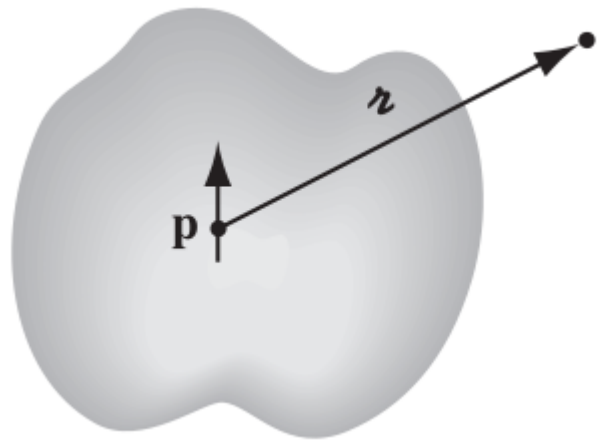


Polarization (\vec{P})

$\vec{P} \equiv$ Dipole moment per unit volume

Field due to a Polarised object

⊗ We have a polarized material.
All the dipoles are pointing in the same direction.

For a single dipole \vec{p} ,

$$V(\vec{r}) = \frac{1}{4\pi\epsilon_0} \frac{\vec{p} \cdot \vec{r}}{r^2}$$

$\vec{r} \equiv$ vector from the dipole to the reference point.

Dipole moment: $\vec{p} = \underline{\vec{p} \, d\tau'}$ for each

infinitesimal volume element dz' .

→ The total potential,

$$V = \frac{1}{4\pi\epsilon_0} \int \frac{\rho(z') \cdot r}{r^2} dz'$$

$$= \frac{1}{4\pi\epsilon_0} \int \rho \cdot \nabla' \left(\frac{1}{r} \right) dz' \quad \Bigg| \quad \nabla' \left(\frac{1}{r} \right) = \frac{-\vec{r}}{r^2}$$

$$\Rightarrow V = \frac{1}{4\pi\epsilon_0} \left[\int \nabla' \cdot \left(\frac{\rho}{r} \right) dz' - \int \frac{1}{r^2} (\nabla \cdot \rho) dz' \right]$$

Using divergence theorem,

$$V = \underbrace{\frac{1}{4\pi\epsilon_0} \oint_S \frac{1}{r} \vec{\rho} \cdot d\vec{a}'}_{\text{Potential due to a surface charge distribution}} - \underbrace{\frac{1}{4\pi\epsilon_0} \int_V \frac{1}{r} (\vec{\nabla}' \cdot \vec{\rho}) d\tau'}_{\text{Potential due to a volume charge distribution}}$$

Potential due to a surface charge distribution

due to a volume charge distribution

Charge density: $\rho_b = \vec{\rho} \cdot \hat{n}$

$$\rho_b = -\vec{\nabla}' \cdot \vec{\rho}$$

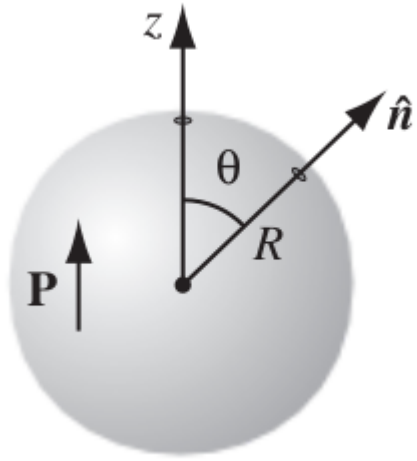
Hence,

$$V(r) = \frac{1}{4\pi\epsilon_0} \oint_S \frac{\rho_b}{r} d\tau' + \frac{1}{4\pi\epsilon_0} \int_V \frac{\rho_b}{r} d\tau'$$

→ Potential of a polarised object is the same as that produced by a volume and a surface charge densities.

There are called bound charges.

\vec{E}_x :



Electric field produced by a uniformly polarized sphere of radius ' R '.

\vec{z} -axis coincides with \vec{p}

→ if \vec{p} is uniform, $\rho_b = 0$

The surface bound charge density,

$$\rho_b = \vec{p} \cdot \hat{n} = p \cos \theta$$

We need to find the potential of a system with surface charge density $\rho_b \cos \theta$

$$V(r, \theta) = \frac{\rho \cos \theta}{3\epsilon_0} \quad \text{for } r \leq R$$

$$= \frac{\rho R^3 \cos \theta}{3\epsilon_0 r^2} \quad \text{for } r \geq R$$

we know, $E = -\frac{\partial V}{\partial r}$

for $r < R$,

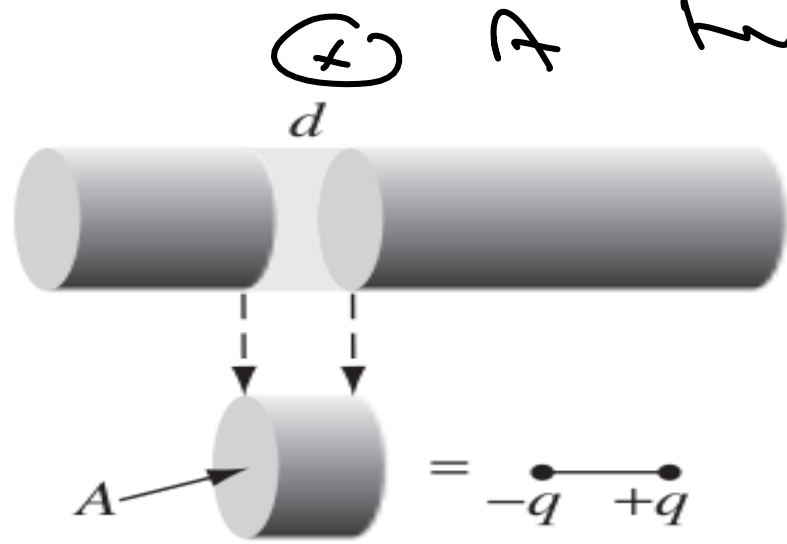
$$E = -\frac{\partial V}{\partial r} = -\frac{\partial}{\partial r} \left(\frac{\rho \cos \theta}{3\epsilon_0} \right) = -\frac{\rho \cos \theta}{3\epsilon_0}$$

Physical interpretation of bound charge.

$$\frac{+}{-} \frac{+}{-} \frac{+}{-} \frac{+}{-} \dots \frac{+}{-} \frac{+}{-} \frac{+}{-} \frac{+}{-} = \frac{+}{-} \frac{+}{-} \rightarrow \text{Bound Charges.}$$

is made up of dipoles consisting of equal amount of +ve and -ve charges.

To calculate the amount of bound charges resulting from a polarisation:
 of dielectric parallel to \vec{P}



On the small chunk,
 the dipole moment
 $= P(A d)$
 \downarrow length of the chunk
 Cross-section
 of the sides

In terms of charges at the
 end, dipole moment $= q d$

$A = A_{\text{end}} \cos \theta$
 \rightarrow The charges
 on the end are
 still the same,
 but

$$q = \frac{q}{A_{\text{end}}} \cos \theta$$

$$= \frac{q \cos \theta}{A}$$

$$= P \cos \theta$$

Then the bound charge

$$P \cdot \hat{n} = \sigma_b$$

$$\Rightarrow \sigma_b = \frac{P \cdot \hat{n}}{1} = \frac{P \cdot \hat{n}}{1}$$

$$\Rightarrow \sigma_b = P \cdot \hat{n}$$

$$\sigma_b = \frac{P}{1} = P = 11.37$$

$$= 11.37$$

The net effect of polarisation is to create a bound charge density $\sigma_b = 11.37$ over the surface of the material.