



Time-Dependent Density Functional Theory Simulations of Coulomb Explosion

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VANDERBILT

Table of Contents

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1. Introduction	3
2. Computational Details	8
2.1. Overview	10
2.2. Formalism - TI	17
2.3. Formalism - TD	31
3. Results & Discussion	41
4. Conclusion	51

TDDFT Simulations of CE

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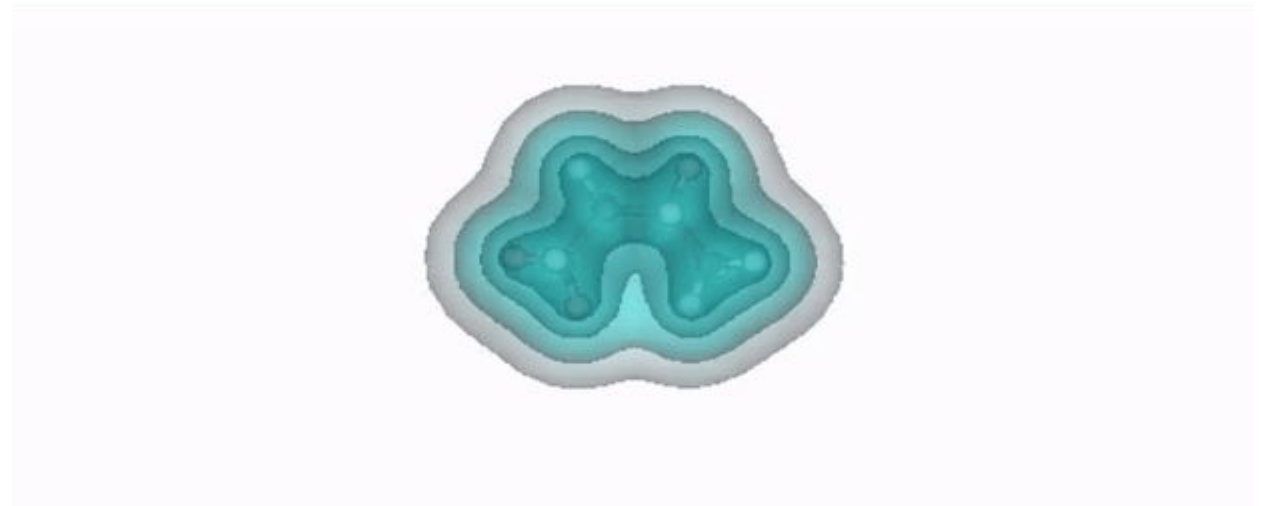
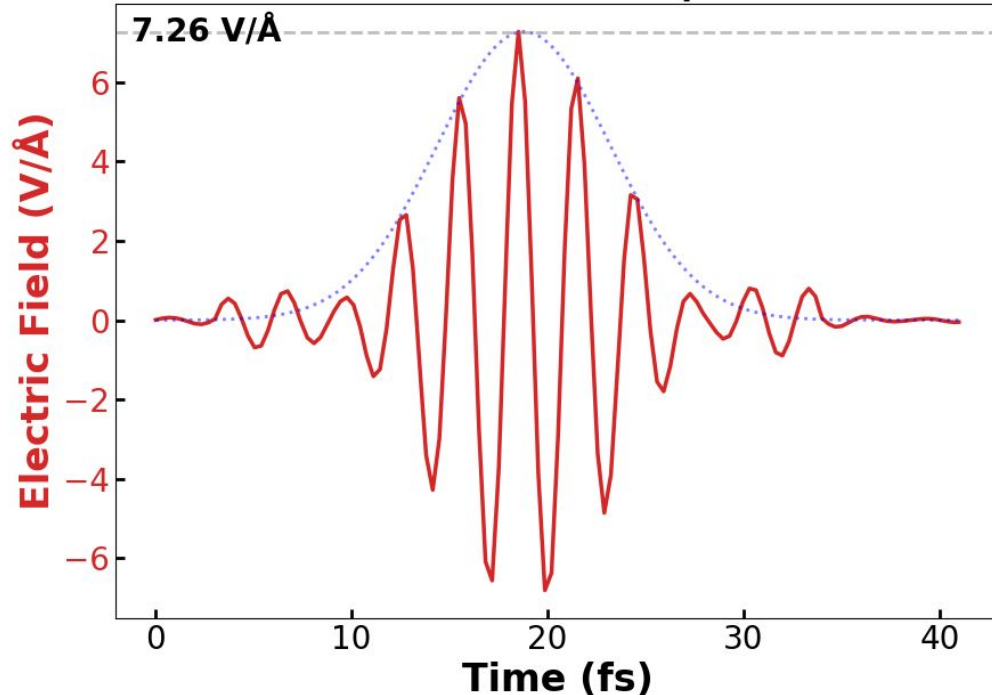


Introduction

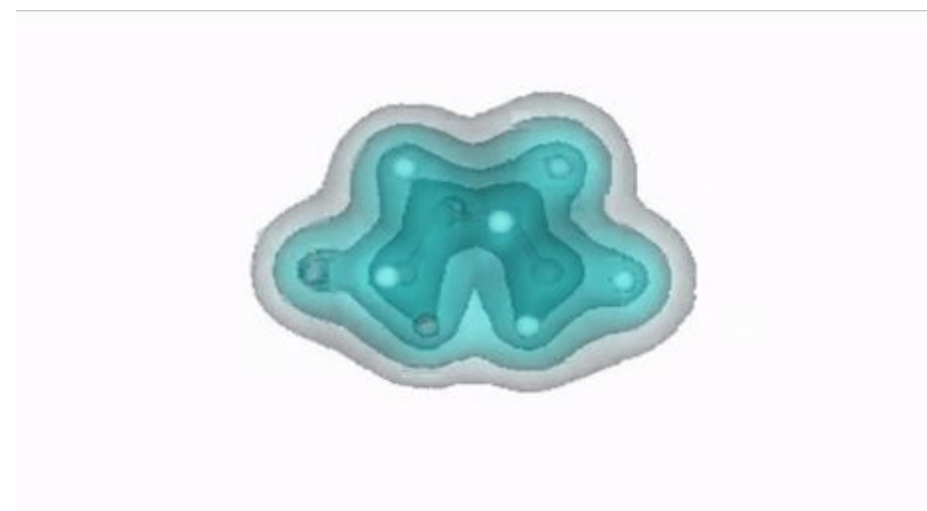
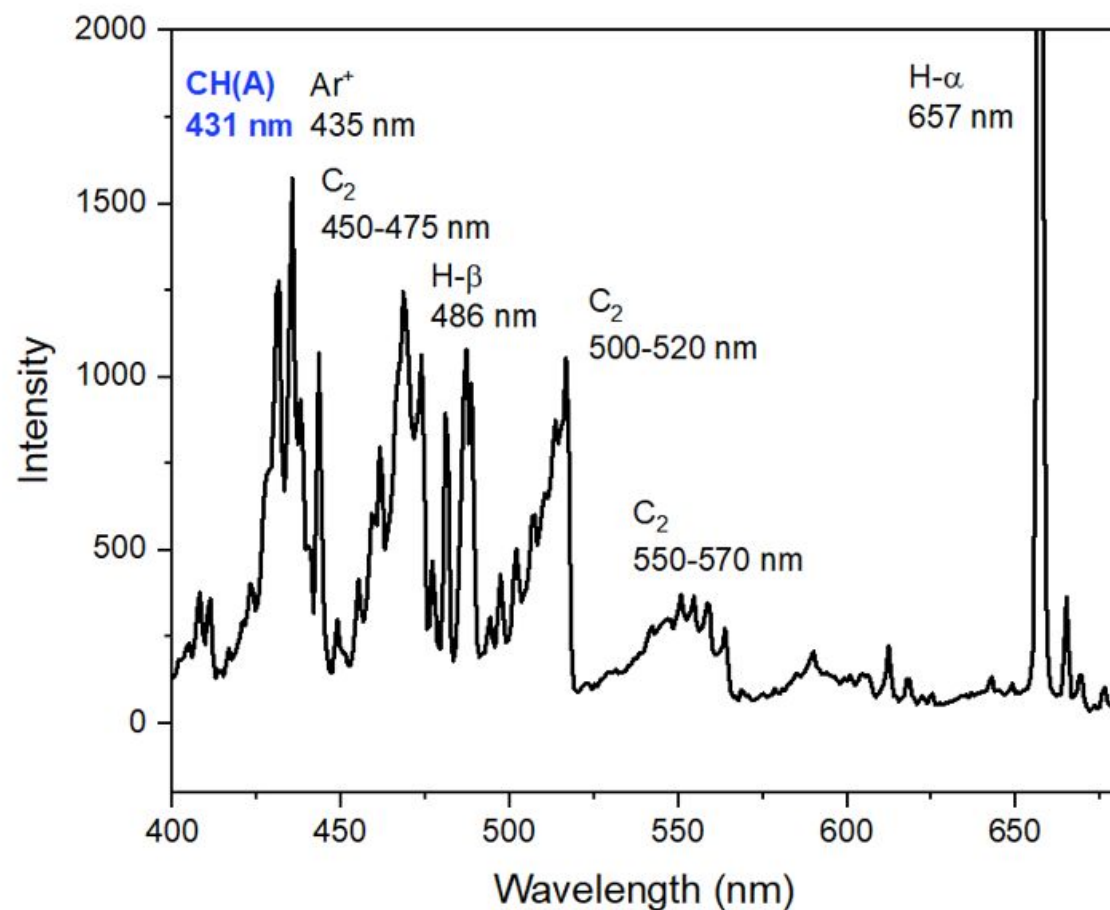
Coulomb Explosion:

- Rapid ionization due to very strong laser pulse
- Coulomb force causes ions to repel

Pulse Used in Butane Coulomb Explosion Simulations

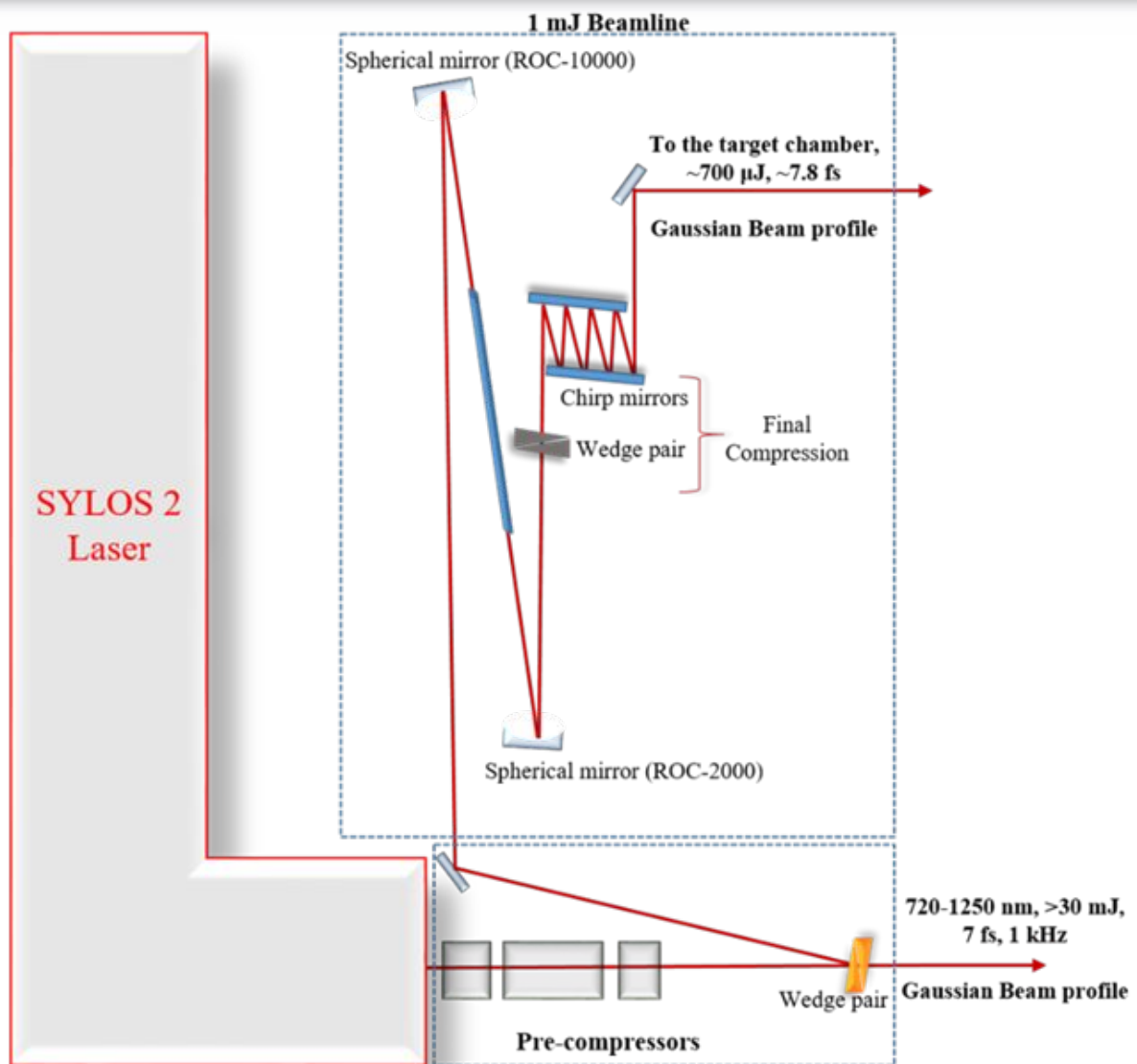


CE can lead to the formation of other molecules. **CH(A) neutral radicals**

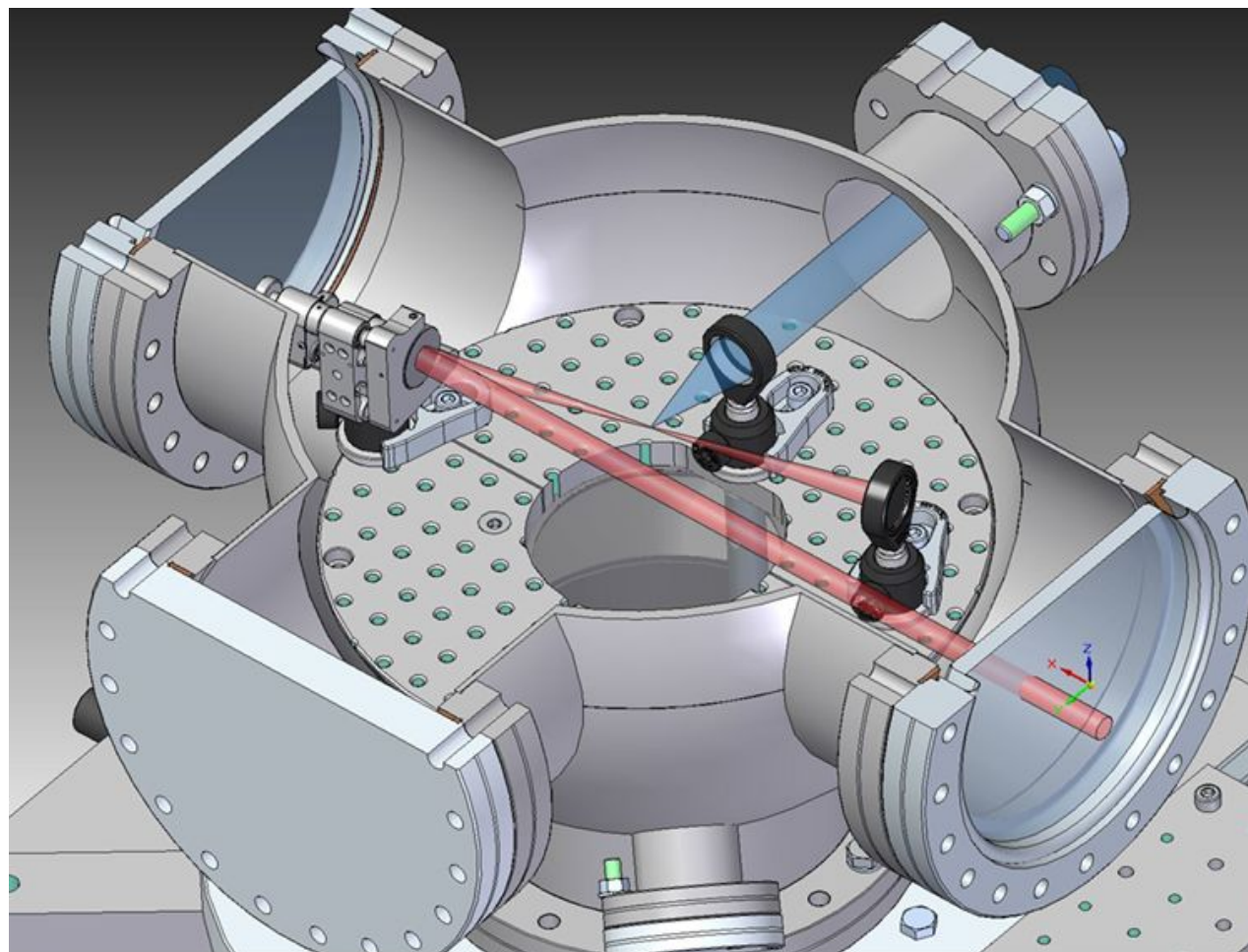


Overview emission spectrum measured from the brightest butane-argon plasma with the CH(A), C₂ emission bands and H atomic lines b, High resolution emission spectra from argon-butane plasma measured at 50 and 250 ns time points (1 cm⁻¹ spectral resolution)

Introduction - The Experiment



Schematic view of the SYLOS 2A laser system and 1 mJ beam line for the photodissociation experiments



The reaction chamber used in the photodissociation experiments (red beam represents the SYLOS beam and the blue one the collected fluorescent light)

Motivation

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- **Show neutral species can be produced from Coulomb explosion**
 - CH neutral radicals - important in many different chemical processes
 - Chemical processes in space, hydrocarbon flames, and plasma-enhanced chemical vapor deposition
- **TDDFT: Match the experimental data with theoretical predictions**
- **Contribute to deeper understanding of:**
 - Radical formation in intense laser fields
 - Various chemical processes in different environments
 - Coulomb Explosion and the production of neutral fragments

TDDFT Simulations of CE

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Computational Details

Computational Details

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Summary of the Process:

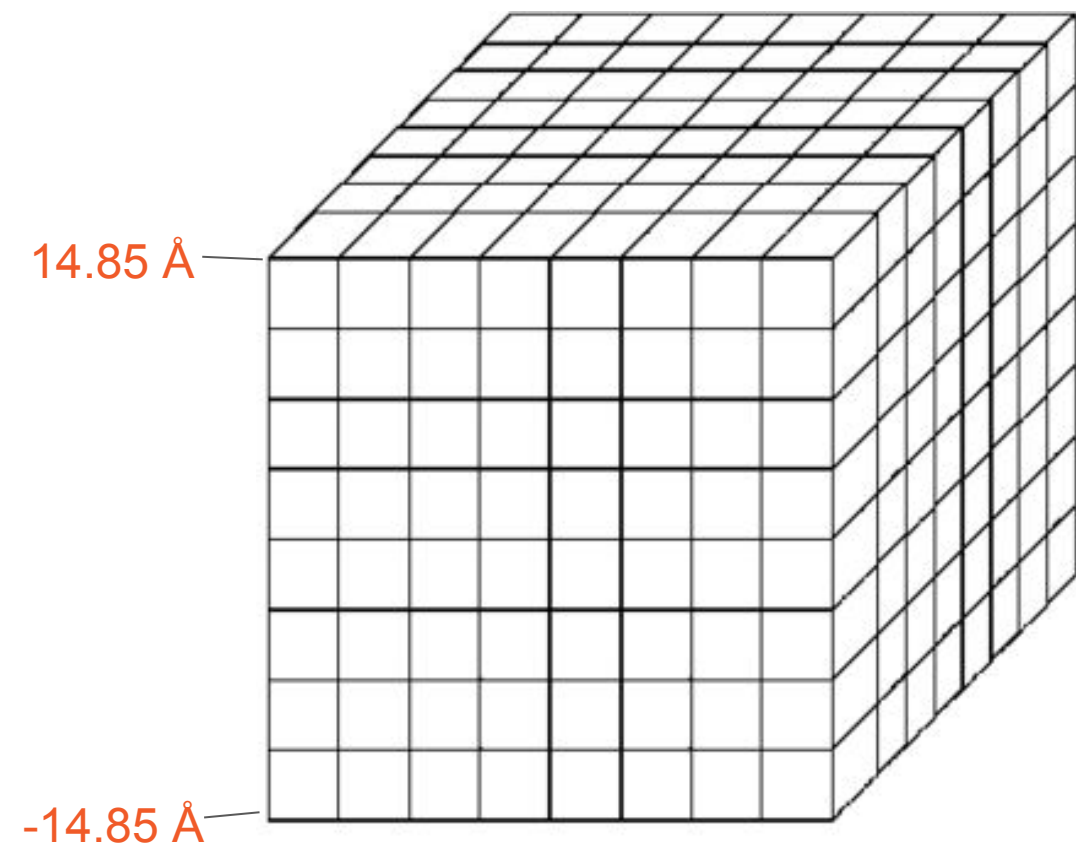
- **Density Functional Theory (DFT)** calculation performed to calculate the **ground state**
 - **Real-space grid used.** Core electrons represented with norm-conserving Troullier-Martin pseudopotentials
- The time-dependent orbitals are propagated using a **fourth-order Taylor expansion of the time-evolution operator**
- Motions of the ions in the simulation are treated **classically** with electron-ion force modeled via Ehrenfest theorem
- A **complex absorbing potential (CAP)** is implemented to prevent reflections of the wave function at the boundary

Computational Details

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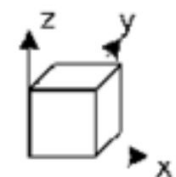


Overview

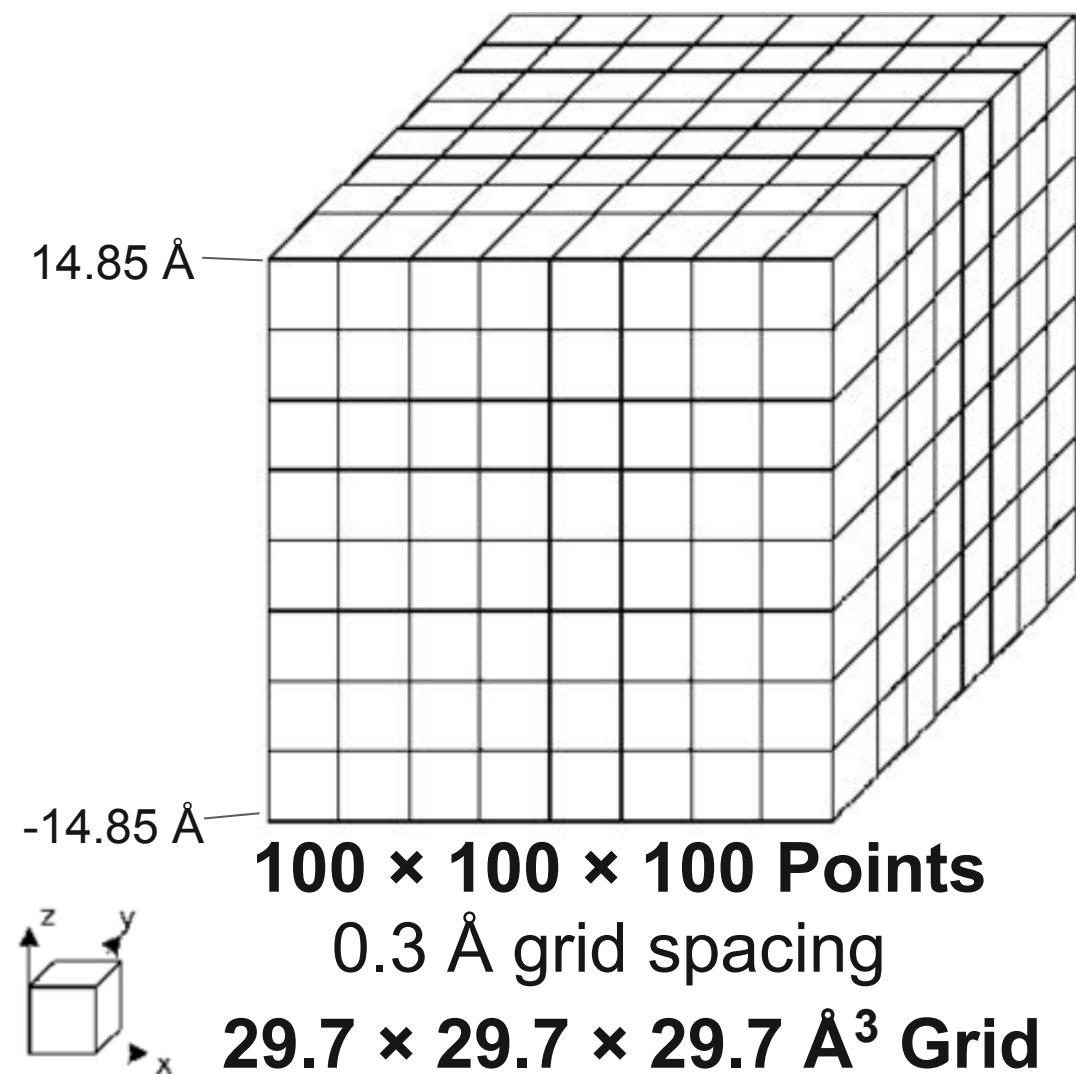


Grid Details and Overview

- Grid is from -14.85 to 14.85 Å on each axis (x,y,z)
- Size is $29.7 \times 29.7 \times 29.7 \text{ Å}^3$
- There are 100 lattice points on every axis
- Grid of $100 \times 100 \times 100 = \mathbf{1,000,000 \text{ points}}$
- Every observable quantity and calculation in the simulation is made from the grid points
- Wave function, matrix operators, electron density
- Distance between each point is **0.3 Å**



$29.7 \times 29.7 \times 29.7 \text{ Å}^3$ Grid

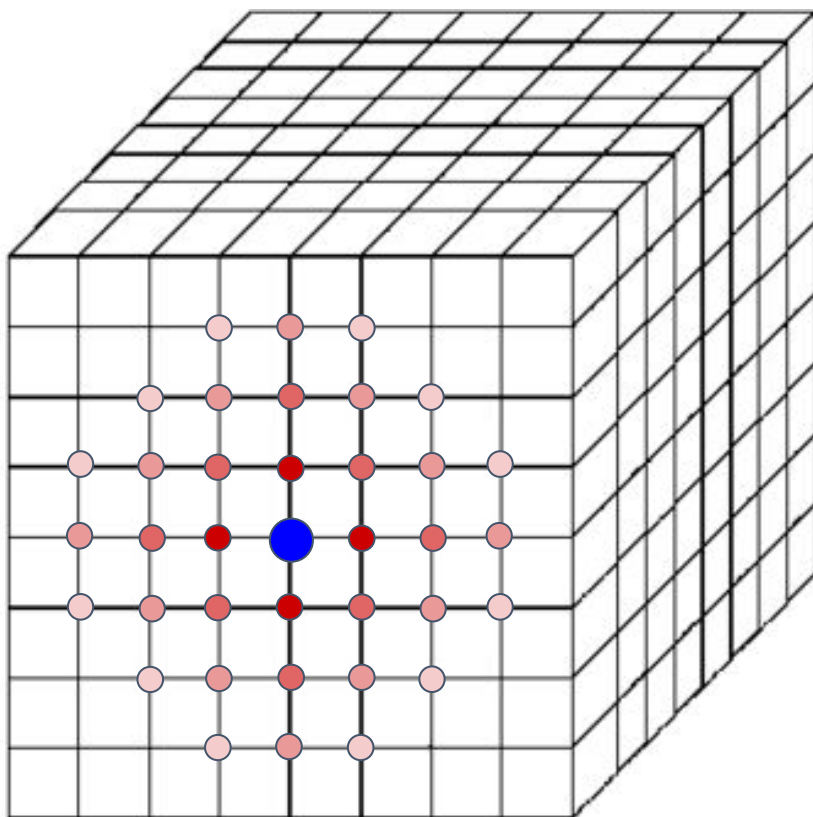


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Overview

Real-space Grid: Example



Example: Using the Grid to calculate number of electrons around a molecule

- Radial “density integration”
- To find the number of electrons N , we need to sum up all of the valence electron density points around a molecule.



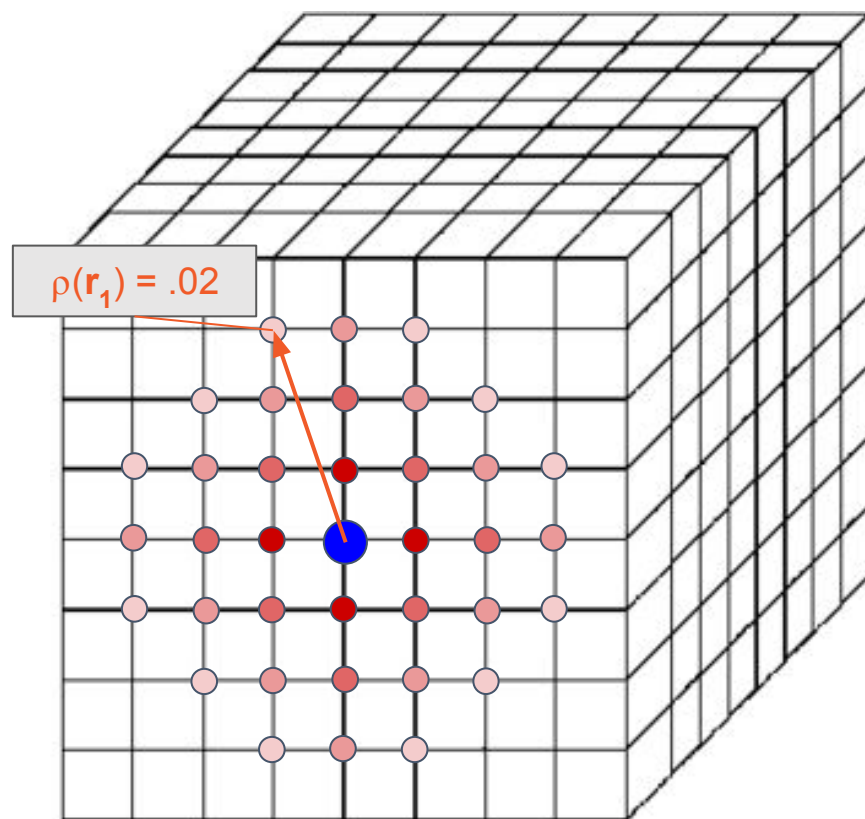
● - Hydrogen



● - Electron Density Point

Overview

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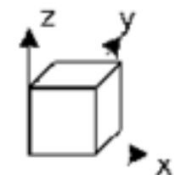
$$N = .02 +$$



- Hydrogen

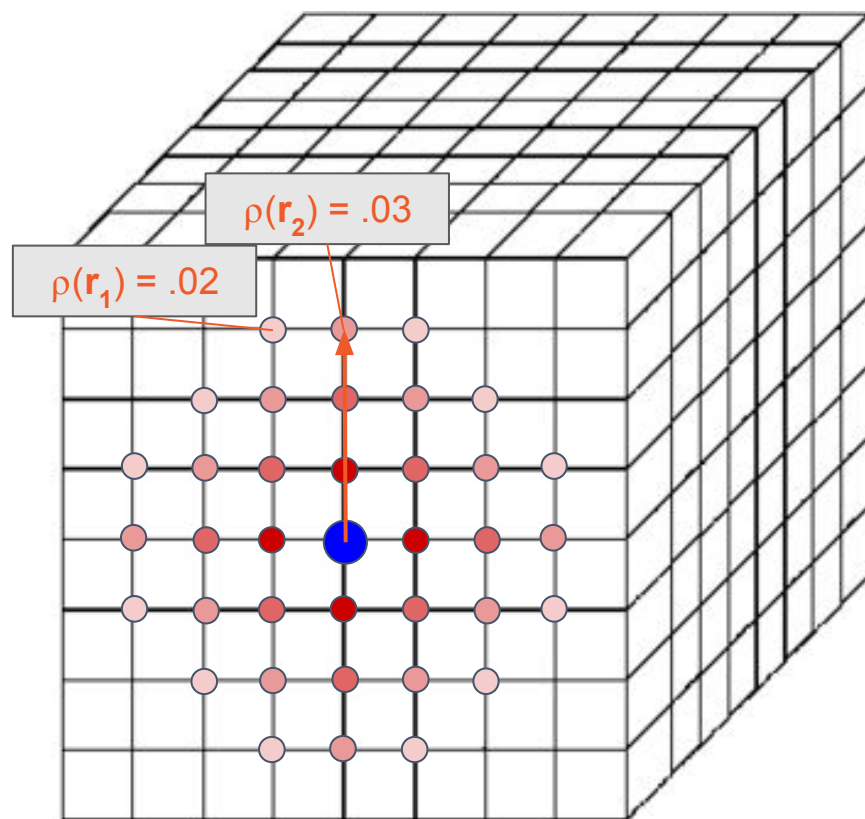


- Electron Density Point



Overview

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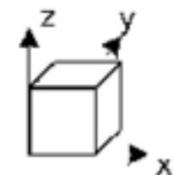
$$N = .02 + .03$$



- Hydrogen

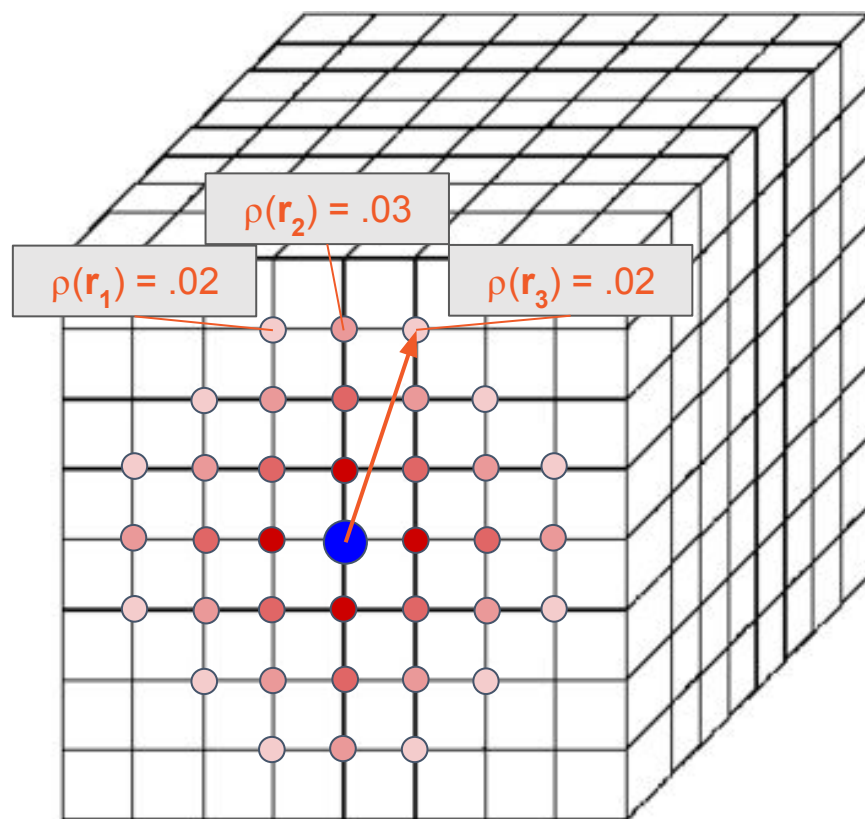


- Electron Density Point



Overview

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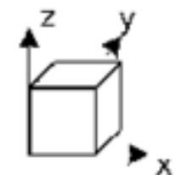
$$N = .02 + .03 + .02$$



- Hydrogen

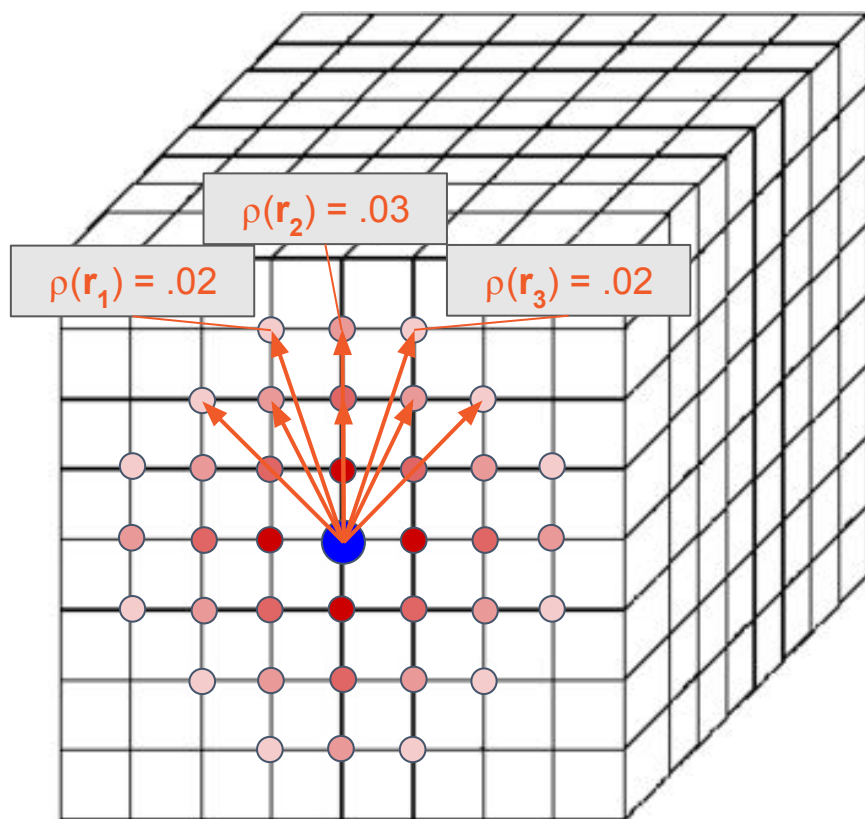


- Electron Density Point



Overview

Real-space Grid: Example



- Hydrogen



- Electron Density Point

Example: Using the Grid to calculate number of electrons around a molecule

- Radial “density integration”
- To find the number of electrons **N**, we need to sum up all of the valence electron density points around a molecule.

$$N = .02 + .03 + .02 + \dots = .70$$

The charge is $1 - .70 = 0.30^+$

(30% probability of being positively charged)



Time-independent calculations

Formalism

DFT Calculation - Kohn Sham

DFT calculation

- Construct Hamiltonian and solve the Kohn-Sham (KS) Equation to find the KS orbitals $\psi_i(\mathbf{r})$

$$H_{\text{KS}}\Psi_i(\mathbf{r}) = E_i\Psi_i(\mathbf{r}), \quad H_{\text{KS}} = -\frac{\hbar^2}{2m}\nabla^2 + V^{\text{KS}}[n(\mathbf{r})]$$

- Solving for the Kohn-Sham orbitals can give us the electron density $n(\mathbf{r})$ as follows:

$$n(\mathbf{r}) = f_k \sum_{i=1}^{N_{\text{occupied}}} |\Psi_i(\mathbf{r})|^2$$

- where f_k is the occupation number of each orbital.
- f_k may take on the values 1 and 2 ($f_k = 2$ is permitted via spin degeneracy)
- The Kohn-Sham potential depends on the density, and yet that is what we are solving for
- **Goal:** solve this equation self-consistently with evolving values for $n(\mathbf{r})$.

Formalism

Kinetic Energy

$$H_{KS}\Psi_i(\mathbf{r}) = E_i\Psi_i(\mathbf{r})$$

$$H_{KS} = -\frac{\hbar^2}{2m}\nabla^2 + V^{KS}[n(\mathbf{r})]$$

Kinetic energy operator

$$-\frac{\hbar^2}{2m}\nabla^2$$

- The Laplacian ∇^2 is approximated using second derivative, 9-point central finite difference

Table 2.2 Higher-order finite difference coefficients for the second derivative

k	C_i^k	$C_{i\pm 1}^k$	$C_{i\pm 2}^k$	$C_{i\pm 3}^k$	$C_{i\pm 4}^k$	$C_{i\pm 5}^k$	$C_{i\pm 6}^k$
1	-2	1					
2	$-\frac{5}{2}$	$\frac{4}{3}$	$-\frac{1}{12}$				
3	$-\frac{49}{18}$	$\frac{3}{2}$	$-\frac{3}{20}$	$\frac{1}{90}$			
4	$-\frac{205}{72}$	$\frac{8}{5}$	$-\frac{1}{5}$	$\frac{8}{315}$	$-\frac{1}{560}$		
5	$-\frac{5269}{1800}$	$\frac{5}{3}$	$-\frac{5}{21}$	$\frac{5}{126}$	$-\frac{5}{1008}$	$\frac{1}{3150}$	
6	$-\frac{5369}{1800}$	$\frac{12}{7}$	$-\frac{15}{56}$	$\frac{10}{189}$	$-\frac{1}{112}$	$\frac{2}{1925}$	$-\frac{1}{16632}$

Look at 9 points:

Current point,
4 to the left,
4 to the right,

to approximate the Laplacian
 $h = 0.3 \text{ \AA}$

$$\phi''(x) = -\frac{1}{560}[\phi(x+4h)+\phi(x-4h)] + \frac{8}{315}[\phi(x+3h)+\phi(x-3h)] - \frac{1}{5}[\phi(x+2h)+\phi(x-2h)] + \frac{8}{5}[\phi(x+h)+\phi(x-h)] - \frac{205}{72}\phi(x)$$

Formalism

Ion potential

$$H_{\text{KS}}\Psi_i(\mathbf{r}) = E_i\Psi_i(\mathbf{r})$$

$$H_{\text{KS}} = -\frac{\hbar^2}{2m}\nabla^2 + V^{\text{KS}}[n(\mathbf{r})]$$

Kohn-Sham potential

$$V_{\text{KS}}(\mathbf{r}) = V_{\text{ion}}(\mathbf{r}) + V_{\text{H}}(\mathbf{r}) + V_{\text{xc}}(\mathbf{r})$$

- V_{ion} - The potential experienced on valence electrons from the ionic core
- Divide the electrons of the atoms into **core electrons** and **valence electrons**
 - Core electron wavefunctions remain essentially unchanged when in different chemical environments
 - Don't solve for core electrons → represent as: **Troullier Martins Pseudopotentials** for the interactions with the valence electrons

$$\hat{V}_{\text{ion}}^{\text{PP}}(\mathbf{r}) = V_{\text{ion,local}}^{\text{PP}}(\mathbf{r}) + \sum_l V_{\text{nonlocal},l}(\mathbf{r})\hat{P}_l,$$

where $V_{\text{ion,local}}^{\text{PP}}(\mathbf{r})$ is the local potential and,

$$V_{\text{nonlocal},l}(\mathbf{r}) = V_{\text{ion},l}^{\text{PP}}(\mathbf{r}) - V_{\text{ion,local}}^{\text{PP}}(\mathbf{r})$$

```
[u.st199789@aces-login1 data]$ ls
ag_pp.dat  au_pp.dat  b_pp.dat  cl_pp.dat  cs_pp.dat  f_pp.dat  he_pp.dat  k_pp.dat  mg_pp.dat  Ng_6  n_pp.dat  pd_pp.dat  rb_pp.dat  s_pp.dat  v_pp.dat
al_pp.dat  be_pp.dat  br_pp.dat  c_pp.dat  cu_pp.dat  ge_pp.dat  h_pp.dat  li_pp.dat  na_pp.dat  ni_pp.dat  o_pp.dat  p_pp.dat  si_pp.dat  ti_pp.dat  w_pp.dat
```

Formalism

Hartree potential

$$H_{\text{KS}}\Psi_i(\mathbf{r}) = E_i\Psi_i(\mathbf{r})$$

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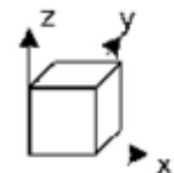
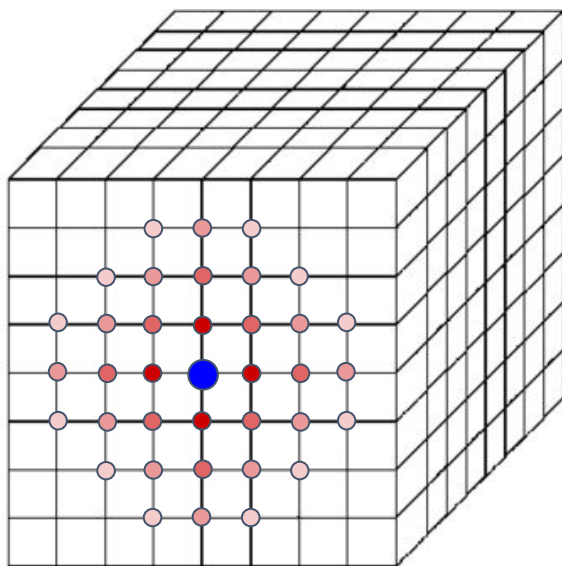
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$$V_{\text{KS}}(\mathbf{r}) = V_{\text{ion}}(\mathbf{r}) + V_{\text{H}}(\mathbf{r}) + V_{\text{xc}}(\mathbf{r})$$

- V_{Hartree} - The electron-electron coulomb interaction

$$V_{\text{H}}(\mathbf{r}) = \int_{\Omega} \frac{n(\mathbf{r}')}{|\mathbf{r} - \mathbf{r}'|} d\mathbf{r}'$$

- The sum of **all** the densities $n(\mathbf{r}')$ at each point \mathbf{r}



Formalism

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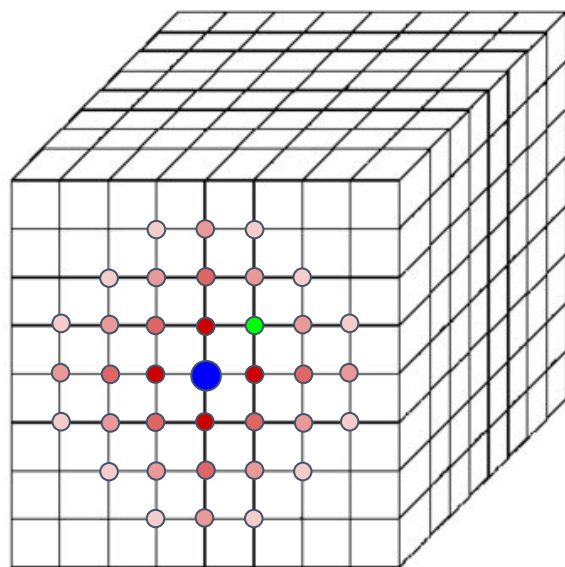
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Example Calculation at \mathbf{r}



Formalism

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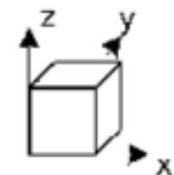
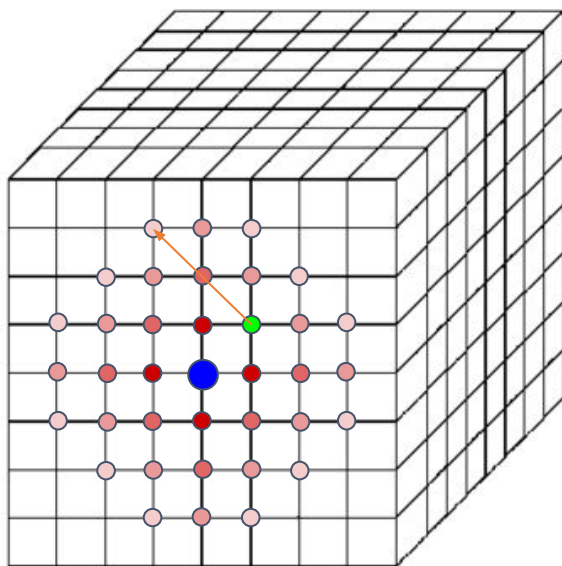
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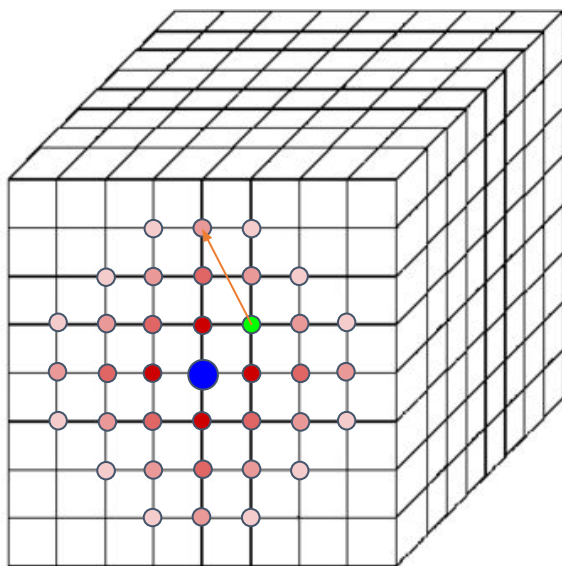
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Formalism

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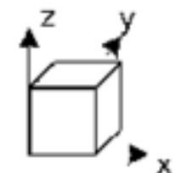
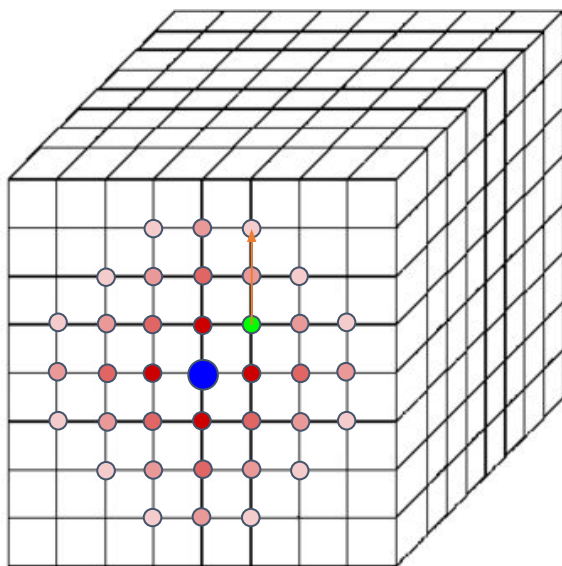
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Example Calculation at \mathbf{r}



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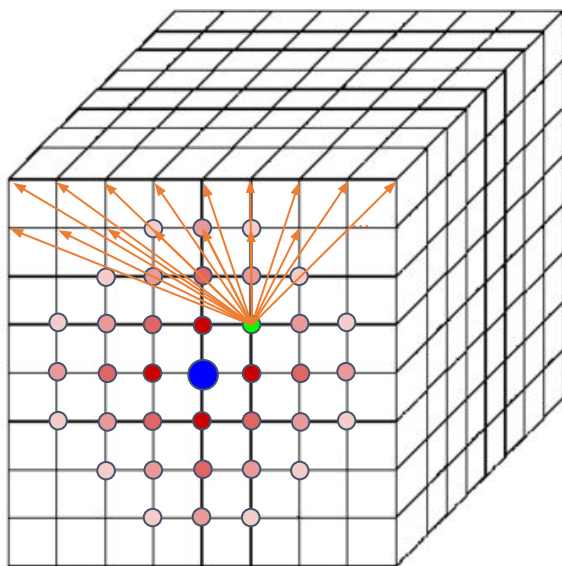
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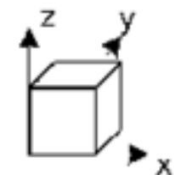
- The sum of **all** the densities $n(\mathbf{r}')$ at each point \mathbf{r}

Example Calculation at \mathbf{r}



The Calculation

- For every single of the 1,000,000 grid points, must sum the resulting effect from 1,000,000 other points
- The work:
 - $1,000,000^2 = (10^6)^2 = \mathbf{10^{12}}$ calculations performed



Formalism

Exchange Correlation

$$H_{\text{KS}}\Psi_i(\mathbf{r}) = E_i\Psi_i(\mathbf{r})$$

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Kohn-Sham potential

$$V_{\text{KS}}(\mathbf{r}) = V_{\text{ion}}(\mathbf{r}) + V_{\text{H}}(\mathbf{r}) + V_{\text{xc}}(\mathbf{r})$$

- V_{xc} - The Exchange and Correlation Potentials – Quantum mechanical electron-electron effects

$$V^{\text{xc}}(\mathbf{r}) = \frac{\delta E_{\text{xc}}[n(\mathbf{r})]}{\delta[n(\mathbf{r})]}$$

- **Exchange:** Energy change that occurs when electrons exchange positions in a system of indistinguishable particles. Comes from Pauli Exclusion principle
- **Correlation:** Includes many-body effects and dynamic correlations beyond the mean-field approximation (Hartree) such as instantaneous positions and momenta of electrons.

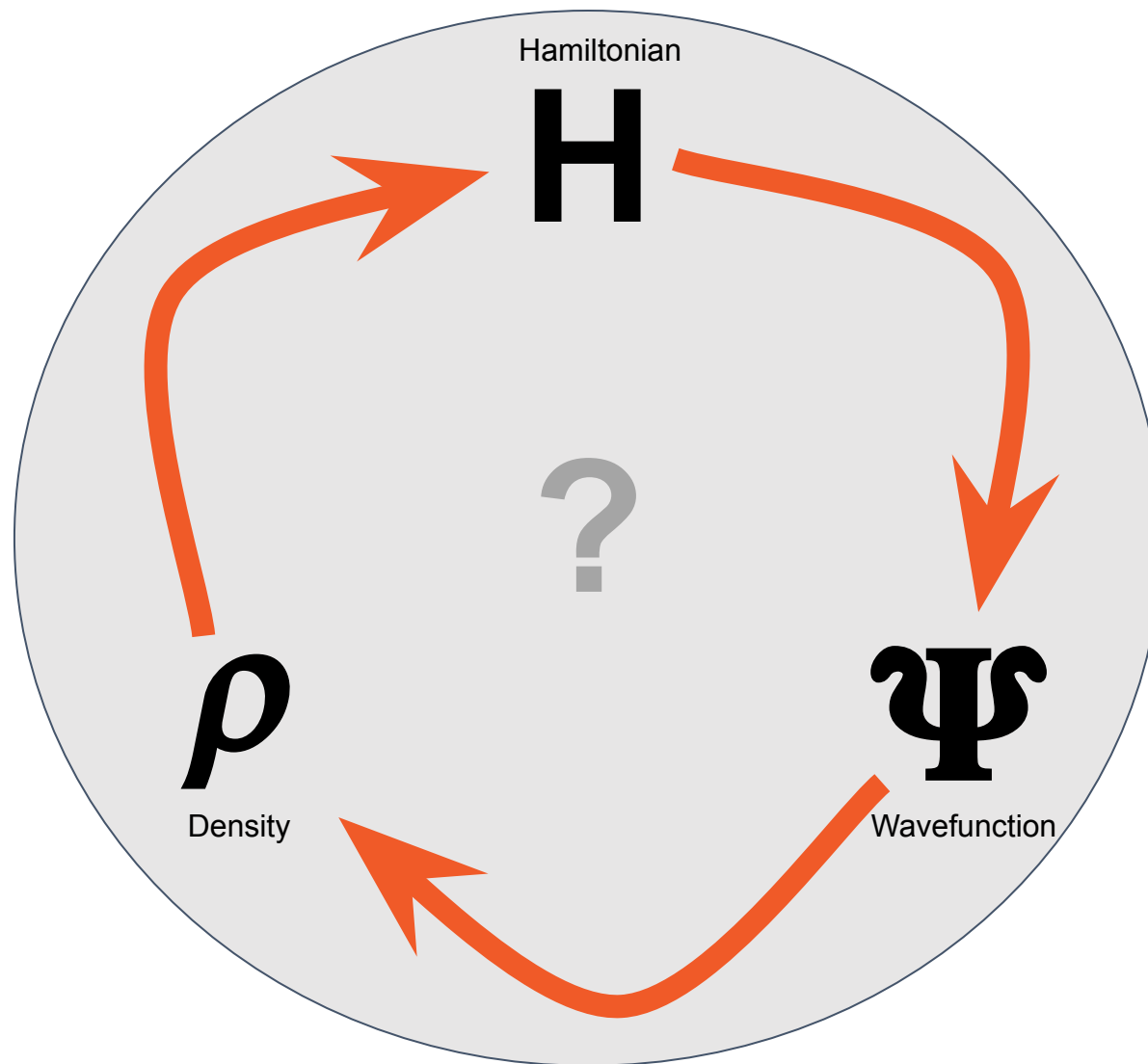
Local Density Approximation: $E_{\text{xc}}^{\text{LDA}} = \int d\mathbf{r} \epsilon[n(\mathbf{r})]n(\mathbf{r}),$

where $\epsilon[n(\mathbf{r})]$ is the exchange-correlation energy per unit volume of a homogeneous electron gas of density $n(\mathbf{r})$.

Resulting LDA potential: $V_{\text{xc}}(\mathbf{r}) = \epsilon_{\text{xc}}(n(\mathbf{r})) + n(\mathbf{r})\frac{d\epsilon_{\text{xc}}}{dn}\big|_{n=n(\mathbf{r})}$

Formalism

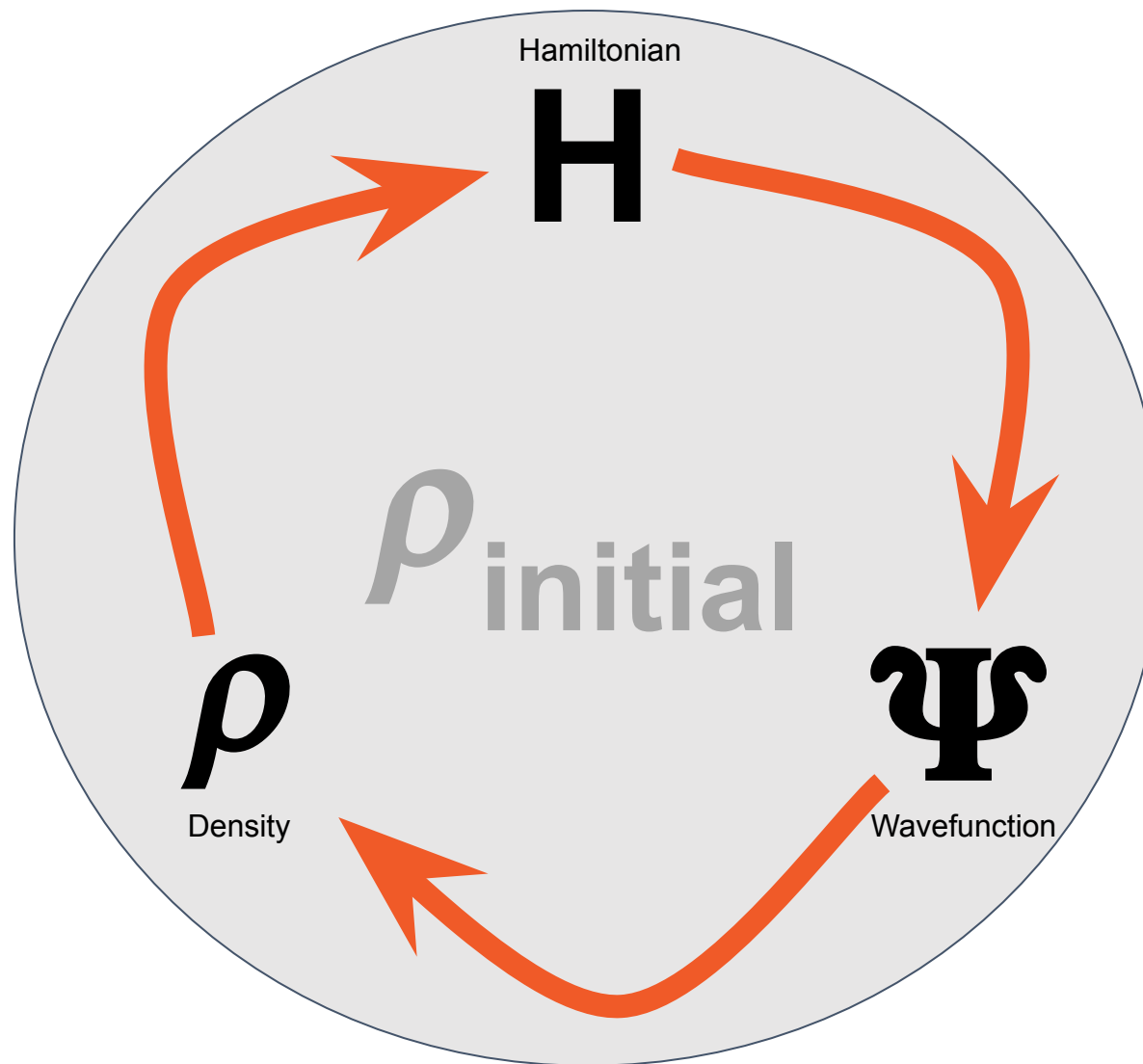
Functionals and Self-Consistency



- The Hamiltonian H is a function of the density ρ
- Need ρ to get H
- Need H to get Ψ
- Need Ψ to get ρ
- *How do we break this loop?*

Formalism

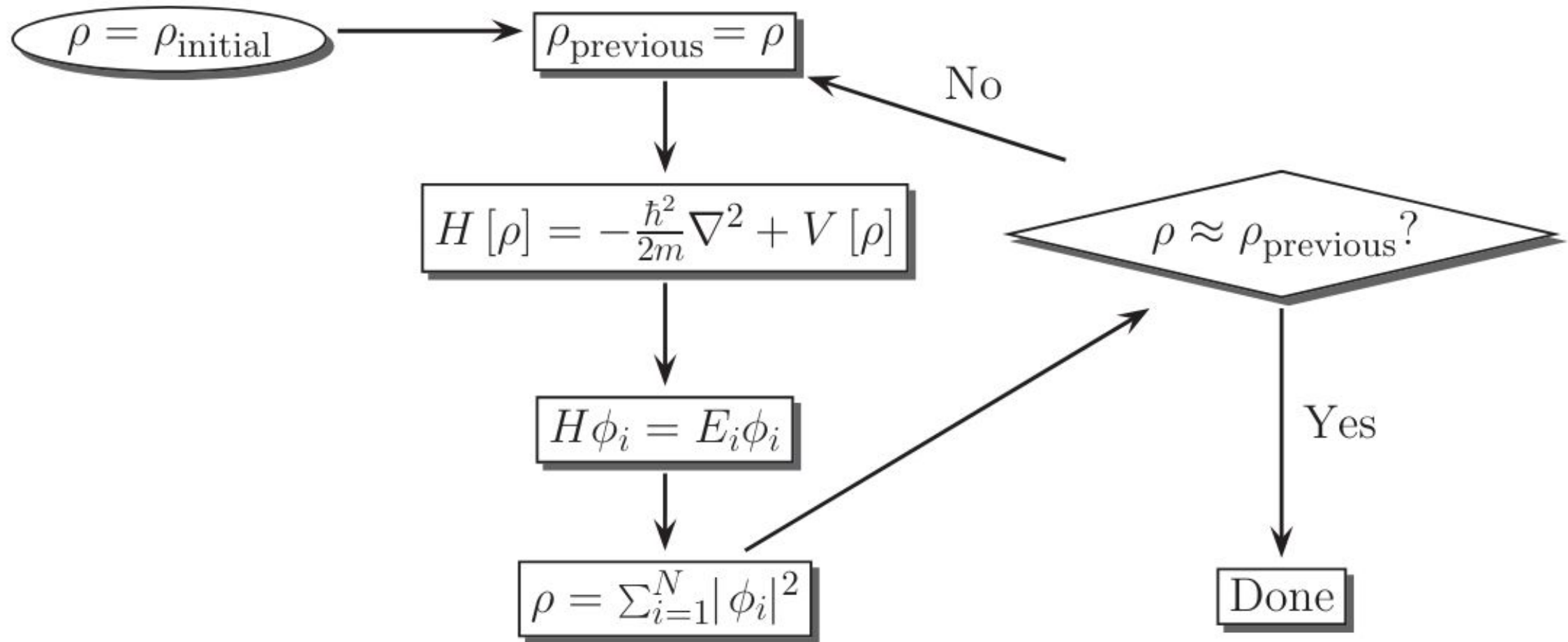
Functionals and Self-Consistency



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- Need ρ to get H
- Need H to get Ψ
- Need Ψ to get ρ
- *How do we break this loop?*
- Initial Guess for ρ

Formalism

DFT Flow Chart



Computational Details

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Time-dependent calculations

Formalism

External Potential

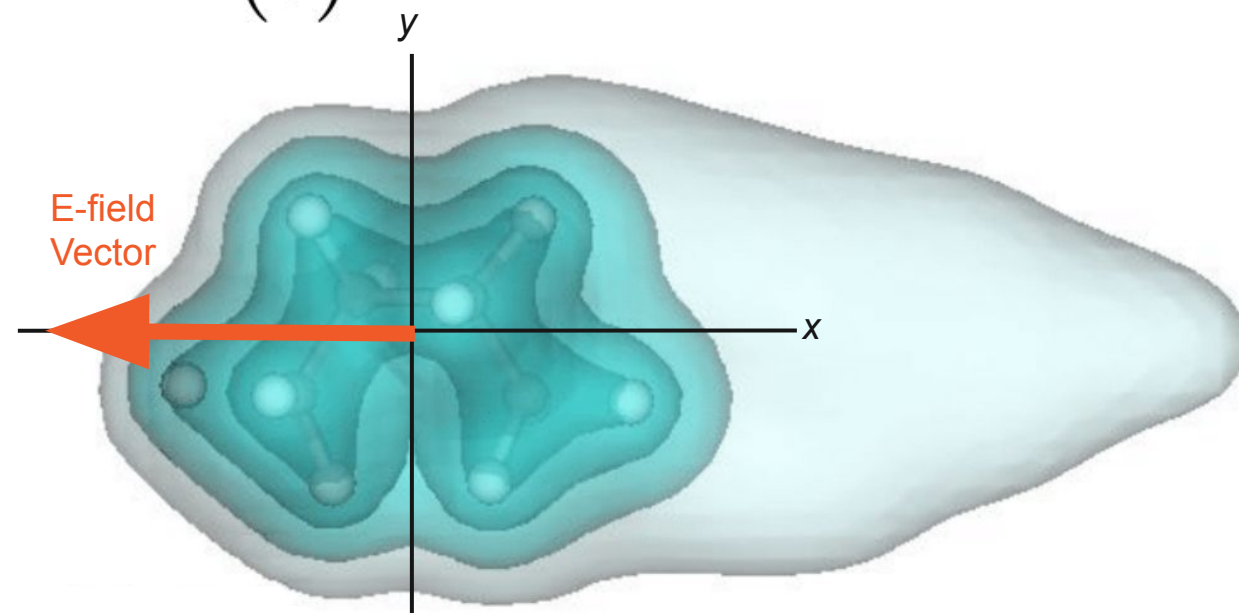
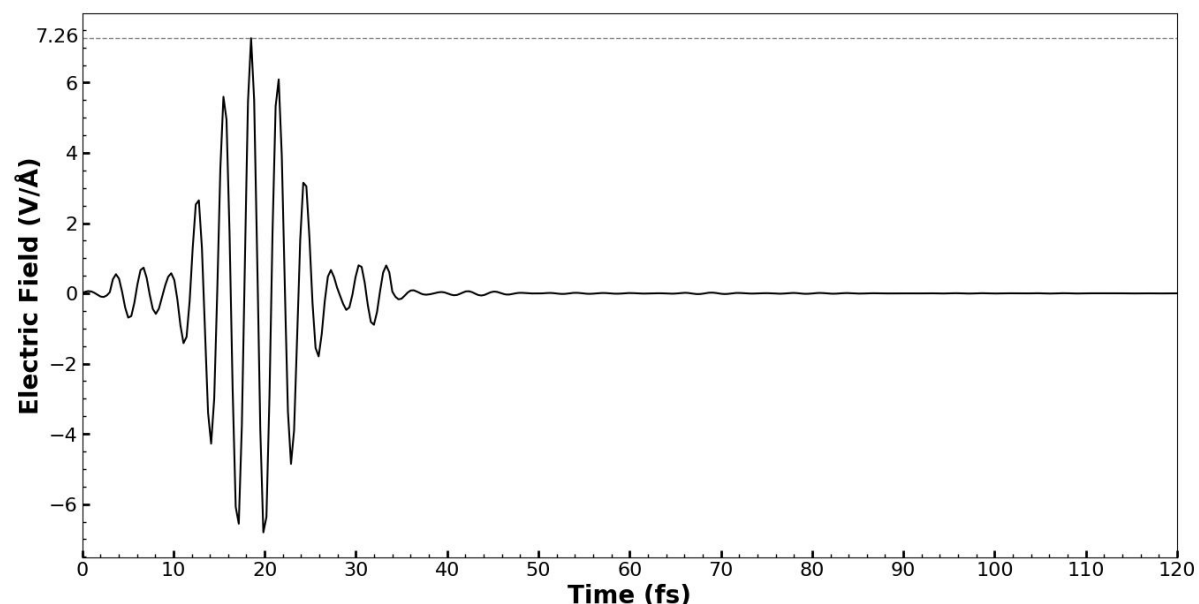
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$$H_{\text{KS}} = -\frac{\hbar^2}{2m}\nabla^2 + V^{\text{KS}}[n(\mathbf{r})]$$

Laser / Electric field potential

- Described using the dipole approximation
- Our laser is polarized in the x-direction—**The E-field strictly exists in the x-axis**

$$V_{\text{laser}} = \mathbf{r} \cdot \mathbf{E}(t)$$



Formalism

External Potential

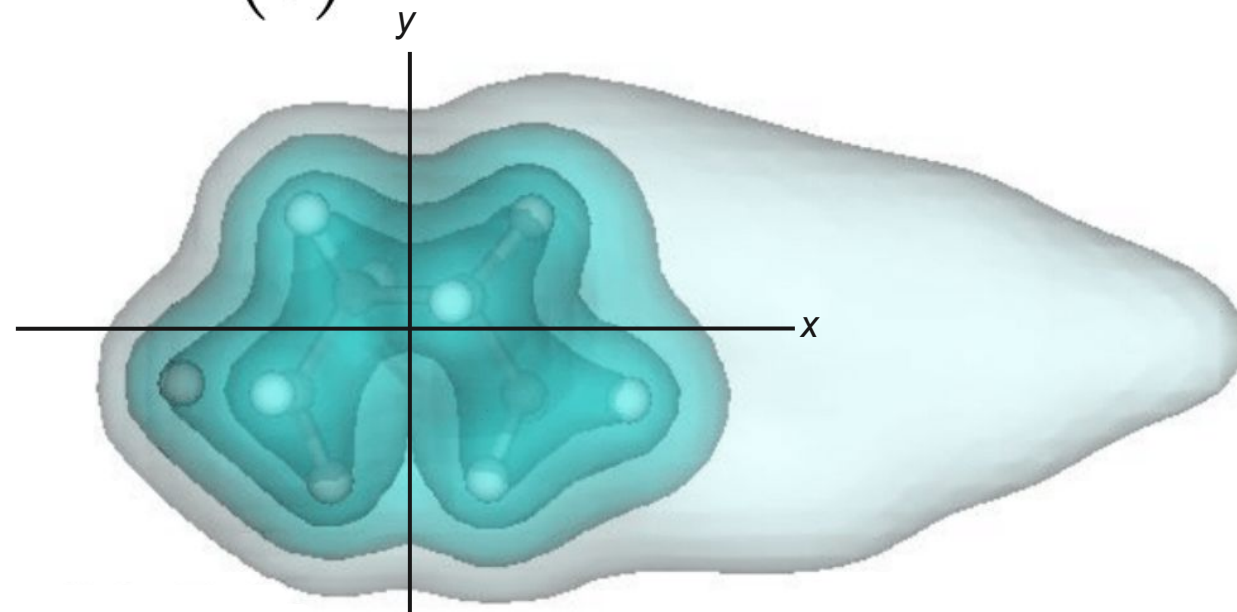
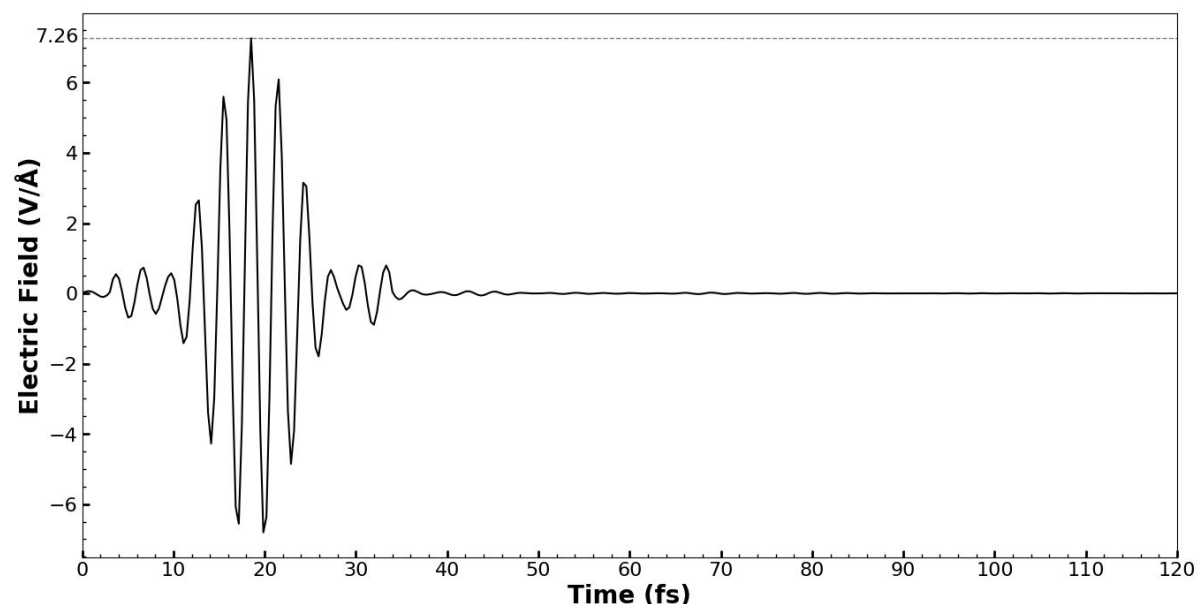
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Laser / Electric field potential

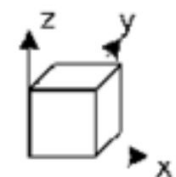
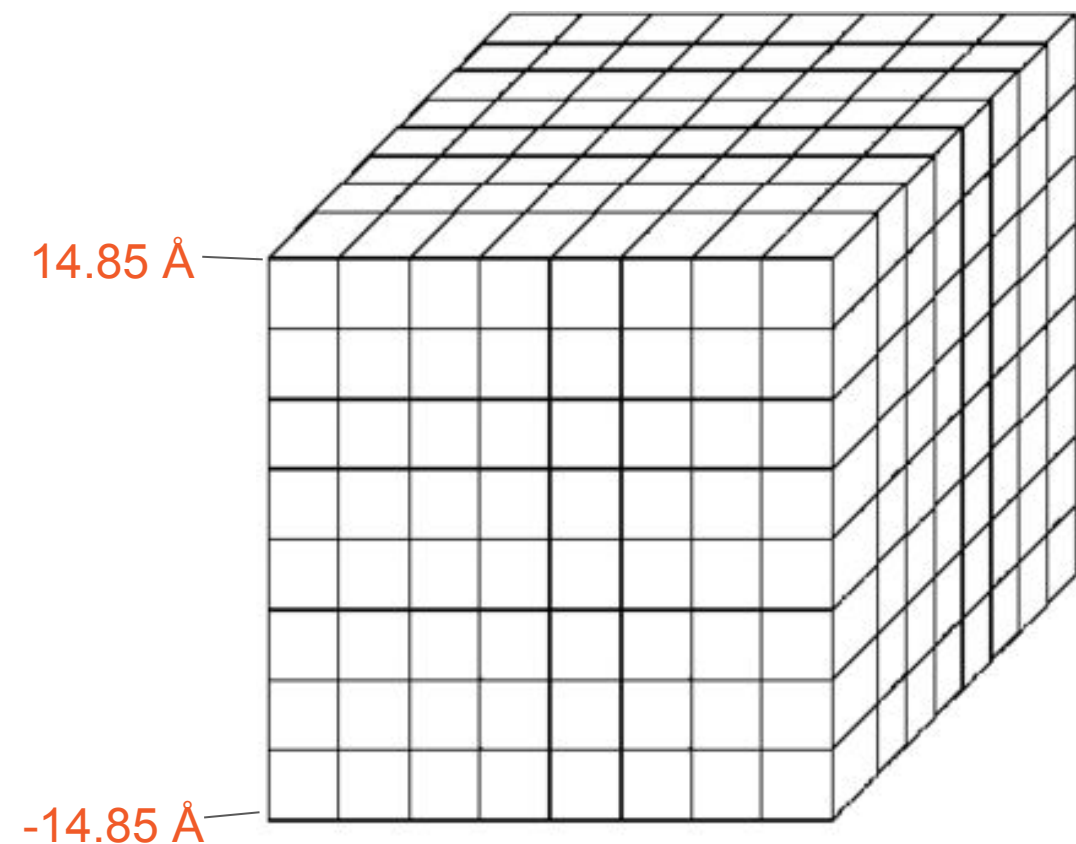
- Described using the dipole approximation
- Our laser is polarized in the x-direction—**The E-field strictly exists in the x-axis**

$$V_{\text{laser}} = \mathbf{r} \cdot \mathbf{E}(t)$$



Grid Details and Overview

- Same grid as in DFT calculation, since the ground state density is used in the TDDFT calculation.
- Grid is from -14.85 to 14.85 Å on each axis (x,y,z)
- Size is $29.7 \times 29.7 \times 29.7 \text{ Å}^3$
- There are 100 lattice points on every axis
- Grid of $100 \times 100 \times 100 = \mathbf{1,000,000 \text{ points}}$
- Distance between each point is $\mathbf{0.3 \text{ Å}}$



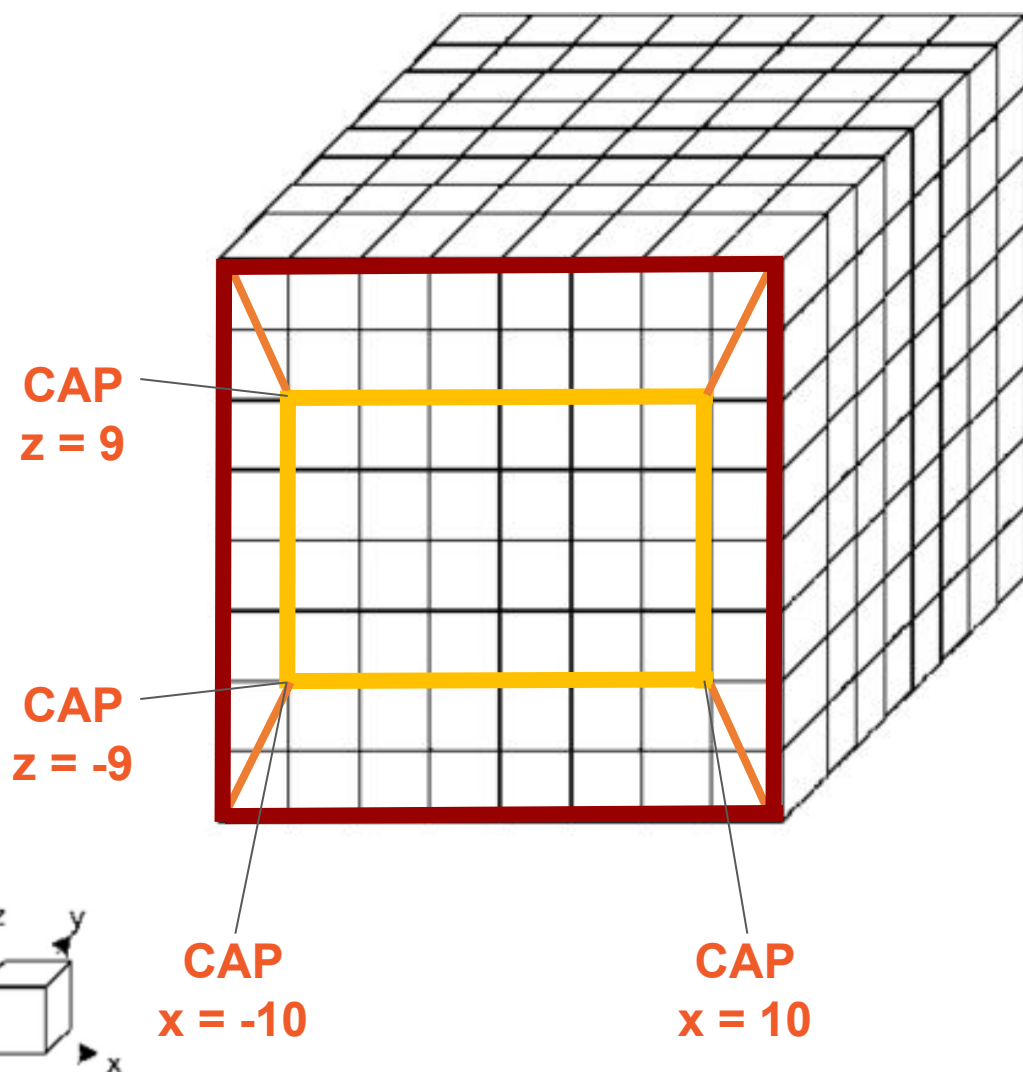
$29.7 \times 29.7 \times 29.7 \text{ Å}^3$ Grid

Overview

Complex Absorbing Potential (CAP)

Complex Absorbing Potential (CAP)

- CAP used to prevent unphysical reflections of the wavefunction
- Implemented as an **imaginary potential** added to the Hamiltonian near the bounds of the simulation
 - *Smooth exponential function to gradually damp wave functions*
- CAP is set at:
 - -10 to 10 Å on the x-axis
 - -9 to 9 Å on the y and z-axis
- **Without the CAP:** Ejected electrons could reflect off the boundary of the grid into the computational domain
 - Like ripples of water reflecting off of walls
- With no CAP, the boundaries of the grid are treated as boundary conditions with effectively infinite potential → **reflections**



Complex Absorbing Potential

- **CAP** is a potential that is added to the Hamiltonian. CAP by **Manolopoulos**

$$-iw(x) = -i \frac{\hbar^2}{2m} \left(\frac{2\pi}{\Delta x} \right)^2 f(y)$$

- \mathbf{x}_1 is the start and \mathbf{x}_2 is the end of the absorbing region,
- $\Delta \mathbf{x} = \mathbf{x}_2 - \mathbf{x}_1$
- \mathbf{m} is the electron's mass

$$f(y) = \frac{4}{c^2} \left(\frac{1}{(1+y)^2} + \frac{1}{(1-y)^2} - 2 \right), \quad y = \frac{(x - x_1)}{\Delta x}$$

- \mathbf{c} is a numerical constant = (2.62)
- Smooth continuous function: **Minimizes transmission and reflection**, maximizes absorption

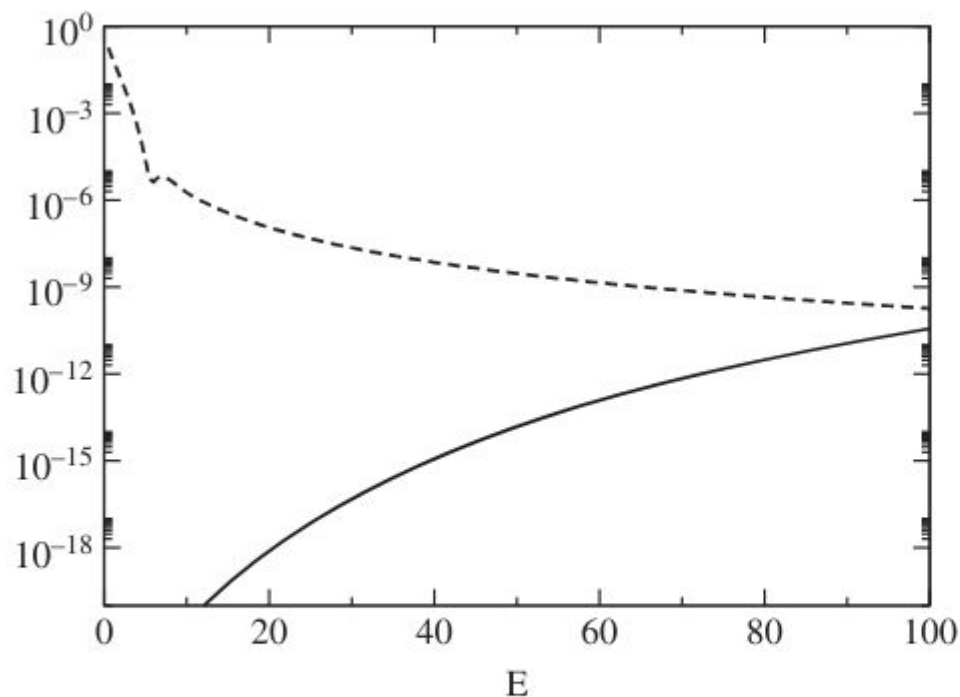
Formalism

Additional info: Other CAPs

Power function CAP

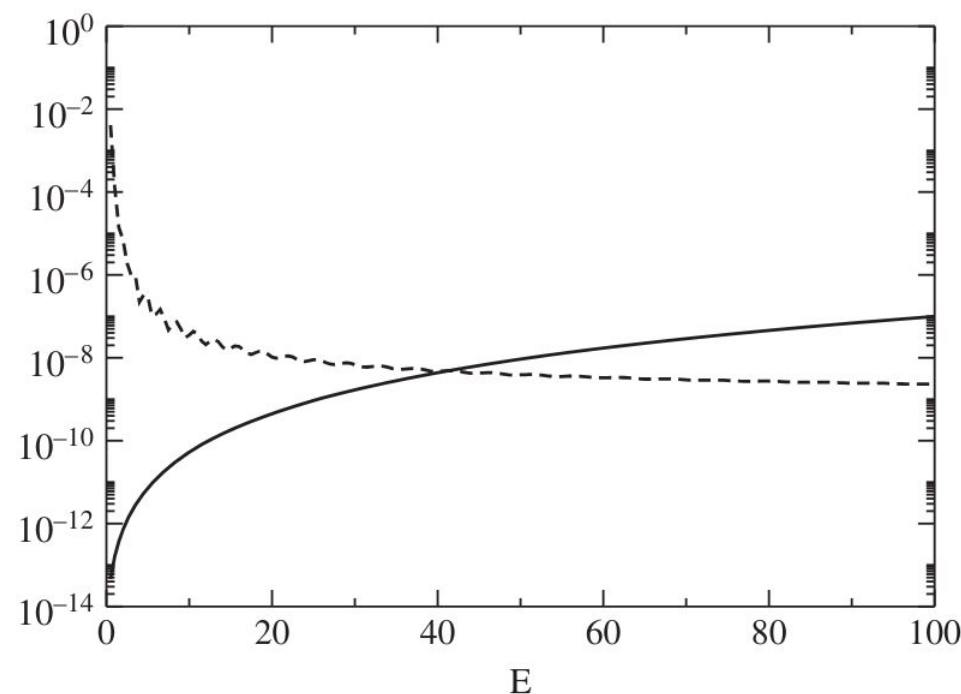
$$W(x) = -i\eta(x - x_K)^n, \quad x_K < x < x_K + L$$

where x_K ($K = L, R$) is the starting point of the complex potential
 L is the range of the complex potential



Manolopoulos CAP

$$-iw(x) = -i \frac{\hbar^2}{2m} \left(\frac{2\pi}{\Delta x} \right)^2 f(y)$$



--- Dashed line:
Reflection Probability

— Solid line:
Transmission Probability

For our simulations: The Manolopoulos CAP best **minimizes reflection and transmission**

Formalism

Time Evolution Operator

Time Evolution Operator

$$\psi(t + \Delta t) = U(t + \Delta t, t)\psi(t)$$

- We use the time evolution operator to propagate the wave function throughout time

$$U(t + \Delta t, t) = e^{-iH_{KS}(t)\Delta t/\hbar}$$

- We have to approximate the exponential operator to use this approach. **4th Order Taylor Expansion**

$$e^{-iH_{KS}\Delta t/\hbar} \approx \sum_{n=0}^N \frac{(-i\Delta t/\hbar)^n H_{KS}^n}{n!}$$

- In our simulations, **$\Delta t = .001 \text{ fs}$** ($1 \times 10^{-18} \text{ s}$) = **1 as**
- Our simulations last 120 fs. So, in a simulation we have **120000 iterations of time propagation.**

Formalism

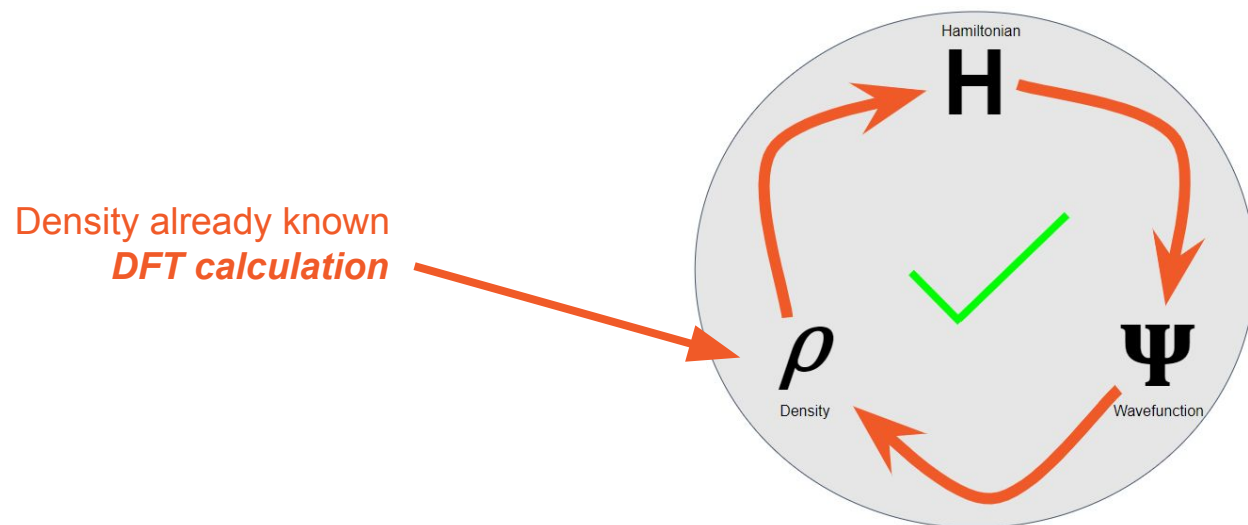
Time Evolution Operator

Time Evolution Operator

$$U(t + \Delta t, t) = e^{-iH\Delta t/\hbar} \approx \sum_{n=0}^N \left(\frac{(-i\Delta t/\hbar)^n H^n}{n!} \right)$$

- Where the Hamiltonian is the $\mathbf{H}(\mathbf{t}) \rightarrow$ the Kohn-Sham Hamiltonian + V_{ext}

$$H(t) = -\frac{\hbar^2}{2m} \nabla^2 + V_{\text{ion}}(\mathbf{r}, t) + V_H[\rho](\mathbf{r}, t) + V_{\text{xc}}[\rho](\mathbf{r}, t) + V_{\text{laser}}(\mathbf{r}, t)$$



Formalism

Ion movement - Classical

Ion movement

- Movement of the ions treated classically. **Newton's second law: $F=ma$**

$$M_i \frac{d^2 \mathbf{R}_i}{dt^2} = Z_i \mathbf{E}_{\text{laser}}(t) + \sum_{j \neq i}^{N_{\text{ions}}} \frac{Z_i Z_j (\mathbf{R}_i - \mathbf{R}_j)}{|\mathbf{R}_i - \mathbf{R}_j|^3} - \nabla_{\mathbf{R}_i} \int V_{\text{ion}}(\mathbf{r}, \mathbf{R}_i) \rho(\mathbf{r}, t) d\mathbf{r}$$

Newton's Law

Coulomb's Law

- Ion \rightarrow Laser
- Ion \rightarrow Ion

Ehrenfest Theorem

- Quantum forces on the ions due to the electrons

- M_i is the mass and Z_i is the charge of the i^{th} ion. N_{ions} is the total number of ions


Formalism

Ion movement - Classical

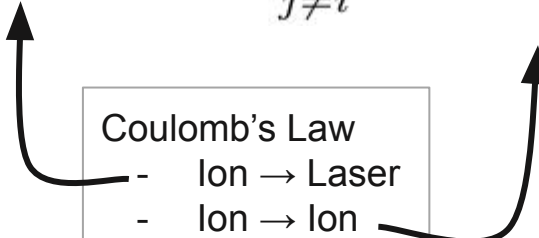
Ion movement

- Movement of the ions treated classically. **Newton's second law: $F=ma$**


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Newton's Law



Coulomb's Law
 - Ion → Laser
 - Ion → Ion



Ehrenfest Theorem
 - Quantum forces on the ions
 due to the electrons

- M_i is the mass and Z_i is the charge of the i^{th} ion. N_{ions} is the total number of ions

Computational Details

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TOWARDS THE SHARP END OF ATTOSCIENCE



Summary of the Process:

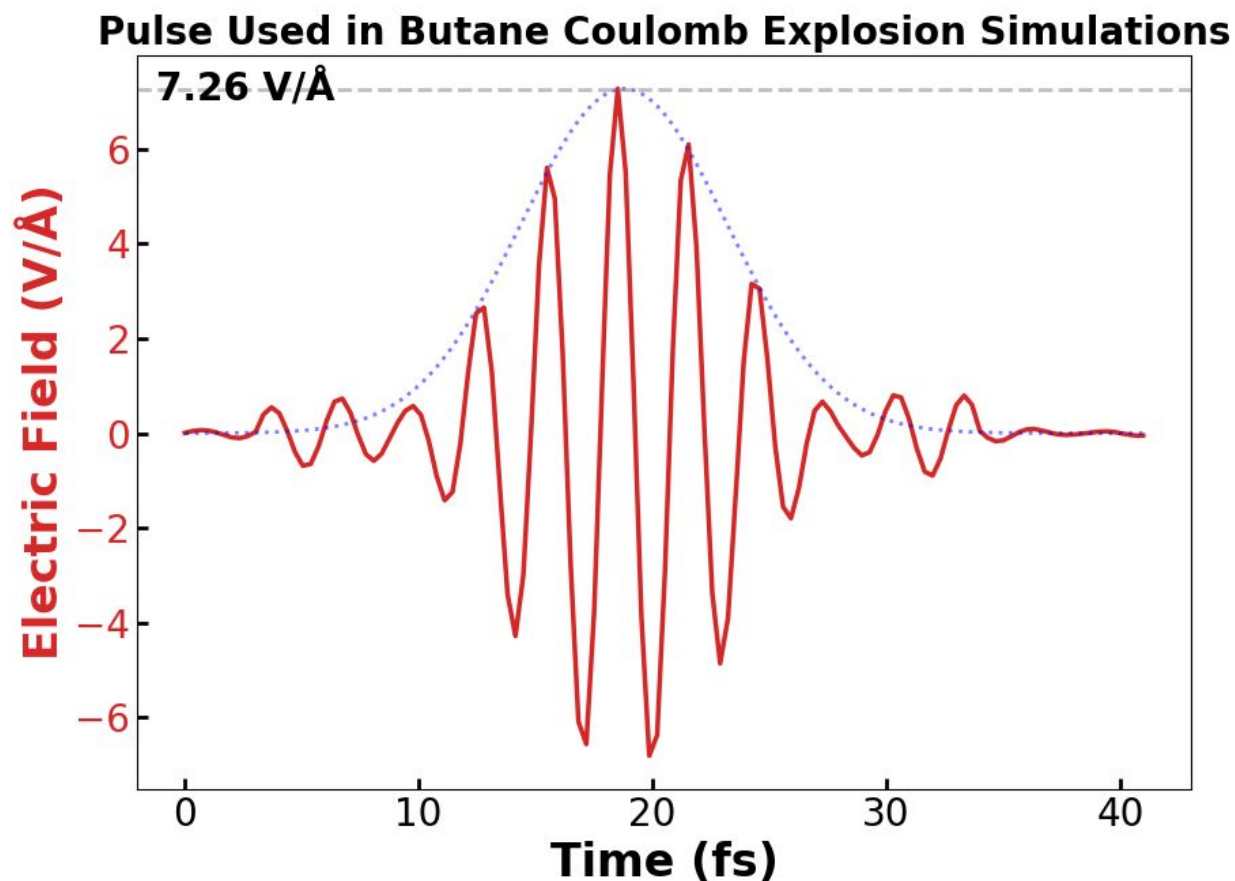
- **Density Functional Theory (DFT)** calculation performed to calculate the **ground state**
 - **Real-space grid used.** Core electrons represented with norm-conserving Troullier-Martin pseudopotentials
- The time-dependent orbitals are propagated using a **fourth-order Taylor expansion of the time-evolution operator**
- Motions of the ions in the simulation are treated **classically** with electron-ion force modeled via Ehrenfest theorem
- A **complex absorbing potential (CAP)** is implemented to prevent reflections of the wave function at the boundary



Results & Discussion

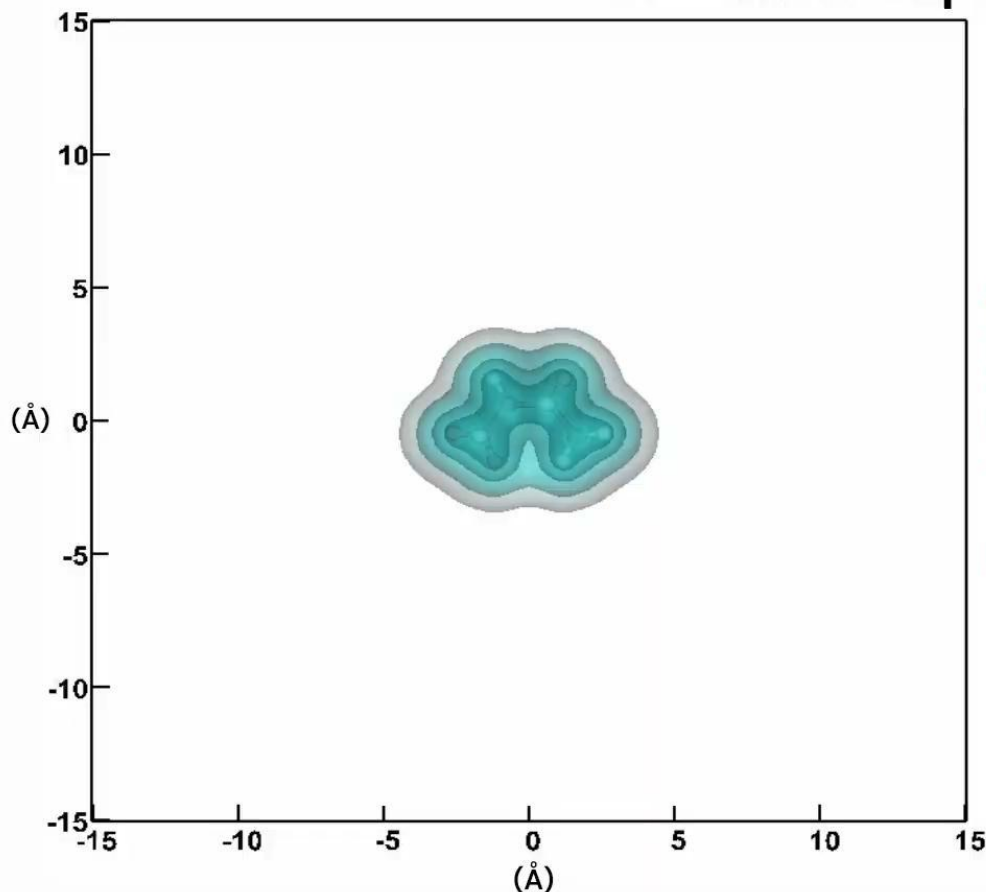
Parameters

- Butane C_4H_{10} Gauche conformation
- Atoms set with random velocities according to *Boltzmann Distribution* at 300 K
- **Laser:**
 - *Intensity:* $7 \times 10^{14} \text{ W/cm}^2$
 - *Peak E-field Strength:* **7.26 V/Å**
 - *GDD:* **0 fs²**
 - *Duration:* **7.8 fs (FWHM)**
 - *Polarization:* **x-direction**
 - *Wavelength:* **890 nm**
 - *Original GDD=0 pulse used in the experiment, but scaled by a factor of 3/16*

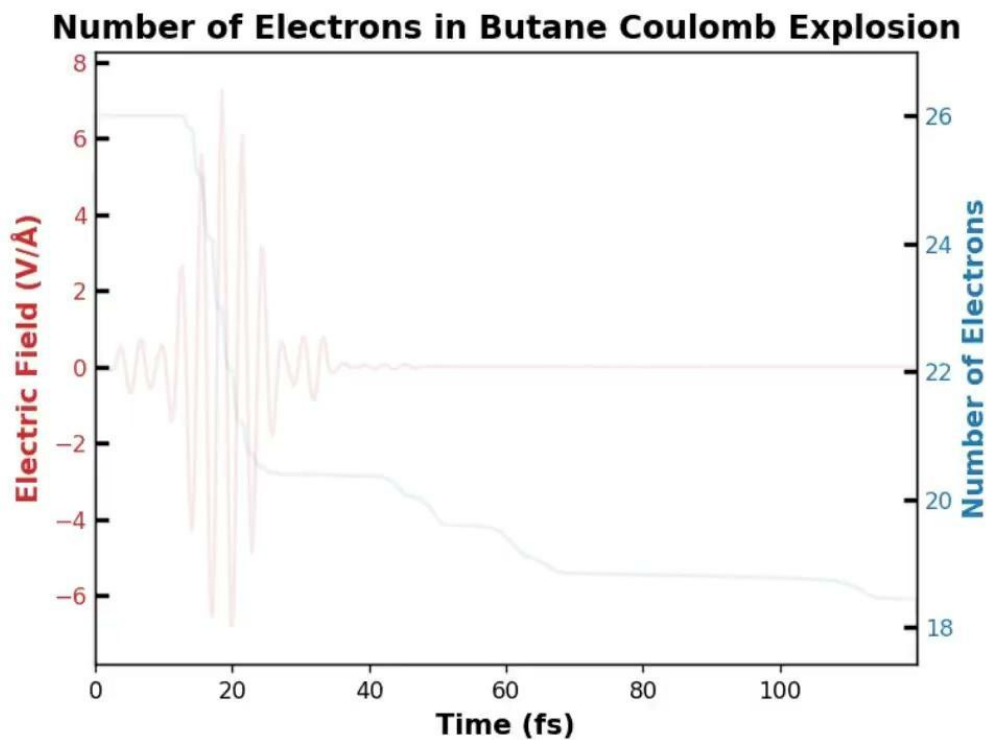


Coulomb Explosion of Butane (C_4H_{10})

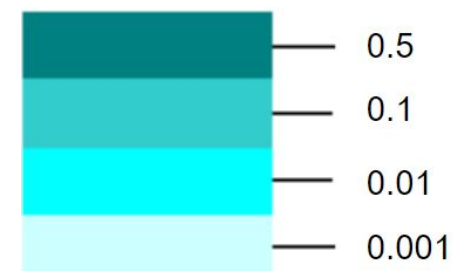
Coulomb Explosion of Butane

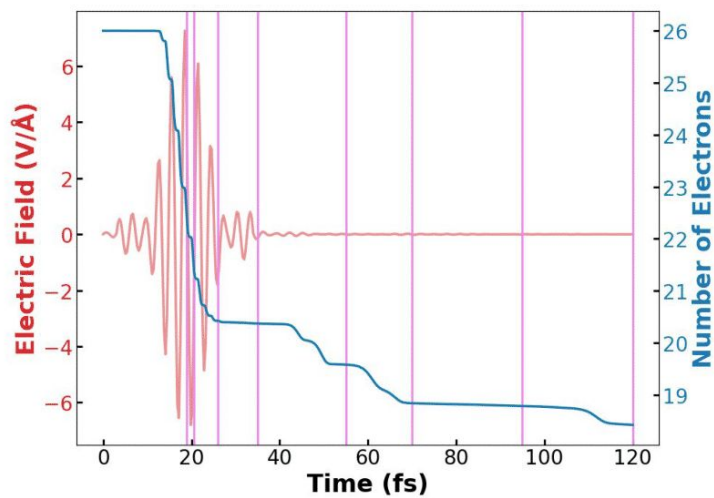
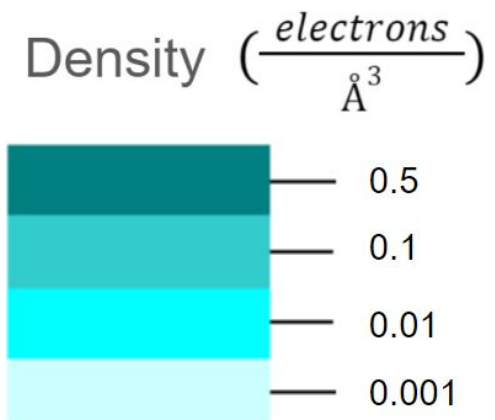
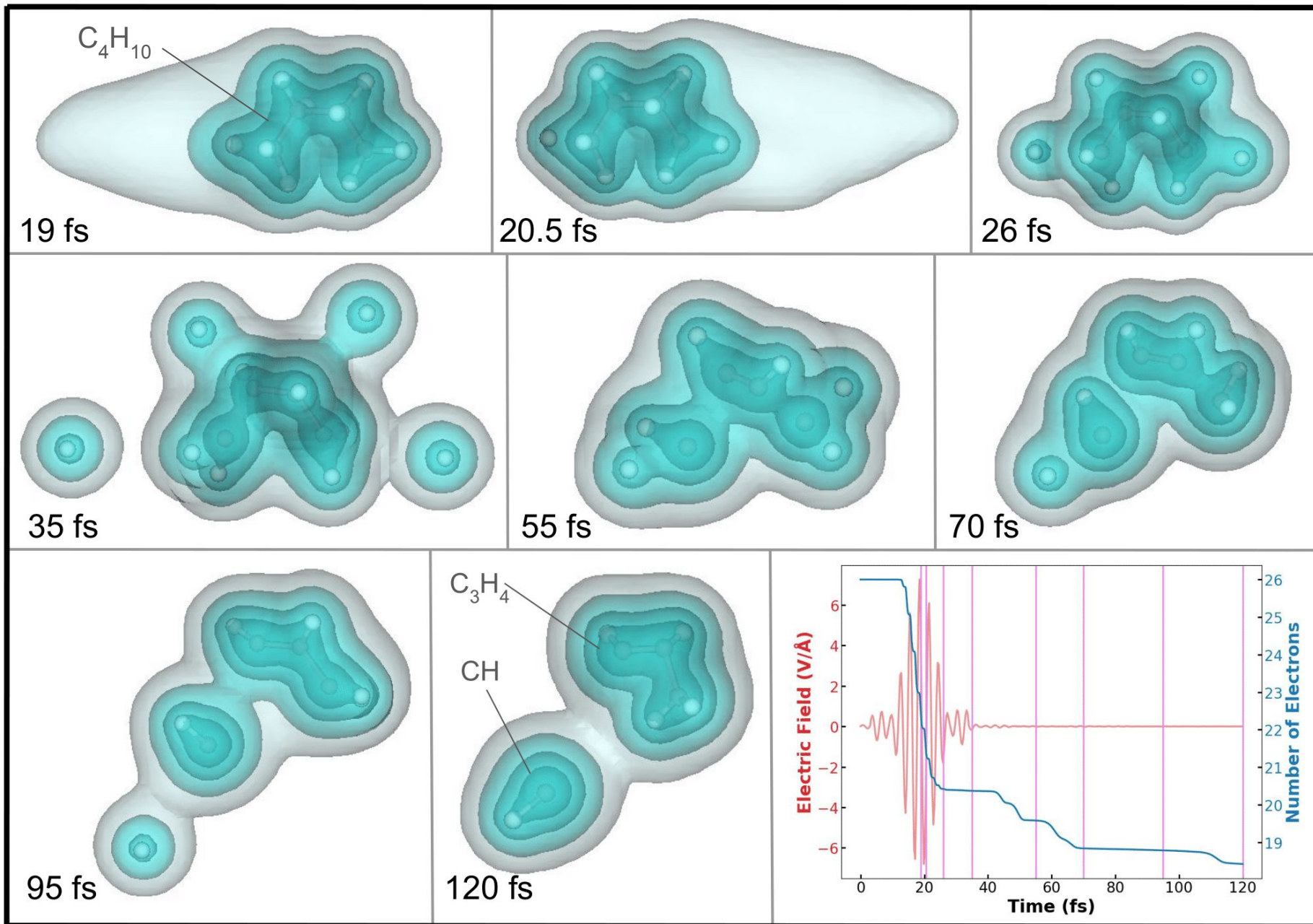
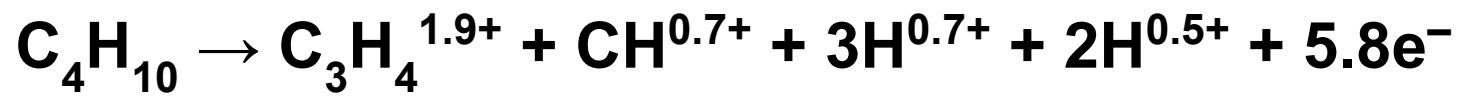


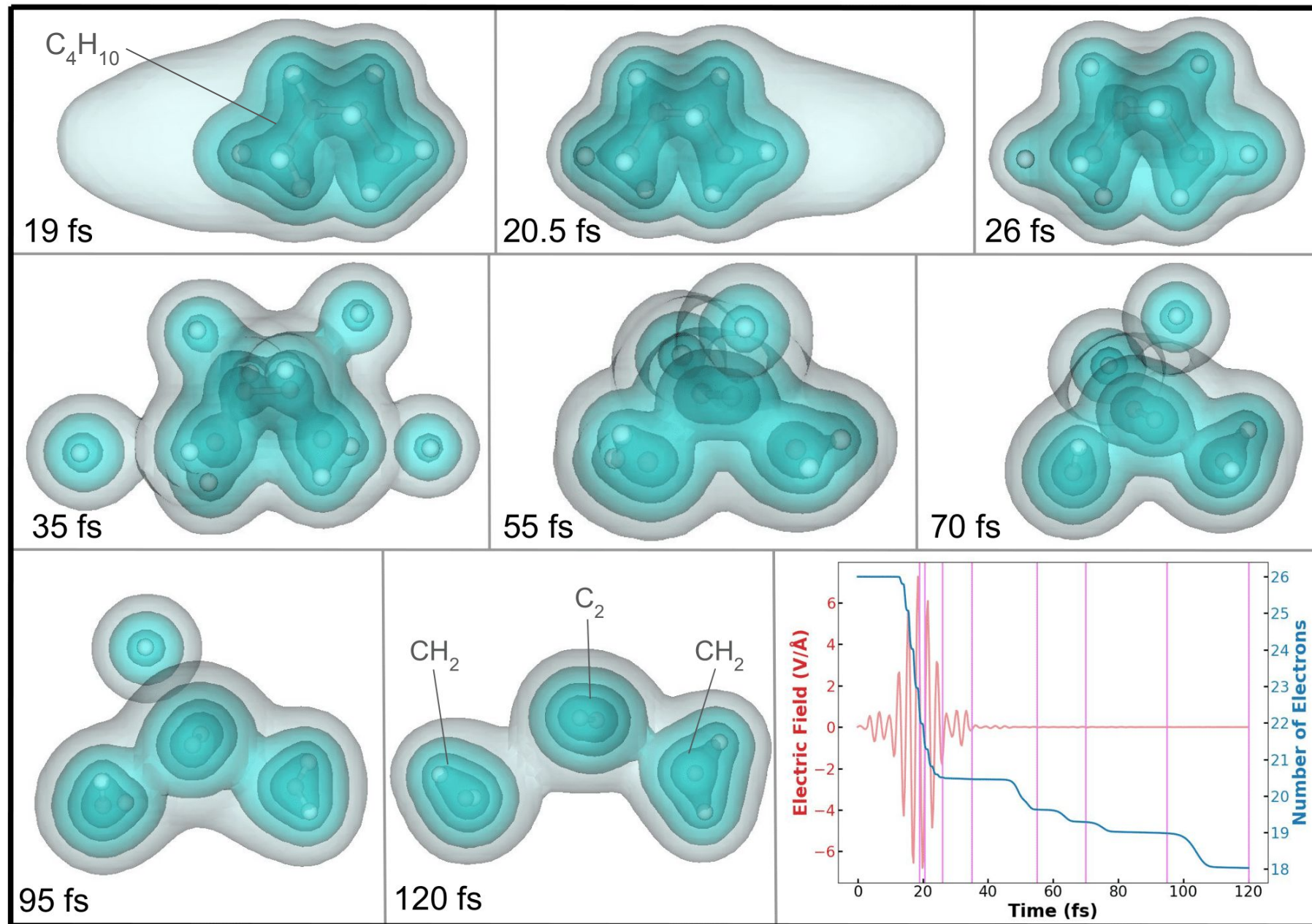
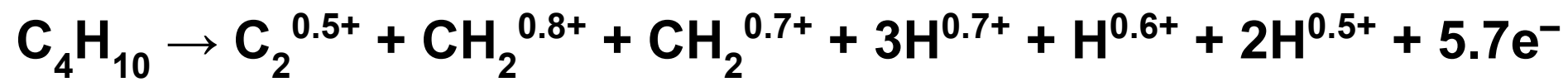
Time: 0 (fs)
E-field: 0.0 (V/\AA)

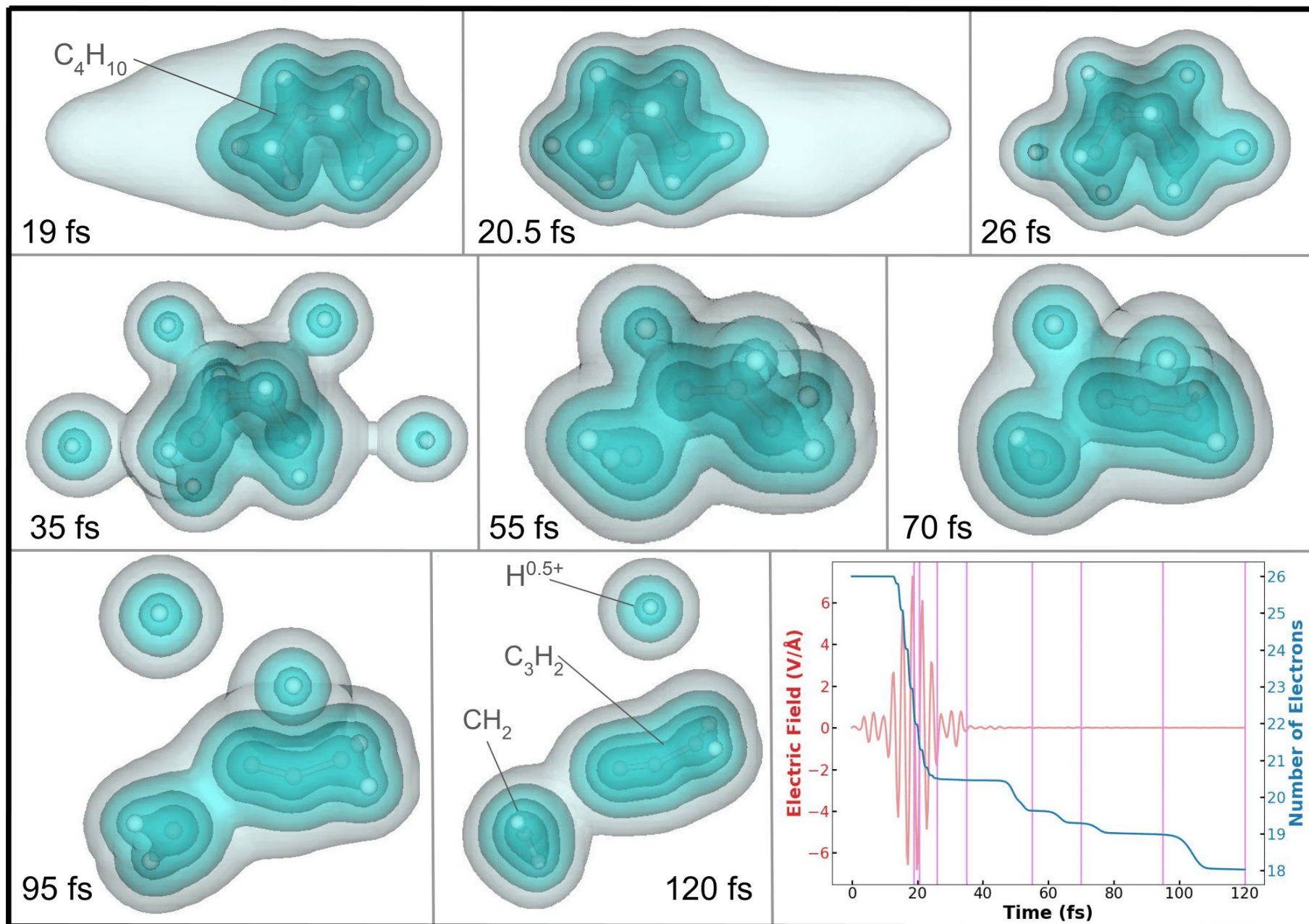
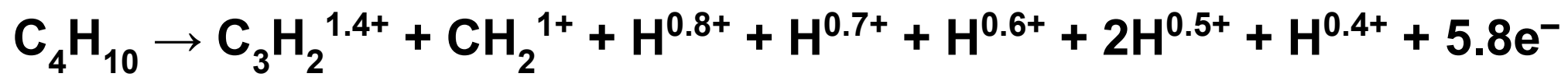


Density ($\frac{\text{electrons}}{\text{\AA}^3}$)

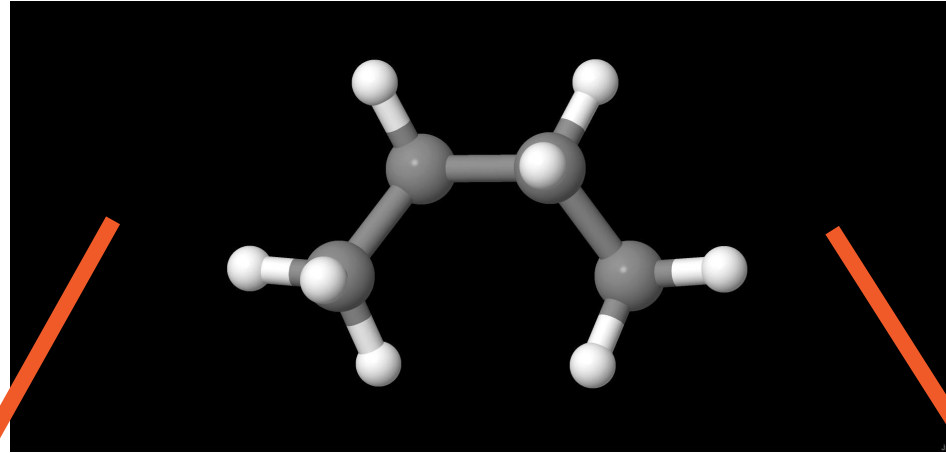






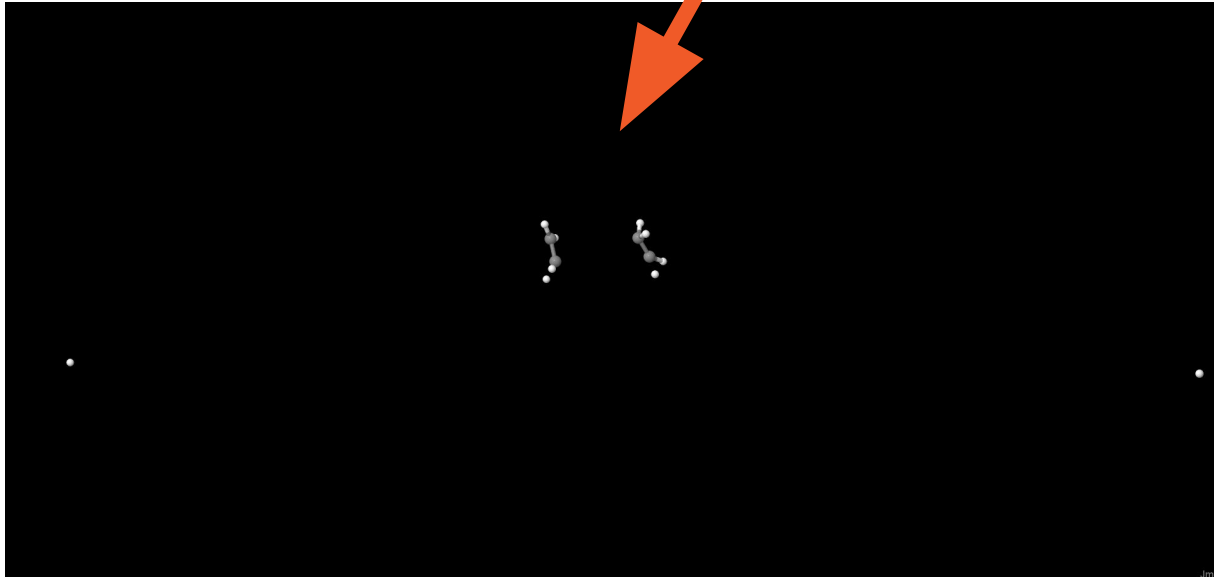


Why reduce the experimental maximum laser intensity ($2 \times 10^{16} \text{ W/cm}^2$) by a factor of $(15/80)^2$ to an intensity of $7 \times 10^{14} \text{ W/cm}^2$?



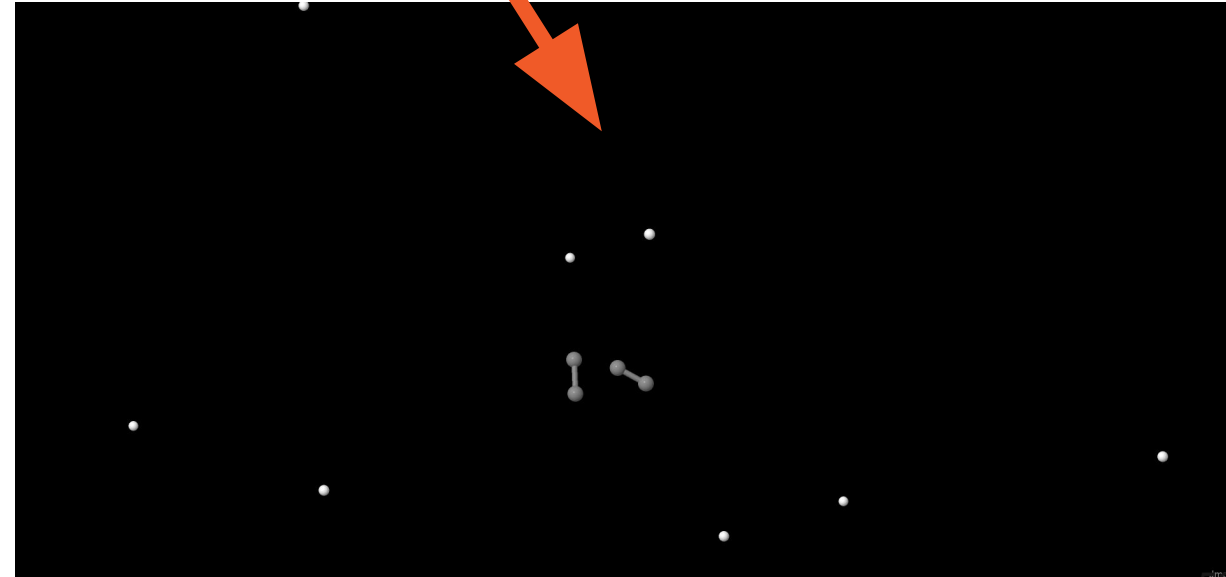
Reduced by factor of $(14/80)^2$

Intensity: 6.125×10^{14}

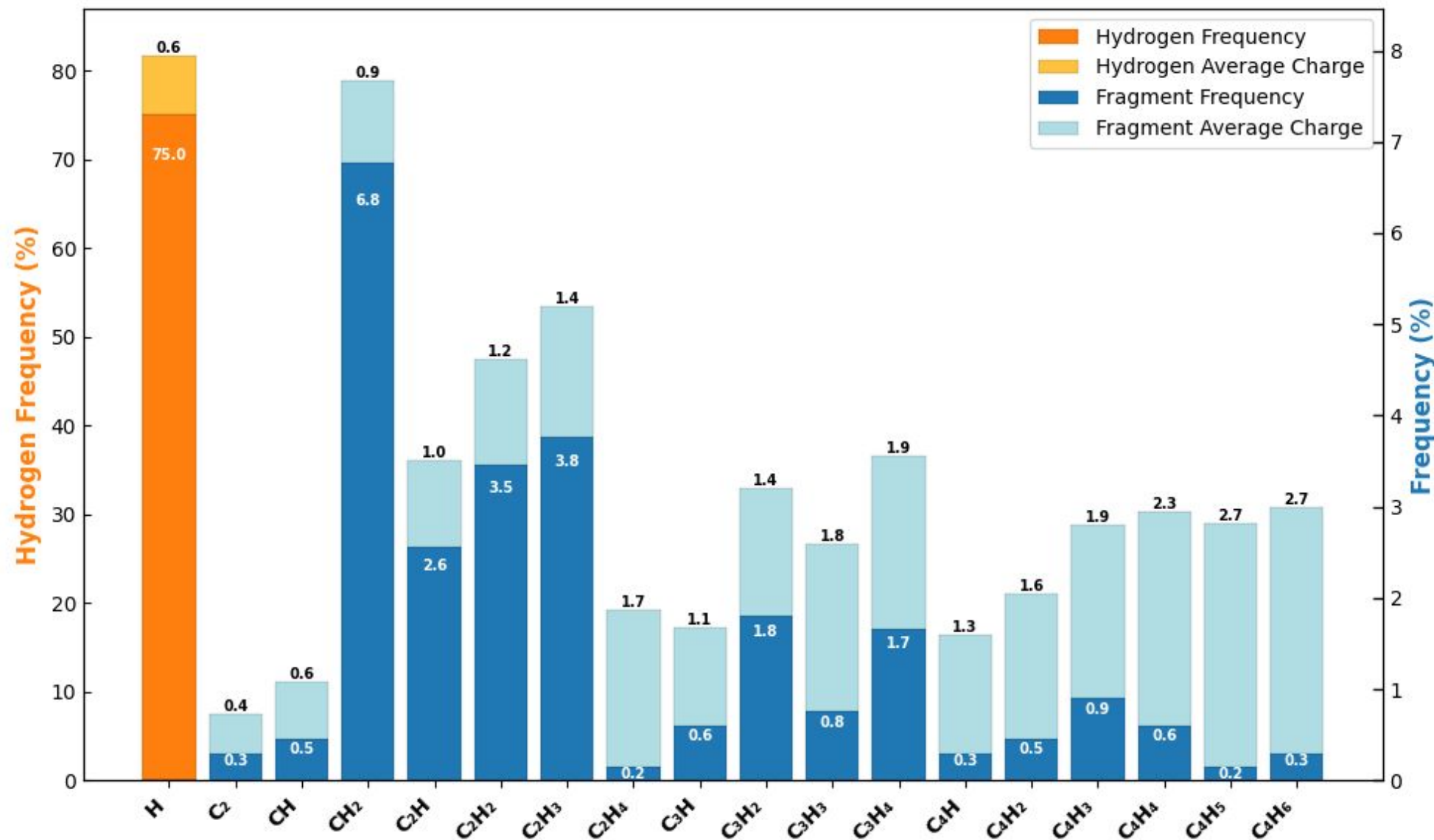


Reduced by factor of $(16/80)^2$

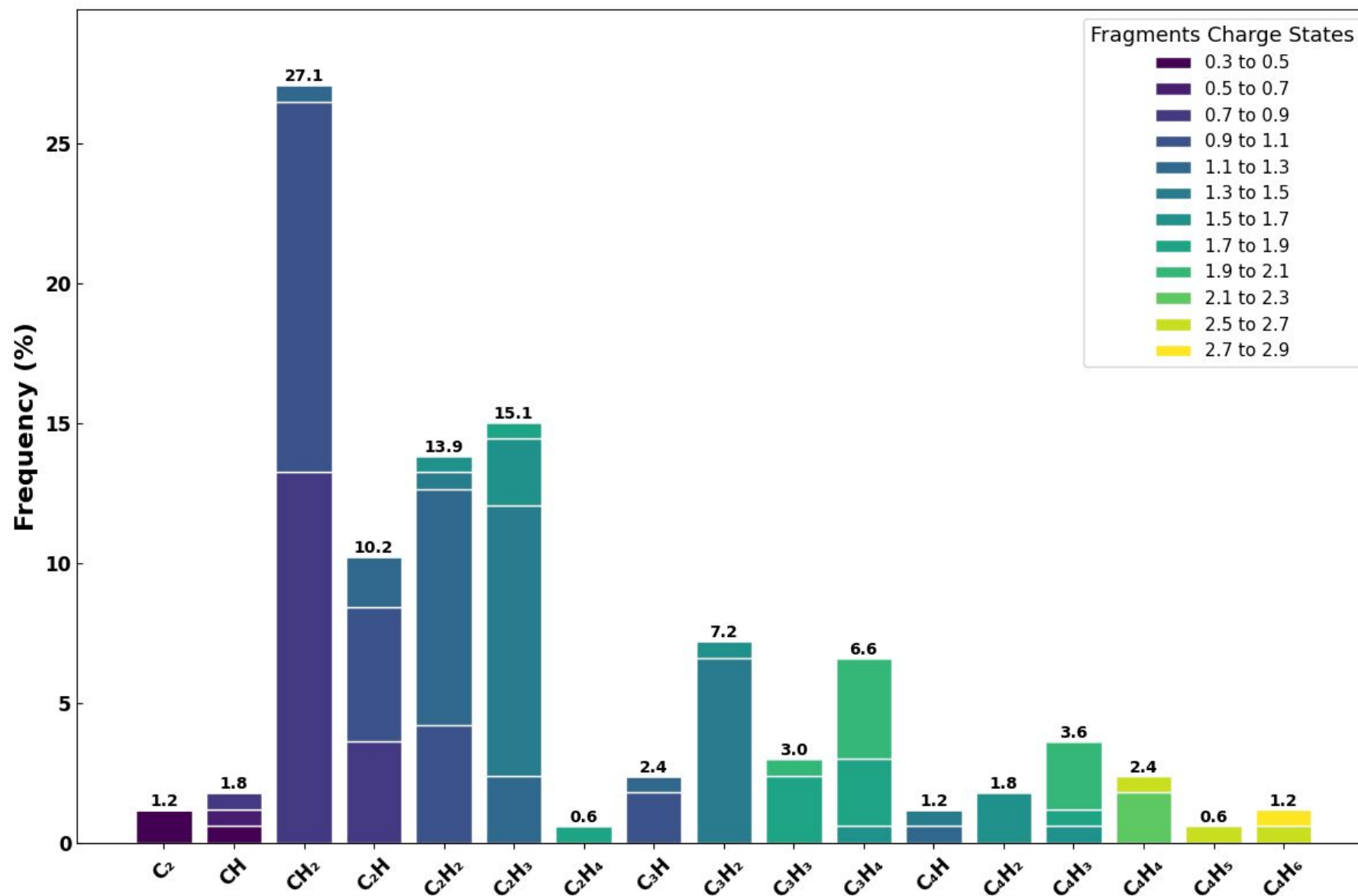
Intensity: 8×10^{14}



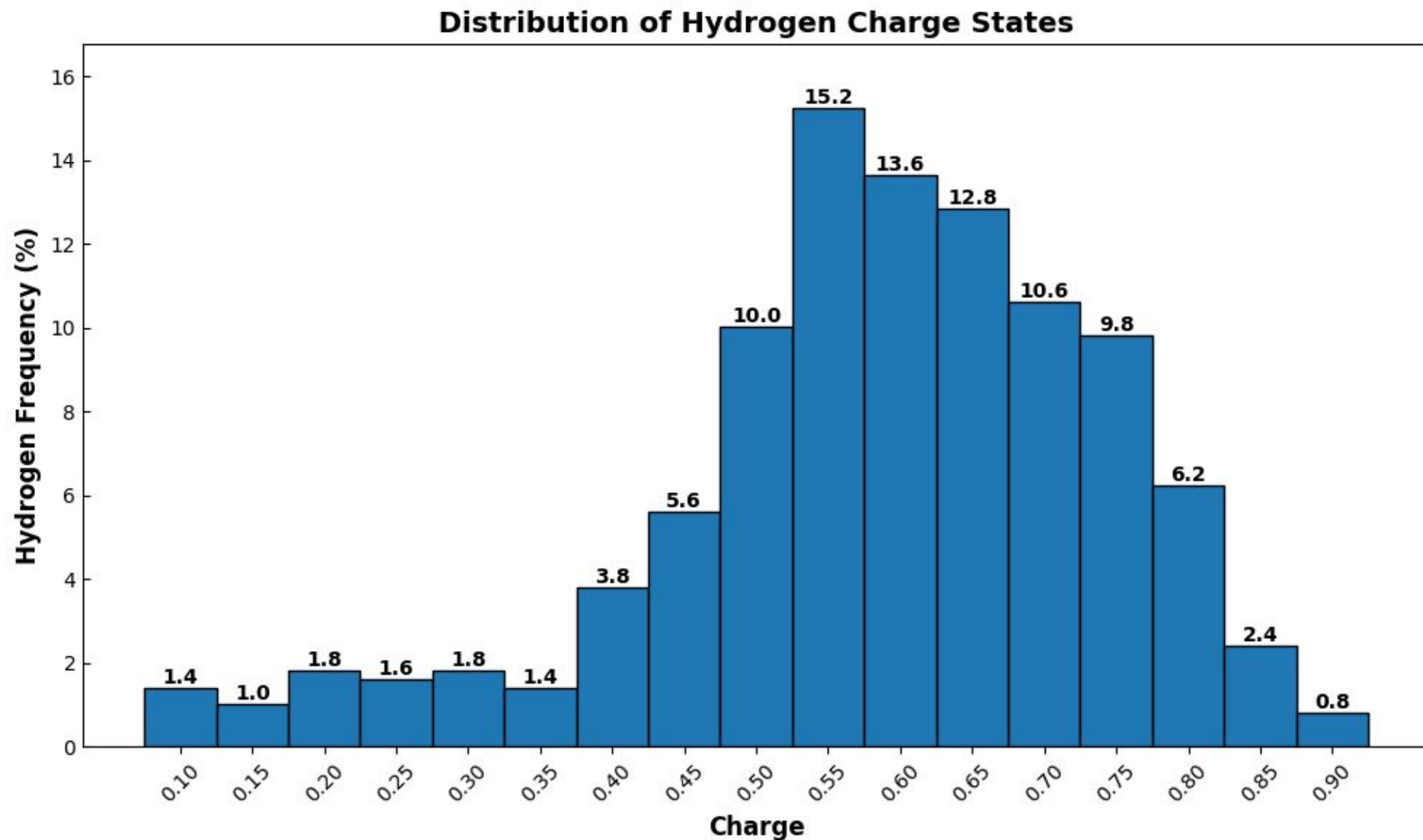
Distribution of Fragments from 88 Butane Coulomb Explosions (Fig. 1)



Distribution of Fragments from 88 Butane Coulomb Explosions (Fig. 2)



Distribution of Charge States from Ejected Hydrogen





Conclusion

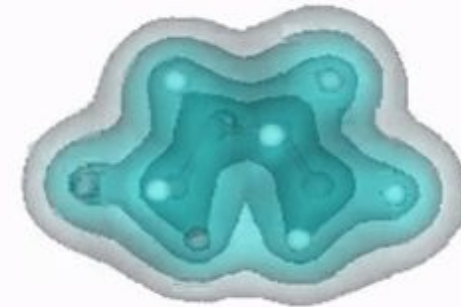
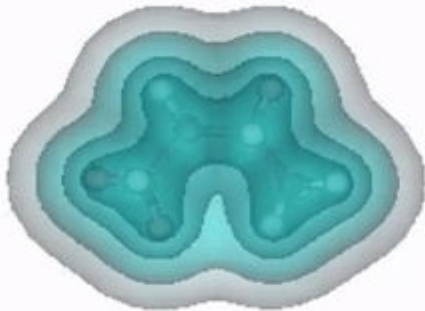
Conclusion

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**TDDFT can be used to accurately predict
Coulomb explosion products.**

**Neutral fragments can be formed
from coulomb explosion**



Currently doing:

- C_4H_{10} with GDD = - 40 fs² pulse
- Methane, ethane, propane with GDD = 0 pulse
- Using supercomputer for the computations
 - ACES (Accelerating Computing for Emerging Sciences)
 - Texas A&M University
- Proton / hydrogen collisions with C_4H_{10} (Other student)

To do:

- Changing direction of laser on molecules
- Other conformations of Butane
- *You tell me!*

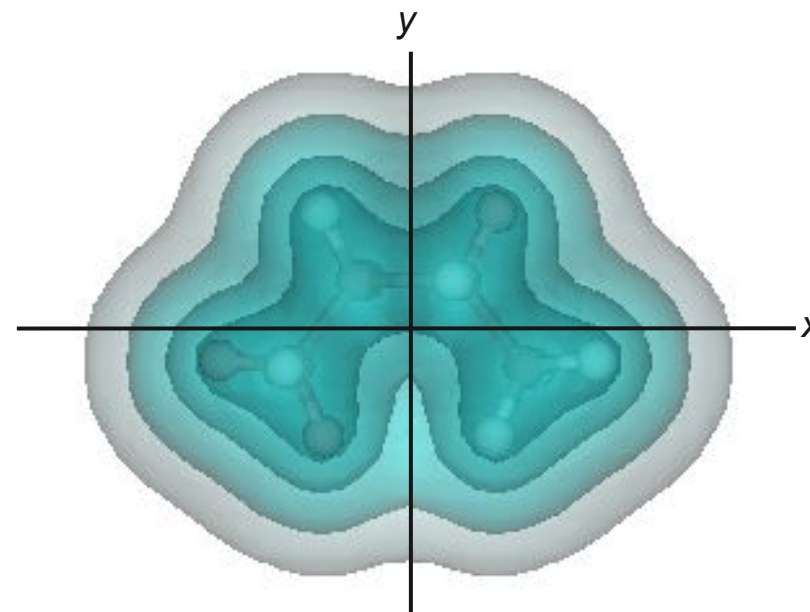


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**THANK YOU FOR YOUR
ATTENTION!**

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European Union
European Regional
Development Fund



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