

Search for chargino and neutralino production in final states with two same-sign leptons, jets and missing transverse momentum at $\sqrt{s} = 13$ TeV with the ATLAS detector

Samuel Lo

The University of Hong Kong

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Introduction

Standard Model : Fundamental Particles

- Standard Model(SM) is the current mainstream theory to describe the electromagnetic force, weak force and strong force.

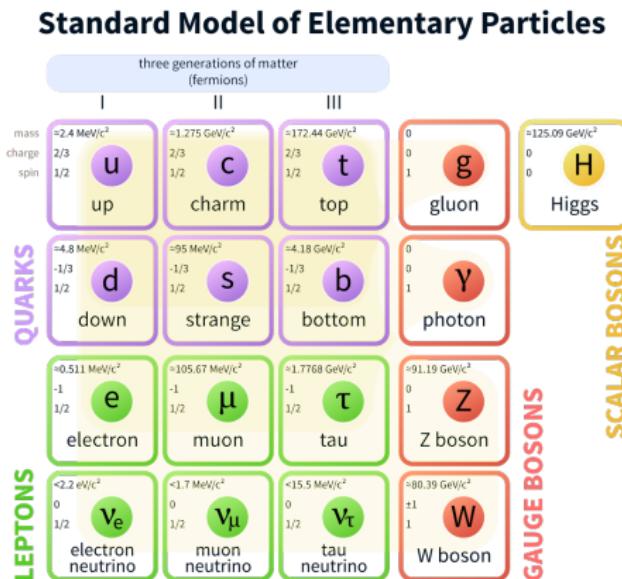


Figure: The “periodic” table for all fundamental particles in SM.

Standard Model : Success

SM gained a huge success in predicting experimental results.

- ▶ With the W boson explaining the beta decay, the existence of the neutral Z boson was predicted and discovered later in 1980s.
- ▶ The existence of a third generation of quarks was predicted in 1970s and top quark was discovered in 1990s.
- ▶ Higgs boson was proposed in 1960s and discovered recently in 2012.

Standard Model : Limitation

- ▶ SM cannot explain the nature of dark matter.
- ▶ The hierarchy problem
 - ▶ Why the weak force is stronger than the gravity by 10^{24} .
 - ▶ The Higgs boson mass is quadratically divergent due to the quantum corrections.
- ▶ The electroweak interaction and the strong interaction cannot be unified in SM.

Supersymmetry

- ▶ Supersymmetry is one of the most promising theory.
- ▶ Supersymmetry(SUSY) is a theoretical extension of the Standard Model.
- ▶ It can solve the hierarchy problem of Higgs mass.
- ▶ It can explain the nature of dark matter.
- ▶ It unifies the electroweak interaction and the strong interaction at $\sim 10^{16}$ GeV.

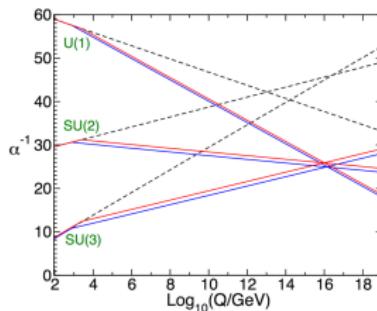


Figure: Black dashed lines are for SM. Red and blue lines are for SUSY.

Supersymmetry

- ▶ It predicts that each particle in the Standard Model has its own partner particle, called the superpartner (SUSY particle).
- ▶ The spin of the superpartner will differ from the Standard Model particle by $1/2$.
- ▶ A symmetry between the fermions and bosons.

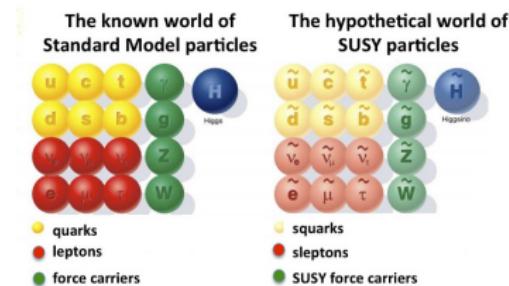


Figure: The particles in Standard Model and their corresponding superpartners and their names.

Supersymmetry : Superpartners

- ▶ In the simple SUSY models, we need two neutral Higgs fields (H_u^0 and H_d^0) and two charged Higgs fields (H_u^+ and H_d^-).
- ▶ The superpartners of 4 neutral bosons γ , Z , H_u^0 and H_d^0 together form 4 mass eigenstates, called neutralinos: $\tilde{\chi}_1^0$, $\tilde{\chi}_2^0$, $\tilde{\chi}_3^0$ and $\tilde{\chi}_4^0$.
- ▶ The superpartners of 4 charged bosons W^+ , W^- , H_u^+ and H_d^- together form two mass eigenstates with electric charge ± 1 , called charginos: $\tilde{\chi}_1^\pm$ and $\tilde{\chi}_2^\pm$.

Type	SM particle	Symbol	Spin	R-parity	Superpartner	Symbol	Spin	R-parity
Fermions	Quark	q	$\frac{1}{2}$	+1	Squark	\tilde{q}	0	-1
	Lepton	l	$\frac{1}{2}$	+1	Slepton	\tilde{l}	0	-1
Gluon	Gluon	g	1	+1	Gluino	\tilde{g}	$\frac{1}{2}$	-1
Neutral EW Bosons	Photon	γ	1	+1	Neutralinos	$\tilde{\chi}_1^0, \tilde{\chi}_2^0, \tilde{\chi}_3^0, \tilde{\chi}_4^0$	$\frac{1}{2}$	-1
	Z Boson	Z	1	+1				
	Neutral Higgs	H_u^0, H_d^0	0	+1				
Charged EW Bosons	W Boson	W^+, W^-	1	+1	Charginos	$\tilde{\chi}_1^\pm, \tilde{\chi}_2^\pm$	$\frac{1}{2}$	-1
	Charged Higgs	H_u^+, H_d^-	0	+1				

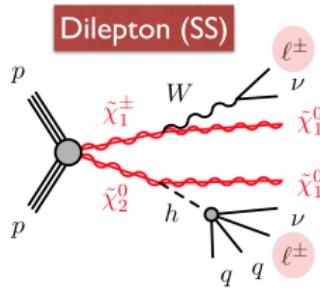
Table: The spin and R-parity for the Standard Model particles and their superpartners.

Supersymmetry : R-parity

- ▶ All SM particles have R-parity +1, while all SUSY particles have R-parity -1.
- ▶ If the R-parity is conserved, the lightest SUSY particle (LSP) cannot decay (stable).
- ▶ If the LSP is neutral and interacts with matter only by the weak interaction and gravity, it could be a candidate for dark matter, for example the lightest neutralinos $\tilde{\chi}_1^0$.
- ▶ The R-parity is assumed to be conserved.
- ▶ Due to the conservation of R-parity, the SUSY particles can only be pair-produced, and will eventually decay into SM particles and the LSP.
- ▶ The lightest neutralino $\tilde{\chi}_1^0$ is assumed to be the LSP.

Our Signal Scenario : Motivation

- ▶ In the recent searches for SUSY particles involved in strong interaction, their masses are suggested to be larger than 1 TeV.
- ▶ In this case, the direct pair production of SUSY electroweak boson $\tilde{\chi}_1^\pm$ and $\tilde{\chi}_2^0$ may be the dominant SUSY production process at the LHC, if their masses are below 1 TeV.
- ▶ Their masses are assumed to be the same, and denoted by $m_{\tilde{\chi}_1^\pm, \tilde{\chi}_2^0}$ (in simplified models).



Our Signal Scenario : Decay Processes

- ▶ $\tilde{\chi}_1^\pm \rightarrow W^\pm + \tilde{\chi}_1^0$, $W^\pm \rightarrow \ell^\pm + \nu$
- ▶ $\tilde{\chi}_2^0 \rightarrow h + \tilde{\chi}_1^0$, $h \rightarrow \ell^\pm + \nu + \text{quarks}$
- ▶ Two same-sign(SS) leptons (electron or muon)
- ▶ A large missing energy is expected
- ▶ Each quark will eventually become a particle shower within a narrow cone, called a jet, by the process of hadronization.

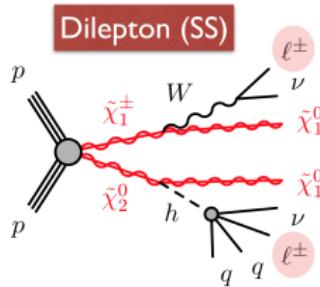


Figure: The Feynman diagram for our Wh same-sign signal scenario.

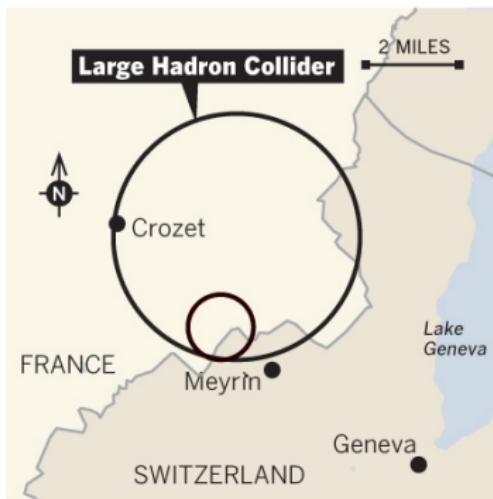
Our Signal Scenario : Sensitive Region

- ▶ If the mass difference ($m_{\tilde{\chi}_1^\pm, \tilde{\chi}_2^0} - m_{\tilde{\chi}_1^0}$) is slightly larger than the Higgs mass, it is called the compressed region.
- ▶ In the compressed region, one of the lepton may have low energy, due to the low momentum of the Higgs boson, and hence it may not be detected.
- ▶ In this case, only 2 out of 3 leptons are detected.
- ▶ Our search channel would be more sensitive than the 3-leptons channel in the compressed region.

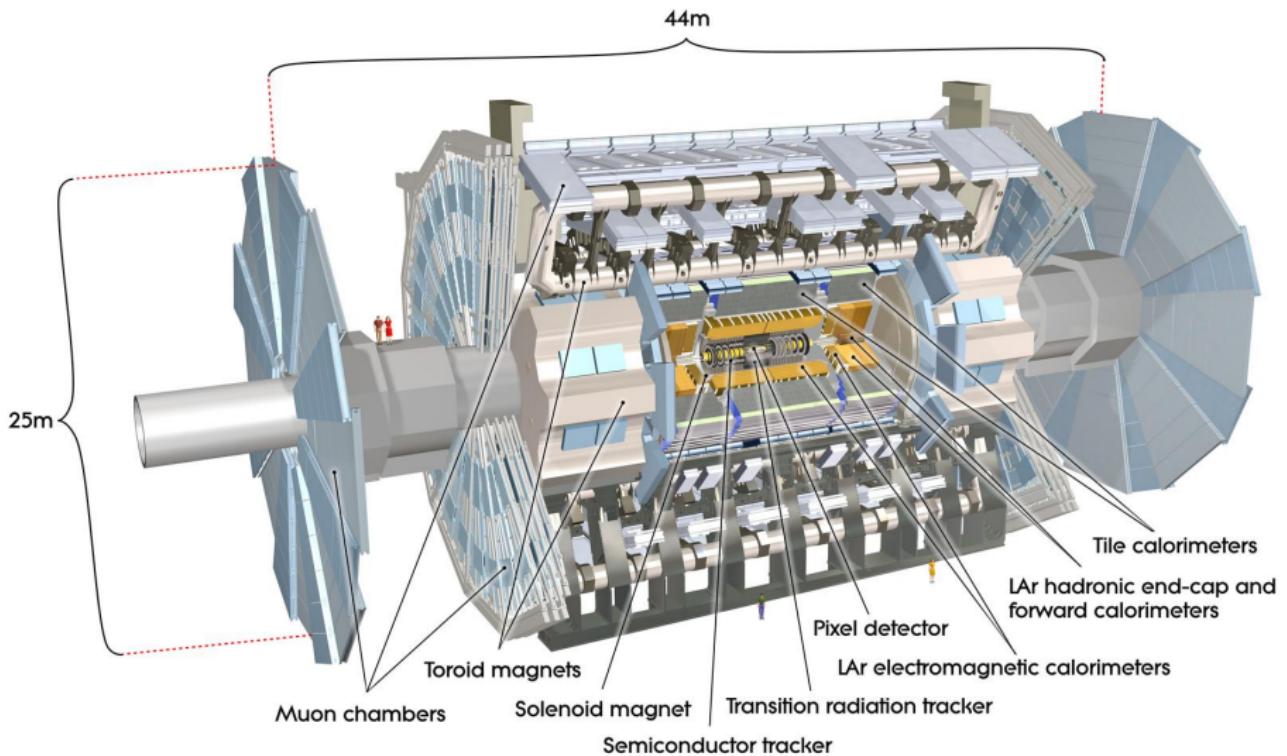
Experimental Setup

Experimental Setup : LHC

- ▶ The Large Hadron Collider (LHC) is the most powerful circular particle accelerator in the world.
- ▶ Its circumference is 27 km.
- ▶ Two beams of protons will be accelerated in opposite direction to centre-of-mass energy 13 TeV.
- ▶ One of the collision point is at the ATLAS detector.

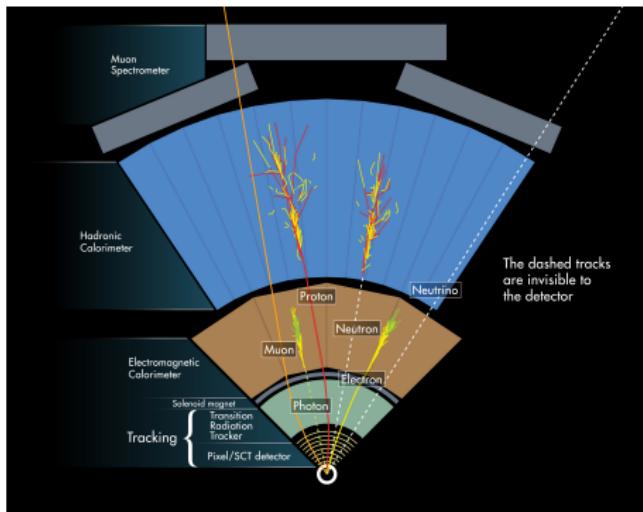


Experimental Setup : ATLAS detector



Experimental Setup : ATLAS detector

- The ATLAS detector consists of 3 main components:
 1. Inner detector: It is a particle tracker. It measures the tracks of charged particles.
 2. Calorimeter: It measures the energy of the particle and stops the particle. It consists of electromagnetic and hadronic calorimeters.
 3. Muon spectrometer: It is a particle tracker for muons.



Basic kinematic variables

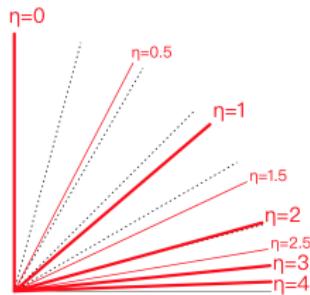
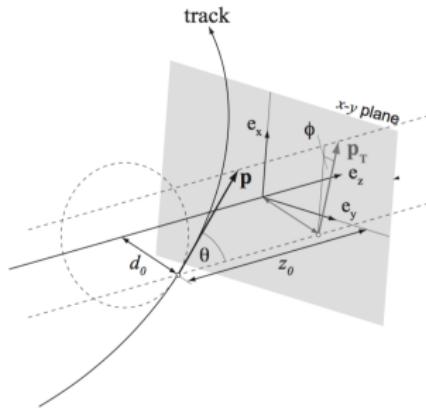
$$p_T = \sqrt{p_x^2 + p_y^2}$$

The pseudorapidity η is defined by

$$\eta = -\ln \left(\tan \frac{\theta}{2} \right)$$

The angle separation ΔR of two particles is defined as:

$$\Delta R = \sqrt{(\Delta\phi)^2 + (\Delta\eta)^2} = \sqrt{(\phi_2 - \phi_1)^2 + (\eta_2 - \eta_1)^2}$$



Challenges of Background Estimation

In order to search for SUSY, we need to understand the SM background very well.

1. Reducible SM background: They are due to the limitation of our detector and the wrong identification of the particle. There are two dominant reducible SM backgrounds, which are estimated by the data-driven method.
 - ▶ Charge flip background
 - ▶ Fake lepton background (dominant background)
2. Irreducible SM background: They are estimated by Monte Carlo (MC)
 - ▶ VV: WZ, WW, ZZ
 - ▶ tt+V
 - ▶ Rare: VVV, multitop, Higgs

Charge Flip Background

Charge Flip Background : Sources

- ▶ The charge flip background is due to the mis-identification of the sign of the charge of a lepton (mainly electron) in the reconstruction.
- ▶ The sign of the charge is determined by the direction of the curvature of the track. The charge may be reversed, or flipped.

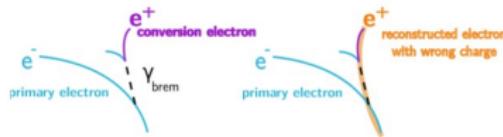


Figure: Bremsstrahlung and γ conversion.

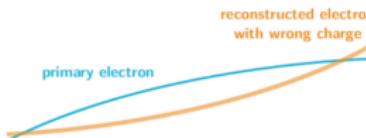


Figure: Track with high velocity
Lo Cheuk Yee

Charge Flip Background : Charge Flip Rate

- ▶ The probability that the charge of an electron is flipped is denoted by the charge-flip rate ϵ_i , where the index i represents the dependency on the p_T and $|\eta|$ of the electron.
- ▶ The value of index i is defined by the index of the following grid splitting.

Variable	Boundary of the grids
p_T (GeV)	25, 60, 90, 130, 150, 1000
$ \eta $	0, 0.50, 1.00, 1.37, 1.52, 1.80, 2.00, 2.47

Table: Binning in p_T and $|\eta|$ for the charge-flip rate ϵ_i .

Charge Flip Background : Charge Flip Rate

- ▶ Most of the same-sign charge-flip background comes from events with two opposite-sign electrons, with one of the charge is flipped.
- ▶ The probability p_{ij} that an OS data event becomes a SS data event (with the leading electron in bin i and the sub-leading electron in bin j) is

$$p_{ij} = (1 - \epsilon_i)\epsilon_j + (1 - \epsilon_j)\epsilon_i$$

- ▶ It is called the charge flip weight.
- ▶ The charge flip background can be estimated by applying the charge flip weights on all OS events in data.
- ▶ The charge flip rates can be measured by the likelihood method described in the next slide.

Charge Flip Background : Likelihood Method

- ▶ We select $Z \rightarrow ee$ events in data by requiring the invariant mass of the two electrons is within the Z mass window cut of 80-100 GeV.
- ▶ The events inside the Z mass window is then subtracted by the non-Zee processes, by using the sideband technique.
- ▶ After the subtraction, the total number of event is denoted by N^{ij} , and the number of events with two same-sign leptons is denoted by N_{SS}^{ij} .
- ▶ The N_{SS}^{ij} is described by a Poisson distribution with the expected value $\lambda = N^{ij} p_{ij}$

$$P(N_{SS}^{ij} | N^{ij}, \epsilon_i, \epsilon_j) = \frac{(N^{ij} p_{ij})^{N_{SS}^{ij}} e^{-N^{ij} p_{ij}}}{N_{SS}^{ij}!}$$

Charge Flip Background : Likelihood Method

- ▶ Converting it to the likelihood function L and taking the negative natural log yields

$$\begin{aligned}-\ln L &= -\ln \prod_{ij} P(N_{SS}^{ij} | N^{ij}, \epsilon_i, \epsilon_j) \\&= -\sum_{ij} \left[N_{SS}^{ij} \ln(N^{ij}[\epsilon_i(1-\epsilon_j) + (1-\epsilon_i)\epsilon_j]) \right. \\&\quad \left. - N^{ij}[\epsilon_i(1-\epsilon_j) + (1-\epsilon_i)\epsilon_j] \right] + \text{constant}\end{aligned}$$

- ▶ where the summation over i and j is taken over all p_T and $|\eta|$ bins of both electrons.
- ▶ By minimizing this likelihood, the charge-flip rate can be estimated.

Charge Flip Background : Result for Charge Flip Rate

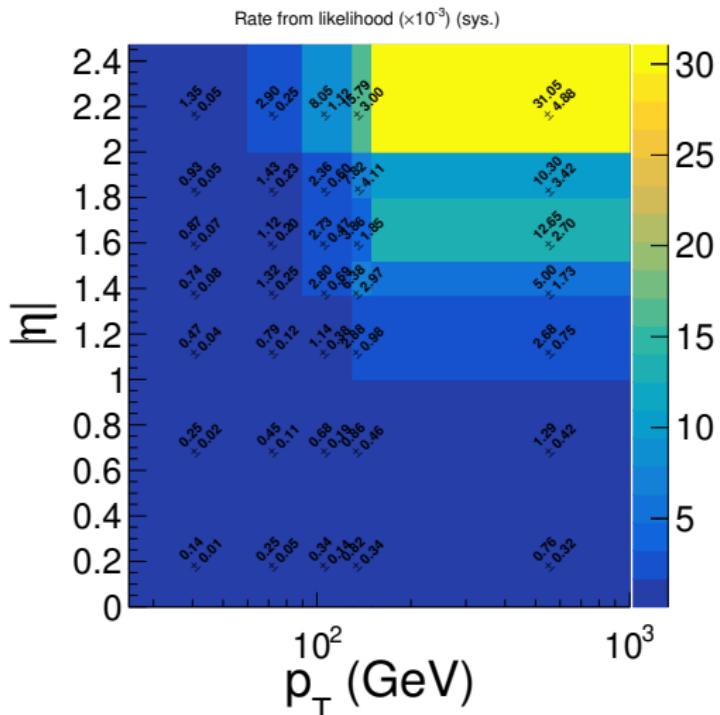


Figure: The measured values of the charge-flip rate $|\epsilon_i|$ in data, with total uncertainties.

Fake Lepton Background

Fake Lepton Background : Sources

The source of the fake lepton is that other object (like jet and photon) is mis-identified as a lepton. They can be mainly classified into 3 types:

- ▶ Heavy-flavor fakes:
 - ▶ leptons from the leptonic decays of heavy-quark (b or c) hadrons.
- ▶ Light-flavor fakes:
 - ▶ leptons from the leptonic decays of light-quark hadrons.
OR
 - ▶ It is due to the mis-identification of light-quark hadrons itself.
- ▶ photon conversion:
 - ▶ leptons from the pair production of a photon

Fake Lepton Background : Matrix Method

- ▶ The fake leptons is estimated by the matrix method.
- ▶ The probability that a real lepton passes the tight selection criteria (i.e. tight lepton) is denoted by the real efficiency ϵ .
- ▶ The probability that a real lepton does not pass the tight selection criteria (i.e. loose lepton) is denoted by $\bar{\epsilon} = 1 - \epsilon$.
- ▶ The probability that a fake lepton passes the tight selection criteria (i.e. tight lepton) is denoted by the fake efficiency f .
- ▶ The probability that a fake lepton does not pass the tight selection criteria (i.e. loose lepton) is denoted by $\bar{f} = 1 - f$.

$$\begin{pmatrix} N_T \\ N_L \end{pmatrix} = \begin{pmatrix} \epsilon & f \\ \bar{\epsilon} & \bar{f} \end{pmatrix} \begin{pmatrix} N_R \\ N_F \end{pmatrix}$$

Fake Lepton Background : Matrix Method

Suppose we know $\begin{pmatrix} N_T \\ N_L \end{pmatrix}$, we can find N_F by inverting the matrix.

$$\begin{pmatrix} 0 \\ N_F \end{pmatrix} = \begin{pmatrix} 0 & 0 \\ 0 & 1 \end{pmatrix} \begin{pmatrix} N_R \\ N_F \end{pmatrix} = \begin{pmatrix} 0 & 0 \\ 0 & 1 \end{pmatrix} \begin{pmatrix} \epsilon & f \\ \bar{\epsilon} & \bar{f} \end{pmatrix}^{-1} \begin{pmatrix} N_T \\ N_L \end{pmatrix}$$

Then, we can find the number of tight lepton due to the fake lepton, N'_T

$$\begin{aligned} \begin{pmatrix} N'_T \\ N'_L \end{pmatrix} &= \begin{pmatrix} \epsilon & f \\ \bar{\epsilon} & \bar{f} \end{pmatrix} \begin{pmatrix} 0 \\ N_F \end{pmatrix} \\ &= \begin{pmatrix} \epsilon & f \\ \bar{\epsilon} & \bar{f} \end{pmatrix} \begin{pmatrix} 0 & 0 \\ 0 & 1 \end{pmatrix} \begin{pmatrix} \epsilon & f \\ \bar{\epsilon} & \bar{f} \end{pmatrix}^{-1} \begin{pmatrix} N_T \\ N_L \end{pmatrix} \end{aligned}$$

Fake Lepton Background : Matrix Method

We can generalize the one-lepton case to the two-leptons case.

$$\begin{pmatrix} N_{TT} \\ N_{TL} \\ N_{LT} \\ N_{LL} \end{pmatrix} = \begin{pmatrix} \epsilon_1 \epsilon_2 & \epsilon_1 f_2 & f_1 \epsilon_2 & f_1 f_2 \\ \epsilon_1 \bar{\epsilon}_2 & \epsilon_1 \bar{f}_2 & f_1 \bar{\epsilon}_2 & f_1 \bar{f}_2 \\ \bar{\epsilon}_1 \epsilon_2 & \bar{\epsilon}_1 f_2 & \bar{f}_1 \epsilon_2 & \bar{f}_1 f_2 \\ \bar{\epsilon}_1 \bar{\epsilon}_2 & \bar{\epsilon}_1 \bar{f}_2 & \bar{f}_1 \bar{\epsilon}_2 & \bar{f}_1 \bar{f}_2 \end{pmatrix} \begin{pmatrix} N_{RR} \\ N_{RF} \\ N_{FR} \\ N_{FF} \end{pmatrix}$$

where ϵ_1 is the probability that a leading real lepton passes the signal selection (i.e. tight lepton), ϵ_2 is the probability that a subleading real lepton passes the signal selection (i.e. tight lepton), etc.

Fake Lepton Background : Matrix Method

Similarly, we can find the number of tight-tight leptons due to the fake leptons, N'_{TT}

$$\begin{pmatrix} N'_{TT} \\ N'_{TL} \\ N'_{LT} \\ N'_{LL} \end{pmatrix} = \begin{pmatrix} \epsilon_1 \epsilon_2 & \epsilon_1 f_2 & f_1 \epsilon_2 & f_1 f_2 \\ \epsilon_1 \bar{\epsilon}_2 & \epsilon_1 \bar{f}_2 & f_1 \bar{\epsilon}_2 & f_1 \bar{f}_2 \\ \bar{\epsilon}_1 \epsilon_2 & \bar{\epsilon}_1 f_2 & \bar{f}_1 \epsilon_2 & \bar{f}_1 f_2 \\ \bar{\epsilon}_1 \bar{\epsilon}_2 & \bar{\epsilon}_1 \bar{f}_2 & \bar{f}_1 \bar{\epsilon}_2 & \bar{f}_1 \bar{f}_2 \end{pmatrix} \begin{pmatrix} 0 & 0 & 0 & 0 \\ 0 & 1 & 0 & 0 \\ 0 & 0 & 1 & 0 \\ 0 & 0 & 0 & 1 \end{pmatrix}^{-1} \begin{pmatrix} N_{TT} \\ N_{TL} \\ N_{LT} \\ N_{LL} \end{pmatrix}$$

Fake Lepton Background : Results of Fake Efficiencies

The method for measuring the fake efficiencies is described in the backup slides.

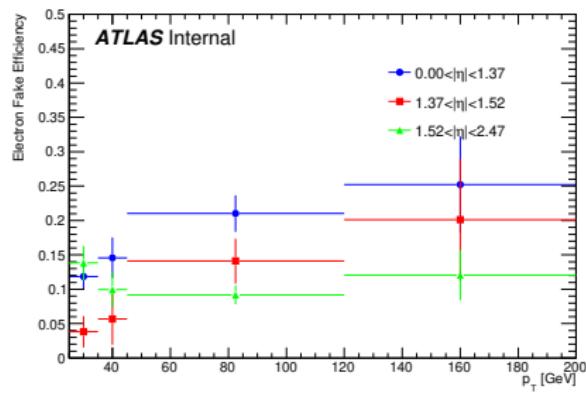


Figure: Electron fake efficiency

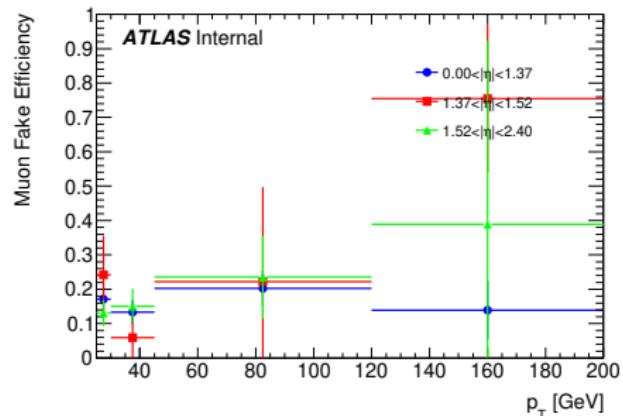


Figure: Muon fake efficiency

Signal Region

Signal Region: Challenges

Estimation of the signal sensitivity

- ▶ Cross section for our signal $\sim 0.1 \text{ pb}$
- ▶ Total luminosity in one year $\sim 10^4 \text{ pb}^{-1}$
- ▶ Expected number of signal events in one year $\sim 10^3$
- ▶ Time interval between each collision = 25 ns
- ▶ Total number of collisions in one year $\sim 10^{15}$
- ▶ The probability to have a signal event $\sim 10^{-12}$
- ▶ Conclusion: the SUSY signal is very rare.

Signal Region : Strategy

In order to extract the rare signal, the following strategies are used.

- ▶ A trigger system are used when taking the data (online), to only store the interested events in the disk.
- ▶ Some discriminant variables (based on the kinematic variables) are defined, to help distinguish the signal and the background.
- ▶ Two dedicated signal regions are defined (SRjet1 and SRjet23) based on the discriminant variables, to maximize the signal sensitivity.

Signal Region : discriminant variables

- ▶ n_{jets} : Number of signal jets:
- ▶ $n_{b\text{-jets}}$: Number of b -jets (jet from b quark).
- ▶ $\Delta\eta_{ll}$: It is designed to estimate the angle between the two leptons.

$$\Delta\eta_{ll} = |\eta_1 - \eta_2|$$

- ▶ m_{eff} : It is designed to estimate the total energy and momentum of the charginos $\tilde{\chi}_1^\pm$ and neutralinos $\tilde{\chi}_2^0$.

$$m_{\text{eff}} = p_T^1 + p_T^2 + E_T^{\text{miss}} + \sum_{\text{signal jets}} p_T$$

- ▶ m_{ll} : The invariant mass of the 4-momentum sum of the two leptons

$$(m_{ll})^2 = (p_1 + p_2)^2$$

Signal Region : discriminant variables

- The missing momentum \mathbf{p}^{miss} due to undetected particles is defined by

$$\mathbf{p}^{\text{miss}} = - \sum_{\text{All detected particles}} \mathbf{p}$$

- E_T^{miss} or MET: The transverse missing energy is designed to estimate the missing energy from the undetected particles.

$$E_T^{\text{miss}} = |\mathbf{p}_T^{\text{miss}}| = \sqrt{(p_x^{\text{miss}})^2 + (p_y^{\text{miss}})^2}$$

Signal Region : discriminant variables

- ▶ m_T : It is designed to reconstruct the mass of the W boson.

$$m_T = \sqrt{2p_T^1 E_T^{\text{miss}} (1 - \cos \Delta\phi)}$$

where $\Delta\phi$ is the azimuthal angle between the leading lepton and the missing transverse momentum.

- ▶ m_{T2} : It is designed to set a lower bound on the masses of charginos $\tilde{\chi}_1^\pm$ and neutralinos $\tilde{\chi}_2^0$.

$$m_{T2} = \min_{\mathbf{q}_T} \left[\max \left(m_T(\mathbf{p}_T^1, \mathbf{q}_T), m_T(\mathbf{p}_T^2, \mathbf{p}_T^{\text{miss}} - \mathbf{q}_T) \right) \right]$$

$$m_T(\mathbf{p}_T, \mathbf{q}_T) = \sqrt{2p_T q_T (1 - \cos \Delta\phi)}$$

Signal Region : discriminant variables

- ▶ m_{lj} or m_{ljj} : It is designed to reconstruct the mass of the Higgs boson.

$$p_{\text{jet-system}} = \begin{cases} p_{\text{jet1}} & \text{for SRjet1 (1 jet)} \\ p_{\text{jet1}} + p_{\text{jet2}} & \text{for SRjet23 (2 or 3 jets)} \end{cases}$$

$$p_{\text{closest-lepton}} = \begin{cases} p_{\ell 1} & \text{if } \Delta R(p_{\ell 1}, p_{\text{jet-system}}) \leq \Delta R(p_{\ell 2}, p_{\text{jet-system}}) \\ p_{\ell 2} & \text{if } \Delta R(p_{\ell 1}, p_{\text{jet-system}}) > \Delta R(p_{\ell 2}, p_{\text{jet-system}}) \end{cases}$$

$$(m_{lj(j)})^2 = (p_{\text{closest-lepton}} + p_{\text{jet-system}})^2$$

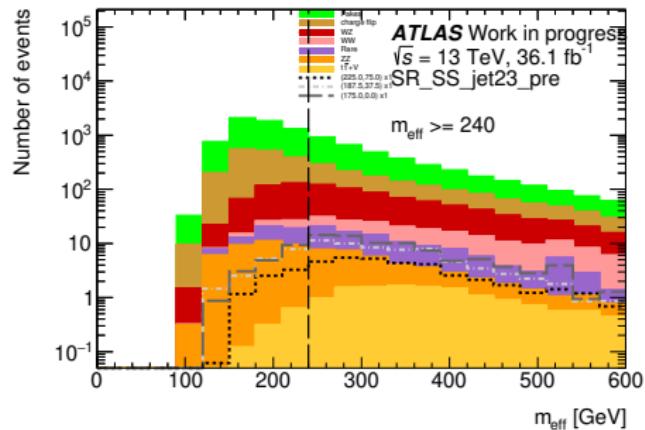
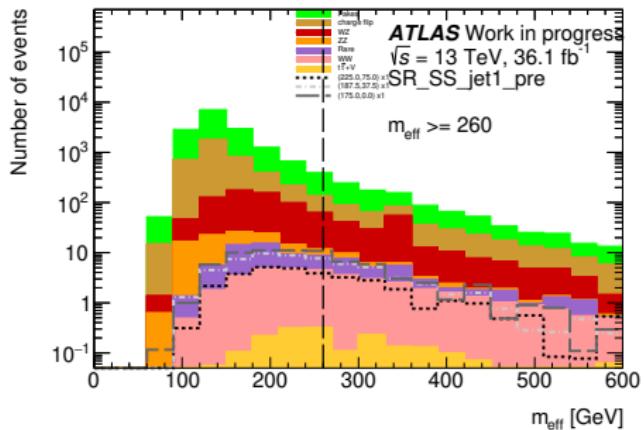
Signal Region

- ▶ Signal Region(SR) is defined by applying a set of selections with discriminant variables, such that the signal sensitivity is high.
- ▶ By trying different sets of selections, the signal region with the highest signal sensitivity can be found. This process is called the signal region optimization.
- ▶ In the SR optimization, the MC signal and background are used to study the expected sensitivity.

Before the signal region optimization, the following pre-selection are applied.

- ▶ Exactly 2 leptons with $p_T > 25$ GeV.
- ▶ Electron: $|\eta| < 2.47$; Muon: $|\eta| < 2.4$
- ▶ The two leptons are same-sign in electric charge (SS).
- ▶ $n_{\text{b-jets}} = 0$, to reduce the SM background from top quark.
- ▶ Two signal regions: SRjet1 with $n_{\text{jets}} = 1$, SRjet23 with $n_{\text{jets}} = 2$ or 3.

Signal Region : Pre-selection plots



Signal Region : Optimized Event Selection

	SRjet1	SRjet23
n_{jets}	1	2 or 3
Leading lepton p_T [GeV]	≥ 25	≥ 25
Sub-leading lepton p_T [GeV]	≥ 25	≥ 25
$ \Delta\eta_{ll} $	< 1.5	-
E_T^{miss} [GeV]	≥ 100	≥ 100
m_{eff} [GeV]	≥ 260	≥ 240
m_T^{l1} [GeV]	≥ 140	≥ 120
m_{lj}/m_{ljj} [GeV]	< 180	< 130
m_{T2} [GeV]	≥ 80	≥ 70

Signal Region : Expected Yields and Sensitivity

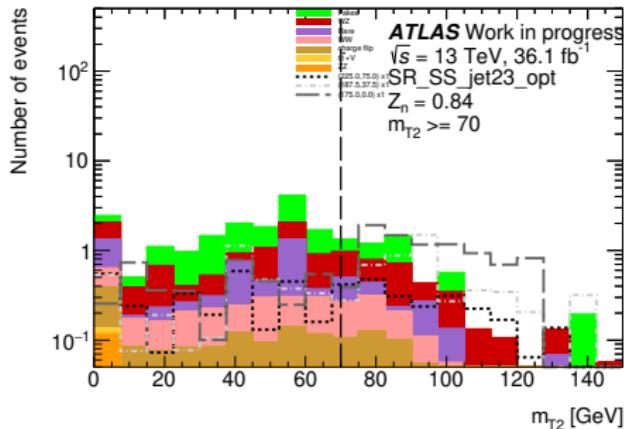
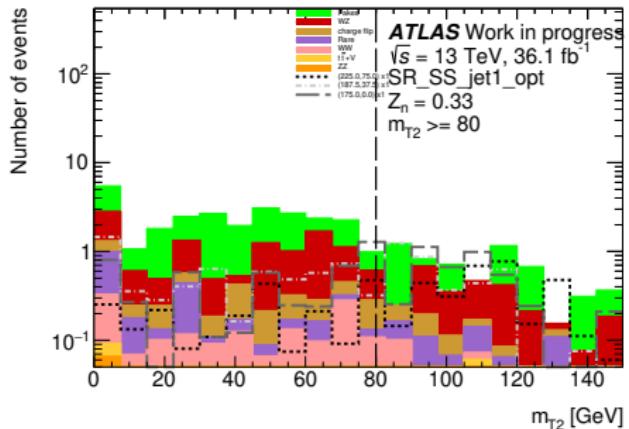
	Number of events	Significance
Fakes	3.295 ± 0.819	
WZ	2.176 ± 0.398	
charge flip	0.472 ± 0.053	
Rare	0.444 ± 0.111	
WW	0.166 ± 0.023	
$t\bar{t} + V$	0.125 ± 0.046	
ZZ	0.055 ± 0.028	
Total BG	6.733 ± 0.921	
(225.0,75.0)	3.33 ± 0.60	0.74
(187.5,37.5)	3.77 ± 0.95	0.86
(175.0,0.0)	4.29 ± 0.73	0.98

Table: Yields in SRjet1

	Number of events	Significance
WZ	1.849 ± 0.273	
Fakes	1.765 ± 0.709	
Rare	0.731 ± 0.195	
WW	0.514 ± 0.037	
charge flip	0.267 ± 0.029	
$t\bar{t} + V$	0.142 ± 0.031	
ZZ	0.067 ± 0.025	
Total BG	5.335 ± 0.787	
(225.0,75.0)	2.35 ± 0.34	0.57
(187.5,37.5)	4.72 ± 0.74	1.26
(175.0,0.0)	8.60 ± 1.51	2.24

Table: Yields in SRjet23

Signal Region : N-1 plots



Validation Region

Validation Region: Introduction

- ▶ The validation region(VR) is designed to ensure that the estimation of background is reliable.
- ▶ Method: The backgrounds are compared with the data, to see whether they agree with each other.
- ▶ One VR is defined for each SR: VRjet1 and VRjet23.

There are three requirements for the validation regions:

- ▶ The signal contribution should be small.
- ▶ The validation region is orthogonal to the corresponding signal region.
- ▶ The background composition is similar to the corresponding signal region.

Validation Region: Method

- ▶ In order to make the VR orthogonal to its SR, E_T^{miss}/m_T cut is inverted.
- ▶ Adjust its lower cut, to have a similar background composition as in its SR.
- ▶ Invert the m_{lj}/m_{ljj} cut, to reduce the signal.
- ▶ Remove m_{T2} and m_{eff} cuts, to increase statistics.

Cut	VRjet1	VRjet23
n_{jets}	1	[2, 3]
$\Delta\eta_{ll}$	< 1.5	—
E_T^{miss} [GeV]	[70, 100]	> 100
m_T [GeV]	> 140	[65, 120]
m_{eff} [GeV]	—	> 240
$m_{lj(j)}$ [GeV]	> 130	> 130
m_{T2} [GeV]	—	—

Validation Region: Yield

Process	VRjet1	VRjet23
Rare	$0.775 \pm 0.389^{+0.661}_{-0.362}$	$2.469 \pm 0.674^{+0.998}_{-0.899}$
$t\bar{t}V$	$0.039 \pm 0.013^{+0.018}_{-0.012}$	$0.959 \pm 0.082^{+0.152}_{-0.146}$
ZZ	$0.298 \pm 0.060^{+0.089}_{-0.063}$	$0.247 \pm 0.045^{+0.113}_{-0.047}$
WZ	$4.909 \pm 0.530^{+0.960}_{-0.899}$	$19.325 \pm 0.643^{+4.393}_{-4.346}$
WW	$0.801 \pm 0.051^{+0.123}_{-0.060}$	$10.477 \pm 0.176^{+0.796}_{-0.726}$
Charge flip	$1.997 \pm 0.128^{+0.260}_{-0.260}$	$2.065 \pm 0.085^{+0.166}_{-0.166}$
Fakes	$8.021 \pm 1.390^{+5.806}_{-5.806}$	$19.990 \pm 2.013^{+13.461}_{-13.461}$
Total BG	$16.839 \pm 1.545^{+5.915}_{-5.912}$	$55.534 \pm 2.228^{+14.396}_{-14.332}$
Data	17	54

Table: The uncertainties include the statistical and systematic uncertainties.

The distribution plots for different variables are in the backup slides.

Results

Results: Observed Events

- The observed events were consistent with the prediction in the SM.

	Number of events
Fakes	3.295 ± 2.100
WZ	2.176 ± 0.398
charge flip	0.472 ± 0.053
Rare	0.444 ± 0.111
WW	0.166 ± 0.023
$t\bar{t} + V$	0.125 ± 0.046
ZZ	0.055 ± 0.028
Total BG	6.733 ± 2.142
Data	2

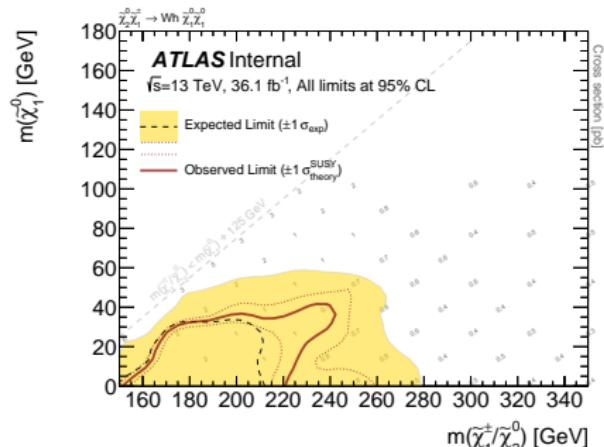
Table: Yields in SRjet1

	Number of events
WZ	1.849 ± 0.273
Fakes	1.765 ± 1.460
Rare	0.731 ± 0.195
WW	0.514 ± 0.037
charge flip	0.267 ± 0.029
$t\bar{t} + V$	0.142 ± 0.031
ZZ	0.067 ± 0.025
Total BG	5.335 ± 1.499
Data	8

Table: Yields in SRjet23

Results: Exclusion Limits

- ▶ New exclusion limits on the mass points were set, by combining the two SR.
 - ▶ This means that the mass points inside the contour line are rejected with 95% confidence level.
 - ▶ The exclusion limits for the masses of $\tilde{\chi}_1^\pm$ and $\tilde{\chi}_2^0$ are extended up to 245 GeV, and the exclusion limits for the mass of $\tilde{\chi}_1^0$ are extended up to 40 GeV.



Conclusion

- ▶ By using the ATLAS data with total integrated luminosity 36.1 fb^{-1} , we searched for the electroweak pair production of a chargino and a neutralino ($p + p \rightarrow \tilde{\chi}_1^\pm + \tilde{\chi}_2^0$), with the Wh channel $\tilde{\chi}_1^\pm \rightarrow \tilde{\chi}_1^0 + W$ and $\tilde{\chi}_2^0 \rightarrow \tilde{\chi}_1^0 + h$, and the same-sign channel.
- ▶ The numbers of observed events in data were consistent with the prediction in the SM.
- ▶ The exclusion limits for the masses of $\tilde{\chi}_1^\pm$ and $\tilde{\chi}_2^0$ are extended up to 245 GeV, and the exclusion limits for the mass of $\tilde{\chi}_1^0$ are extended up to 40 GeV, with 95% confidence level.

Acknowledgments

- ▶ Thank my supervisor Yanjun Tu for her guidance and suggestions on my analysis.
- ▶ Thank Dongliang Zhang for teaching me all technical stuff on my analysis.
- ▶ Thank Daniela Paredes for her help and suggestions on my analysis.

Thank You

Backup

My Contribution

- ▶ All the study and results I show here are made by me.
- ▶ Some of them are selected for publication.
- ▶ Others are cross-check to the work by the post-doc in the group.

Standard Model : Fundamental Interaction

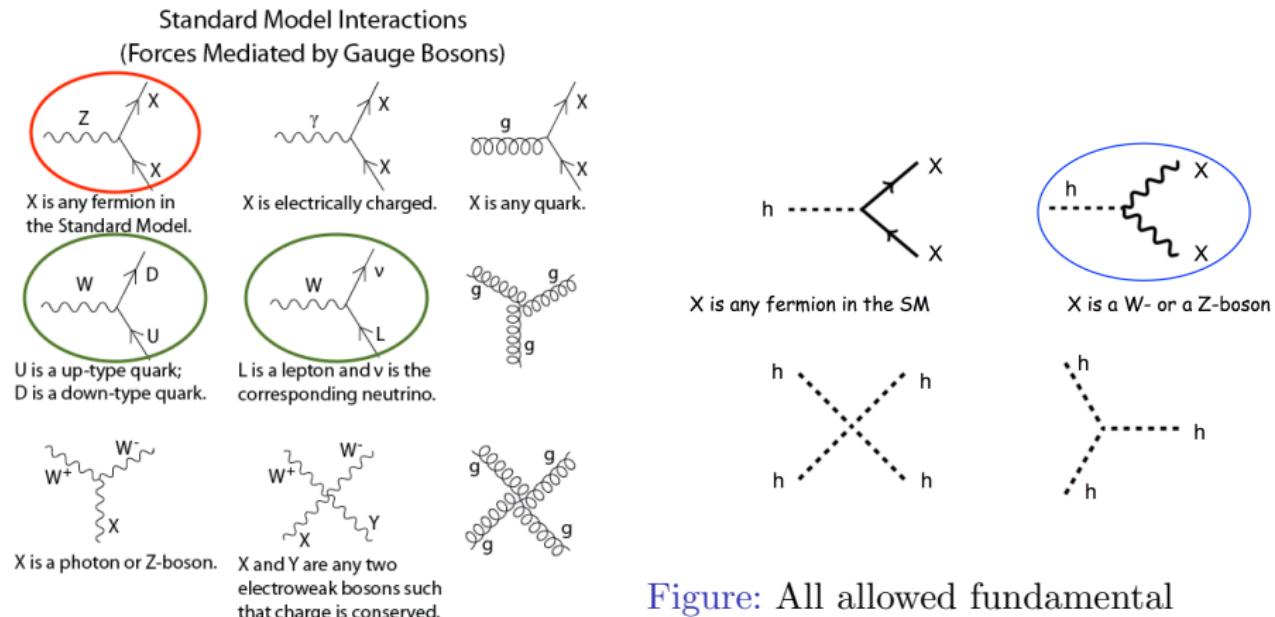


Figure: All allowed fundamental higgs-related Feynman vertices in SM.

Figure: All allowed fundamental Feynman vertices in SM, except higgs-related vertices.

Supersymmetry : R-parity

- ▶ The baryon number B is defined by $\frac{1}{3}(n_q - n_{\bar{q}})$, where n_q is the number of quarks and $n_{\bar{q}}$ is the number of anti-quarks.
- ▶ The lepton number L is defined by $n_l - n_{\bar{l}}$, where n_l is the number of leptons and $n_{\bar{l}}$ is the number of anti-leptons.
- ▶ In SM, $(B - L)$ is conserved. But in SUSY, it is no longer conserved.
- ▶ To keep $(B - L)$ conservation and prevent the proton decay, the R-parity P_R is introduced.

$$P_R = (-1)^{3(B-L)-2s}$$

where s is the spin.

Charge Flip Background : Likelihood Method

Likelihood Method : Method

- ▶ The probability that the charge of an electron is flipped is denoted by the charge-flip rate ϵ_i , where the index i represents the dependency on the p_T and $|\eta|$ of the electron.
- ▶ The value of index i is defined by the index of the following grids in the table.

Variable	Boundary of the bins
p_T (GeV)	25, 60, 90, 130, 150, 1000
$ \eta $	0, 0.50, 1.00, 1.37, 1.52, 1.80, 2.00, 2.47

Table: Binning in p_T and $|\eta|$ for the charge-flip rate ϵ_i .

- ▶ The probability p_{ij} that an OS data event becomes a SS data event (with the leading lepton in bin i and the subleading lepton in bin j) is

$$p_{ij} = (1 - \epsilon_i)\epsilon_j + (1 - \epsilon_j)\epsilon_i$$

Likelihood Method : Definition of Control Region

A control region is defined.

- ▶ Trigger requirement
- ▶ exactly 2 signal leptons
- ▶ 2 electrons
- ▶ $25 < p_T < 1000$
- ▶ $|\eta| < 2.47$

Likelihood Method : Method

- ▶ The charge flip rate ϵ_i is measured by using the data in the control region, within the Z mass window cut of 80-100 GeV.
- ▶ The events inside the Z mass window is then subtracted by the non-Zee processes, by using the sideband technique.
- ▶ After the subtraction, the total number of event is denoted by N^{ij} , and the number of events with two same-sign leptons is denoted by N_{SS}^{ij} .
- ▶ The expected value of N_{SS}^{ij} is then given by $\lambda = N^{ij} p_{ij}$. And, if N_{SS}^{ij} is described by a Poisson distribution, then

$$P(N_{SS}^{ij}|N^{ij}, \epsilon_i, \epsilon_j) = \frac{(N^{ij} p_{ij})^{N_{SS}^{ij}} e^{-N^{ij} p_{ij}}}{N_{SS}^{ij}!}$$

Likelihood Method : Method

- ▶ Converting it to the likelihood function L and taking the negative natural log yields

$$\begin{aligned} -\ln L &= -\ln \prod_{ij} P(N_{SS}^{ij} | N^{ij}, \epsilon_i, \epsilon_j) \\ &= -\sum_{ij} \left[N_{SS}^{ij} \ln(N^{ij}[\epsilon_i(1-\epsilon_j) + (1-\epsilon_i)\epsilon_j]) \right. \\ &\quad \left. - N^{ij}[\epsilon_i(1-\epsilon_j) + (1-\epsilon_i)\epsilon_j] \right] + \text{constant} \end{aligned}$$

- ▶ where the summation over i and j is taken over all p_T and $|\eta|$ bins of both electrons.
- ▶ By minimizing this likelihood, the charge-flip rate can be estimated.

Likelihood Method : Background Subtraction

Both for data and MC

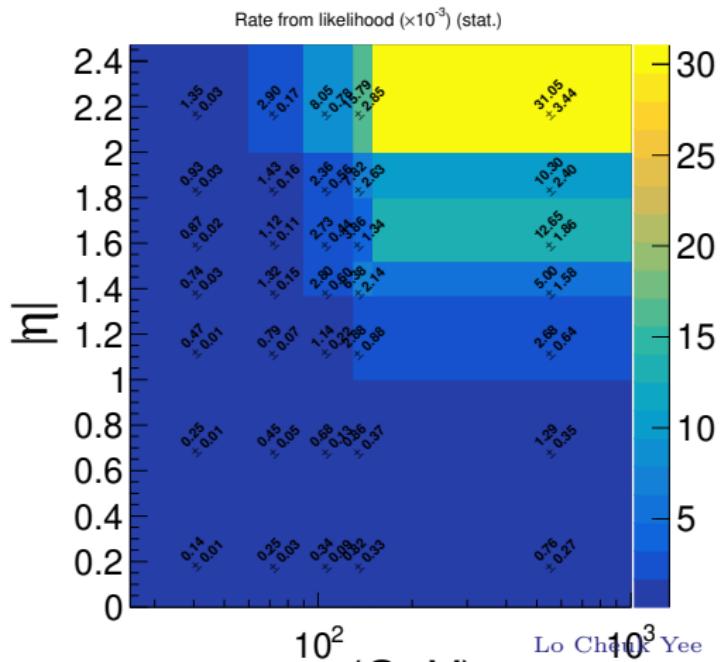
- ▶ Central region: 80 - 100 GeV
- ▶ Sideband region: 60 - 80 and 100 - 120
(Sideband width 20 GeV)

$$N_{Central} = N_{Central} - \frac{20}{20 + 20} N_{SB}$$

- ▶ Background subtraction is done both on N^{ij} and N_{SS}^{ij}

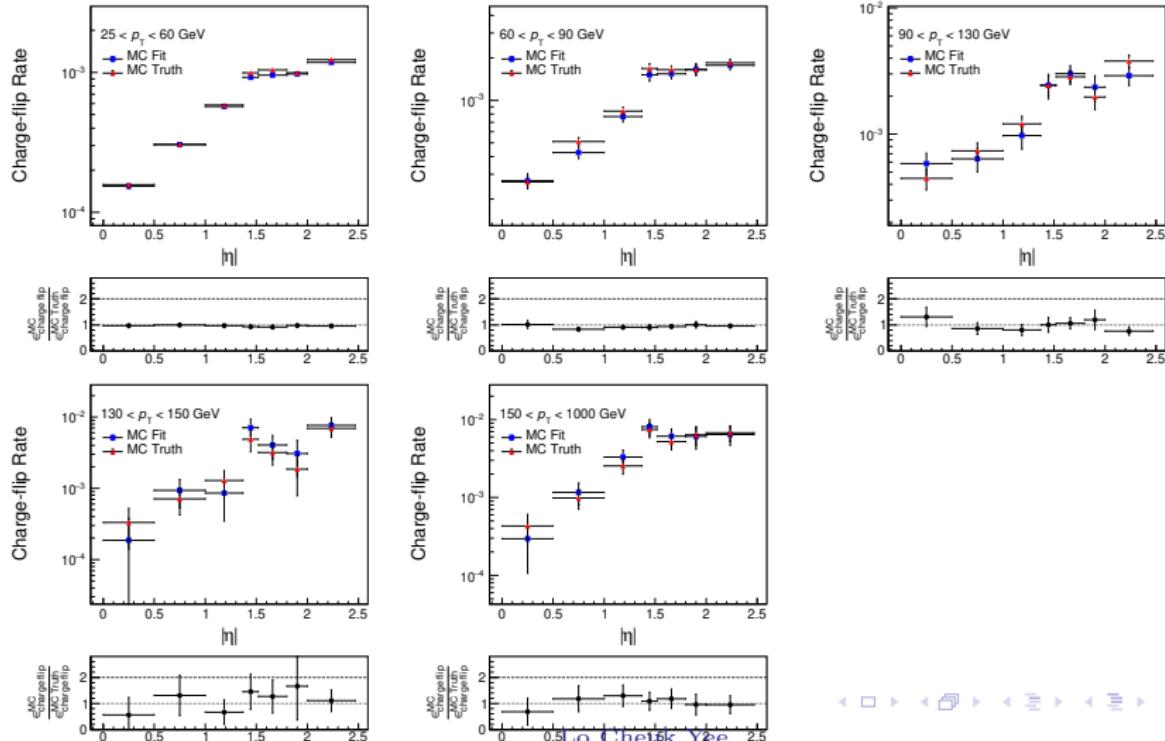
Likelihood Method : Result for Charge Flip Rate (statistical error only)

- ▶ Results by likelihood method from data, $\epsilon_{\text{lik,data}}$
- ▶ The error is statistical only, from likelihood method, $\sigma_{\text{lik,data}}$.



Likelihood Method : MC Truth Validation

- ▶ Results by likelihood method from MC (blue points)
 - ▶ Results by MC truth (red points)

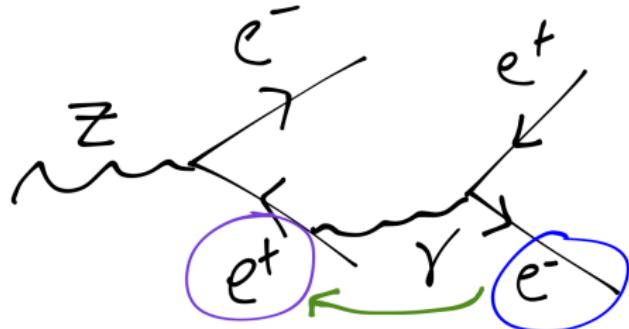


Likelihood Method : MC Truth Matching

- ▶ It is a complicated part, because we do not know the original charge of the reconstructed electron for some cases.
- ▶ The reconstructed electrons match the truth particle by the smallest ΔR ($\Delta R < 0.1$).
- ▶ No match for reconstructed electrons if no any truth particle is inside $\Delta R < 0.1$. (cannot find the original charge)
- ▶ If the truth particle is matched, the truth particle is not electron. (fake lepton, not charge flip)

Likelihood Method : MC Truth Matching

- ▶ If the truth particle is electron, the (grand)mother particle is not Z Boson. (ignore it)
- ▶ If (grand)mother particle is Z Boson, the daughter of the Z Boson is not electron. (ignore it)
- ▶ If the daughter of the Z Boson is electron, the charge of that electron is the original charge.



Likelihood Method : Systematic uncertainties due to likelihood method

For each bin, the systematic uncertainties due to likelihood method are estimated by the difference between the likelihood method and the MC truth method.

For MC,

$$\sigma_{\text{truth,MC}} = |\epsilon_{\text{lik,MC}} - \epsilon_{\text{MC truth}}|$$

For data,

$$\sigma_{\text{truth,data}} = \epsilon_{\text{lik,data}} \times \frac{\sigma_{\text{truth,MC}}}{\epsilon_{\text{lik,MC}}}$$

Likelihood Method : Systematic uncertainties due to background subtraction

Nominal:

- ▶ Central region: 80 - 100 GeV, Sideband width: 20 GeV

Variations:

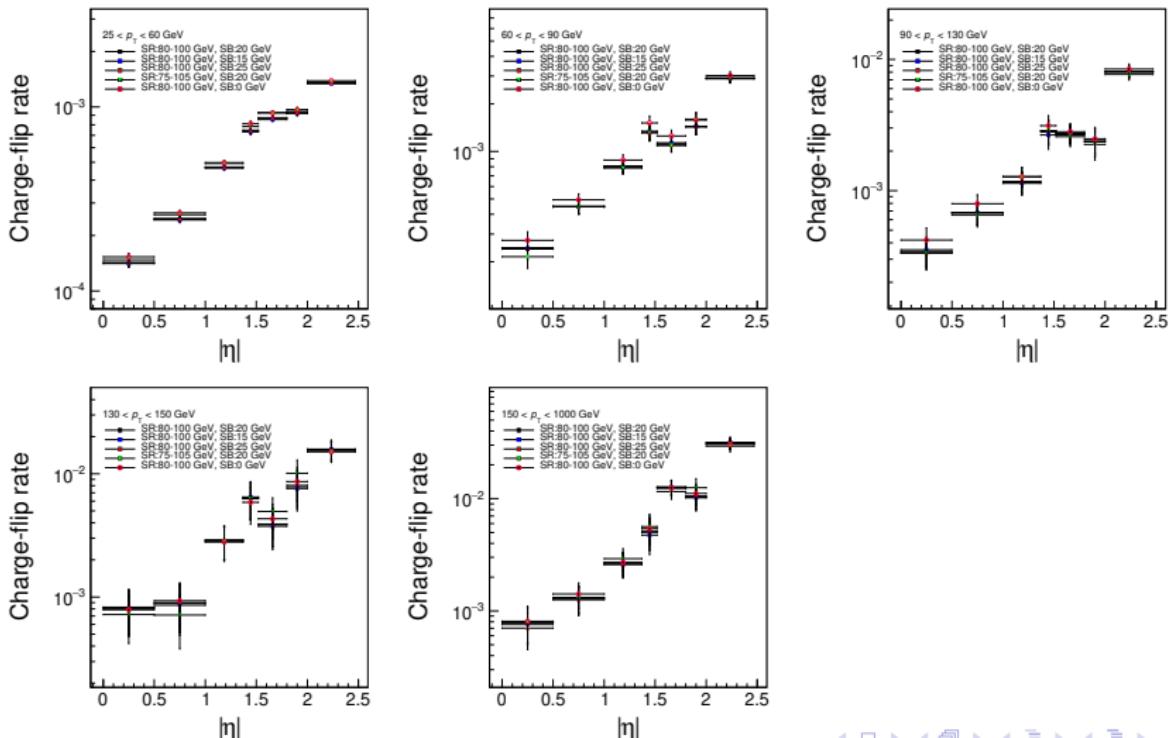
- ▶ Central region: 80 - 100 GeV, Sideband width: 15 GeV
- ▶ Central region: 80 - 100 GeV, Sideband width: 25 GeV
- ▶ Central region: 75 - 105 GeV, Sideband width: 20 GeV
- ▶ Central region: 80 - 100 GeV, no background subtraction

For each bin, the largest deviation from the nominal among these variations is the systematic uncertainty due to background subtraction.

$$\sigma_{\text{bgk}} = \max\{|\sigma_{\text{nominal}} - \sigma_{\text{variation}}|\}$$

Likelihood Method : Result for different variation

► Results for different variations



Likelihood Method : Total uncertainties

$$\sigma_{\text{sys}} = \sqrt{\sigma_{\text{truth}}^2 + \sigma_{\text{bgk}}^2}$$

$$\sigma_{\text{tot}} = \sqrt{\sigma_{\text{lik}}^2 + \sigma_{\text{sys}}^2}$$

Likelihood Method : MC Validation Plots

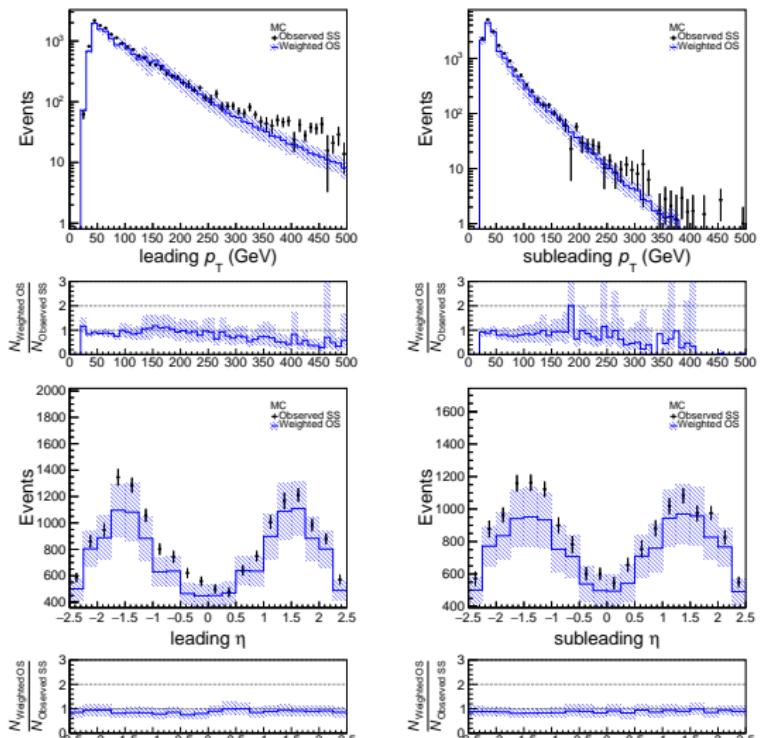


Figure: The comparison between the weighted OS events and the SS events, with different variables.

Likelihood Method : MC Validation Plots

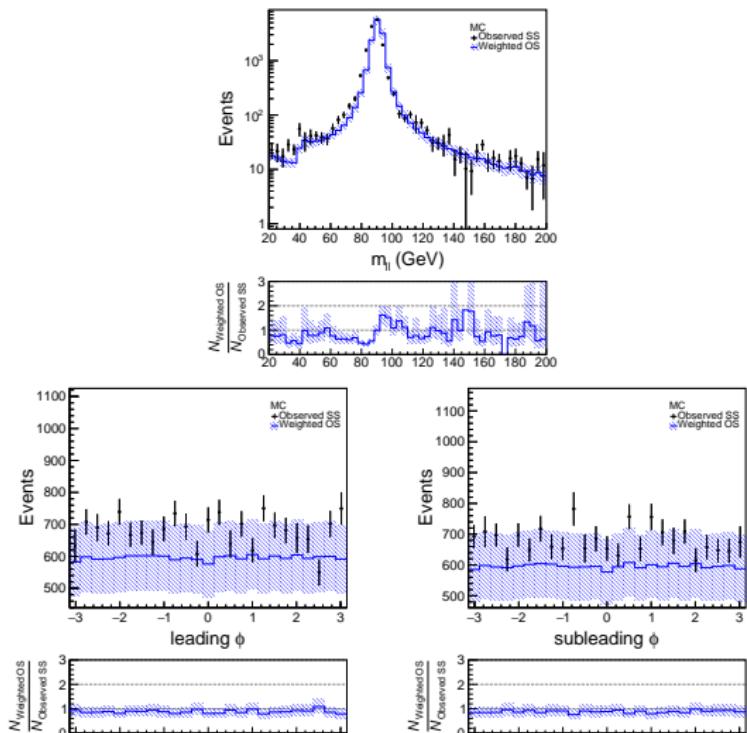


Figure: The comparison between the weighted OS events and the SS events, with different variables.

Fake Lepton Background : Matrix Method

Fake Lepton Background : Matrix Method

We can generalize the one-lepton case to the two-leptons case.

$$\begin{pmatrix} N_{TT} \\ N_{TL} \\ N_{LT} \\ N_{LL} \end{pmatrix} = \begin{pmatrix} \epsilon_1 \epsilon_2 & \epsilon_1 f_2 & f_1 \epsilon_2 & f_1 f_2 \\ \epsilon_1 \bar{\epsilon}_2 & \epsilon_1 \bar{f}_2 & f_1 \bar{\epsilon}_2 & f_1 \bar{f}_2 \\ \bar{\epsilon}_1 \epsilon_2 & \bar{\epsilon}_1 f_2 & \bar{f}_1 \epsilon_2 & \bar{f}_1 f_2 \\ \bar{\epsilon}_1 \bar{\epsilon}_2 & \bar{\epsilon}_1 \bar{f}_2 & \bar{f}_1 \bar{\epsilon}_2 & \bar{f}_1 \bar{f}_2 \end{pmatrix} \begin{pmatrix} N_{RR} \\ N_{RF} \\ N_{FR} \\ N_{FF} \end{pmatrix}$$

where ϵ_1 is the probability that a leading real lepton passes the signal selection (i.e. tight lepton), ϵ_2 is the probability that a subleading real lepton passes the signal selection (i.e. tight lepton), etc.

Fake Lepton Background : Matrix Method

We can find N_{RF} , N_{FR} , N_{FF} by inverting the matrix.

$$\begin{pmatrix} 0 \\ N_{RF} \\ N_{FR} \\ N_{FF} \end{pmatrix} = \begin{pmatrix} 0 & 0 & 0 & 0 \\ 0 & 1 & 0 & 0 \\ 0 & 0 & 1 & 0 \\ 0 & 0 & 0 & 1 \end{pmatrix} \begin{pmatrix} N_{RR} \\ N_{RF} \\ N_{FR} \\ N_{FF} \end{pmatrix}$$

$$= \begin{pmatrix} 0 & 0 & 0 & 0 \\ 0 & 1 & 0 & 0 \\ 0 & 0 & 1 & 0 \\ 0 & 0 & 0 & 1 \end{pmatrix} \begin{pmatrix} \epsilon_1\epsilon_2 & \epsilon_1f_2 & f_1\epsilon_2 & f_1f_2 \\ \epsilon_1\bar{\epsilon}_2 & \epsilon_1\bar{f}_2 & f_1\bar{\epsilon}_2 & f_1\bar{f}_2 \\ \bar{\epsilon}_1\epsilon_2 & \bar{\epsilon}_1f_2 & \bar{f}_1\epsilon_2 & \bar{f}_1f_2 \\ \bar{\epsilon}_1\bar{\epsilon}_2 & \bar{\epsilon}_1\bar{f}_2 & \bar{f}_1\bar{\epsilon}_2 & \bar{f}_1\bar{f}_2 \end{pmatrix}^{-1} \begin{pmatrix} N_{TT} \\ N_{TL} \\ N_{LT} \\ N_{LL} \end{pmatrix}$$

Fake Lepton Background : Matrix Method

Then, we can find the number of tight-tight leptons due to the fake leptons, N'_{TT}

$$\begin{pmatrix} N'_{TT} \\ N'_{TL} \\ N'_{LT} \\ N'_{LL} \end{pmatrix} = \begin{pmatrix} \epsilon_1\epsilon_2 & \epsilon_1f_2 & f_1\epsilon_2 & f_1f_2 \\ \epsilon_1\bar{\epsilon}_2 & \epsilon_1\bar{f}_2 & f_1\bar{\epsilon}_2 & f_1\bar{f}_2 \\ \bar{\epsilon}_1\epsilon_2 & \bar{\epsilon}_1f_2 & \bar{f}_1\epsilon_2 & \bar{f}_1f_2 \\ \bar{\epsilon}_1\bar{\epsilon}_2 & \bar{\epsilon}_1\bar{f}_2 & \bar{f}_1\bar{\epsilon}_2 & \bar{f}_1\bar{f}_2 \end{pmatrix} \begin{pmatrix} 0 \\ N_{RF} \\ N_{FR} \\ N_{FF} \end{pmatrix}$$

$$= \begin{pmatrix} \epsilon_1\epsilon_2 & \epsilon_1f_2 & f_1\epsilon_2 & f_1f_2 \\ \epsilon_1\bar{\epsilon}_2 & \epsilon_1\bar{f}_2 & f_1\bar{\epsilon}_2 & f_1\bar{f}_2 \\ \bar{\epsilon}_1\epsilon_2 & \bar{\epsilon}_1f_2 & \bar{f}_1\epsilon_2 & \bar{f}_1f_2 \\ \bar{\epsilon}_1\bar{\epsilon}_2 & \bar{\epsilon}_1\bar{f}_2 & \bar{f}_1\bar{\epsilon}_2 & \bar{f}_1\bar{f}_2 \end{pmatrix} \begin{pmatrix} 0 & 0 & 0 & 0 \\ 0 & 1 & 0 & 0 \\ 0 & 0 & 1 & 0 \\ 0 & 0 & 0 & 1 \end{pmatrix}$$

$$\begin{pmatrix} \epsilon_1\epsilon_2 & \epsilon_1f_2 & f_1\epsilon_2 & f_1f_2 \\ \epsilon_1\bar{\epsilon}_2 & \epsilon_1\bar{f}_2 & f_1\bar{\epsilon}_2 & f_1\bar{f}_2 \\ \bar{\epsilon}_1\epsilon_2 & \bar{\epsilon}_1f_2 & \bar{f}_1\epsilon_2 & \bar{f}_1f_2 \\ \bar{\epsilon}_1\bar{\epsilon}_2 & \bar{\epsilon}_1\bar{f}_2 & \bar{f}_1\bar{\epsilon}_2 & \bar{f}_1\bar{f}_2 \end{pmatrix}^{-1} \begin{pmatrix} N_{TT} \\ N_{TL} \\ N_{LT} \\ N_{LL} \end{pmatrix}$$

Method for measuring real efficiencies

Control Region for Real Efficiencies

It uses the Z tag-and-probe method to measure the real efficiencies of electrons and muons.

Pre-selection:

- ▶ Trigger requirement
- ▶ exactly 2 baseline leptons
- ▶ same flavour and opposite charge ($Z \rightarrow ee$, $Z \rightarrow \mu\mu$)
- ▶ $80 < m_{ll} < 100$
- ▶ $p_T > 20$
- ▶ $|\eta| < 2.47$ for electron, $|\eta| < 2.5$ for muon

Requirements for tag lepton

- ▶ $p_T > 25$
- ▶ pass signal requirement

Permutation of tag and probe

Check the probe lepton:

Tight lepton: pass signal requirement

Loose lepton: pass baseline requirement (include tight lepton)

4 possible cases: (slightly simplified)

- ▶ TL: fill loose hist by the subleading lepton
- ▶ LT: fill loose hist by the leading lepton
- ▶ TT: fill both loose and tight hist by the both leading and subleading lepton
- ▶ LL: do nothing

The real eff is calculated by

$$\epsilon = \frac{N_{\text{tight}}}{N_{\text{loose}}}$$

Fake Lepton Background : Results of Real Efficiencies

The method for measuring the real efficiencies is described in the backup slides.

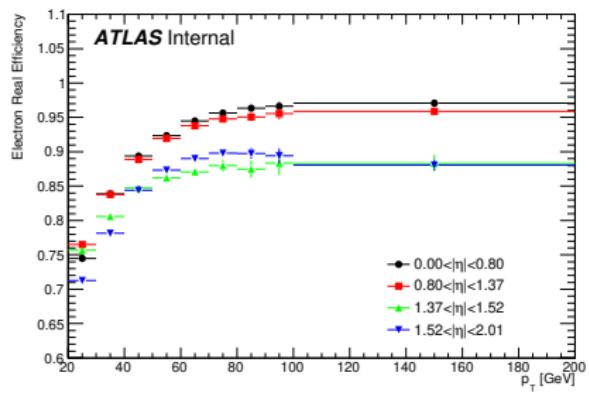


Figure: Electron real efficiency

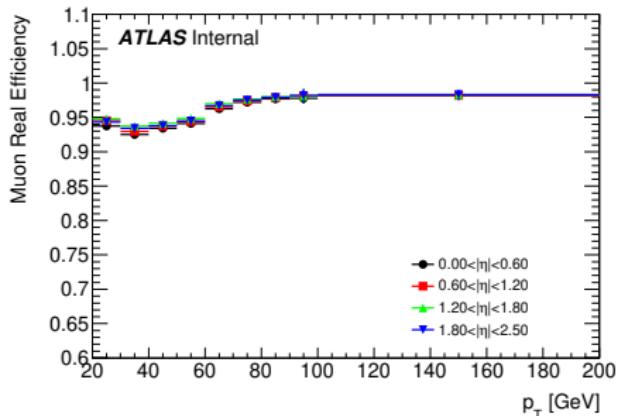


Figure: Muon real efficiency

Method for measuring fake efficiencies

Control Region for fake eff for muon

It uses the heavy-flavour fake-enriched region to measure the fake efficiencies of electrons and muons.

For muon fake eff, it uses the muon-muon events.

Pre-selection for muon-muon events:

- ▶ Trigger requirement
- ▶ exactly 2 baseline muons
- ▶ SS
- ▶ at least 1 jet
- ▶ at least 1 b-jet
- ▶ $p_T > 20$
- ▶ $|\eta| < 2.5$

Requirements for the tag lepton

- ▶ $p_T > 40$
- ▶ pass signal requirement

Control Region for fake eff for electron

For electron fake eff, it uses the muon-electron events.

Pre-selection for muon-electron events:

- ▶ Trigger requirement
- ▶ exactly 1 baseline electron and 1 baseline muon
- ▶ SS
- ▶ at least 1 jet
- ▶ at least 1 b-jet
- ▶ $p_T > 20$
- ▶ $|\eta| < 2.5$

Requirements for the tag lepton

- ▶ muon
- ▶ $p_T > 40$
- ▶ pass signal requirement

Calculate the fake eff

The fake eff is calculated by

$$\epsilon = \frac{N_{\text{tight}}^{\text{data}} - N_{\text{tight}}^{\text{prompt MC}}}{N_{\text{loose}}^{\text{data}} - N_{\text{loose}}^{\text{prompt MC}}}$$

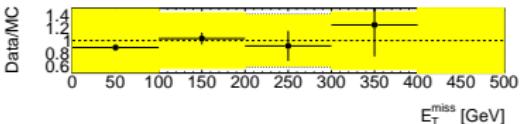
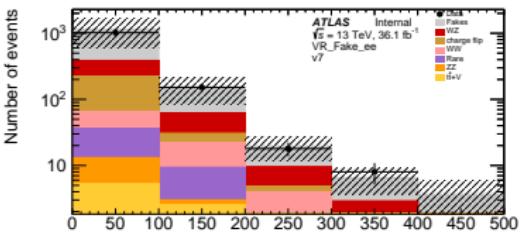
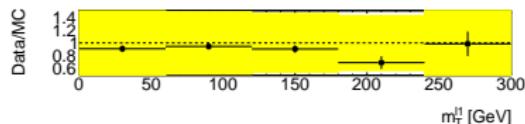
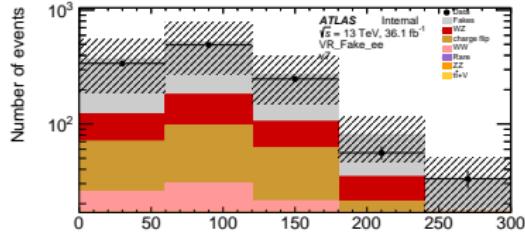
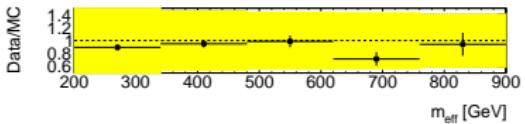
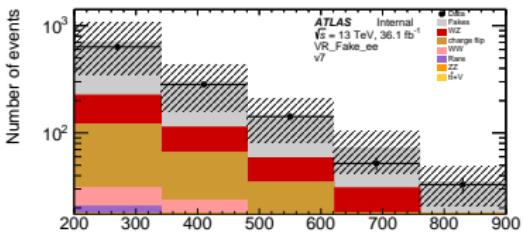
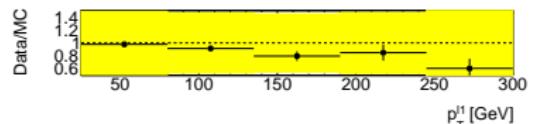
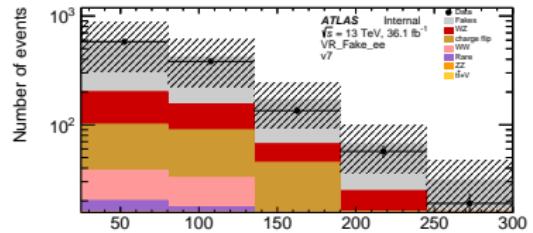
The prompt lepton is estimated from MC by using the truth information. The prompt lepton is defined by the MCTruthClassifier, with the ParticleType = IsoElectron or IsoMuon.

Validation plots for fake lepton background

ee channel

	Number of events	
Fakes	$860.738787 \pm 14.359660 \pm 630.863451$ (9663)	
WZ	201.069385 ± 3.094538 (39857)	
charge flip	167.954322 ± 1.110819 (86972)	
WW	45.176809 ± 0.401964 (21140)	
Rare	30.757392 ± 3.403960 (1268)	
ZZ	8.351345 ± 0.345416 (6466)	
$t\bar{t} + V$	8.346022 ± 0.254296 (4158)	
Total BG	$1322.394062 \pm 15.130843 \pm 630.863451$ (169524)	
Data	1201	

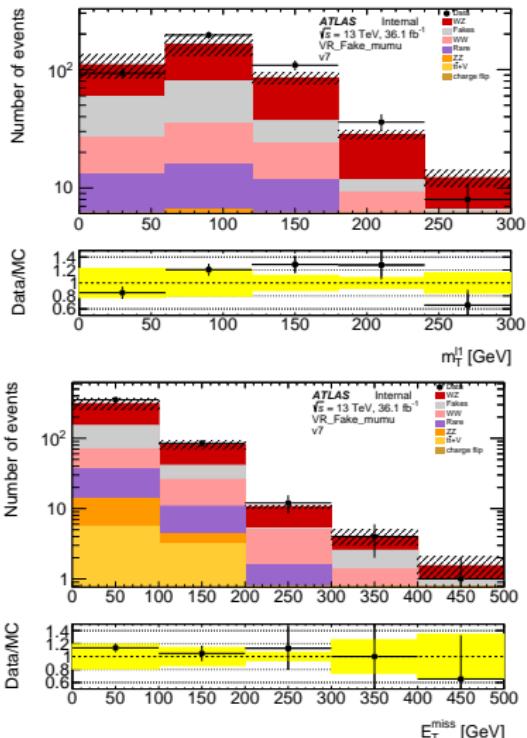
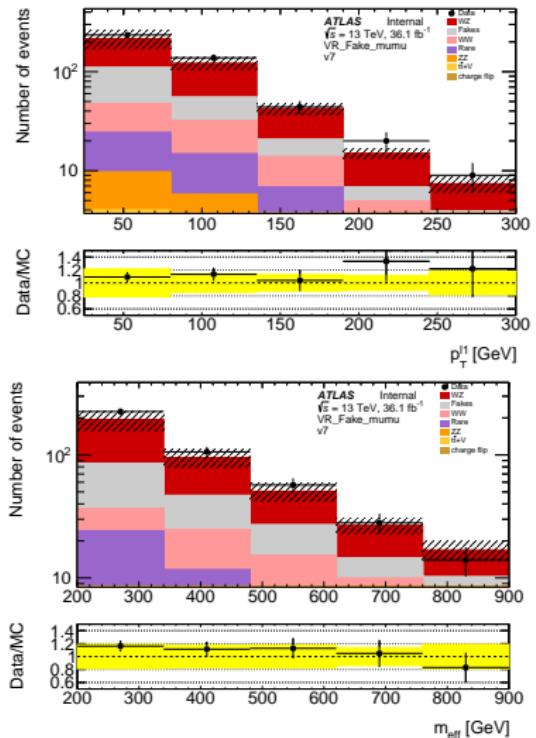
ee channel



mumu channel

	Number of events	
WZ	209.219829 ± 2.994737 (42458)	
Fakes	$96.203916 \pm 5.013796 \pm 70.510979$ (1160)	
WW	53.675707 ± 0.409851 (25092)	
Rare	30.460547 ± 4.359663 (1361)	
ZZ	9.474511 ± 0.534342 (7674)	
$t\bar{t} + V$	9.369165 ± 0.284944 (4801)	
charge flip	0.000000 ± 0.000000 (0)	
Total BG	$408.403674 \pm 7.324477 \pm 70.510979$ (82546)	
Data	454	

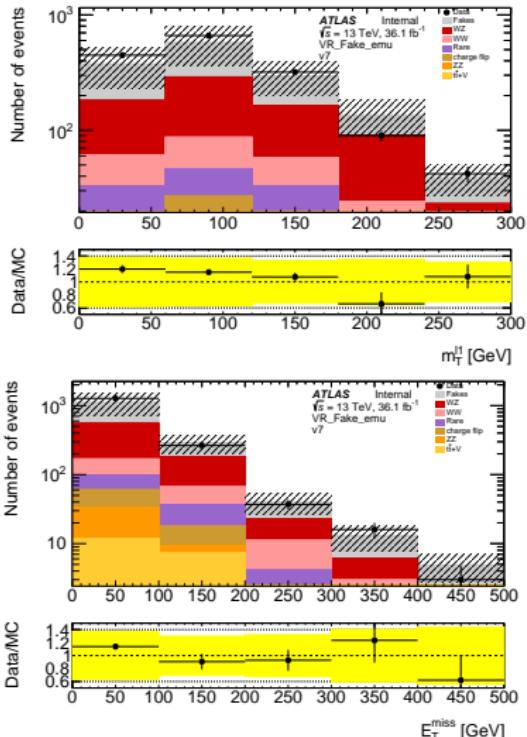
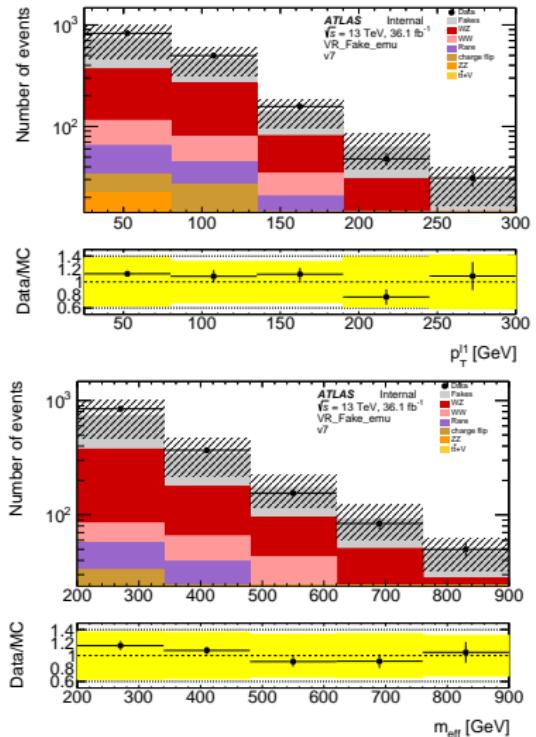
mumu channel



emu channel

	Number of events	
Fakes	$702.995427 \pm 13.784203 \pm 515.248212$ (8544)	
WZ	516.828375 ± 33.696304 (94571)	
WW	110.280002 ± 0.578575 (51615)	
Rare	58.742111 ± 4.005608 (3037)	
charge flip	36.901490 ± 0.479301 (36998)	
ZZ	23.391749 ± 0.482387 (16698)	
$t\bar{t} + V$	20.739252 ± 0.424107 (10526)	
Total BG	$1469.878407 \pm 36.639692 \pm 515.248212$ (221989)	
Data	1593	

emu channel



Pre-selection plots

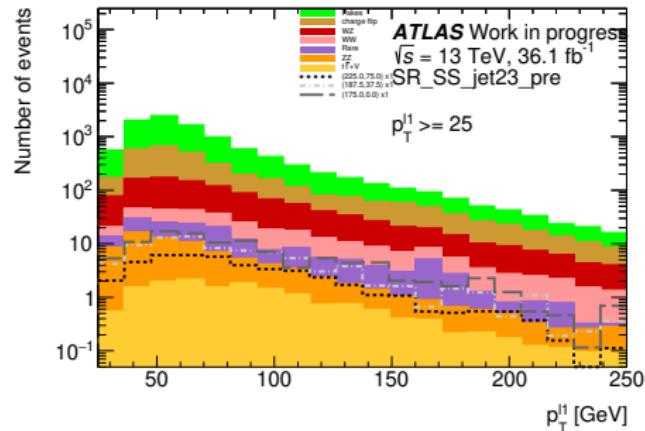
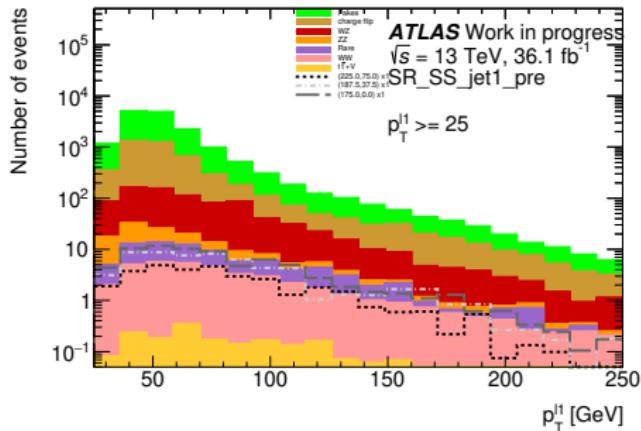
Yields in pre-selection

Number of events	
Fakes	$11580.965417 \pm 45.775452$ (107219)
charge flip	3622.756937 ± 2.810481 (3375964)
WZ	691.333566 ± 34.238864 (68626)
ZZ	66.945860 ± 1.049840 (21749)
Rare	47.371873 ± 4.082859 (1662)
WW	38.128141 ± 0.388221 (15666)
$t\bar{t} + V$	1.971267 ± 0.138488 (671)
Total BG	$16049.473060 \pm 57.389305$ (3591557)
(225.0,75.0)	33.08 ± 1.57 (676)
(187.5,37.5)	62.09 ± 3.13 (568)
(175.0,0.0)	74.30 ± 3.57 (556)

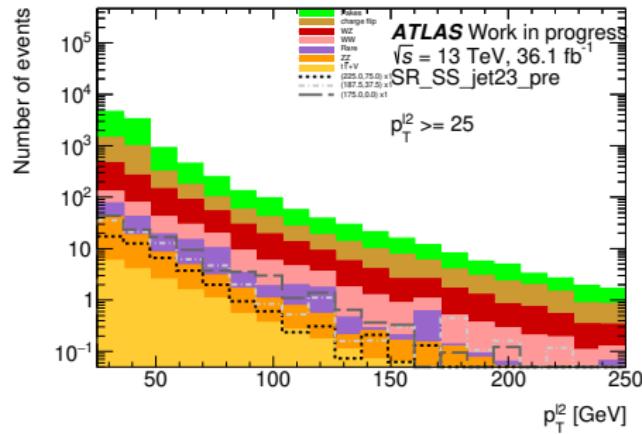
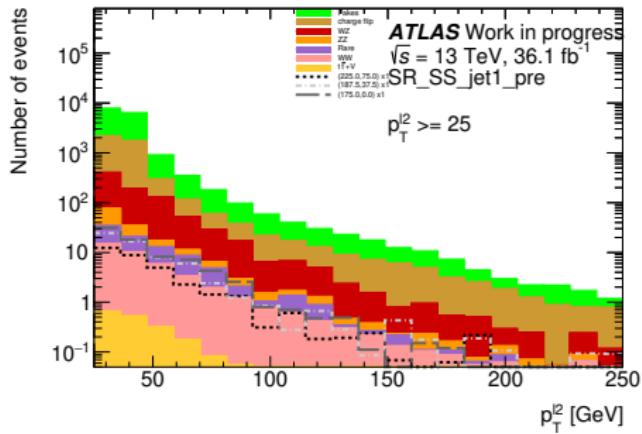
Table: Yields in SRjet1

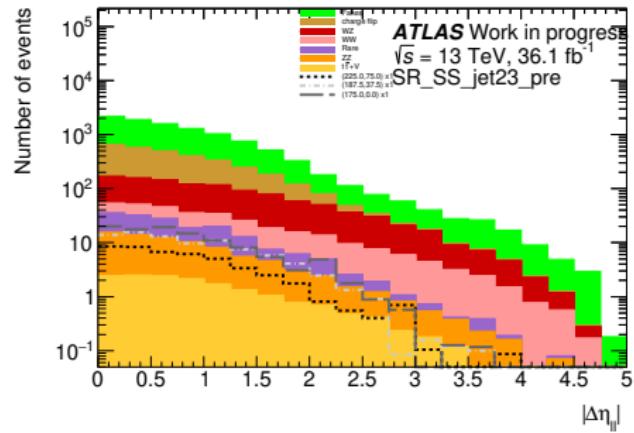
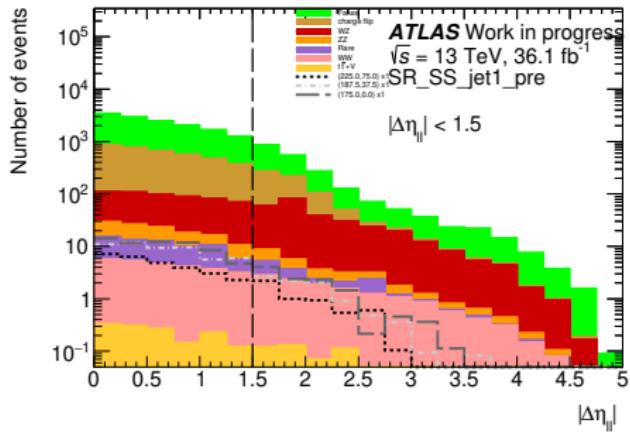
Number of events	
Fakes	$6848.284248 \pm 36.074557$ (66096)
charge flip	2074.427691 ± 2.680118 (1679673)
WZ	776.282836 ± 5.978666 (147413)
WW	153.455294 ± 0.679752 (72175)
Rare	92.181843 ± 6.338615 (3011)
ZZ	62.932966 ± 1.198123 (35945)
$t\bar{t} + V$	16.375388 ± 0.359456 (6082)
Total BG	$10023.940266 \pm 37.235816$ (2010395)
(225.0,75.0)	45.01 ± 1.76 (937)
(187.5,37.5)	86.02 ± 3.84 (754)
(175.0,0.0)	107.96 ± 5.14 (743)

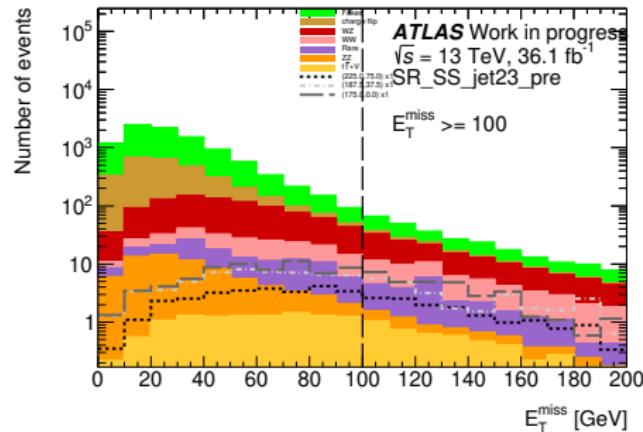
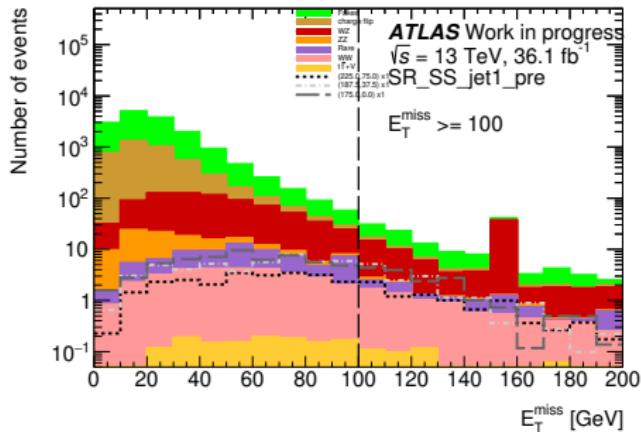
Table: Yields in SRjet23

p_T^{l1} 

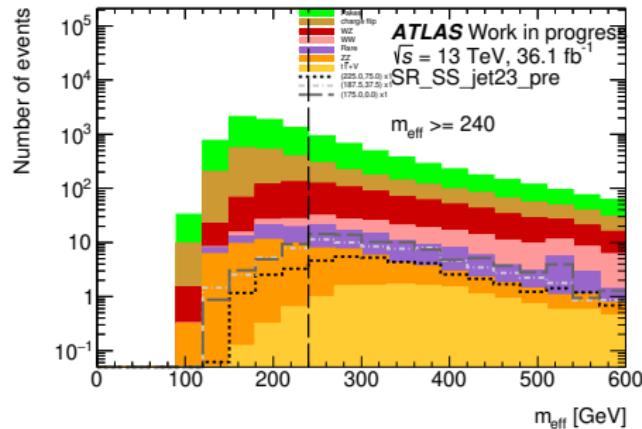
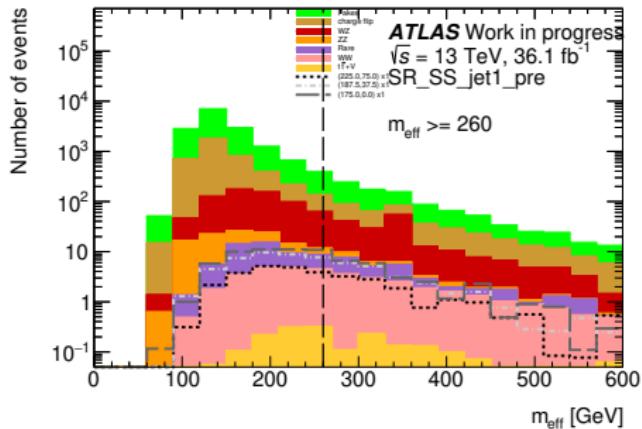
p_T^{l2}



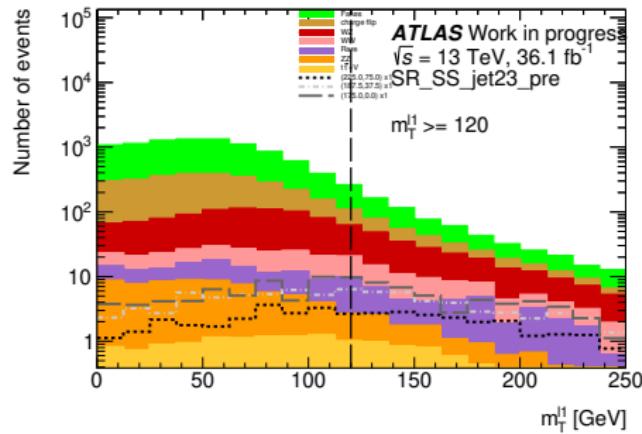
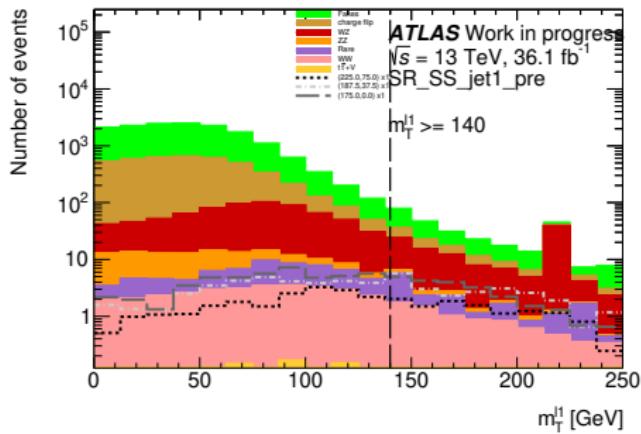
$|\Delta\eta_{ll}|$ 

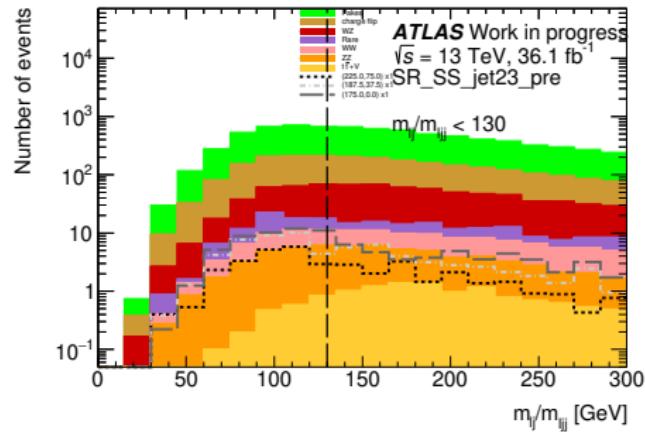
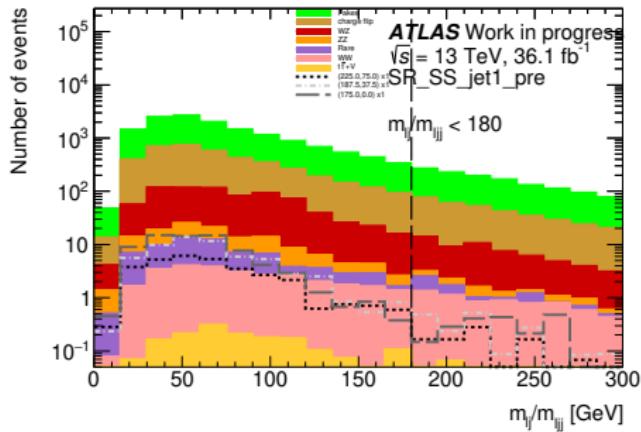
E_T^{miss} 

m_{eff}

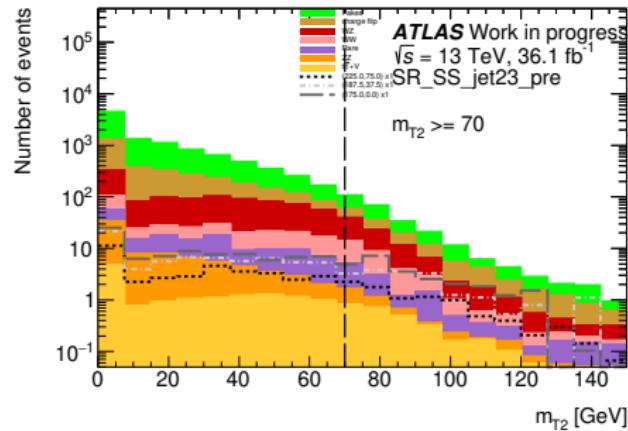
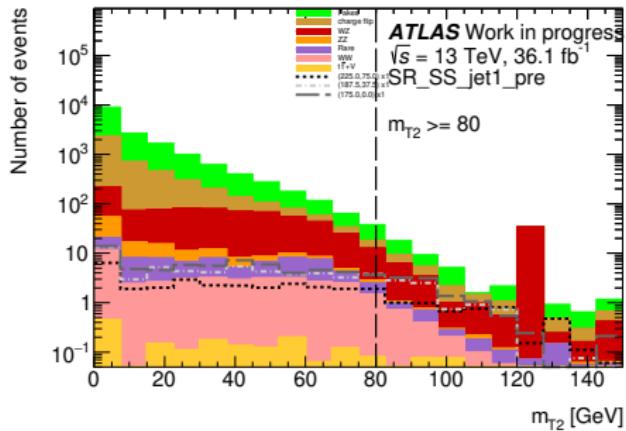


$$m_{\mathrm{T}}^{l_1}$$



m_{lj}/m_{ljj}


$$m_{T2}$$



N-1 plots

Yields in SR

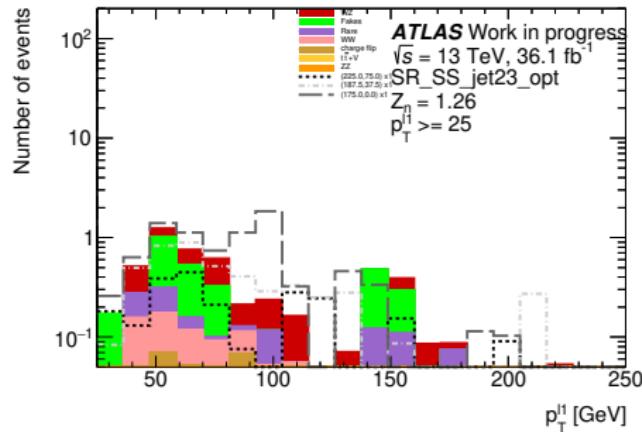
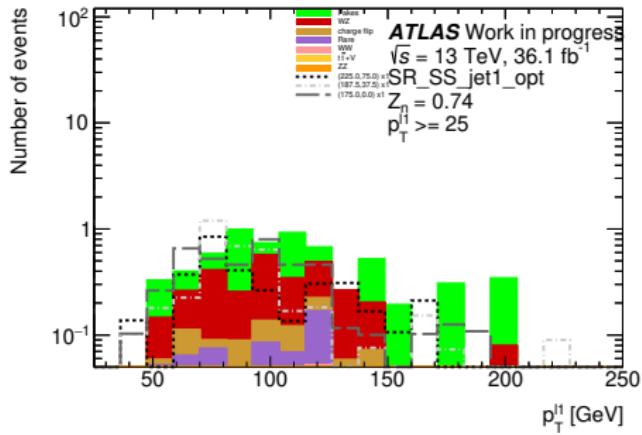
	Number of events	Significance
Fakes	3.295434 ± 0.819175 (24)	
WZ	2.176292 ± 0.397871 (257)	
charge flip	0.472185 ± 0.052737 (265)	
Rare	0.443536 ± 0.111301 (56)	
WW	0.165504 ± 0.023455 (67)	
$t\bar{t} + V$	0.124898 ± 0.046227 (36)	
ZZ	0.055330 ± 0.027713 (31)	
Total BG	6.733178 ± 0.920855 (736)	
(225.0,75.0)	3.33 ± 0.60 (60)	0.74
(187.5,37.5)	3.77 ± 0.95 (31)	0.86
(175.0,0.0)	4.29 ± 0.73 (36)	0.98

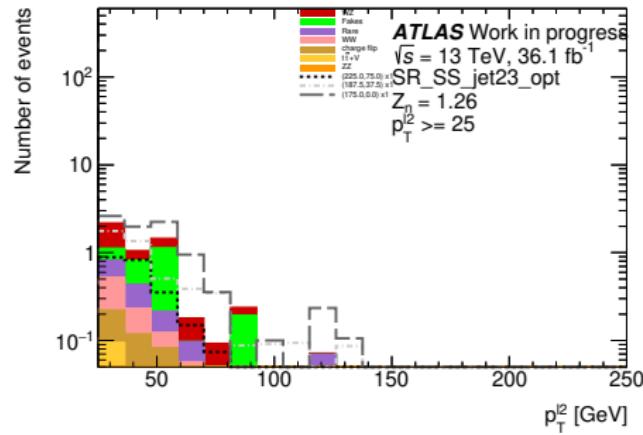
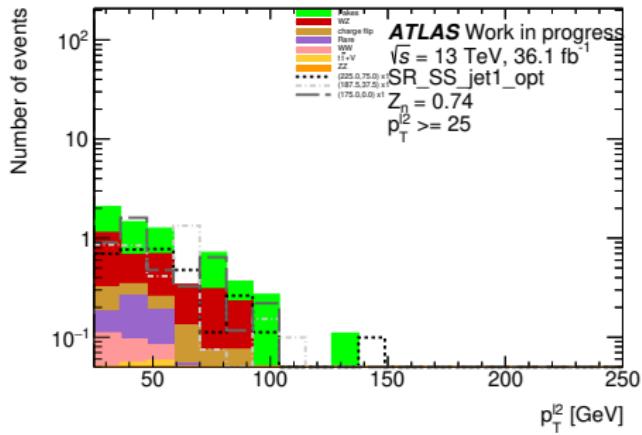
Table: Yields in SRjet1

	Number of events	Significance
WZ	1.848838 ± 0.273025 (319)	
Fakes	1.764693 ± 0.709215 (20)	
Rare	0.731015 ± 0.195114 (48)	
WW	0.514488 ± 0.036878 (235)	
charge flip	0.267485 ± 0.029086 (274)	
$t\bar{t} + V$	0.141679 ± 0.030772 (67)	
ZZ	0.067226 ± 0.024995 (24)	
Total BG	5.335425 ± 0.787005 (987)	
(225.0,75.0)	2.35 ± 0.34 (58)	0.57
(187.5,37.5)	4.72 ± 0.74 (47)	1.26
(175.0,0.0)	8.60 ± 1.51 (58)	2.24

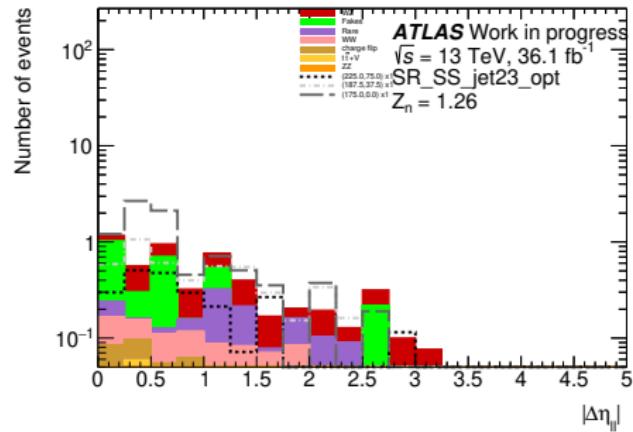
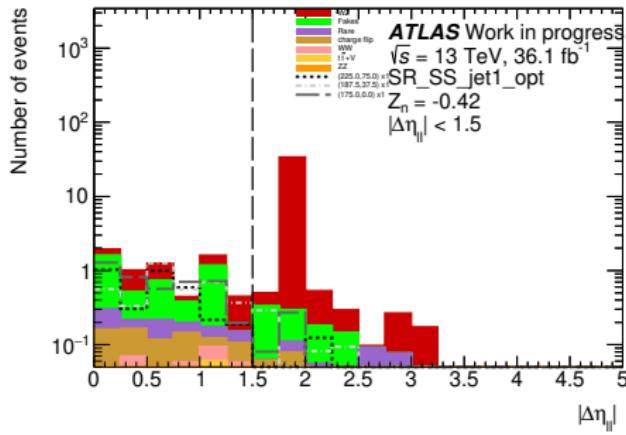
Table: Yields in SRjet23

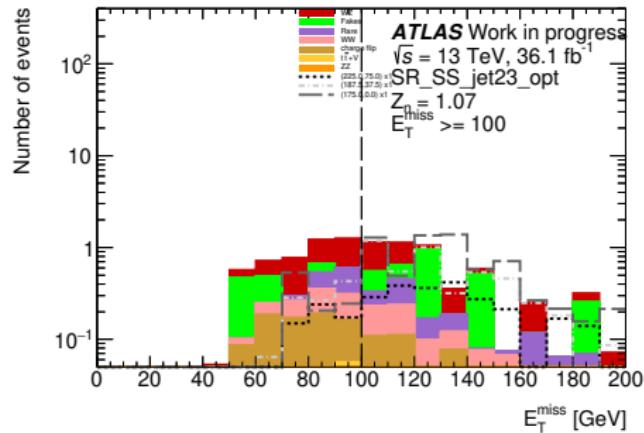
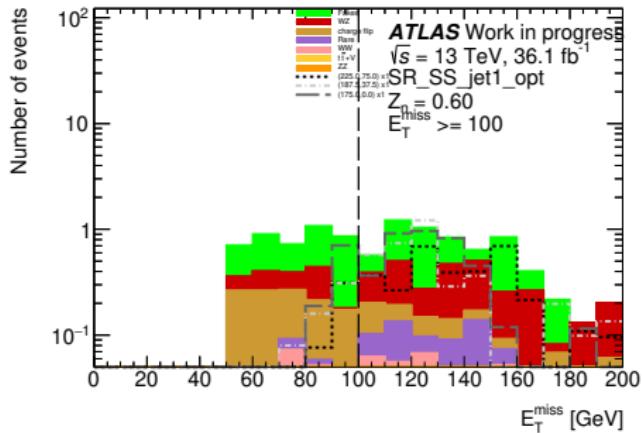
$$p_T^{l1}$$



p_T^{l2} 

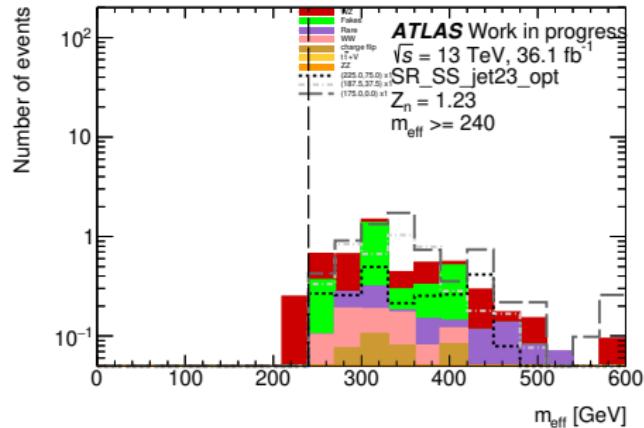
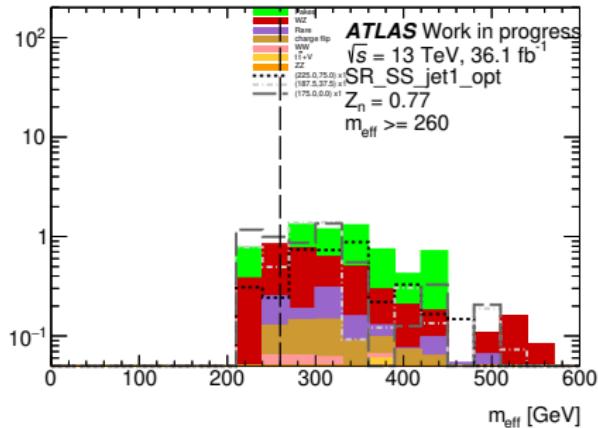
$$|\Delta\eta_{ll}|$$



E_T^{miss} 

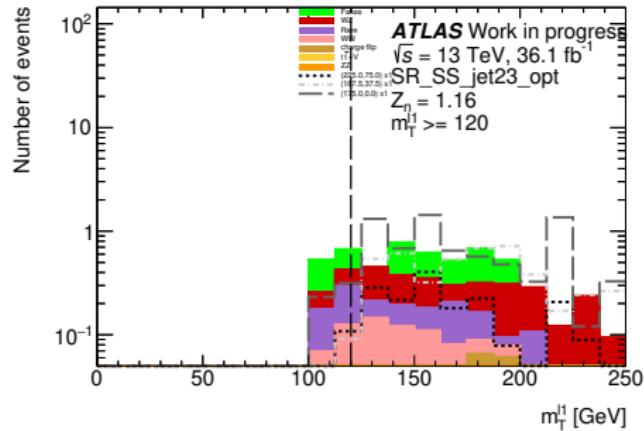
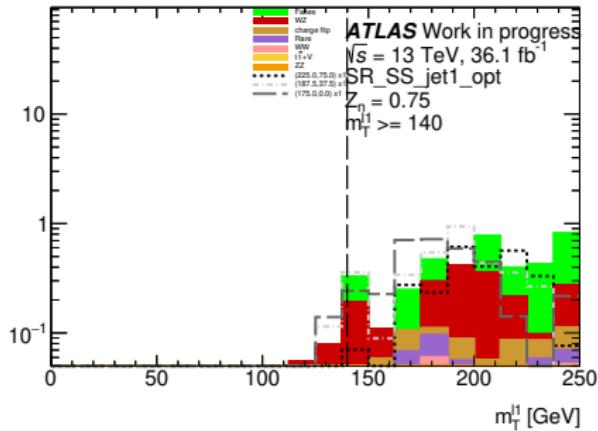
m_{eff}

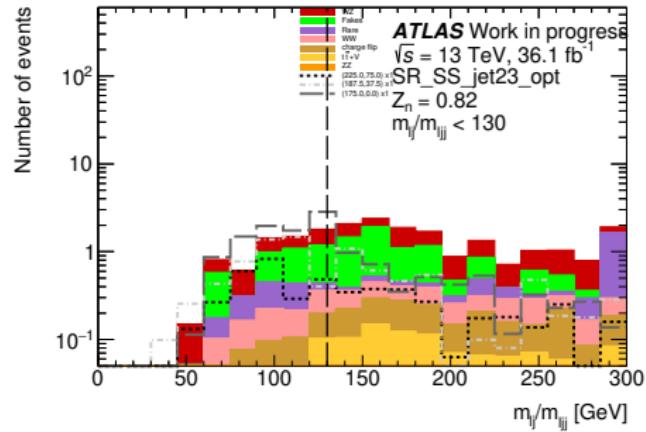
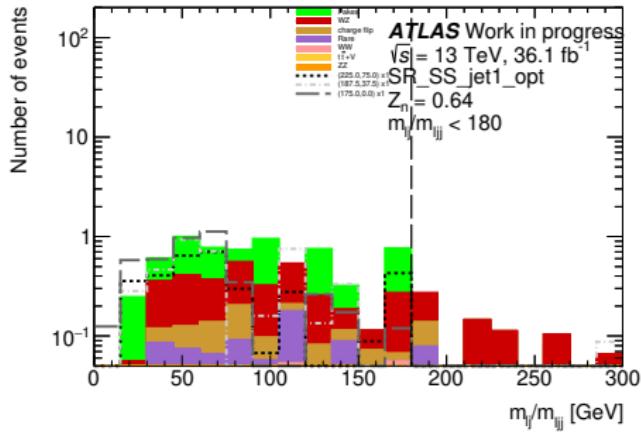
Number of events

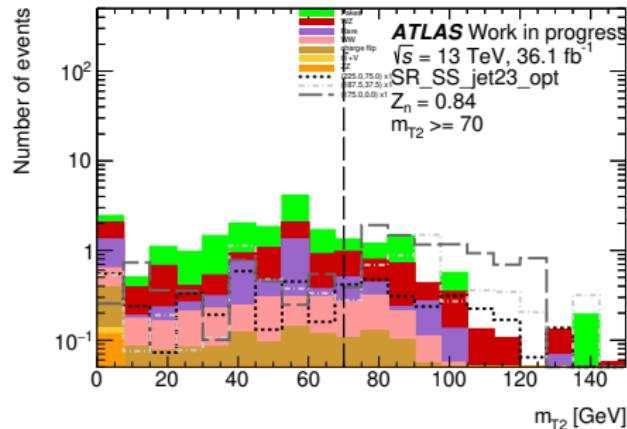
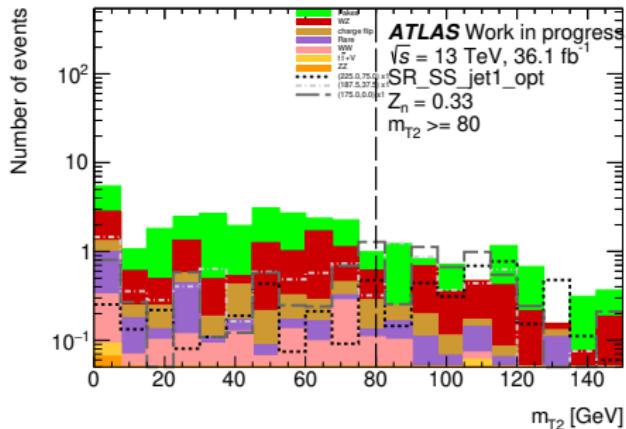


m_T^{l1}

Number of events



m_{lj}/m_{ljj}


m_{T2} 

Validation Region

Validation Region: Check for Signal Contribution

The signal contribution is below 12% and 10% for VRjet1 and VRjet23 respectively.

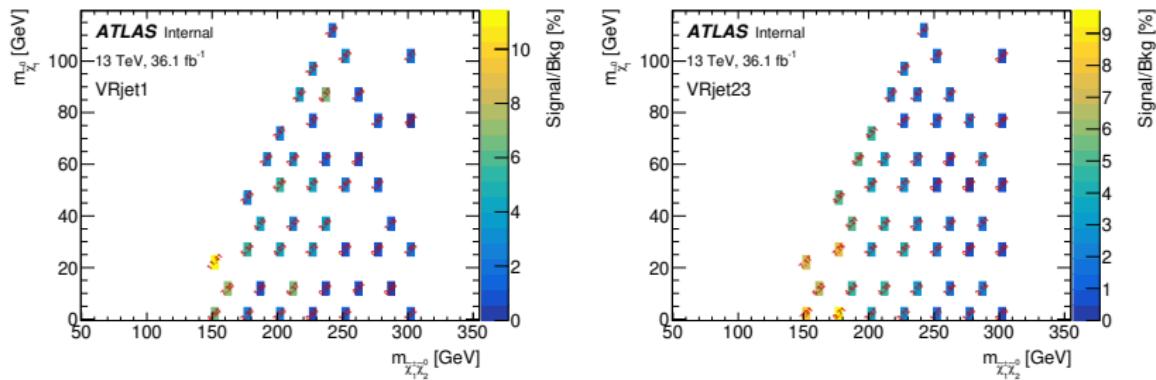


Figure: Percentages of signal contribution are shown for the VRjet1 (left) and VRjet23 (right).

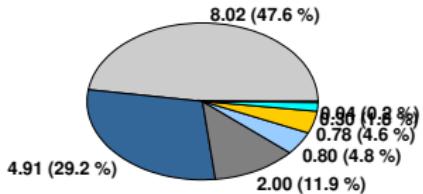
Validation Region: Check for Background Composition

The background compositions for VRjet1 and SRjet1 are similar.

ATLAS Internal
13 TeV, 36.1 fb⁻¹

VRjet1

Rare ZZ WW Charge flip Fakes



ATLAS Internal
13 TeV, 36.1 fb⁻¹

SRjet1

Rare ZZ WW Charge flip Fakes

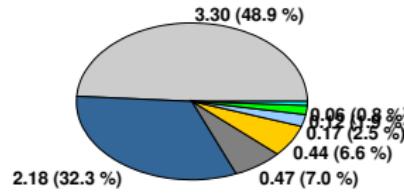


Figure: Background composition for the VRjet1 (left) and SRjet1 (right).

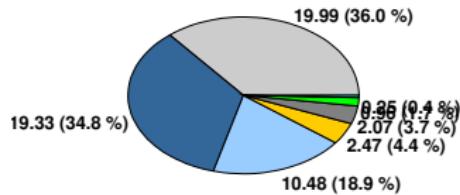
Validation Region: Check for Background Composition

The background compositions for VRjet23 and SRjet23 are similar.

ATLAS Internal
13 TeV, 36.1 fb⁻¹

VRjet23

Rare $t\bar{t}V$
 ZZ WZ
 WW Charge flip
Fakes



ATLAS Internal
13 TeV, 36.1 fb⁻¹

SRjet23

Rare $t\bar{t}V$
 ZZ WZ
 WW Charge flip
Fakes

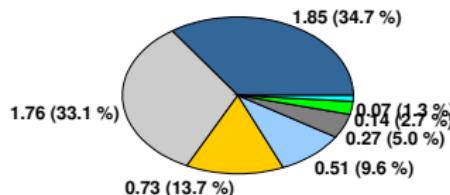
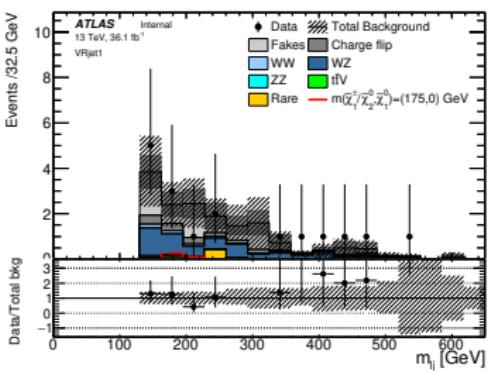
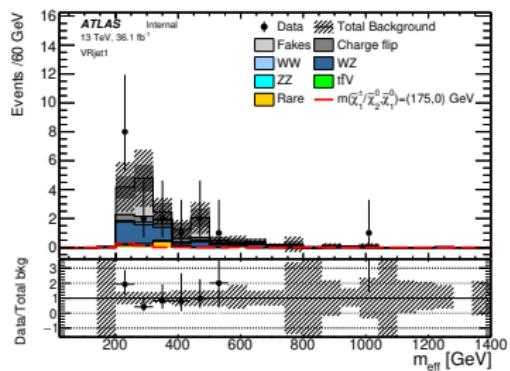
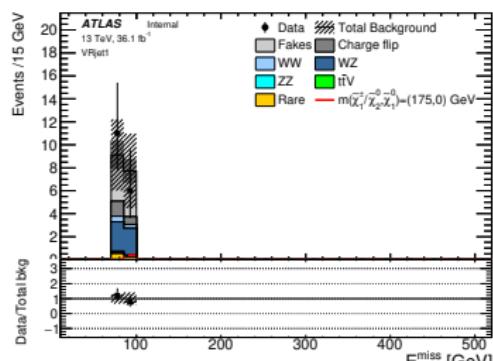
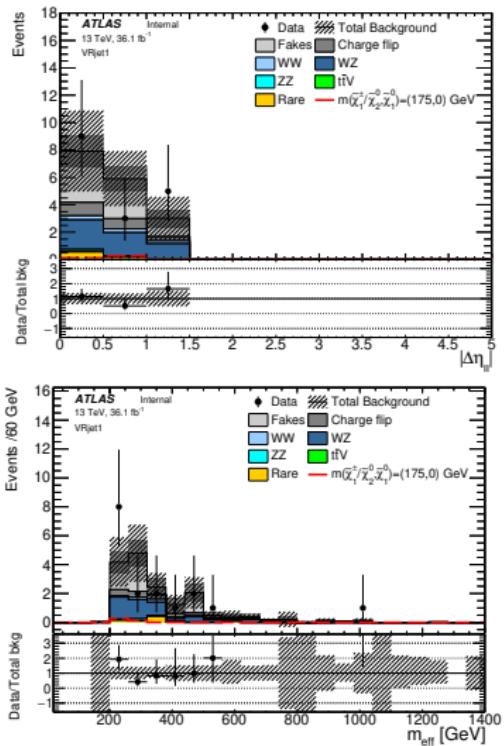
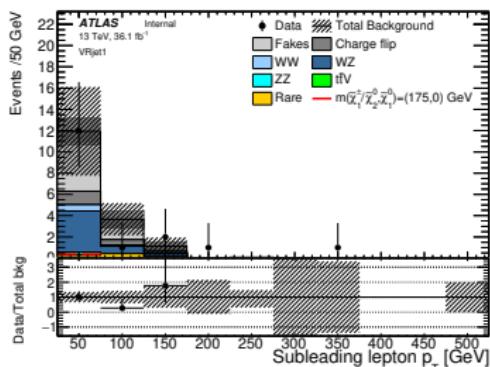
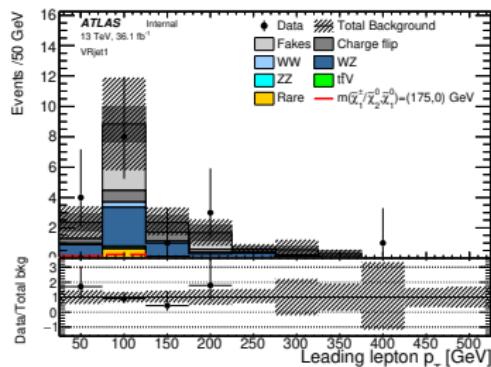
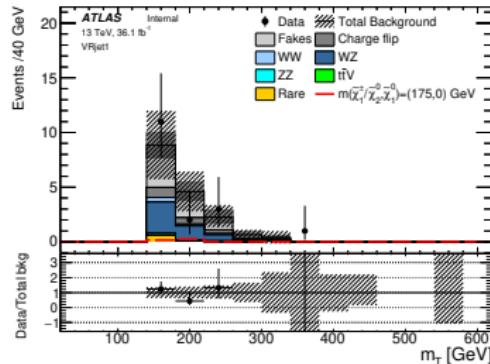
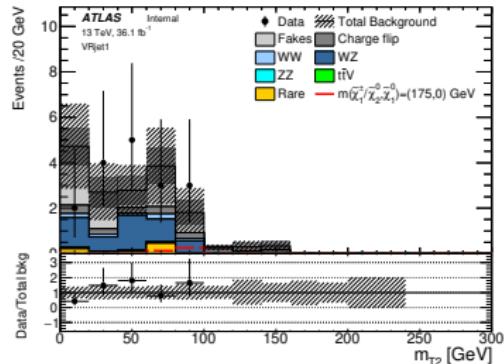


Figure: Background composition for the VRjet23 (left) and SRjet23 (right).

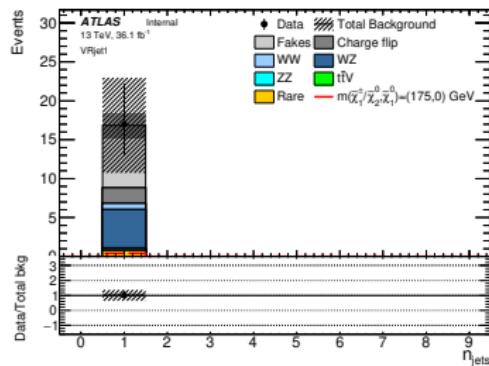
Validation Region : Plots for different variables (VRjet1)



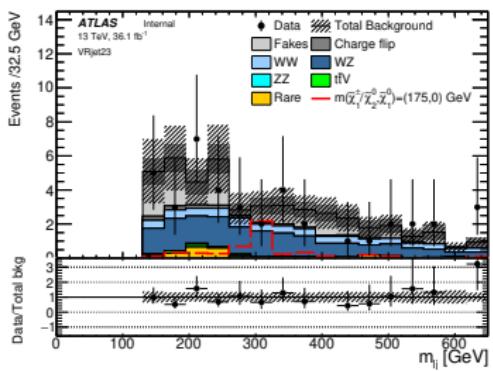
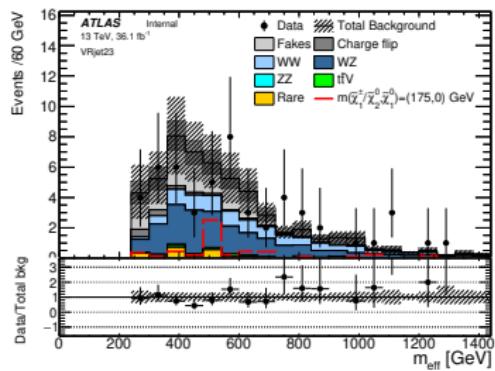
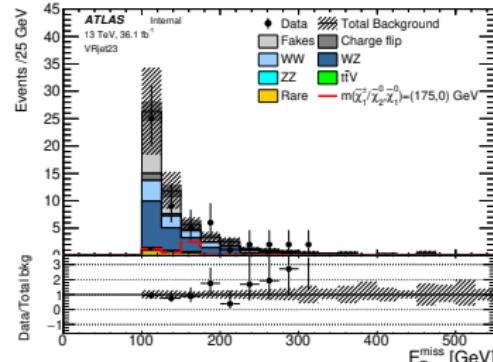
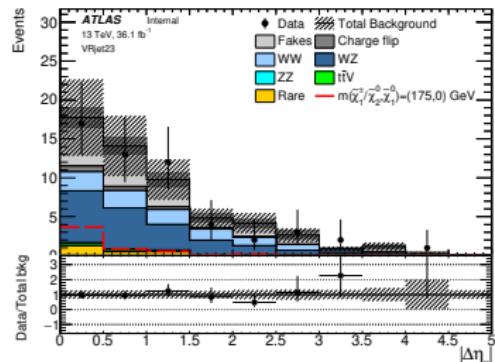
Validation Region : Plots for different variables (VRjet1)



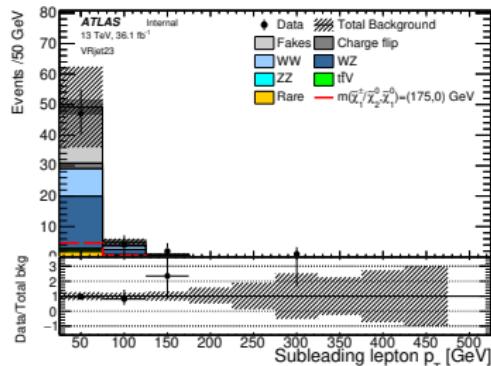
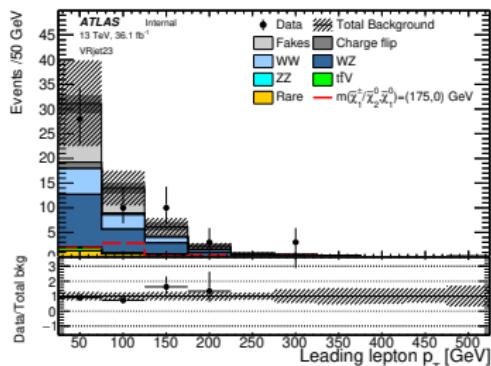
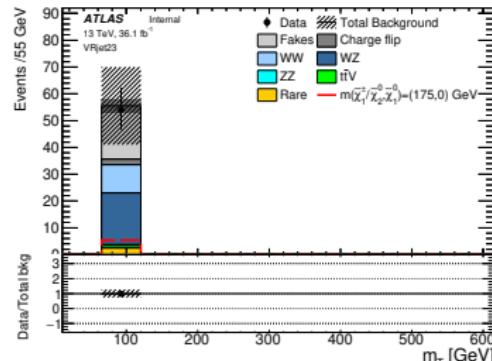
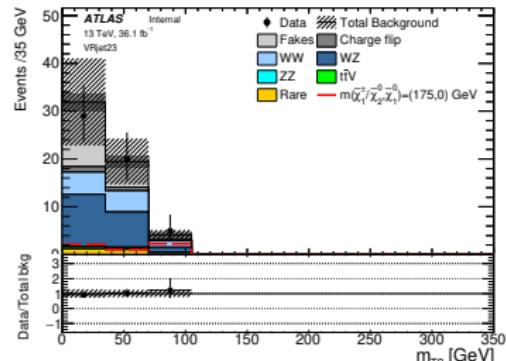
Validation Region : Plots for different variables (VRjet1)



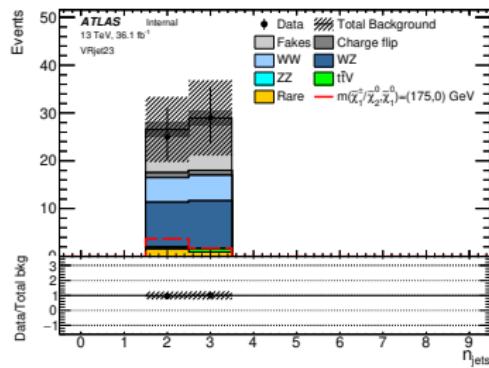
Validation Region : Plots for different variables (VRjet23)



Validation Region : Plots for different variables (VRjet23)

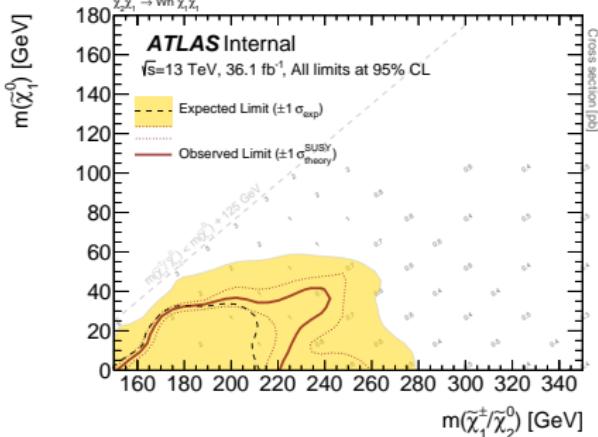
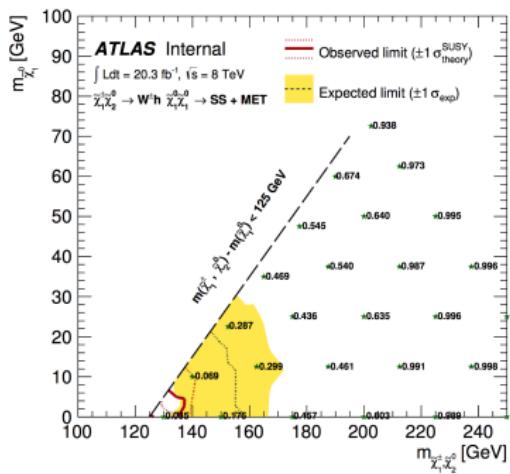


Validation Region : Plots for different variables (VRjet23)

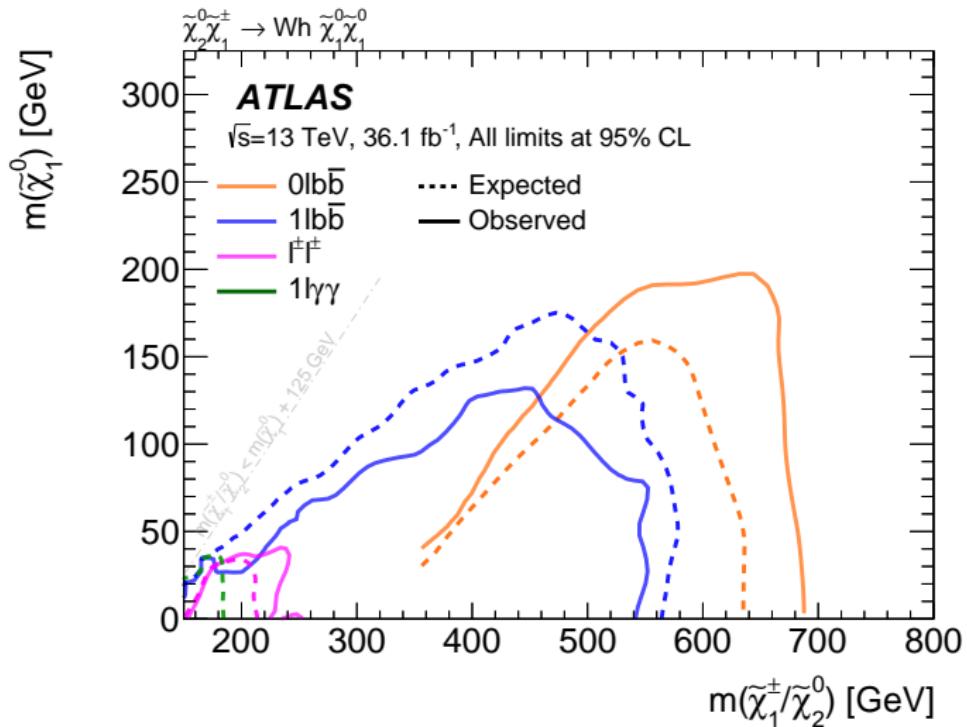


Exclusion Limits

Exclusion Limits : Compare with run 1 analysis



Exclusion Limits : Compare with other channels



Results: Observed Events (systematics)

- The dominant systematic uncertainty is the systematic uncertainty from the fake lepton background.

	Number of events
Total BG	$6.74 \pm 2.60 \pm 2.17$
Data	2

Table: Yields in SRjet1

	Number of events
Total BG	$5.33 \pm 2.31 \pm 1.59$
Data	8

Table: Yields in SRjet23