

Noise prediction for serrated trailing-edges

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Outline

Motivation

Why is TE noise important?

Introduction to TE noise generation and control

TE noise generation

TE noise control

Inaccurate existing model

Analytical formulation

The mathematical model

Fourier transformation and iterative-solving procedure

Far-field sound

Results

FEM validation

Model results

Comparison with Howe's model

Noise reduction mechanisms

Conclusion

TE noise problems are important

Figure 1: Applications where TE noise is important

1 2 3

¹Fig(a): sites.google.com/site/flightdeckathome/liveatc

²Fig(b): blog.journals.cambridge.org/2013/01/wind-turbine-syndrome-fact-or-fiction

³Fig(c): www.aliexpress.com/promotion/electronic_computer-fan-noise-promotion.html

TE noise problems are important

- ▶ TE noise of an approaching aircraft



(a) An approaching aircraft

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(a) An approaching aircraft



(b) Wind turbines

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- ▶ TE noise of wind turbines
- ▶ TE noise of rotating fans



(a) An approaching aircraft



(b) Wind turbines



(c) A rotating fan

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When the turbulent boundary layer convects past the TE, the non-radiating pressure fluctuation is scattered into sound capable of propagating to the far-field.



Figure 2: TE noise generation by edge-scattering

TE noise reduction techniques

Figure 3: TE noise reduction techniques

4 5 6 7

⁴Fig(a): T.Geyer *et al* 2010

⁵Fig(b): Michaela Herr *et al* 2005

⁶Fig(c): Gruber's PhD thesis 2012

⁷Theory: (a) and (b) Jaworski and Peake 2013, Lorlna Ayton (c) Howe 1991

TE noise reduction techniques

- ▶ Porous airfoil
- ▶ Brush-type TE



Figure 3: TE noise reduction techniques

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TE noise reduction techniques

- ▶ Porous airfoil
- ▶ Brush-type TE
- ▶ Serrated TEs

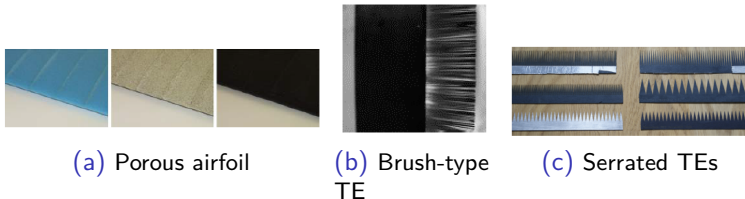


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Howe's model

Howe's model significantly overpredicts the noise reduction capability of serrated TEs.

Figure 4: Comparison of experiment and Howe's model

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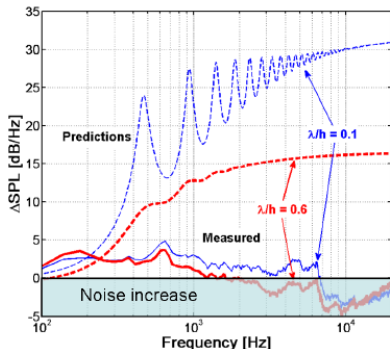


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The mathematical model

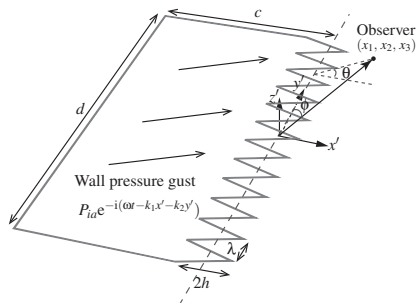


Figure 5: The schematic of a flat plate with a serrated TE

The following wave equation needs to be solved (Roger and Moreau 2013)

$$\begin{aligned} \left(\beta^2 + H'^2(y) \right) \frac{\partial^2 P}{\partial x^2} + \frac{\partial^2 P}{\partial y^2} + \frac{\partial^2 P}{\partial z^2} - 2H'(y) \frac{\partial^2 P}{\partial x \partial y} + \\ (2iM_0 k - H''(y)) \frac{\partial P}{\partial x} + k^2 P = 0, \end{aligned} \quad (1)$$

Fourier transformation

Making use of Fourier transformation

$$P(x, y, z) = \sum_{-\infty}^{\infty} P_n(x, z) e^{ik_{2n}y}, \quad (2)$$

where, $k_{2n} = k_2 + 2n\pi/\lambda$,

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where, $k_{2n} = k_2 + 2n\pi/\lambda$, the wave equation reduces to

$$\mathcal{D}\mathbf{P} = \mathbf{A}\mathbf{P} + \mathbf{B}\frac{\partial\mathbf{P}}{\partial x}, \quad (3)$$

where,

$$\mathcal{D} = \left\{ (\beta^2 + \sigma^2) \frac{\partial^2}{\partial x^2} + \frac{\partial^2}{\partial z^2} + 2ikM_0 \frac{\partial}{\partial x} \right\}. \quad (4)$$

The iterative-solving procedure

$\mathbf{P}^{(0)}$ is obtained by solving

$$\mathcal{D}\mathbf{P} = \mathbf{A}\mathbf{P}. \quad (5)$$

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A solution sequence

$$\mathbf{P}^{(0)}, \mathbf{P}^{(1)}, \mathbf{P}^{(2)}, \mathbf{P}^{(3)} \dots$$

Far-field sound

The far-field sound is obtained by evaluating the surface integral based on the theories of Kirchoff and Curle.

$$p_f(\mathbf{x}, \omega) = \frac{-i\omega x_3}{4\pi c_0 S_0^2} \iint_s \Delta P(x', y') e^{-ikR} \mathrm{d}x' \mathrm{d}y', \quad (8)$$

where $S_0^2 = x_1^2 + \beta^2(x_2^2 + x_3^2)$, and R takes the following form:

$$R = \frac{M_0(x_1 - x') - S_0}{\beta^2} + \frac{x_1 x' + x_2 y' \beta^2}{\beta^2 S_0}, \quad (9)$$

where, ΔP denotes the pressure jump across the flat plate.

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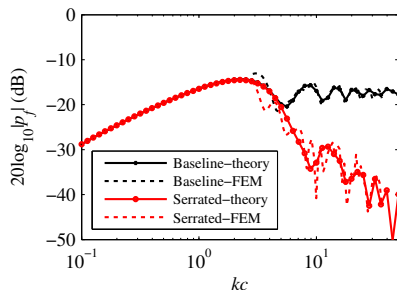
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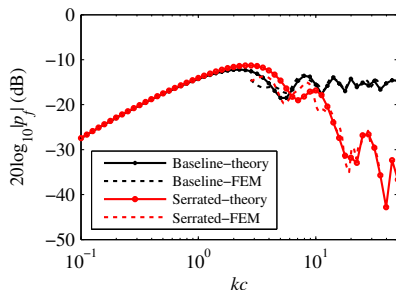
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FEM validation

For wide serrations



(a) $M_0 = 0.1$

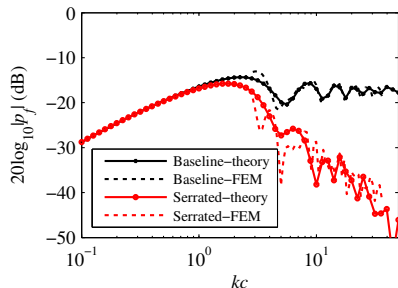


(b) $M_0 = 0.2$

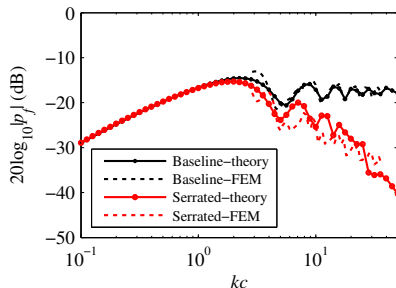
Figure 6: SPL at 90° above the trailing-edge in the mid-span plane with $x_3 = 1$ due to a wall pressure gust of frequency ω with $k_2 = 0$, parameters of the serrations are $\lambda/h = 6$, $h/c = 0.025$.

FEM validation cont.

For narrow serrations,



(a) $\lambda/h = 3$

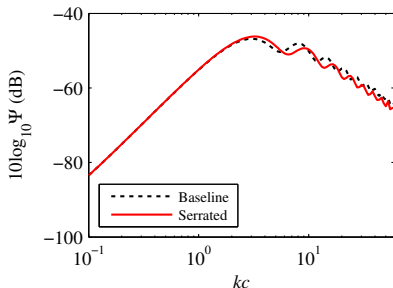


(b) $\lambda/h = 1$

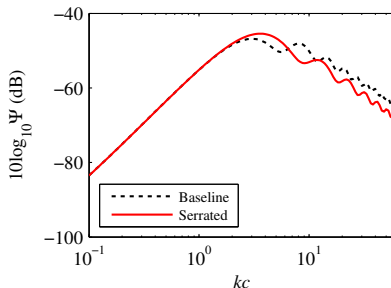
Figure 7: SPL at 90° above the trailing-edge in the mid-span plane with $x_3 = 1$ due to a wall pressure gust of frequency ω with $k_2 = 0$, parameters of the serrations are $h/c = 0.05$ with $M_0 = 0.1$.

Model results

For wide serrations,



(a) $\lambda/h = 8$

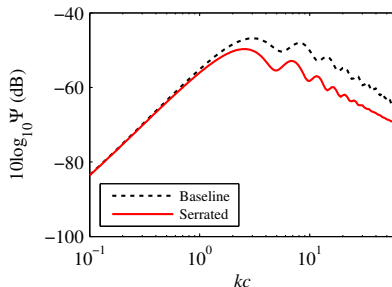


(b) $\lambda/h = 4$

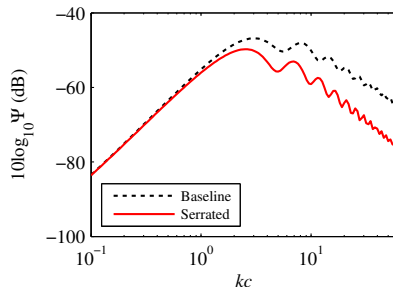
Figure 8: The normalized spectrum for straight and serrated trailing-edges, $h/c = 0.025$, $M_0 = 0.1$, the observer is at 90° above the trailing-edge in the mid-span plane with $x_3 = 1\lambda/h = 8$, $\lambda/h = 4$.

Model results cont.

For narrow serrations,



(a) $\lambda/h = 0.4$



(b) $\lambda/h = 0.2$

Figure 9: The normalized spectrum for straight and serrated trailing-edges, $h/c = 0.05$, $M_0 = 0.1$, the observer is at 90° above the trailing-edge in the mid-span plane with $x_3 = 1$.

Comparison with Howe's model

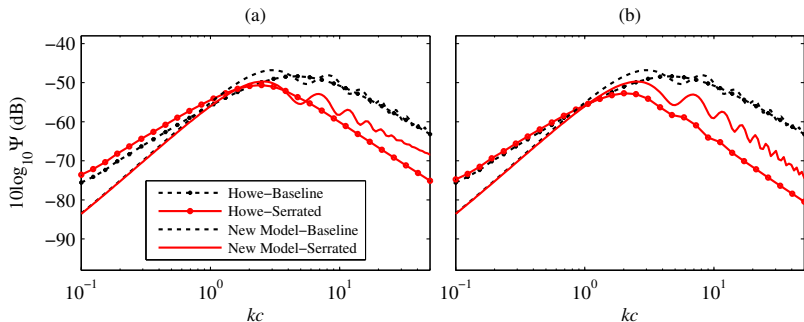
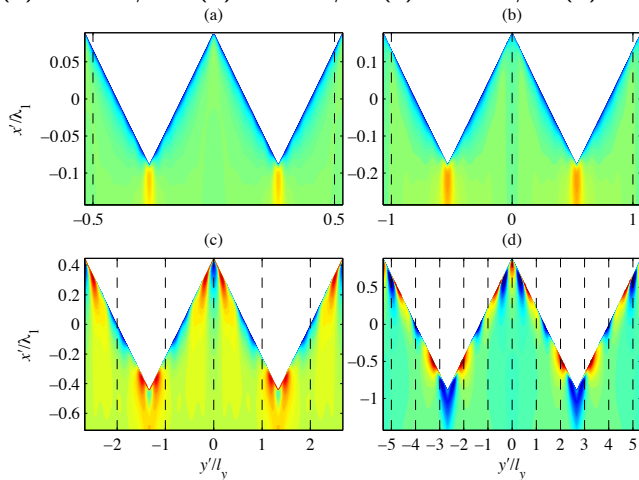


Figure 10: The normalized spectrum of Howe's model and the new model, $h/c = 0.05$, $M_0 = 0.1$, the observer is at 90° above the trailing-edge in the mid-span plane with $x_3 = 1$. (a) $\lambda/h = 0.4$, $h/c = 0.05$, $M_0 = 0.1$ (b) $\lambda/h = 0.2$, $h/c = 0.05$, $M_0 = 0.1$

Noise reduction mechanism

$$\sigma = 5$$

(a) $k\lambda = \pi/10$, (b) $k\lambda = \pi/5$, (c) $k\lambda = \pi/2$, (d) $k\lambda = \pi$



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1. Compared to Howe's model, the presented model includes the convection effect of the mean flow, and can better agrees with experiments.
2. It is found that the destructive interference of the scattered pressure is the cause of sound reduction.
3. The approach used in this model can be used for other serrations. Future work on optimizing the serration profiles can be done.

Thank You!