), with the Wh channel  $\tilde{\chi}_1^\pm \rightarrow \tilde{\chi}_1^0 + W$  and  $\tilde{\chi}_2^0 \rightarrow \tilde{\chi}_1^0 + h$ , and the same-sign channel.
- ▶ The numbers of observed events in data were consistent with the prediction in the SM.
- ▶ The exclusion limits for the masses of  $\tilde{\chi}_1^\pm$  and  $\tilde{\chi}_2^0$  are extended up to 245 GeV, and the exclusion limits for the mass of  $\tilde{\chi}_1^0$  are extended up to 40 GeV, with 95% confidence level.

# Thank You

# Backup

# Supersymmetry : R-parity

- ▶ The baryon number  $B$  is defined by  $\frac{1}{3}(n_q - n_{\bar{q}})$ , where  $n_q$  is the number of quarks and  $n_{\bar{q}}$  is the number of anti-quarks.
- ▶ The lepton number  $L$  is defined by  $n_l - n_{\bar{l}}$ , where  $n_l$  is the number of leptons and  $n_{\bar{l}}$  is the number of anti-leptons.
- ▶ In SM,  $(B - L)$  is conserved. But in SUSY, it is no longer conserved.
- ▶ To keep  $(B - L)$  conservation and prevent the proton decay, the R-parity  $P_R$  is introduced.

$$P_R = (-1)^{3(B-L)-2s}$$

where  $s$  is the spin.

# Charge Flip Background : Likelihood Method

## Likelihood Method : Method

- ▶ The probability that the charge of an electron is flipped is denoted by the charge-flip rate  $\epsilon_i$ , where the index i represents the dependency on the  $p_T$  and  $|\eta|$  of the electron.
- ▶ The value of index i is defined by the index of the following grids in the table.

Variable	Boundary of the bins
$p_T$ (GeV)	25, 60, 90, 130, 150, 1000
$ \eta $	0, 0.50, 1.00, 1.37, 1.52, 1.80, 2.00, 2.47

Table: Binning in  $p_T$  and  $|\eta|$  for the charge-flip rate  $\epsilon_i$ .

- ▶ The probability  $p_{ij}$  that an OS data event becomes a SS data event (with the leading lepton in bin  $i$  and the subleading lepton in bin  $j$ ) is

$$p_{ij} = (1 - \epsilon_i)\epsilon_j + (1 - \epsilon_j)\epsilon_i$$

# Likelihood Method : Definition of Control Region

A control region is defined.

- ▶ Trigger requirement
- ▶ exactly 2 signal leptons
- ▶ 2 electrons
- ▶  $25 < p_T < 1000$
- ▶  $|\eta| < 2.47$

# Likelihood Method : Method

- ▶ The charge flip rate  $\epsilon_i$  is measured by using the data in the control region, within the Z mass window cut of 80-100 GeV.
- ▶ The events inside the Z mass window is then subtracted by the non-Zee processes, by using the sideband technique.
- ▶ After the subtraction, the total number of event is denoted by  $N^{ij}$ , and the number of events with two same-sign leptons is denoted by  $N_{SS}^{ij}$ .
- ▶ The expected value of  $N_{SS}^{ij}$  is then given by  $\lambda = N^{ij} p_{ij}$ . And, if  $N_{SS}^{ij}$  is described by a Poisson distribution, then

$$P(N_{SS}^{ij}|N^{ij}, \epsilon_i, \epsilon_j) = \frac{(N^{ij} p_{ij})^{N_{SS}^{ij}} e^{-N^{ij} p_{ij}}}{N_{SS}^{ij}!}$$

# Likelihood Method : Method

- ▶ Converting it to the likelihood function  $L$  and taking the negative natural log yields

$$\begin{aligned} -\ln L &= -\ln \prod_{ij} P(N_{SS}^{ij} | N^{ij}, \epsilon_i, \epsilon_j) \\ &= -\sum_{ij} \left[ N_{SS}^{ij} \ln(N^{ij}[\epsilon_i(1-\epsilon_j) + (1-\epsilon_i)\epsilon_j]) \right. \\ &\quad \left. - N^{ij}[\epsilon_i(1-\epsilon_j) + (1-\epsilon_i)\epsilon_j] \right] + \text{constant} \end{aligned}$$

- ▶ where the summation over  $i$  and  $j$  is taken over all  $p_T$  and  $|\eta|$  bins of both electrons.
- ▶ By minimizing this likelihood, the charge-flip rate can be estimated.

# Likelihood Method : Background Subtraction

Both for data and MC

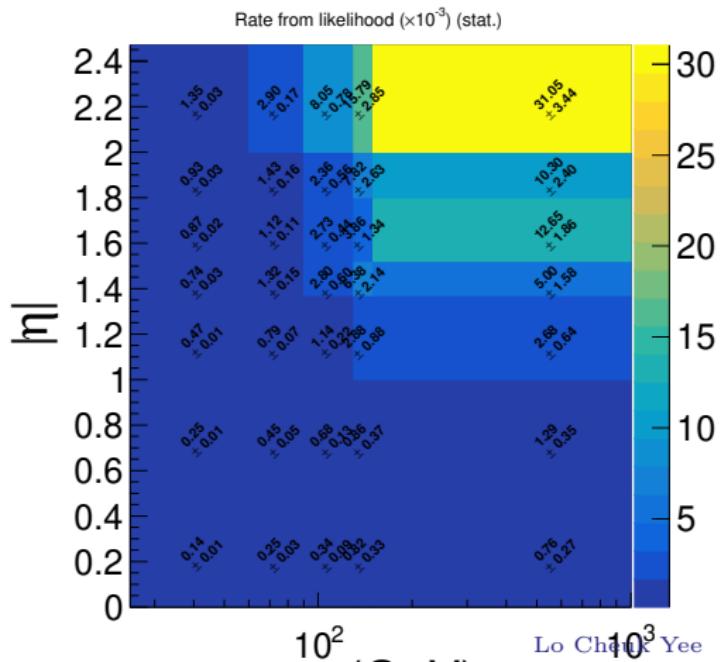
- ▶ Central region: 80 - 100 GeV
- ▶ Sideband region: 60 - 80 and 100 - 120  
(Sideband width 20 GeV)

$$N_{Central} = N_{Central} - \frac{20}{20 + 20} N_{SB}$$

- ▶ Background subtraction is done both on  $N^{ij}$  and  $N_{SS}^{ij}$

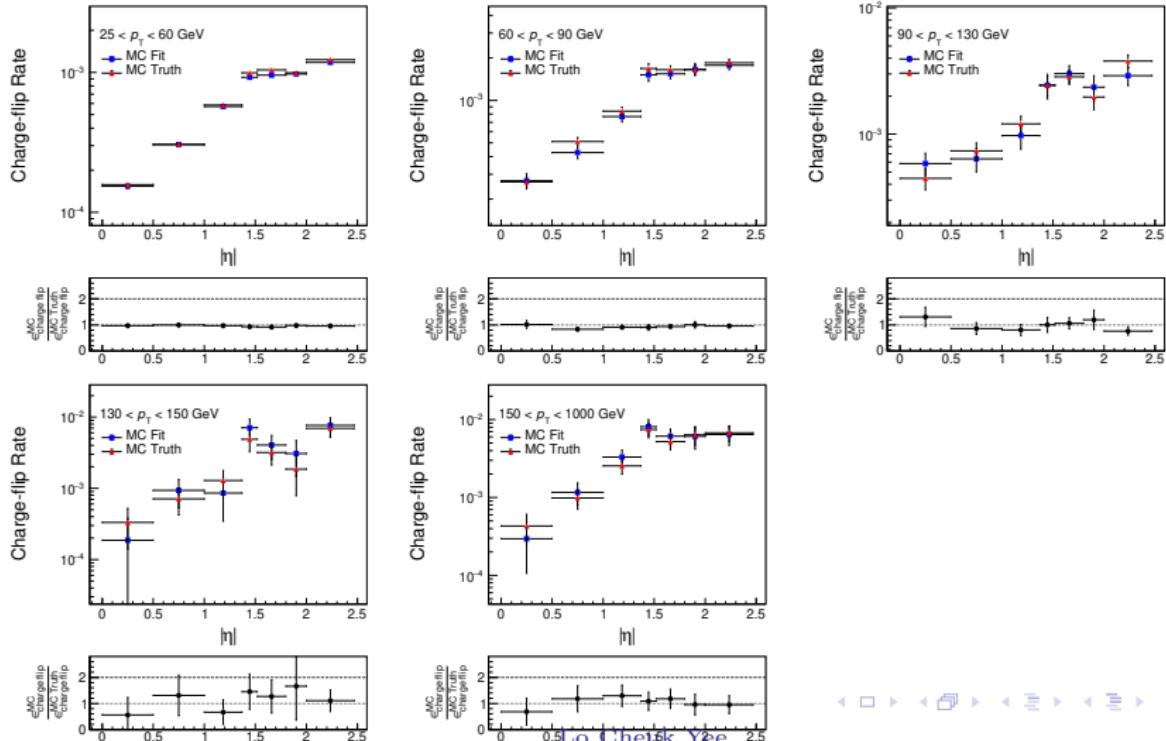
# Likelihood Method : Result for Charge Flip Rate (statistical error only)

- ▶ Results by likelihood method from data,  $\epsilon_{\text{lik,data}}$
- ▶ The error is statistical only, from likelihood method,  $\sigma_{\text{lik,data}}$ .



# Likelihood Method : MC Truth Validation

- ▶ Results by likelihood method from MC (blue points)
- ▶ Results by MC truth (red points)

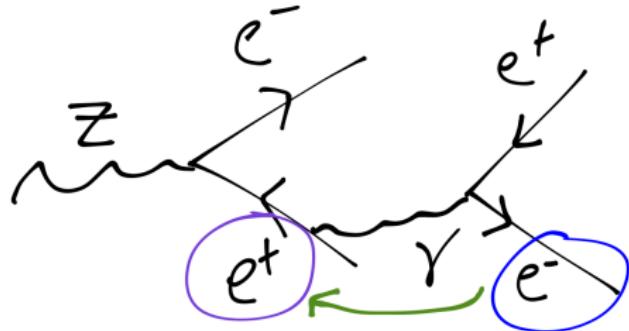


# Likelihood Method : MC Truth Matching

- ▶ It is a complicated part, because we do not know the original charge of the reconstructed electron for some cases.
- ▶ The reconstructed electrons match the truth particle by the smallest  $\Delta R$  ( $\Delta R < 0.1$ ).
- ▶ No match for reconstructed electrons if no any truth particle is inside  $\Delta R < 0.1$ . (cannot find the original charge)
- ▶ If the truth particle is matched, the truth particle is not electron. (fake lepton, not charge flip)

# Likelihood Method : MC Truth Matching

- ▶ If the truth particle is electron, the (grand)mother particle is not Z Boson. (ignore it)
- ▶ If (grand)mother particle is Z Boson, the daughter of the Z Boson is not electron. (ignore it)
- ▶ If the daughter of the Z Boson is electron, the charge of that electron is the original charge.



# Likelihood Method : Systematic uncertainties due to likelihood method

For each bin, the systematic uncertainties due to likelihood method are estimated by the difference between the likelihood method and the MC truth method.

For MC,

$$\sigma_{\text{truth,MC}} = |\epsilon_{\text{lik,MC}} - \epsilon_{\text{MC truth}}|$$

For data,

$$\sigma_{\text{truth,data}} = \epsilon_{\text{lik,data}} \times \frac{\sigma_{\text{truth,MC}}}{\epsilon_{\text{lik,MC}}}$$

# Likelihood Method : Systematic uncertainties due to background subtraction

Nominal:

- ▶ Central region: 80 - 100 GeV, Sideband width: 20 GeV

Variations:

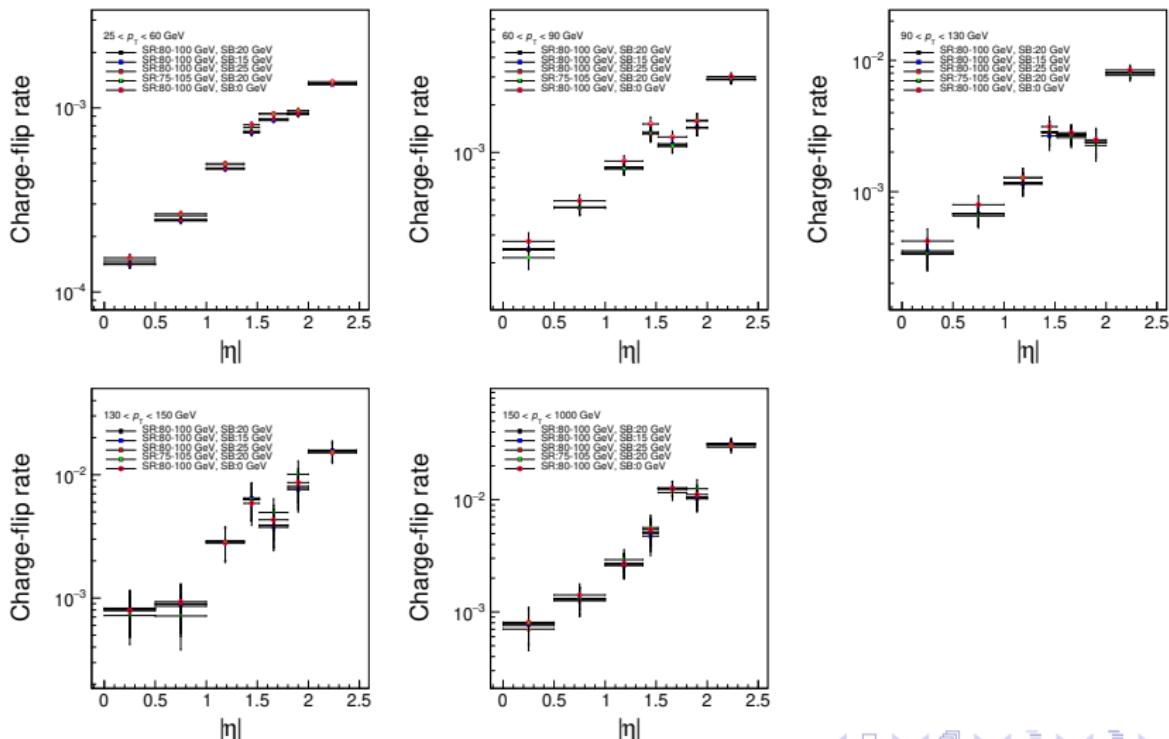
- ▶ Central region: 80 - 100 GeV, Sideband width: 15 GeV
- ▶ Central region: 80 - 100 GeV, Sideband width: 25 GeV
- ▶ Central region: 75 - 105 GeV, Sideband width: 20 GeV
- ▶ Central region: 80 - 100 GeV, no background subtraction

For each bin, the largest deviation from the nominal among these variations is the systematic uncertainty due to background subtraction.

$$\sigma_{\text{bgk}} = \max\{|\sigma_{\text{nominal}} - \sigma_{\text{variation}}|\}$$

# Likelihood Method : Result for different variation

## ► Results for different variations

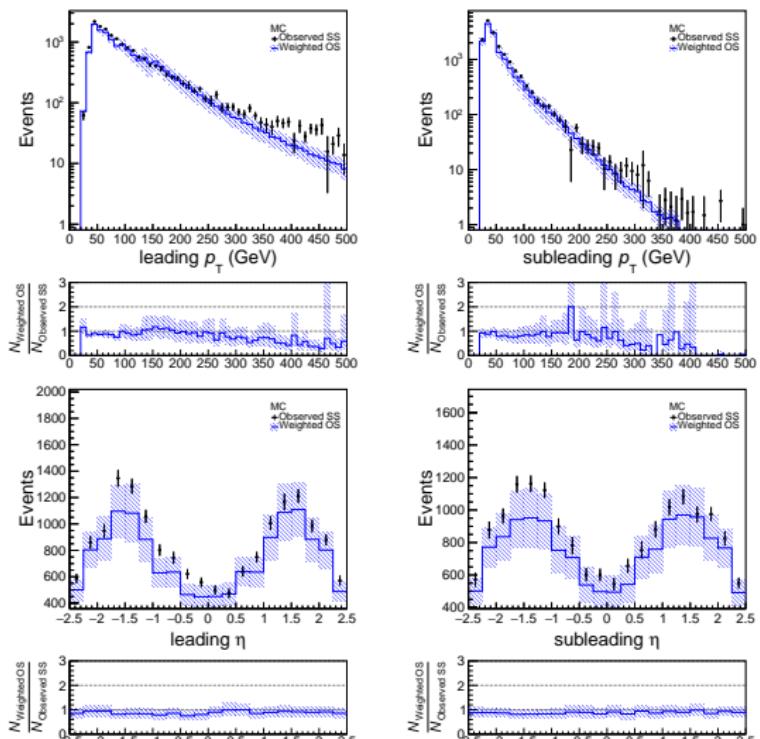


# Likelihood Method : Total uncertainties

$$\sigma_{\text{sys}} = \sqrt{\sigma_{\text{truth}}^2 + \sigma_{\text{bgk}}^2}$$

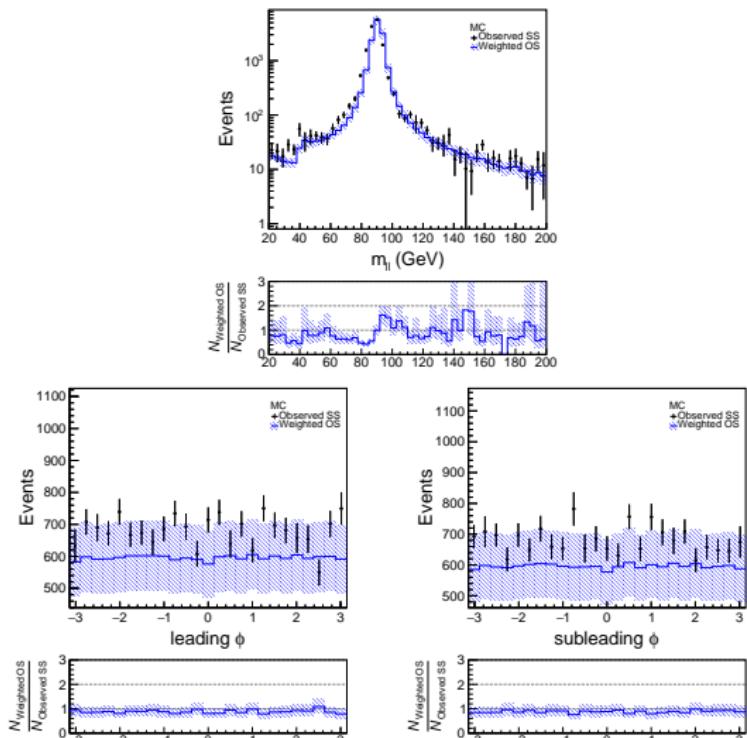
$$\sigma_{\text{tot}} = \sqrt{\sigma_{\text{lik}}^2 + \sigma_{\text{sys}}^2}$$

# Likelihood Method : MC Validation Plots



**Figure:** The comparison between the weighted OS events and the SS events, with different variables.

# Likelihood Method : MC Validation Plots



**Figure:** The comparison between the weighted OS events and the SS events, with different variables.

# Fake Lepton Background : Matrix Method

# Fake Lepton Background : Matrix Method

We can generalize the one-lepton case to the two-leptons case.

$$\begin{pmatrix} N_{TT} \\ N_{TL} \\ N_{LT} \\ N_{LL} \end{pmatrix} = \begin{pmatrix} \epsilon_1 \epsilon_2 & \epsilon_1 f_2 & f_1 \epsilon_2 & f_1 f_2 \\ \epsilon_1 \bar{\epsilon}_2 & \epsilon_1 \bar{f}_2 & f_1 \bar{\epsilon}_2 & f_1 \bar{f}_2 \\ \bar{\epsilon}_1 \epsilon_2 & \bar{\epsilon}_1 f_2 & \bar{f}_1 \epsilon_2 & \bar{f}_1 f_2 \\ \bar{\epsilon}_1 \bar{\epsilon}_2 & \bar{\epsilon}_1 \bar{f}_2 & \bar{f}_1 \bar{\epsilon}_2 & \bar{f}_1 \bar{f}_2 \end{pmatrix} \begin{pmatrix} N_{RR} \\ N_{RF} \\ N_{FR} \\ N_{FF} \end{pmatrix}$$

where  $\epsilon_1$  is the probability that a leading real lepton passes the signal selection (i.e. tight lepton),  $\epsilon_2$  is the probability that a subleading real lepton passes the signal selection (i.e. tight lepton), etc.

# Fake Lepton Background : Matrix Method

We can find  $N_{RF}$ ,  $N_{FR}$ ,  $N_{FF}$  by inverting the matrix.

$$\begin{pmatrix} 0 \\ N_{RF} \\ N_{FR} \\ N_{FF} \end{pmatrix} = \begin{pmatrix} 0 & 0 & 0 & 0 \\ 0 & 1 & 0 & 0 \\ 0 & 0 & 1 & 0 \\ 0 & 0 & 0 & 1 \end{pmatrix} \begin{pmatrix} N_{RR} \\ N_{RF} \\ N_{FR} \\ N_{FF} \end{pmatrix}$$

$$= \begin{pmatrix} 0 & 0 & 0 & 0 \\ 0 & 1 & 0 & 0 \\ 0 & 0 & 1 & 0 \\ 0 & 0 & 0 & 1 \end{pmatrix} \begin{pmatrix} \epsilon_1\epsilon_2 & \epsilon_1f_2 & f_1\epsilon_2 & f_1f_2 \\ \epsilon_1\bar{\epsilon}_2 & \epsilon_1\bar{f}_2 & f_1\bar{\epsilon}_2 & f_1\bar{f}_2 \\ \bar{\epsilon}_1\epsilon_2 & \bar{\epsilon}_1f_2 & \bar{f}_1\epsilon_2 & \bar{f}_1f_2 \\ \bar{\epsilon}_1\bar{\epsilon}_2 & \bar{\epsilon}_1\bar{f}_2 & \bar{f}_1\bar{\epsilon}_2 & \bar{f}_1\bar{f}_2 \end{pmatrix}^{-1} \begin{pmatrix} N_{TT} \\ N_{TL} \\ N_{LT} \\ N_{LL} \end{pmatrix}$$

# Fake Lepton Background : Matrix Method

Then, we can find the number of tight-tight leptons due to the fake leptons,  $N'_{TT}$

$$\begin{pmatrix} N'_{TT} \\ N'_{TL} \\ N'_{LT} \\ N'_{LL} \end{pmatrix} = \begin{pmatrix} \epsilon_1\epsilon_2 & \epsilon_1f_2 & f_1\epsilon_2 & f_1f_2 \\ \epsilon_1\bar{\epsilon}_2 & \epsilon_1\bar{f}_2 & f_1\bar{\epsilon}_2 & f_1\bar{f}_2 \\ \bar{\epsilon}_1\epsilon_2 & \bar{\epsilon}_1f_2 & \bar{f}_1\epsilon_2 & \bar{f}_1f_2 \\ \bar{\epsilon}_1\bar{\epsilon}_2 & \bar{\epsilon}_1\bar{f}_2 & \bar{f}_1\bar{\epsilon}_2 & \bar{f}_1\bar{f}_2 \end{pmatrix} \begin{pmatrix} 0 \\ N_{RF} \\ N_{FR} \\ N_{FF} \end{pmatrix}$$

$$= \begin{pmatrix} \epsilon_1\epsilon_2 & \epsilon_1f_2 & f_1\epsilon_2 & f_1f_2 \\ \epsilon_1\bar{\epsilon}_2 & \epsilon_1\bar{f}_2 & f_1\bar{\epsilon}_2 & f_1\bar{f}_2 \\ \bar{\epsilon}_1\epsilon_2 & \bar{\epsilon}_1f_2 & \bar{f}_1\epsilon_2 & \bar{f}_1f_2 \\ \bar{\epsilon}_1\bar{\epsilon}_2 & \bar{\epsilon}_1\bar{f}_2 & \bar{f}_1\bar{\epsilon}_2 & \bar{f}_1\bar{f}_2 \end{pmatrix} \begin{pmatrix} 0 & 0 & 0 & 0 \\ 0 & 1 & 0 & 0 \\ 0 & 0 & 1 & 0 \\ 0 & 0 & 0 & 1 \end{pmatrix}$$

$$\begin{pmatrix} \epsilon_1\epsilon_2 & \epsilon_1f_2 & f_1\epsilon_2 & f_1f_2 \\ \epsilon_1\bar{\epsilon}_2 & \epsilon_1\bar{f}_2 & f_1\bar{\epsilon}_2 & f_1\bar{f}_2 \\ \bar{\epsilon}_1\epsilon_2 & \bar{\epsilon}_1f_2 & \bar{f}_1\epsilon_2 & \bar{f}_1f_2 \\ \bar{\epsilon}_1\bar{\epsilon}_2 & \bar{\epsilon}_1\bar{f}_2 & \bar{f}_1\bar{\epsilon}_2 & \bar{f}_1\bar{f}_2 \end{pmatrix}^{-1} \begin{pmatrix} N_{TT} \\ N_{TL} \\ N_{LT} \\ N_{LL} \end{pmatrix}$$

# Method for measuring real efficiencies

# Control Region for Real Efficiencies

It uses the Z tag-and-probe method to measure the real efficiencies of electrons and muons.

Pre-selection:

- ▶ Trigger requirement
- ▶ exactly 2 baseline leptons
- ▶ same flavour and opposite charge ( $Z \rightarrow ee$ ,  $Z \rightarrow \mu\mu$ )
- ▶  $80 < m_{ll} < 100$
- ▶  $p_T > 20$
- ▶  $|\eta| < 2.47$  for electron,  $|\eta| < 2.5$  for muon

Requirements for tag lepton

- ▶  $p_T > 25$
- ▶ pass signal requirement

# Permutation of tag and probe

Check the probe lepton:

Tight lepton: pass signal requirement

Loose lepton: pass baseline requirement (include tight lepton)

4 possible cases: (slightly simplified)

- ▶ TL: fill loose hist by the subleading lepton
- ▶ LT: fill loose hist by the leading lepton
- ▶ TT: fill both loose and tight hist by the both leading and subleading lepton
- ▶ LL: do nothing

The real eff is calculated by

$$\epsilon = \frac{N_{\text{tight}}}{N_{\text{loose}}}$$

# Fake Lepton Background : Results of Real Efficiencies

The method for measuring the real efficiencies is described in the backup slides.

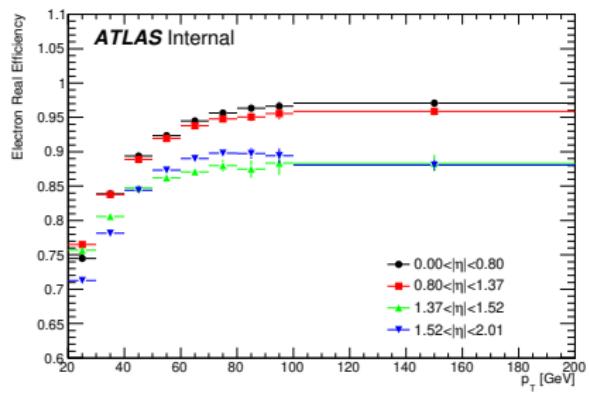


Figure: Electron real efficiency

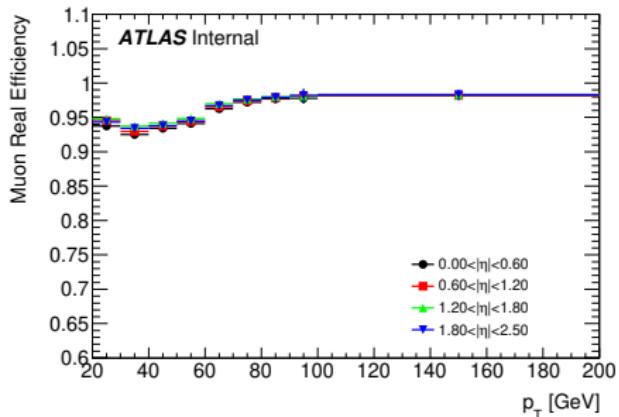


Figure: Muon real efficiency

# Method for measuring fake efficiencies

## Control Region for fake eff for muon

It uses the heavy-flavour fake-enriched region to measure the fake efficiencies of electrons and muons.

For muon fake eff, it uses the muon-muon events.

Pre-selection for muon-muon events:

- ▶ Trigger requirement
- ▶ exactly 2 baseline muons
- ▶ SS
- ▶ at least 1 jet
- ▶ at least 1 b-jet
- ▶  $p_T > 20$
- ▶  $|\eta| < 2.5$

Requirements for the tag lepton

- ▶  $p_T > 40$
- ▶ pass signal requirement

# Control Region for fake eff for electron

For electron fake eff, it uses the muon-electron events.

Pre-selection for muon-electron events:

- ▶ Trigger requirement
- ▶ exactly 1 baseline electron and 1 baseline muon
- ▶ SS
- ▶ at least 1 jet
- ▶ at least 1 b-jet
- ▶  $p_T > 20$
- ▶  $|\eta| < 2.5$

Requirements for the tag lepton

- ▶ muon
- ▶  $p_T > 40$
- ▶ pass signal requirement

# Calculate the fake eff

The fake eff is calculated by

$$\epsilon = \frac{N_{\text{tight}}^{\text{data}} - N_{\text{tight}}^{\text{prompt MC}}}{N_{\text{loose}}^{\text{data}} - N_{\text{loose}}^{\text{prompt MC}}}$$

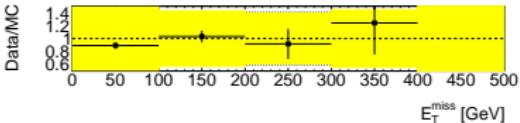
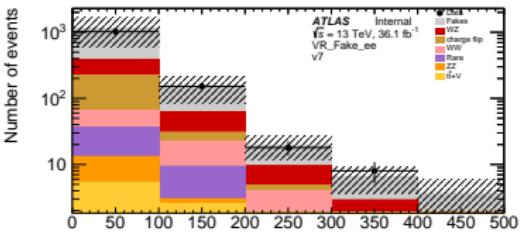
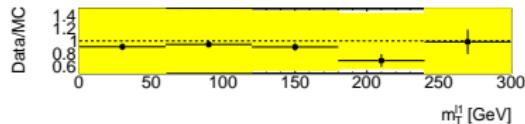
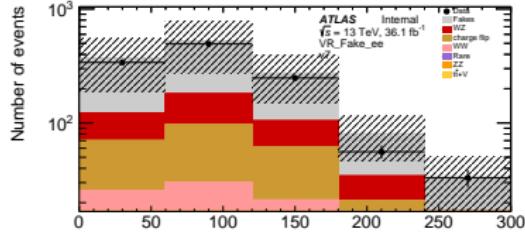
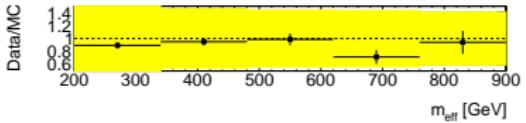
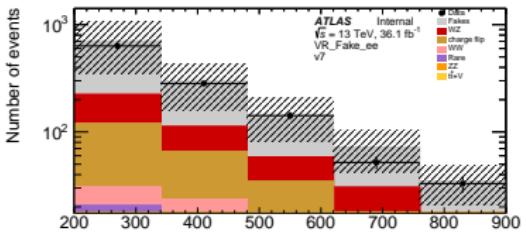
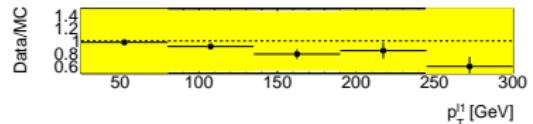
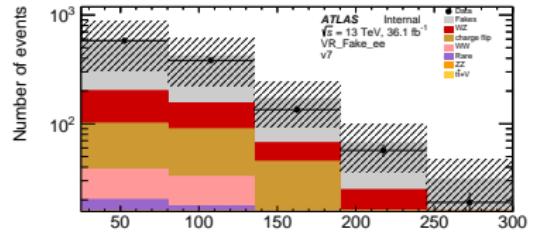
The prompt lepton is estimated from MC by using the truth information. The prompt lepton is defined by the MCTruthClassifier, with the ParticleType = IsoElectron or IsoMuon.

# Validation plots for fake lepton background

## ee channel

	Number of events	
Fakes	$860.738787 \pm 14.359660 \pm 630.863451$ (9663)	
WZ	$201.069385 \pm 3.094538$ (39857)	
charge flip	$167.954322 \pm 1.110819$ (86972)	
WW	$45.176809 \pm 0.401964$ (21140)	
Rare	$30.757392 \pm 3.403960$ (1268)	
ZZ	$8.351345 \pm 0.345416$ (6466)	
$t\bar{t} + V$	$8.346022 \pm 0.254296$ (4158)	
Total BG	$1322.394062 \pm 15.130843 \pm 630.863451$ (169524)	
Data	1201	

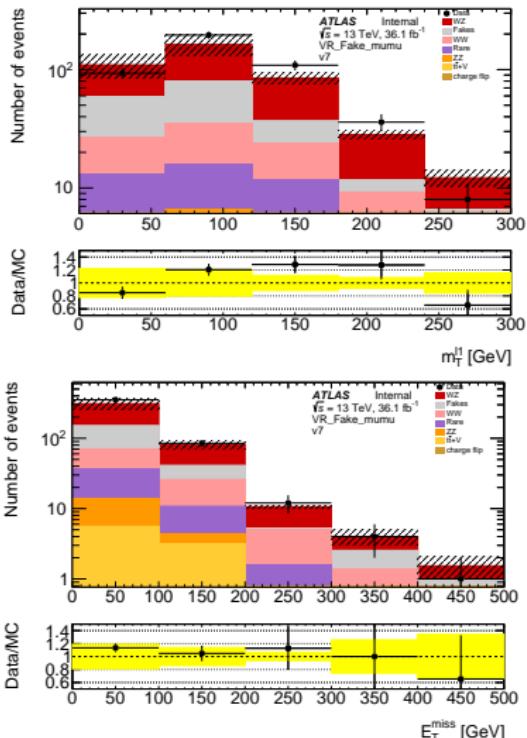
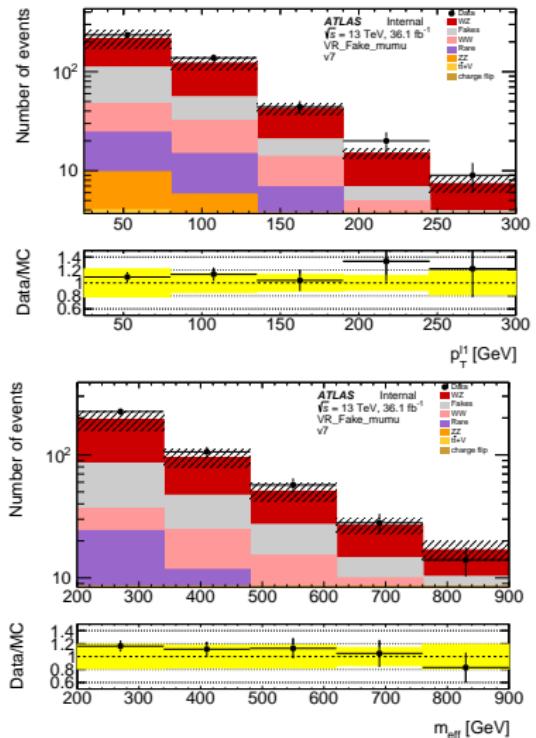
## ee channel



# mumu channel

	Number of events	
WZ	$209.219829 \pm 2.994737$ (42458)	
Fakes	$96.203916 \pm 5.013796 \pm 70.510979$ (1160)	
WW	$53.675707 \pm 0.409851$ (25092)	
Rare	$30.460547 \pm 4.359663$ (1361)	
ZZ	$9.474511 \pm 0.534342$ (7674)	
$t\bar{t} + V$	$9.369165 \pm 0.284944$ (4801)	
charge flip	$0.000000 \pm 0.000000$ (0)	
Total BG	$408.403674 \pm 7.324477 \pm 70.510979$ (82546)	
Data	454	

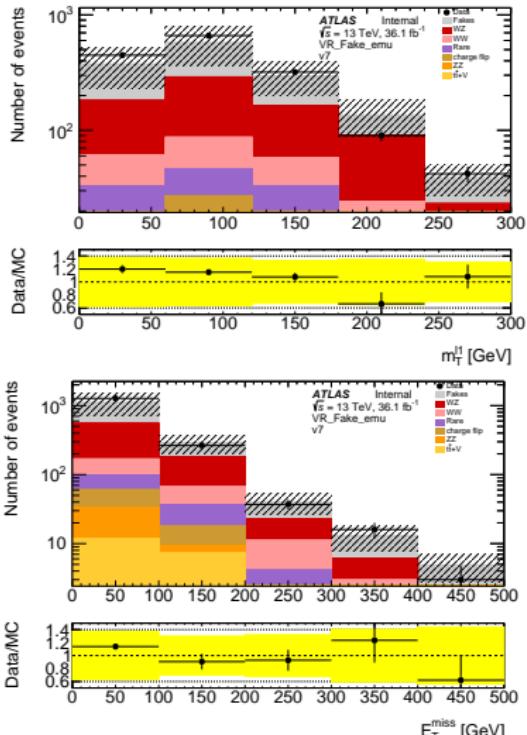
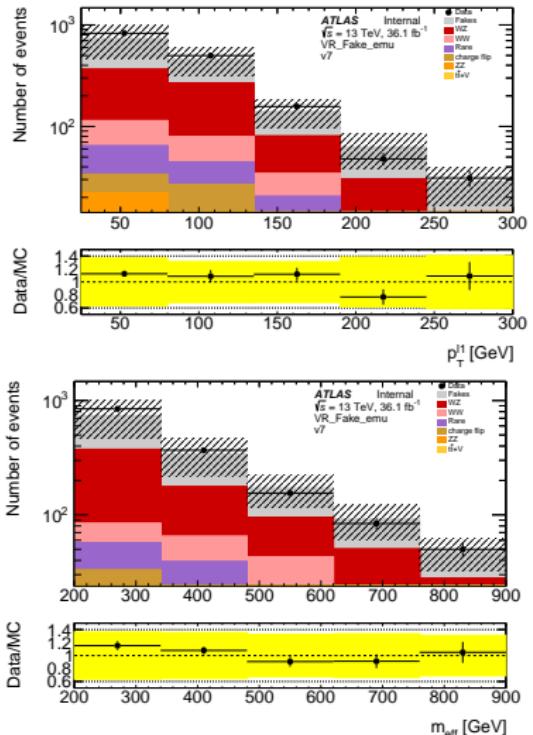
# mumu channel



## emu channel

	Number of events	
Fakes	$702.995427 \pm 13.784203 \pm 515.248212$ (8544)	
WZ	$516.828375 \pm 33.696304$ (94571)	
WW	$110.280002 \pm 0.578575$ (51615)	
Rare	$58.742111 \pm 4.005608$ (3037)	
charge flip	$36.901490 \pm 0.479301$ (36998)	
ZZ	$23.391749 \pm 0.482387$ (16698)	
$t\bar{t} + V$	$20.739252 \pm 0.424107$ (10526)	
Total BG	$1469.878407 \pm 36.639692 \pm 515.248212$ (221989)	
Data	1593	

# emu channel



# Pre-selection plots

# Yields in pre-selection

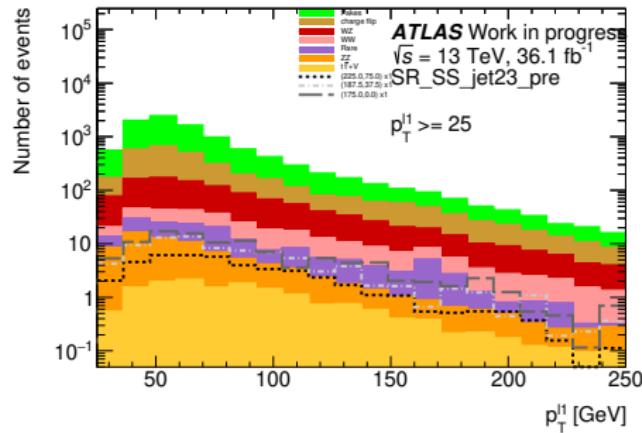
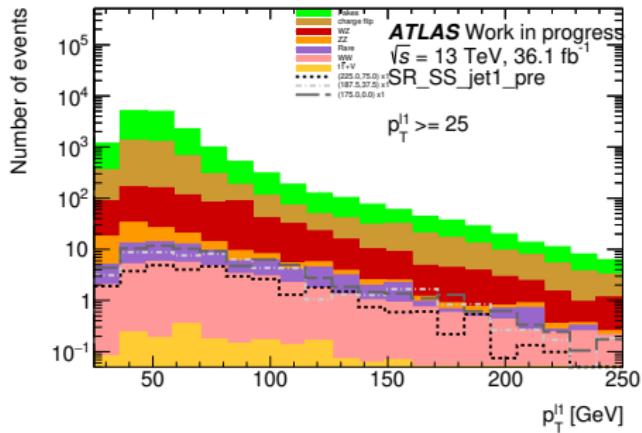
Number of events	
Fakes	$11580.965417 \pm 45.775452$ (107219)
charge flip	$3622.756937 \pm 2.810481$ (3375964)
WZ	$691.333566 \pm 34.238864$ (68626)
ZZ	$66.945860 \pm 1.049840$ (21749)
Rare	$47.371873 \pm 4.082859$ (1662)
WW	$38.128141 \pm 0.388221$ (15666)
$t\bar{t} + V$	$1.971267 \pm 0.138488$ (671)
Total BG	$16049.473060 \pm 57.389305$ (3591557)
(225.0,75.0)	$33.08 \pm 1.57$ (676)
(187.5,37.5)	$62.09 \pm 3.13$ (568)
(175.0,0.0)	$74.30 \pm 3.57$ (556)

Table: Yields in SRjet1

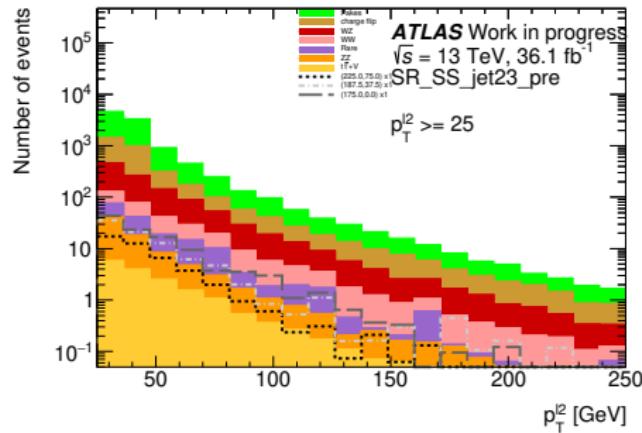
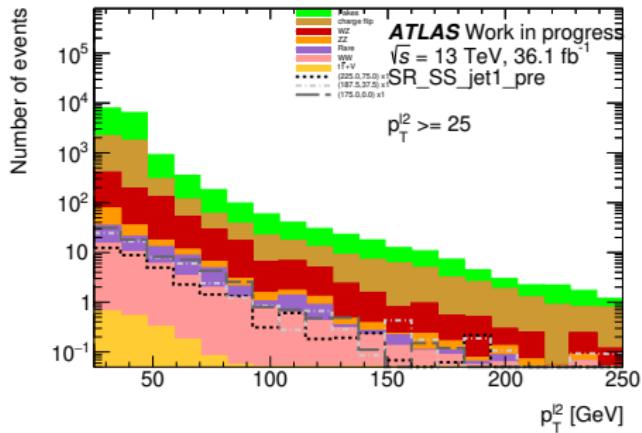
Number of events	
Fakes	$6848.284248 \pm 36.074557$ (66096)
charge flip	$2074.427691 \pm 2.680118$ (1679673)
WZ	$776.282836 \pm 5.978666$ (147413)
WW	$153.455294 \pm 0.679752$ (72175)
Rare	$92.181843 \pm 6.338615$ (3011)
ZZ	$62.932966 \pm 1.198123$ (35945)
$t\bar{t} + V$	$16.375388 \pm 0.359456$ (6082)
Total BG	$10023.940266 \pm 37.235816$ (2010395)
(225.0,75.0)	$45.01 \pm 1.76$ (937)
(187.5,37.5)	$86.02 \pm 3.84$ (754)
(175.0,0.0)	$107.96 \pm 5.14$ (743)

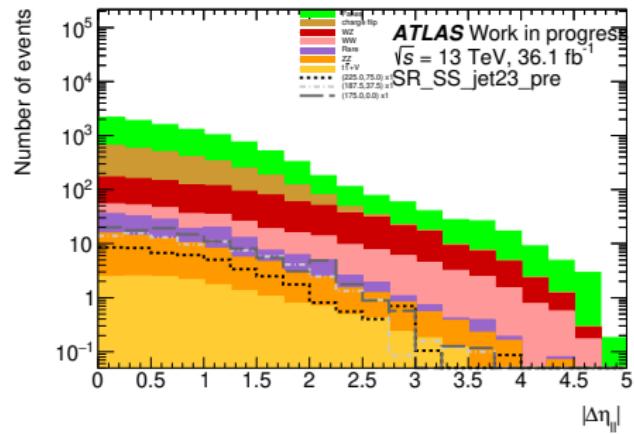
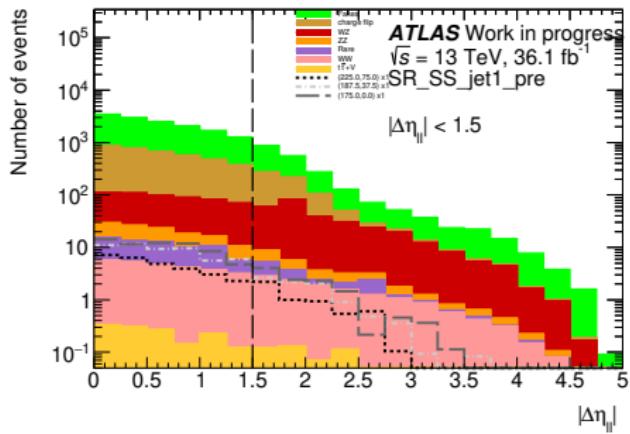
Table: Yields in SRjet23

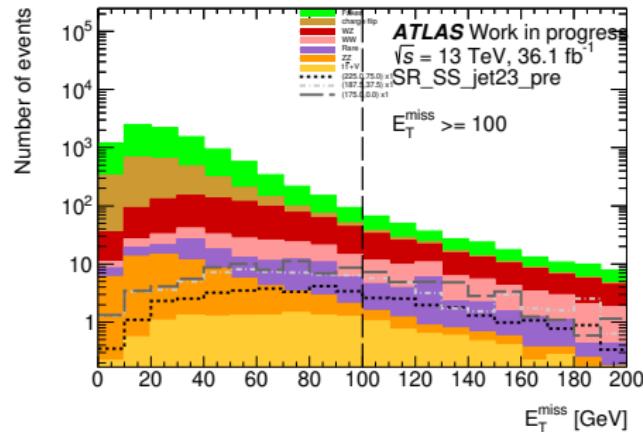
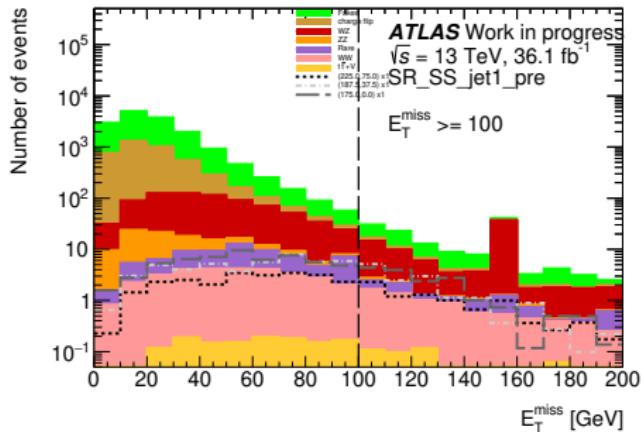
$$p_T^{l1}$$



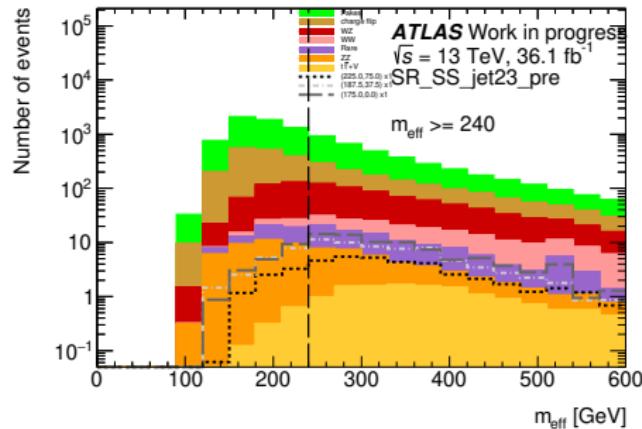
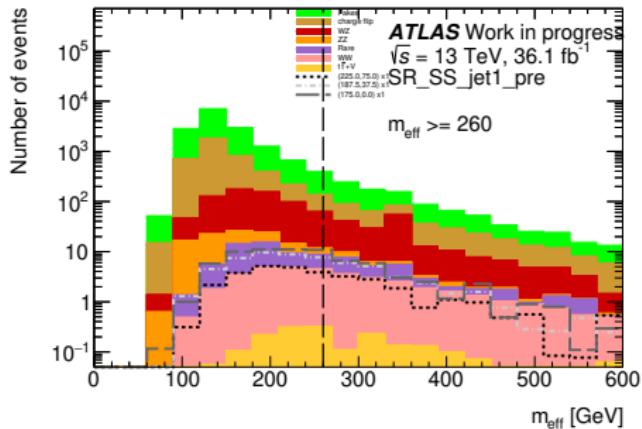
p<sub>T</sub><sup>l2</sup>



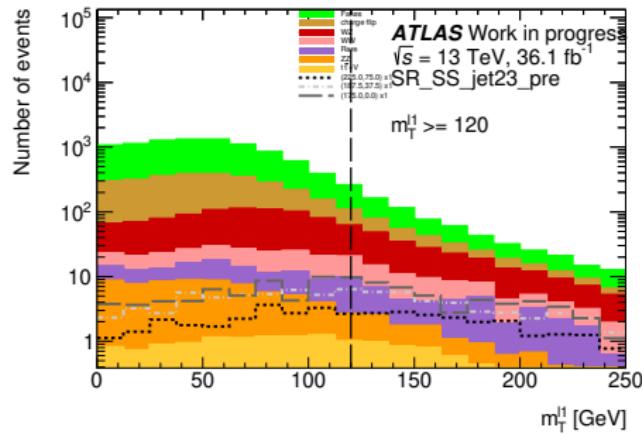
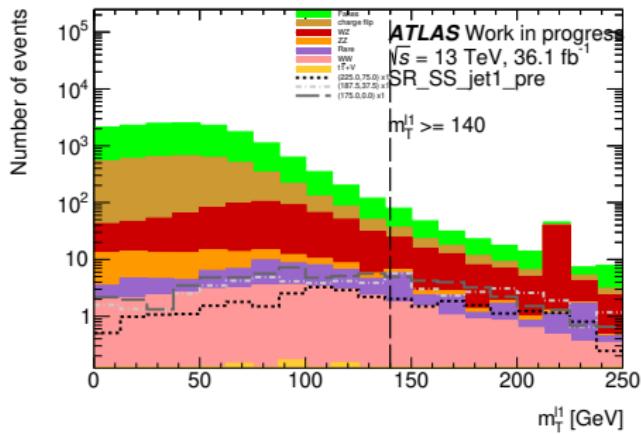
$|\Delta\eta_{ll}|$ 

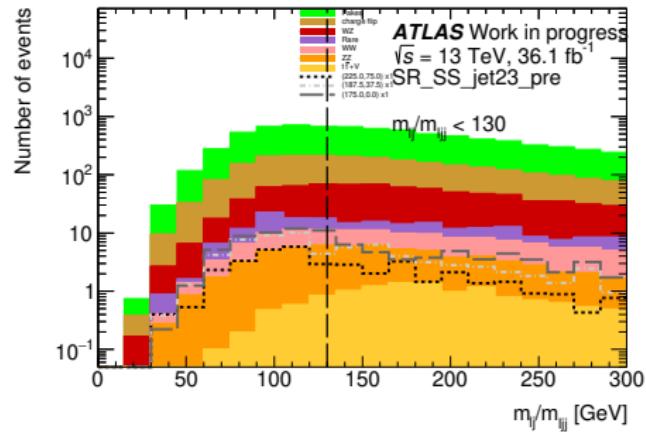
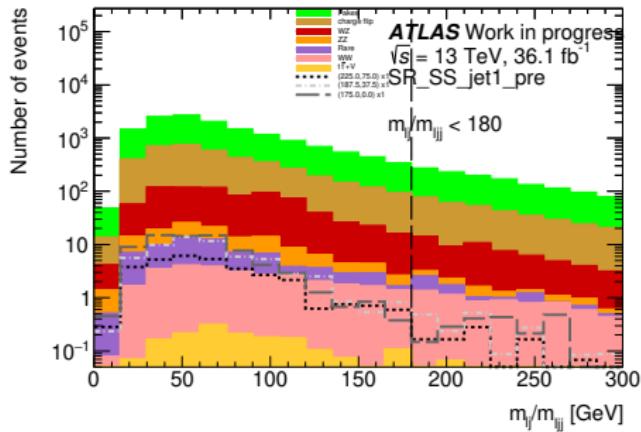
$E_T^{\text{miss}}$ 

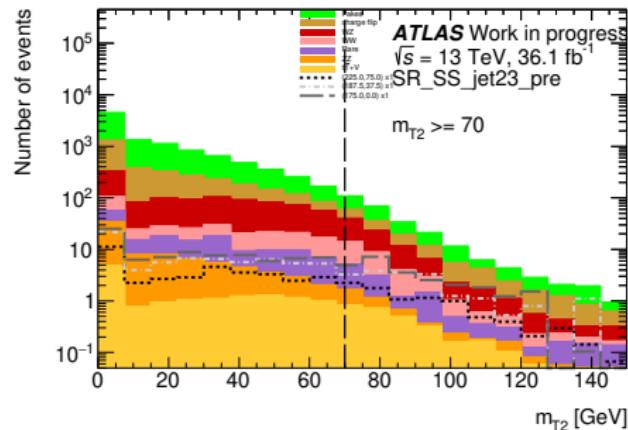
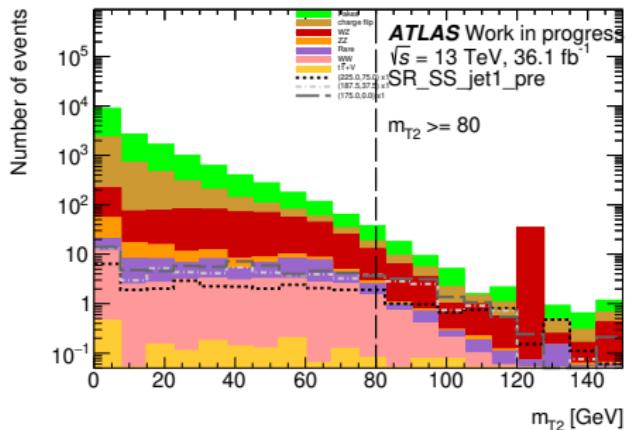
$m_{\text{eff}}$



$$m_{\mathrm{T}}^{l_1}$$



$m_{lj}/m_{ljj}$ 


$m_{T2}$ 

# N-1 plots

# Yields in SR

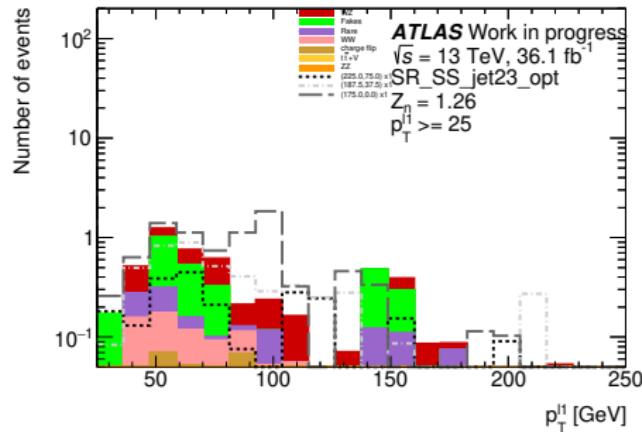
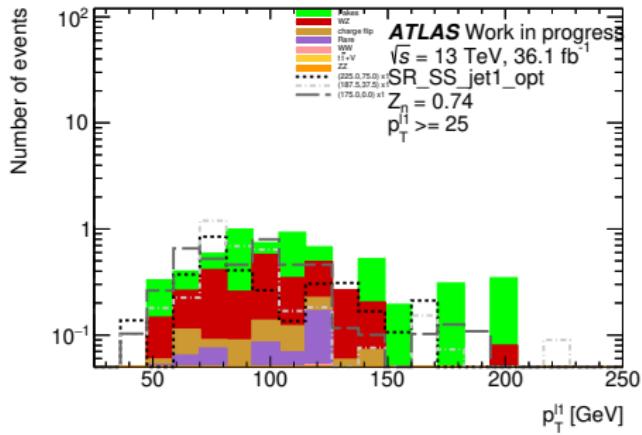
	Number of events	Significance
Fakes	$3.295434 \pm 0.819175$ (24)	
WZ	$2.176292 \pm 0.397871$ (257)	
charge flip	$0.472185 \pm 0.052737$ (265)	
Rare	$0.443536 \pm 0.111301$ (56)	
WW	$0.165504 \pm 0.023455$ (67)	
$t\bar{t} + V$	$0.124898 \pm 0.046227$ (36)	
ZZ	$0.055330 \pm 0.027713$ (31)	
Total BG	$6.733178 \pm 0.920855$ (736)	
(225.0,75.0)	$3.33 \pm 0.60$ (60)	<b>0.74</b>
(187.5,37.5)	$3.77 \pm 0.95$ (31)	0.86
(175.0,0.0)	$4.29 \pm 0.73$ (36)	0.98

Table: Yields in SRjet1

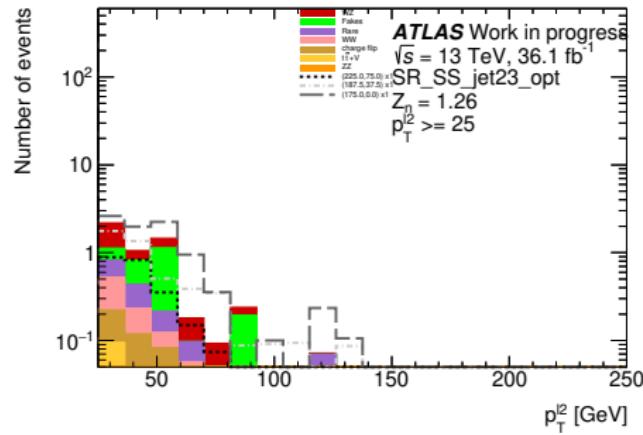
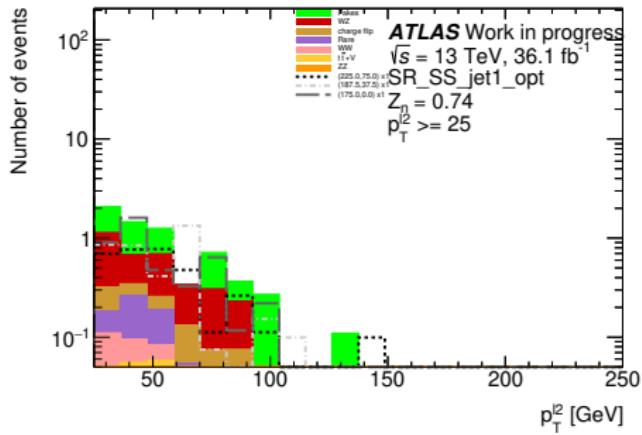
	Number of events	Significance
WZ	$1.848838 \pm 0.273025$ (319)	
Fakes	$1.764693 \pm 0.709215$ (20)	
Rare	$0.731015 \pm 0.195114$ (48)	
WW	$0.514488 \pm 0.036878$ (235)	
charge flip	$0.267485 \pm 0.029086$ (274)	
$t\bar{t} + V$	$0.141679 \pm 0.030772$ (67)	
ZZ	$0.067226 \pm 0.024995$ (24)	
Total BG	$5.335425 \pm 0.787005$ (987)	
(225.0,75.0)	$2.35 \pm 0.34$ (58)	0.57
(187.5,37.5)	$4.72 \pm 0.74$ (47)	<b>1.26</b>
(175.0,0.0)	$8.60 \pm 1.51$ (58)	2.24

Table: Yields in SRjet23

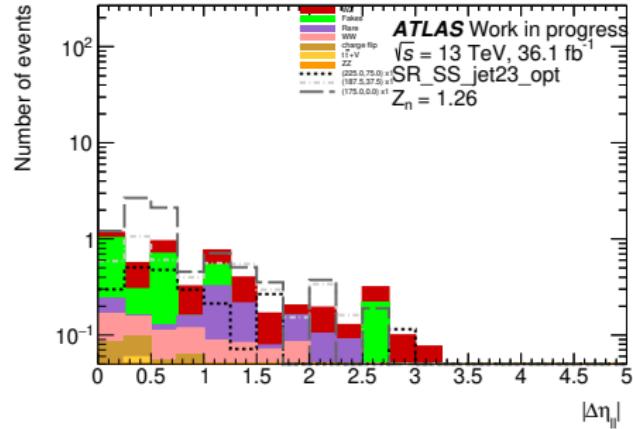
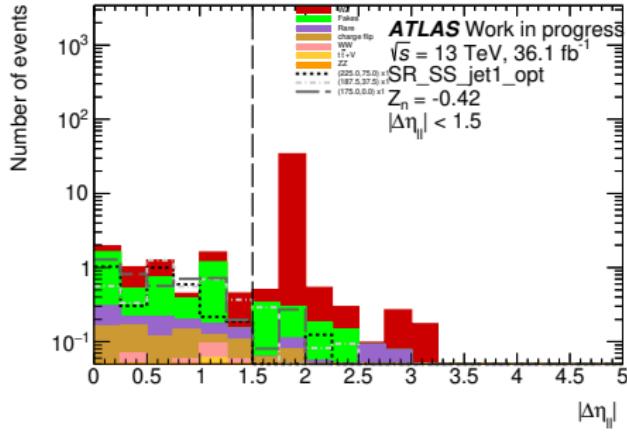
$$p_T^{l1}$$



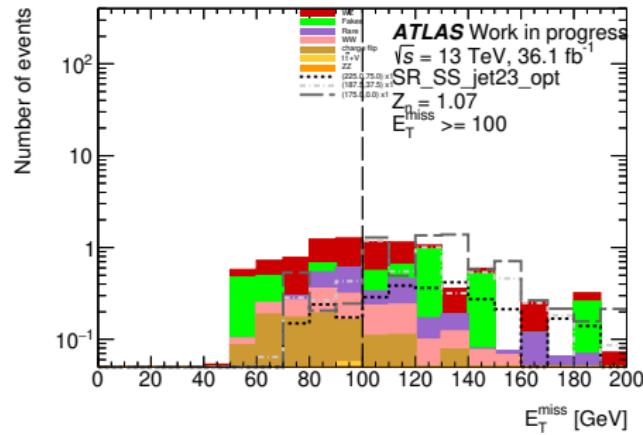
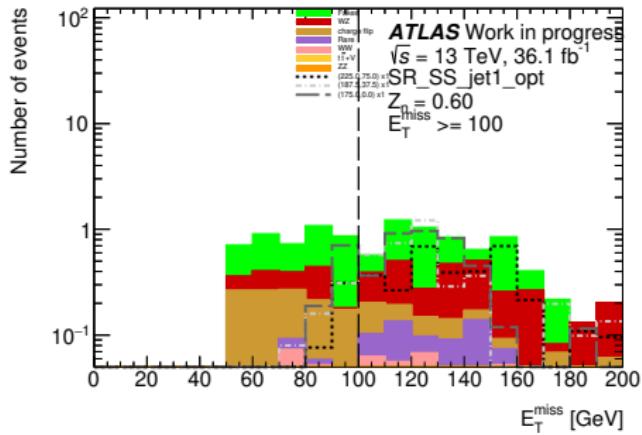
$$p_T^{l2}$$



$$|\Delta\eta_{ll}|$$

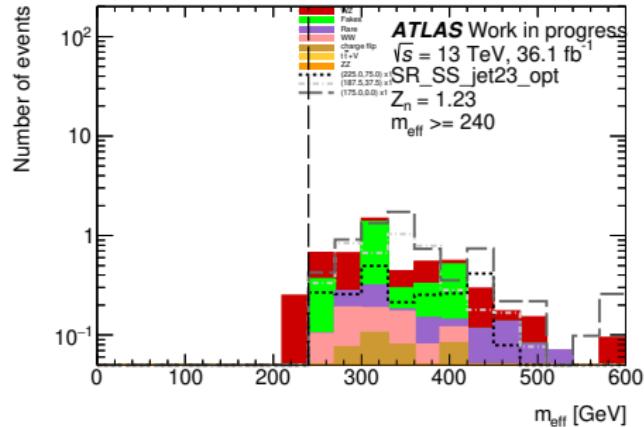
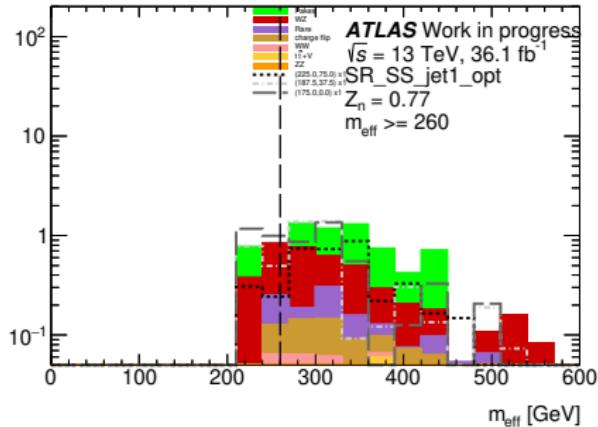


$E_T^{\text{miss}}$



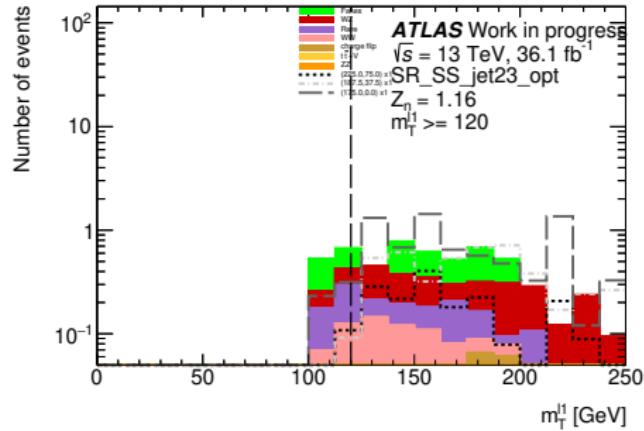
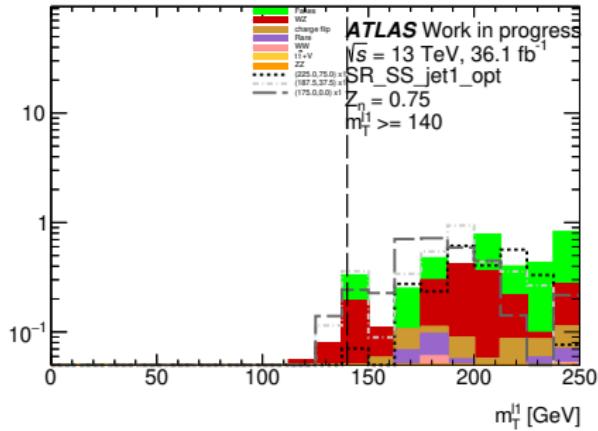
$m_{\text{eff}}$ 

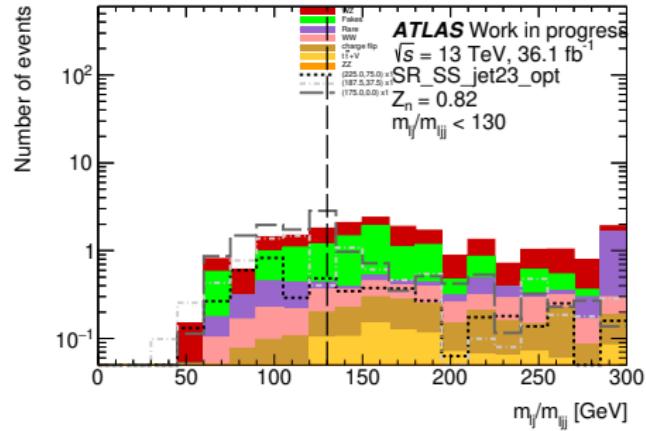
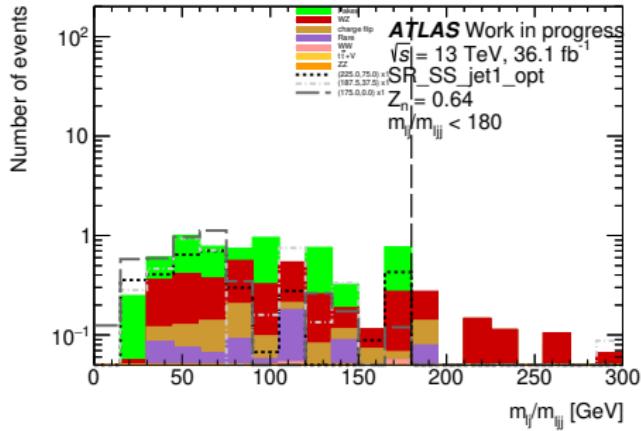
Number of events



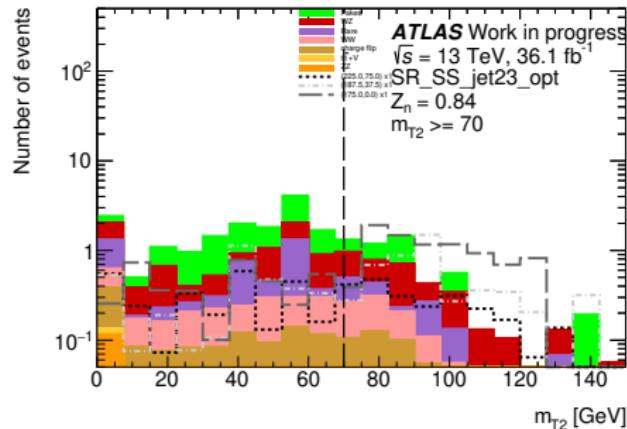
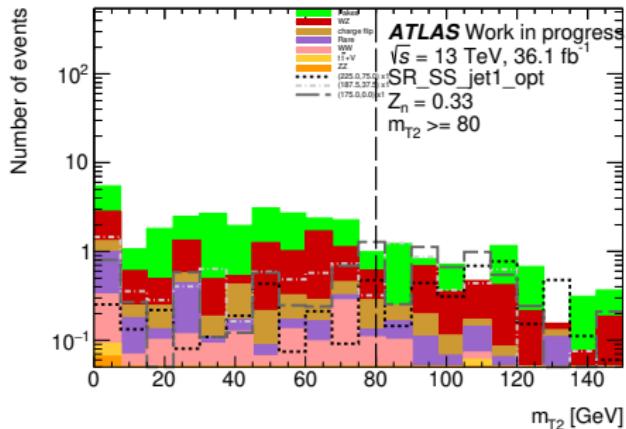
$m_T^{l1}$ 

Number of events



$m_{lj}/m_{ljj}$ 


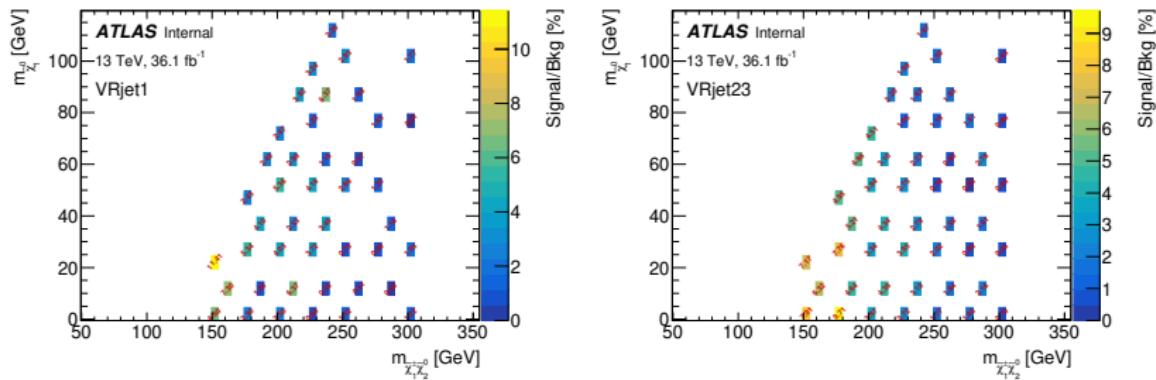
$$m_{T2}$$



# Validation Region

# Validation Region: Check for Signal Contribution

The signal contribution is below 12% and 10% for VRjet1 and VRjet23 respectively.



**Figure:** Percentages of signal contribution are shown for the VRjet1 (left) and VRjet23 (right).

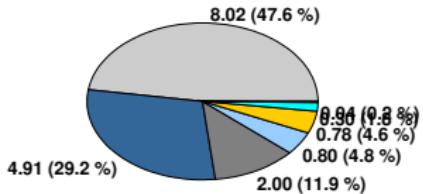
# Validation Region: Check for Background Composition

The background compositions for VRjet1 and SRjet1 are similar.

**ATLAS Internal**  
13 TeV, 36.1 fb<sup>-1</sup>

VRjet1

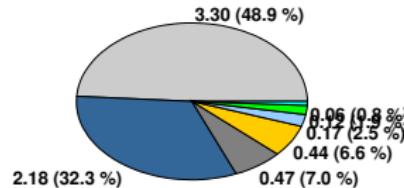
Rare ZZ WW Charge flip Fakes



**ATLAS Internal**  
13 TeV, 36.1 fb<sup>-1</sup>

SRjet1

Rare ZZ WW Charge flip Fakes



**Figure:** Background composition for the VRjet1 (left) and SRjet1 (right).

# Validation Region: Check for Background Composition

The background compositions for VRjet23 and SRjet23 are similar.

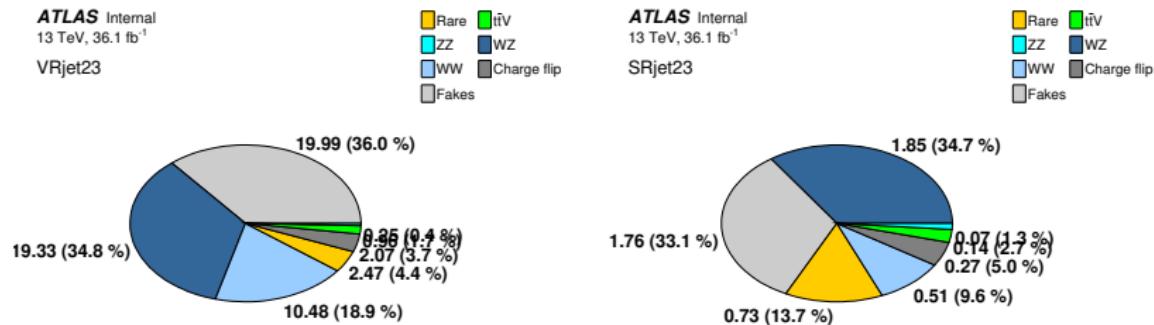
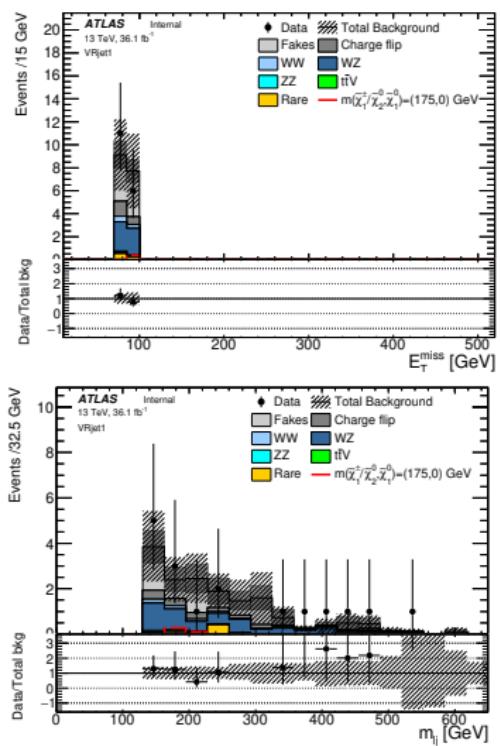
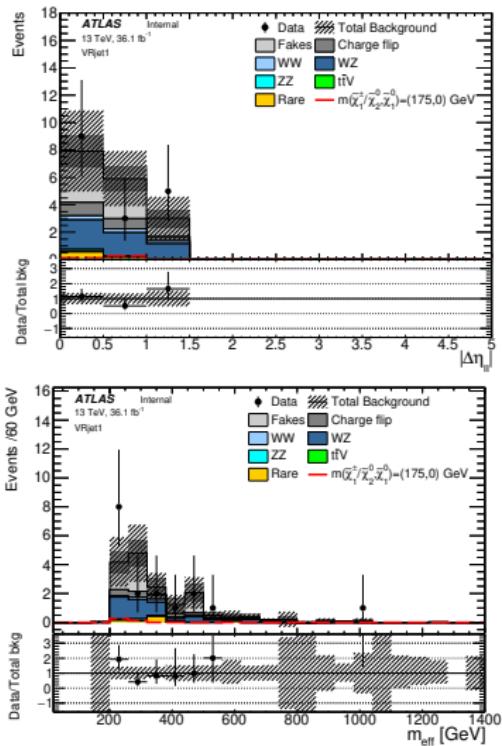
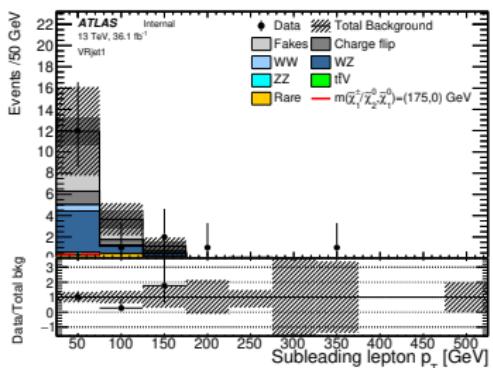
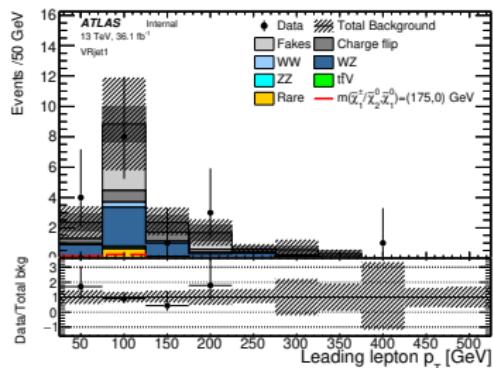
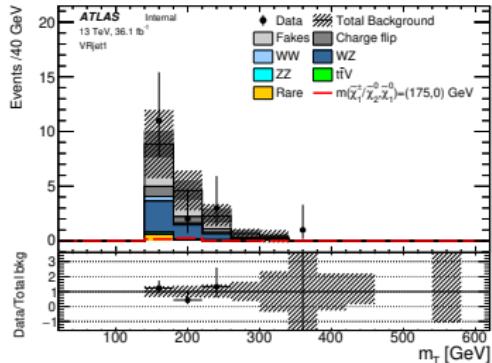
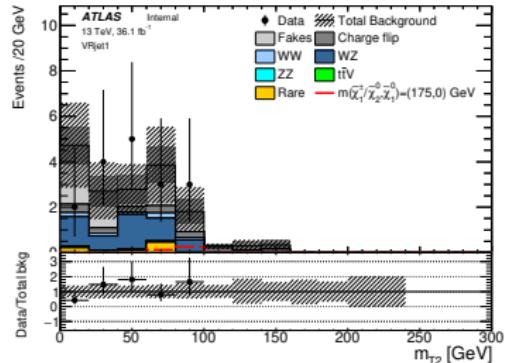


Figure: Background composition for the VRjet23 (left) and SRjet23 (right).

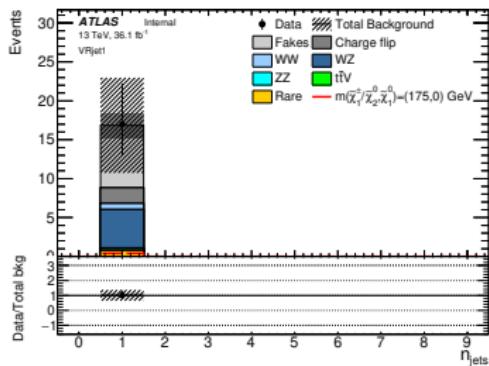
# Validation Region : Plots for different variables (VRjet1)



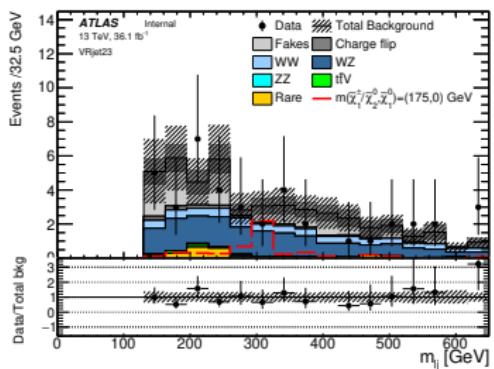
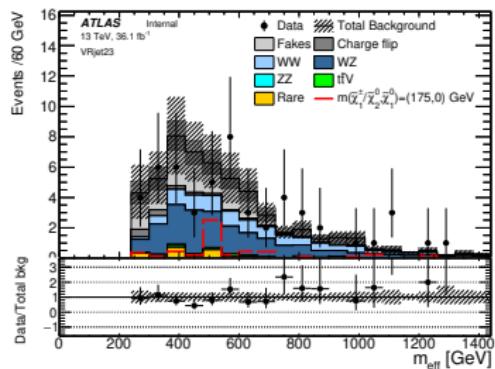
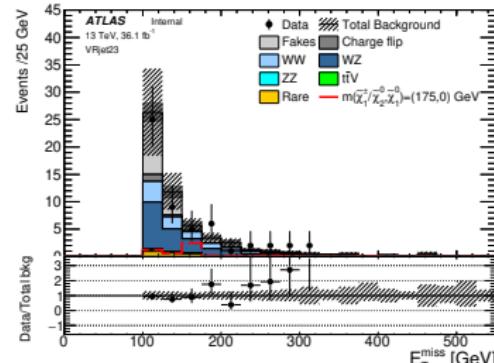
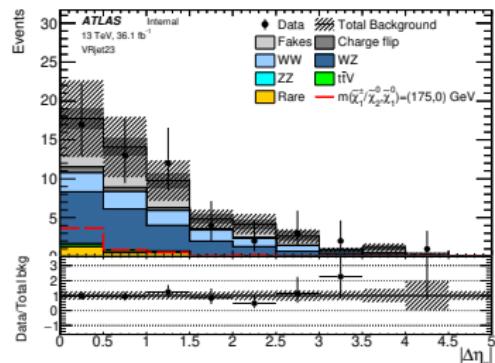
# Validation Region : Plots for different variables (VRjet1)



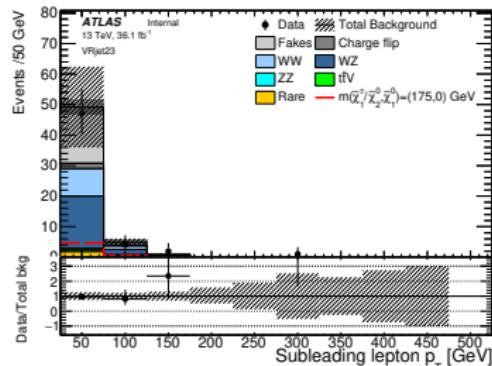
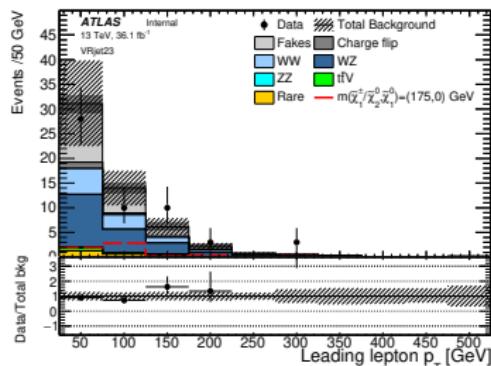
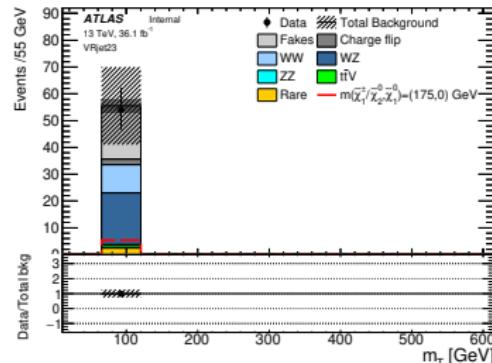
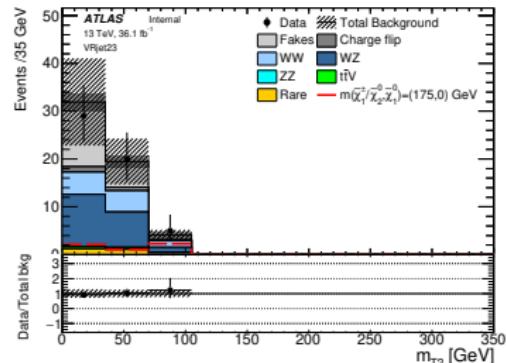
# Validation Region : Plots for different variables (VRjet1)



## Validation Region : Plots for different variables (VRjet23)



# Validation Region : Plots for different variables (VRjet23)



# Validation Region : Plots for different variables (VRjet23)

