# 1 Hamiltonian and Lagrangian Systems

# 1.1 Symplectic Geometry

**Definition 1** Let M be a manifold and  $\omega \in \Omega^2(M)$ . Then define the isomorphism  $\flat : \mathfrak{X}(M) \to \mathfrak{X}^*(M)$ ;  $X \mapsto X^{\flat} = i_X \omega$ , and the map  $\sharp$  be its inverse.

**Theorem 1 (Darboux)** Suppose  $\omega$  is a nondegenerate two-form on a 2n-manifold. Then  $d\omega = 0$  iff there is a chart  $(U, \phi)$  around each point m such that  $\phi(m) = 0$  and  $\omega|_U$  is canonical.

**Definition 2** A sympletic form on a manifold M is a nondegenerate, closed two-form  $\omega$  on M. A Symplectic Manifold is a manifold equipped with a symplectic form. The associated volume form is  $\Omega_{\omega} = [(-1)^{[n/2]}/n!]\omega^n$ . The charts in which the symplectic form takes the canonical form are called symplectic charts, and the coordinate functions are called canonical coordinates.

**Definition 3** If  $(M, \omega)$  and  $(N, \rho)$  are symplectic manifolds, a  $C^{\infty}$  map between them that preserves the symplectic structure is called a **canonical transformation**.

**Proposition 1** A canonical transformation has determinant 1 and is a local diffeomorphism.

**Theorem 2** Let  $M = T^*Q$ , with  $\tau_Q^* : M \to Q$  and  $T\tau_Q^* : TQ \to TM$ . Let  $\alpha_q \in M$  and  $\omega_{\alpha_q} \in T_{\alpha_q}M$ . Then let  $\theta_{alpha_q} : T_{\alpha_q}M \to \mathbb{R} : \omega_{\alpha_q} \mapsto \langle \alpha_q, T\tau_Q^*(\omega_{\alpha_q}) \rangle$ , and  $\theta_0 : \alpha_q \mapsto \theta_{\alpha_q}$ . Then  $\omega_0 = -d\theta_0$  is symplectic and the forms  $\omega_0$  and  $\theta_0$  are called the **canonical forms**.

The canonical forms, given in the usual coordinates of a cotangent bundle, take the form:

$$\theta_0 = \sum p_i dq^i$$

$$\omega = \sum dq^i \wedge dp_i$$

The canonical one-form can be thought of as a 'formal adjoint' to the projection operator:

$$\langle \theta(\alpha_q), w_{\alpha_q} \rangle = \langle T \tau_Q^* w_{\alpha_q}, \alpha_q \rangle$$

# 1.2 Hamiltonain Vector Fields and Poisson Brackets

**Definition 4** On a symplectic manifold, given a function  $H: M \to \mathbb{R}$ , the **Hamiltonian Vector Field** associate to the function is a the vector field  $X_H$  satisfying  $\omega(X_H, Y) = \langle dH, Y \rangle$ , or that  $i_{X_H} \omega = dH$ .

**Proposition 2** H is constant along the flow of  $X_H$ .

**Proposition 3** Along a Hamiltonian flow, the symplectic form is conserved.

**Definition 5** A vector field X is **locally Hamiltonian** if for every point, there is a neighborhood U of m such that  $X|_{U}$  is Hamiltonian

### Proposition 4 TFAE:

- 1. X is locally Hamiltonian
- 2.  $\mathcal{L}_X \omega = 0$
- 3. The flow of X consists of canonical transformations

**Remark 1** Locally Hamiltonian vector fields for a Lie subalgebra of  $\mathfrak{X}(M)$ . Globally Hamiltonian vector fields are locally Hamiltonian, but the other way around requires  $H^1(M) = 0$ .

**Definition 6** Let  $\alpha, \beta \in \mathfrak{X}^*(M)$ . Then the Poisson Bracket of alpha and  $\beta$  is the one-form  $-[\alpha^{\sharp}, \beta^{\sharp}]^{\flat}$ 

**Definition 7** Let M be a sympletic manifold and  $f,g:M\to\mathbb{R}$ , then the Poisson bracket of f and g is  $\{f,g\}=-i_{X_f}i_{X_g}\omega$ .

### Proposition 5

$$\{f,g\} = -\mathcal{L}_{X_f}g = \mathcal{L}_{X_g}f$$

Which mean the Poisson bracket is a derivation over f and g individually.

#### Corollary 1

- 1.  $\frac{d}{dt}(f \circ F_t^{X_H}) = \{f \circ F_t^{X_H}, H\}$
- 2.  $d\{f,g\} = \{df,dg\}$

**Definition 8** The Lagrange Bracket of two vector fields is the function  $[[X,Y]] = \omega(X,Y)$  and the Lagrange bracket of a chart is a matrix formed from the Lagrange bracket of each coordinate vector.

**Theorem 3** Let  $(u, \varphi)$  be a chart on a symplectic manifold. Then

- 1.  $\omega|_U = \sum [[u^i, u^j]] du^i \wedge du^j$
- 2. In a symplectic chart, the matrix  $\omega_{ij}$  takes the off-diagonal block matrix form of a almost-complex structure.

3. If 
$$f(q,p) = (Q,P)$$
, then  $[[Q,P]] = \sum_{i=1}^{\infty} \left( \frac{\partial q^i}{\partial Q} \frac{\partial p^i}{\partial P} - \frac{\partial q^i}{\partial P} \frac{\partial p^i}{\partial Q} \right)$ 

4.  $[[q,p]] \circ f^{-1} = [[Q,P]]$ 

**Theorem 4** If X is a locally Hamiltonian vector field, and the pushforward of the canonical coordinates by the flow is denoted  $(Q_t, P_t)$ , then  $[[Q_t, P_t]] \circ F_t^X = [[q, p]]$ 

# 1.3 Integral Invariants, Energy Surfaces, and Stability

**Definition 9** An invariant form for a vector field is one whose Lie derivative is zero.

**Proposition 6** Let X be a vector field and  $\alpha, \beta$  invariant forms of it. Then

- 1.  $i_X \alpha$  is invariant
- 2.  $d\alpha$  is invariant
- 3.  $\mathcal{L}_X \gamma$  is closed  $\iff$   $d\gamma$  is invariant
- 4.  $\alpha \wedge \beta$  is invariant

**Definition 10**  $\alpha$  is relatively invariant  $\iff \mathcal{L}_X \alpha$  is closed.

**Definition 11**  $A_X$  is the algebra of all invariant forms of X,  $\mathcal{R}_X$  the relatively invariant forms of X,  $\mathcal{C}$  the closed forms of  $\Omega(M)$  and  $\mathcal{E}$  the exact forms.

**Theorem 5** The following sequences are exact:

1. 
$$0 \to \mathcal{A}_X \xrightarrow{i} \Omega(M) \xrightarrow{\mathcal{L}_X} \Omega(M) \xrightarrow{\pi} \Omega(M) / Im(\mathcal{L}_X) \to 0$$

2. 
$$0 \to \mathcal{C} \xrightarrow{i} \mathcal{R}_X \xrightarrow{d} \mathcal{A}_X \xrightarrow{\pi} \mathcal{A}_X / (\mathcal{E} \cap \mathcal{A}_X) \to 0$$

Let  $\Sigma_e$  be a connected component of  $H^{-1}(e)$ , where e is a regular value of H.

**Theorem 6** There is a volume element  $\mu_e$  invariant on  $\Sigma_e$  invariant under  $X|\Sigma_e$ 

**Definition 12**  $V \subset M$  is a submanifold is an invariant manifold of a vector field if the vector field is tangent to V at every point.

**Definition 13** Let  $f_k: M \to \mathbb{R}$  be constants of motion for a Hamiltonian system  $X_H$ , and let  $\vec{F} = (f_1, \ldots, f_n): M \to \mathbb{R}^k$ , and c a regular value of  $\vec{F}$ , and let  $\Sigma_c = \vec{F}^{-1}(c)$ . Then  $\Sigma_c$  is an invariant manifold of  $X_H$  of codimension n and there is an invariance volume  $\mu_c$  defined on  $\Sigma_c$ .

### 1.4 Lagrangian Systems

**Definition 14** Let f be any map between vector bundles E, F over the same base space. Then the **Fiber Derivative** of the function f is the function  $\mathbf{F}f$ :  $E \to L(E, F)$ ;  $e \mapsto Df(e)$ .

**Proposition 7** Let  $L: TQ \to \mathbb{R}$ . Then  $FL: TQ \to T^*Q$  is smooth and fiber-preserving.

**Definition 15** Let  $\omega_0$  be the canonical symplectic form on  $T^*Q$  and let  $L: TQ \to \mathbb{R}$ . Then the **Lagrange two-form** is  $\omega_L = (FL)^*\omega_0$ 

**Definition 16** Let Q be a manifold and L a function on the tangent bundle. Then L is a regular Lagrangian if every point is a regular point of FL

**Definition 17** Given  $L: TQ \to \mathbb{R}$ , define the action  $A: TQ \to \mathbb{R}$  by  $A(v) = \langle \mathbf{F}L(v), v \rangle$  and the energy E = A - L. A Lagrangian vector field for L is a vector field  $X_L$  s.t.  $i_{X_L}\omega_L = dE$ .

**Theorem 7** Let  $X_L$  be a Lagrangian vector field for L, then in a chart, the integral curves (u(t), v(t)) satisfy Lagrange's Equations:

$$\frac{d}{dt}u(t) = v(t)$$

$$\frac{d}{dt}\left(\langle D_2L(u(t),v(t)),w\rangle\right) = \langle D_1L(u(t),v(t)),w\rangle$$

 $\forall w \in TQ.$ 

**Theorem 8** Let L and  $\tilde{L}$  be regular Lagrangians, and  $X_L, X_{\tilde{L}}$  be their respective vector fields. Then TFAE:

1. 
$$L = \tilde{L} + \alpha + C$$
,  $d\alpha = 0$ 

2. 
$$X_L = X_{\tilde{L}} \, \& \, \omega_L = \omega_{\tilde{L}}$$

The set of closed one-forms on Q form the 'gauge group' of Lagrangians, i.e. Lagrangians can be transformed without changing the dynamics.

## 1.5 The Legendre Transformation

**Definition 18** L is a hyperregular Lagrangian if  $FL: TQ \to T^*Q$  is a diffeomorphism.

**Theorem 9** Let L be a hyperregular Lagrangian on Q and let  $H = E \circ (\mathbf{F}L)^{-1}$ :  $T^*Q \to \mathbb{R}$ , where E is the energy of L. Then  $\mathbf{F}L$  conjugates the flow  $X_L$  to  $X_H$ .

Theorem 10  $FH = (FL)^{-1}$ 

**Corollary 2** Hyperregular Hamiltonians and Lagrangians correspond bijectively by their fiber derivatives.

### 1.6 Variational Principles in Mechanics

**Definition 19** The path space between two points is defined as  $\Omega(q_1, q_2, [a, b]) = \{c : [a, b] \to Q | c \text{ is a } C^2 \text{ curve}, c(a) = q_1; c(b) = q_2\}$ 

**Proposition 8** The tangent space of the path space is  $T_c\Omega(q_1, q_2, [a, b]) = \{v : [a, b] \to TQ | \pi_Q(v) = c, v(a) = 0, v(b) = 0\}$ 

**Theorem 11** A function satisfies the Euler-Lagrange equations iff the resulting curve is a critical point of the action functional.

**Theorem 12** (Euler-Lagrange-Jacobi-Maupertuis Principle of Least Action) Let  $c_0(t)$  be a base integral curve of  $X_L$ ,  $q_1 = c_0(a)$ ;  $q_2 = c_0(b)$ , and e be the energy of  $c_0(t)$  and be a regular value of a e. Let A be the accumulated (integrate) action along a path. Then dA(c) = 0, and the converse holds.