

# Operators

## Linear Operator

$$F(a\Phi_1 + b\Phi_2) = a \cdot F\Phi_1 + b \cdot F\Phi_2 \quad \forall \Phi_1, \Phi_2$$

## Eigenvalue, Eigenfunction

$$Fu_n = f_n u_n$$

$u_n$  is a eigen function to the operator  $F$  with corresponding eigenvalue  $f_n$ .

## Hermitian Operator

$$\int (Hu) * v \, d^3r = \int u * Hv \, d^3r, \quad \forall u, v$$

A hermitian operator has real eigenvalues and corresponding eigenfunctions can be choosen to be orthonormal. Practically all operators in quantum mechanics are linear and hermitian.

## Eigenfunction Expansion

$$\psi(\mathbf{r}) = \sum_n a_n m \cdot u_n(\mathbf{r}), \quad a_n = \int u_n * \psi \cdot d^3r$$

## Expansion Postulate

At a measurement of an observable  $F$  on a system described by a wavefunction  $\psi$  only eigenvalues of the operator  $F$  can be found. The probability of the result  $F = f_n$  is given by

$$P(F = f_n) = \left| \int u_n * \psi \, d^3r \right|^2, \quad Fu_n = f_n u_n$$

## Momentum Operators

$$L^2 = -\hbar^2 \left[ \frac{1}{\sin^2 \theta} \frac{\partial^2}{\partial \varphi^2} + \frac{1}{\sin \theta} \frac{\partial}{\partial \theta} \left( \sin \theta \frac{\partial}{\partial \theta} \right) \right]$$

$$L_z = \frac{\hbar}{i} \frac{\partial}{\partial \varphi}$$

$L^2$  and  $L_z$  have normalized eigenfunctions  $\Upsilon_l^m(\theta, \varphi)$  for which it holds that:

$$L^2 \Upsilon_l^m = \hbar^2 l(l+1) \Upsilon_l^m$$

$$L_z \Upsilon_l^m = m \hbar \Upsilon_l^m$$

$l$	$m$	$\Upsilon_l^m(\theta, \varphi)$
0	0	$\Upsilon_0^0 = \frac{1}{\sqrt{4\pi}}$
1	0	$\Upsilon_1^0 = \sqrt{\frac{3}{4\pi}} \cos \theta$
1	$\pm 1$	$\Upsilon_1^{\pm 1} = \pm \sqrt{\frac{3}{8\pi}} \sin \theta e^{\pm i\varphi}$
2	0	$\Upsilon_2^0 = \sqrt{\frac{5}{16\pi}} (3 \cos^2 \theta - 1)$
2	$\pm 1$	$\Upsilon_2^{\pm 1} = \pm \sqrt{\frac{15}{8\pi}} \sin \theta \cos \theta e^{\pm i\varphi}$
2	$\pm 2$	$\Upsilon_2^{\pm 2} = \sqrt{\frac{15}{32\pi}} \sin^2 \theta e^{\pm 2i\varphi}$

## Commutators and Momentum Operators

$$\epsilon_{ijk} = \begin{cases} 1 & ijk \text{ even} \\ -1 & ijk \text{ odd} \\ 0 & \text{otherwise} \end{cases}$$

$$[x_i, p_j] = i\hbar \cdot \delta_{ij}$$

$$[x_i, L_j] = i\hbar \cdot \epsilon_{ijk} \cdot x_k$$

$$[L_i, L_j] = i\hbar \cdot \epsilon_{ijk} \cdot L_k$$

$$[x_i, x_j] = [p_i, p_j] = 0$$

$$[p_i, L_j] = i\hbar \cdot \epsilon_{ijk} \cdot p_k$$

$$J_+ = J_x + iJ_y$$

$$J_- = J_x - iJ_y$$

$$J_{\pm}J_{\mp} = J^2 - J_z^2 \pm \hbar \cdot J_z$$

$$[J_+, J_-] = 2\hbar \cdot J_z$$

$$[J_z, J_{\pm}] = \pm \hbar \cdot J_{\pm}$$

$$J_+\phi_{j,m} = \sqrt{(j-m)(j+m+1)} \cdot \hbar \cdot \phi_{j,m+1}$$

$$J_-\phi_{j,m} = \sqrt{(j+m)(j-m+1)} \cdot \hbar \cdot \phi_{j,m-1}$$

$$\Upsilon_l^l(\theta, \varphi) = (-1)^l \sqrt{\frac{2l+1}{4\pi} \frac{(2l)!}{2^{2l}(l!)^2}} \cdot \sin^l \theta \cdot e^{il\varphi}$$