
Problem 1

```
function f = ffunctgravgradsc02(t,x,IMoIbody,norbit)
%
% Copyright (c) 2019 Mark L. Psiaki. All rights reserved.
%
% This function implements a nonlinear dynamic model
% of a rigid-body spacecraft that includes the
% gravity-gradient torques produced by a  $1/r^2$ 
% gravity field for a spacecraft that is flying
% in a circular orbit. This particular model
% uses quaternion to parameterize the transformation
% from local-level orbit following coordinates
% to the body-fixed coordinates in which
% the moment-of-inertia matrix IMoIbody is
% defined so that the transformation from
% local-level coordinates to body-fixed coordinates
% is defined by the orthonormal rotation matrix:
%
%   R = R(q)
%
% Inputs:
%
%   t                The time, in seconds, at which f is
%                   to be computed.
%
%   x                = [q1;q2;q3;q4;omegabody1;omegabody2;
%                   omegabody3], the 7-by-1 state
%                   vector of this system. The first four
%                   elements give the non-dimensional unit-
%                   normalized attitude quaternion for the
%                   rotation from local-level coordinates
%                   to spacecraft body coordinates. The last
%                   three elements give the body rotation rate
%                   vector relative to inertial coordinates
%                   in radians/second and expressed in
%                   the same body-fixed axes in which
%                   IMoIbody is defined.
%
%   IMoIbody         The 3-by-3 symmetric spacecraft
%                   moment-of-inertia matrix about the
%                   spacecraft center of mass and given
%                   in body-fixed coordinates in units
%                   of kg-m^2.
%
%   norbit            The orbital motion rate in
%                   radians/sec. This is known as
%                   the mean motion in Keplerian
%                   orbital dynamics parlance.
%
% Outputs:
```

```

%
%      f      =
%      [q1dot;q2dot;q3dot;q4dot;omegabody1dot;...
%              omegabody2dot;omegabody3dot],
%              the 7-by-1 vector that contains the
%              computed time derivatives of the state
%              from the kinematics and dynamics models
%              of the spacecraft. f(1:4,1) is given
%              in 1/second units. f(5:7,1) is given in
%              radians/second^2.
%

%
% Determine the rotation matrix from local-level coordinates
% to body-fixed coordinates.
%
q = x(1:4);
qnorm = q*(1/sqrt(sum(q.^2)));
R = rotmatquaternion(qnorm);
%
% Determine the rotation rate of the body-fixed coordinates
% relative to local-level coordinates and given along
% body axes.
%
omegavec = x(5:7);
deltaomegavec = omegavec - R*[0;-norbit;0];
%
% Determine the quaternion time rate of change from
% the quaternion kinematics model.
%
Omegamat = [0 deltaomegavec(3) -deltaomegavec(2)
deltaomegavec(1);...
            -deltaomegavec(3) 0 deltaomegavec(1)
deltaomegavec(2);...
            deltaomegavec(2) -deltaomegavec(1) 0
deltaomegavec(3);...
            -deltaomegavec(1) -deltaomegavec(2) -deltaomegavec(3)
0];
qdot = 0.5*Omegamat*q;
%
% Compute the unit direction vector from the Earth to
% the spacecraft and given in spacecraft body coordinates:
%
rhaticmvec = R*[0 0 -1]';
%
% Compute the gravity-gradient torque in body coordinates.
%
IMoI_rhaticmvec = IMoIbody*rhaticmvec;
Tgravgradvec = 3*(norbit^2)*cross(rhaticmvec,IMoI_rhaticmvec);
%
% Compute the angular velocity rate using Euler's equation.
%
hvec = IMoIbody*omegavec;
omegavecdot = IMoIbody\((cross(-omegavec,hvec)+Tgravgradvec);

```

```
%  
% Assemble the computed state time derivative elements  
% into the output vector.  
%  
f = [qdot;omegavecdot];
```

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Problem 2

```
%script_simgravgradsc11.m
%
% Copyright (c) 2019 Mark L. Psiaki. All rights reserved.
%
% This Matlab script simulates the torque-free motion of
% an axi-symmetric spinning satellite.
%
% This script makes a plot of the
% angular momentum time history in body-fixed
% coordinates and in inertial
% coordinates. It also makes plots
% of the time histories of the 3 body-axis spin-
% rate vector elements.
%
% Clear the Matlab workspace.
%
clear;clc;close all;
%
% Set up the simulation parameters.
%
Itr = 60;
Ispin = 100;
IMoIbody = diag([Itr;Itr;Ispin]);
omegabody0 = [(-0.13*(2*pi/60));(0.07*(2*pi/60));(2*pi/60)];
norbit = 0; % eliminate gravity-gradient torque and
            % rotation of the reference frame relative to which
            % x(1:4,1) defines the attitude quaternion so that
            % it becomes an inertial reference frame rather
            % than a non-inertial orbit-following local-level
            % reference frame.
q0 = [0;0;0;1];
x0 = [q0;omegabody0];
%
% Define the aircraft dynamics function handle
% in a form that is suitable for input to ode45.m.
%
ffunctode45 = @(tdum,xdum) ...
    ffunctgravgradsc02(tdum,xdum,IMoIbody,norbit);
%
% Define the time span of the simulation, computing outputs
% every 0.5. This time span should be large enough
% to see several spins periods and several nutation periods.
%
tspan = ((0:900)')*0.5;
%
% Set up numerical integration options for ode45.m
% in a way that uses a tighter relative tolerance than
% is normally used.
%
optionsode45 = odeset('RelTol',1.e-10);
```

```

%
% Call ode45.m in order to perform numerical integration.
%
tic
[thist,xhist] = ode45(ffunctode45,tspan,x0,optionsode45);
timetosim = toc;
%
% Compute the angular momentum vector time history in
% inertial coordinates.
%
tic
N = size(thist,1);
hvecbodyhist = zeros(N,3);
hvecinertialhist = zeros(N,3);
for k = 1:N
    xk = xhist(k,:)';
    hvecbodyk = IMoIbody*xk(5:7,1);
    hvecbodyhist(k,:) = hvecbodyk';
    qk = xk(1:4,1);
    qknorm = qk*(1/sqrt(sum(qk.^2)));
    Rk = rotmatquaternion(qknorm);
    hvecinertialk = (Rk')*hvecbodyk;
    hvecinertialhist(k,:) = hvecinertialk';
end
timetohvecinertial = toc
clear k xk hvecbodyk qk qknorm Rk hvecinertialk
%
% Compute the nutation frequency.
%
omeganut = abs(Ispin-Itr)*omegabody0(3)/Itr;
%
% Compute the theoretical body-axis spin vector component
% time histories that are valid for this axially-symmetric
% spacecraft.
%
signItrminusIspin = sign(Itr - Ispin);
omegabody1hist = omegabody0(1)*cos(omeganut*thist)
+omegabody0(2)*signItrminusIspin*sin(omeganut*thist);
omegabody2hist = omegabody0(2)*cos(omeganut*thist) -
omegabody0(1)*signItrminusIspin*sin(omeganut*thist);
omegabody3hist = ones(N,1)*omegabody0(3);
%
% Plot the body-axes angular momentum time history.
%
figure(1)
hold off
plot(thist,hvecbodyhist,'LineWidth',2)
set(get(gcf,'CurrentAxes'),'FontSize',16)
grid
xlabel('Time (sec)')
ylabel('Angular Momentum (N-m-sec)')
title(['Body-Axes Angular Momentum','...
' script\_simgravgradsc11.m'])

```

```

    legend('h1 body','h2 body','h3 body')
%
% Plot the inertial angular momentum time history.
%
figure(2)
hold off
plot(thist,hvecinertialhist,'LineWidth',2)
set(get(gcf,'CurrentAxes'),'FontSize',16)
grid
xlabel('Time (sec)')
ylabel('Angular Momentum (N-m-sec)')
title(['Inertial Angular Momentum','...
        ' script\_simgravgradsc11.m'])
legend('h1 ECIF','h2 ECIF','h3 ECIF')
%
% Plot the body-axis spin vector component time histories,
% both from the numerical integration and the theoretical
% values.
%
figure(3)
hold off
plot(thist,xhist(:,5:7),'-','LineWidth',2)
set(get(gcf,'CurrentAxes'),'FontSize',16)
hold on
plot(thist,[omegabody1hist,omegabody2hist,omegabody3hist],'.',...
        'MarkerSize',10)
grid
xlabel('Time (sec)')
ylabel('Angular Velocity (radians/sec)')
title(['Body-Axis Angular Velocity','...
        ' script\_simgravgradsc11.m'])
legend('omegabody1 sim','omegabody2 sim','omegabody3 sim',...
        'omegabody1 theory','omegabody2 theory','omegabody3 theory')
%
% Save the results.
%
textcommands = ['These data have been generated by the',...
                ' commands in script_simgravgradsc11.m'];
save simgravgradsc11

format long
xfinal = xhist(end,:)'
qfinalmag = norm(xfinal(1:4,1))

disp('To verify that angular momentum is conserved, we subtract
the initial value of inertial angular momentum from the rest of
its time history and compute the norm. The closer this is to zero the
better the conservation of angular momentum is. This would imply zero
external torque')
disp('The norm on variation of inertial angular momentum is:')
disp(norm(hvecinertialhist(:,1)-
hvecinertialhist(1,1))+norm(hvecinertialhist(:,2)-
hvecinertialhist(1,2))+norm(hvecinertialhist(:,3)-
hvecinertialhist(1,3)))

```

```
disp('The norm of the difference in the time histories between  
theory and simulated angular velocities is:')  
disp(norm(omegabody1hist-xhist(:,5))+norm(omegabody2hist-  
xhist(:,6))+norm(omegabody3hist-xhist(:,7)))
```

```
timetohvecinertial =
```

```
0.0105829000000000
```

```
xfinal =
```

```
0.076778204795434  
-0.041341864928188  
-0.984339512015986  
0.153186692608215  
-0.013613444068028  
0.007330204591270  
0.104719755119660
```

```
qfinalmag =
```

```
0.999997240113923
```

To verify that angular momentum is conserved, we subtract the initial value of inertial angular momentum from the rest of its time history and compute the norm. The closer this is to zero the better the conservation of angular momentum is. This would imply zero external torque

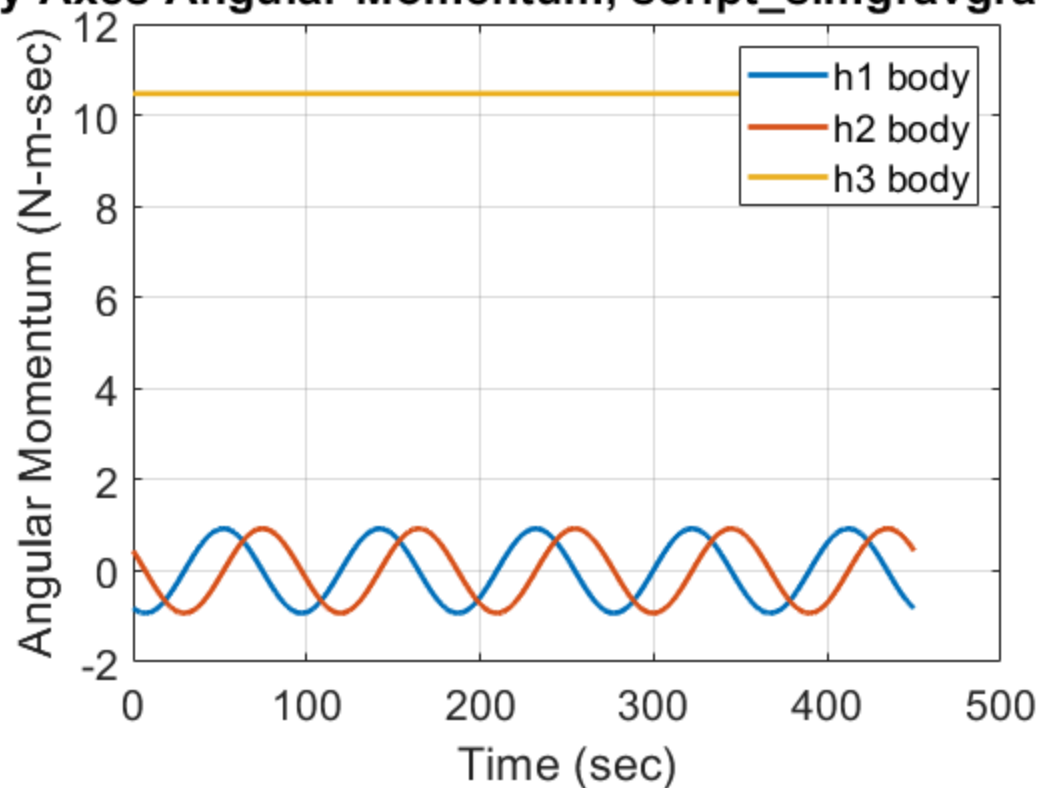
The norm on variation of inertial angular momentum is:

```
1.320999550676519e-04
```

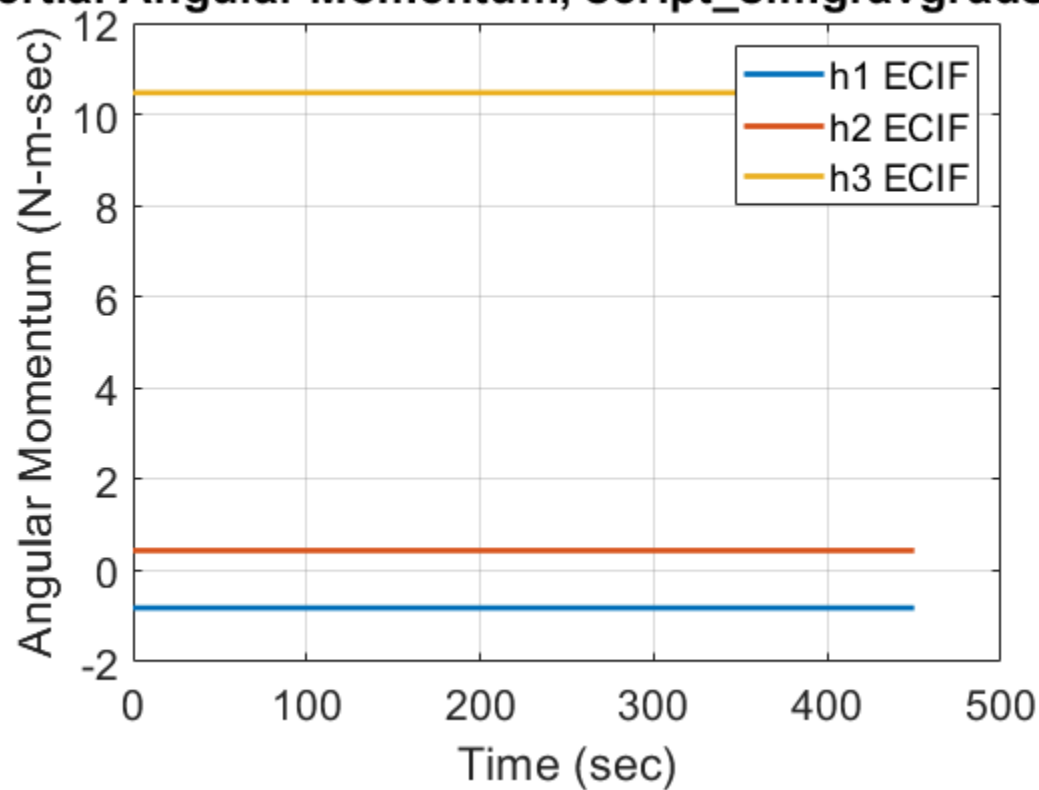
The norm of the difference in the time histories between theory and simulated angular velocities is:

```
5.332291859654702e-06
```

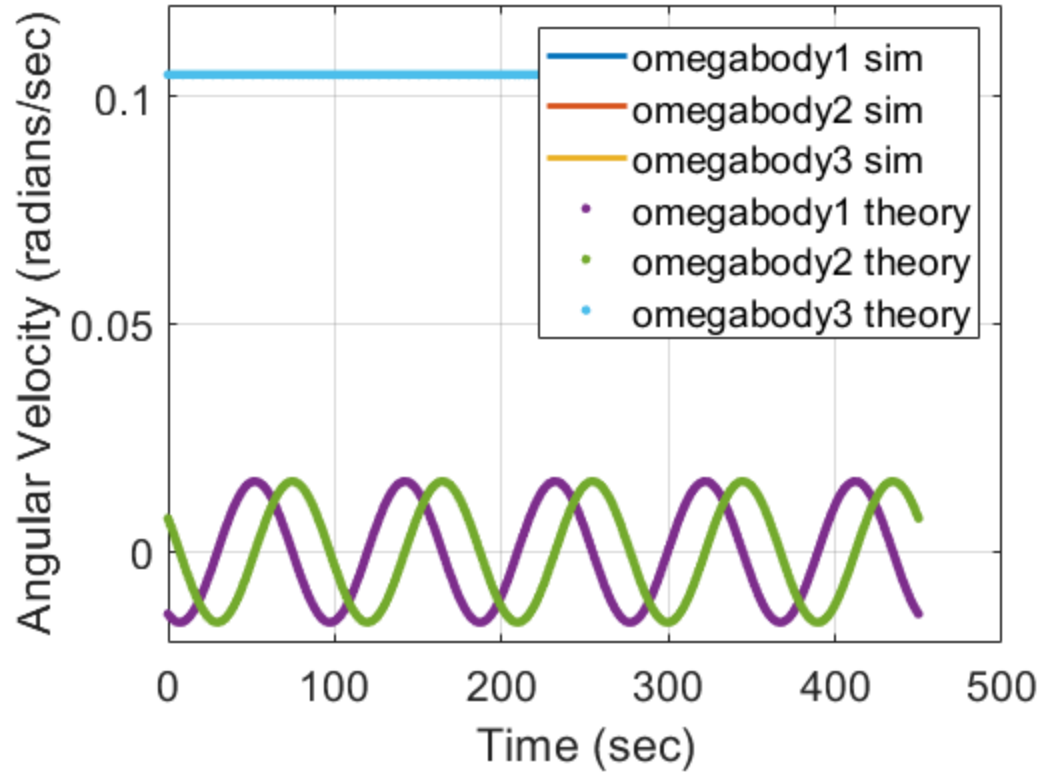
ly-Axes Angular Momentum, script_simgravgradsc



ierial Angular Momentum, script_simgravgradsc1



Body-Axis Angular Velocity, script_simgravgradsc'



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Problem 3

```
%script_simgravgradsc13.m
%
% Copyright (c) 2019 Mark L. Psiaki. All rights reserved.
%
% This Matlab script simulates the motion of a
% nadir-pointing satellite that is acting
% under the influence of a gravity-gradient torque
% in a circular Low-Earth Orbit (LEO) that has an
% orbital period of 6000 sec. The initial conditions
% start with zero roll, pitch, and yaw angles relative
% to the principal axes, but with non-zero initial roll,
% pitch, and yaw rates relative to the principal axes.
%
% The principal axes are not exactly body axes, which is
% why the point about which the satellite attitude
% oscillates is not exactly zero for the roll, pitch, and
% yaw angles.
%
% This script also makes plots of the roll, pitch, and
% yaw attitude time history.
%
% Clear the Matlab workspace.
%
clear;clc;close all;
%
% Load simulation parameters IMoIbody, norbit, omegabody0
% and q0 from the file simgravgradsc13_data.mat.
%
load simgravgradsc13_data
%
% Set up the initial state. This initial state vector
% has been chosen with a knowledge of the principal
% axes coordinates, and q0 has been chosen so that
% the principal axes are exactly aligned with
% the local-level orbit-following coordinate
% system so that, had omegabody0 also been chosen
% correctly, q(t) and omegabody(t) would have
% remained constant. A slightly perturbed
% omegabody0 has been chosen to produce motion.
% This allows the stability of the gravity-gradient
% system to be tested.
%
x0 = [q0;omegabody0];
%
% Define the aircraft dynamics function handle
% in a form that is suitable for input to ode45.m.
%
ffunctode45 = @(tdum,x dum) ...
    ffunctgravgradsc02(tdum,x dum,IMoIbody,norbit);
%
```

```

% Define the time span of the simulation, computing outputs
% 1000 times per orbit for 6 orbits.
%
Torbit = 2*pi/norbit;
tspan = (0:6000)'.*(Torbit/1000);
%
% Set up numerical integration options for ode45.m
% in a way that uses a tighter relative tolerance than
% is normally used.
%
optionsode45 = odeset('RelTol',1.e-10);
%
% Call ode45.m in order to perform numerical integration.
%
tic
[thist,xhist] = ode45(ffunctode45,tspan,x0,optionsode45);
timetosim = toc;
%
% Determine the 3-1-2 Euler angle time histories. The function
% yawpitchrollcalc02.m has been designed with the 3-1-2
% assumption specifically in mind. Note that the
% computation of qknorm via normalization is included
% in order to remove any small errors in the quaternion
% unit normalization that may have built up due
% to numerical intergration error of the quaternion
% kinematics.
%
N = size(thist,1);
phihist = zeros(N,1);
thetahist = zeros(N,1);
psihist = zeros(N,1);
for k = 1:N
    qk = xhist(k,1:4)';
    qknorm = qk*(1/sqrt(sum(qk.^2)));
    [phik,thetak,psik] = yawpitchrollcalc02(qknorm);
    phihist(k,1) = phik;
    thetahist(k,1) = thetak;
    psihist(k,1) = psik;
end
clear k qk qknorm phik thetak psik
%
% Plot the roll, pitch, and yaw time histories.
% An analysis of this case indicates that there should
% be 8.05 pitch angle oscillations, 4.61 oscillations
% of one of the coupled roll-yaw modes, and 11.4
% oscillations of the other coupled roll-yaw modes.
% Of course, the actual roll, pitch, and yaw
% principal axes differ slightly from the
% body axes used in this simulation. Therefore,
% the roll, pitch, and yaw angle time histories plotted
% below are not pure principal axes quantities.
% This fact causes slight additional coupling
% between modes beyond the theoretical coupling of the
% two roll/yaw modes of oscillation.

```

```

%
figure(1)
hold off
plot(thist*(1/Torbit),[phihist,thetahist,psihist]*(180/pi),...
     'LineWidth',2)
hold on
plot(thist*(1/Torbit),ones(N,1)*...
     ([phihist(1,1),thetahist(1,1),psihist(1,1)]*(180/pi)),...
     ':','LineWidth',1.5)
hold off
set(get(gcf,'CurrentAxes'),'FontSize',10)
grid
xlabel('Time (orbits)')
ylabel('Euler Angle (deg)')
title(['3-1-2 Euler Angle Time Histories','...
      ' script\_simgravgradsc13.m'])
legend('Roll','Pitch','Yaw',...
      'Roll Equilibrium','Pitch Equilibrium',...
      'Yaw Equilibrium')

%
% Save the results.
%
textcommands = ['These data have been generated by the',...
               ' commands in script_simgravgradsc13.m'];
save simgravgradsc13
format long
xfinal = xhist(end,:)
qfinalmag = norm(xfinal(1:4,1))

disp('The system produces a neutrally stable response, evident
from the time histories showing non increasing crest and troughs in
attitude over 6 orbits')

```

```
xfinal =
```

```

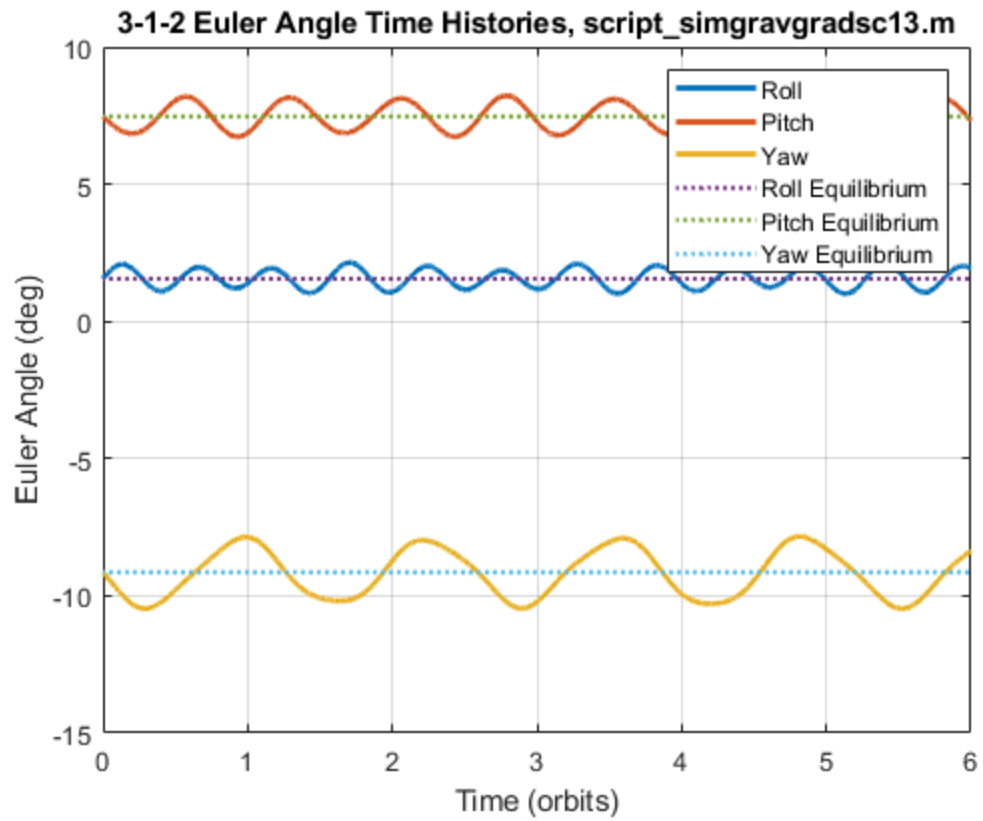
0.021414146457538
0.062360917783887
-0.071803040606404
0.995237366114879
0.000136192378780
-0.001052654303743
0.000062483571553

```

```
qfinalmag =
```

```
1.000000270643444
```

The system produces a neutrally stable response, evident from the time histories showing non increasing crest and troughs in attitude over 6 orbits



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Q4) the Euler's Eq

$$\frac{d\vec{h}}{dt} = \vec{T}$$
 where $\vec{h} \rightarrow$ Angular momentum
 $\vec{T} \rightarrow$ External torque
 for \vec{h}^b , Angular momentum Expressed in
 a body frame Rotating as $\vec{\omega}^b$, the Euler Eq
 is as follows

$$\frac{d\vec{h}^b}{dt} + \vec{\omega}^b \times \vec{h}^b = \vec{T}^b$$

for a Torque-Free case $\vec{T}^b = \begin{bmatrix} 0 \\ 0 \\ 0 \end{bmatrix}$

Expanding $\vec{h}^b = I_{MOI}^b \vec{\omega}^b$

where I_{MOI}^b is a principle Frame of the
 body

$$I_{MOI}^b = \begin{bmatrix} I_{tr1} & 0 & 0 \\ 0 & I_{tr2} & 0 \\ 0 & 0 & I_{spin} \end{bmatrix}$$

So, the torque-free rigid body equations become:

$$\underbrace{\begin{bmatrix} I_{tr1} & 0 & 0 \\ 0 & I_{tr2} & 0 \\ 0 & 0 & I_{spin} \end{bmatrix}}_{I_{MOI}^b} \underbrace{\begin{bmatrix} \dot{\omega}_1 \\ \dot{\omega}_2 \\ \dot{\omega}_3 \end{bmatrix}}_{\dot{\vec{\omega}}^b} + \underbrace{\begin{bmatrix} \omega_1 \\ \omega_2 \\ \omega_3 \end{bmatrix}}_{\vec{\omega}^b} \times \left(\underbrace{\begin{bmatrix} I_{tr1} & 0 & 0 \\ 0 & I_{tr2} & 0 \\ 0 & 0 & I_{spin} \end{bmatrix}}_{I_{MOI}^b} \underbrace{\begin{bmatrix} \omega_1 \\ \omega_2 \\ \omega_3 \end{bmatrix}}_{\vec{\omega}^b} \right) = \underbrace{\begin{bmatrix} 0 \\ 0 \\ 0 \end{bmatrix}}_{\vec{T}=0}$$

further simplifying,

$$\begin{bmatrix} I_{TR1} \dot{\omega}_1 \\ I_{TR2} \dot{\omega}_2 \\ I_{spin} \dot{\omega}_3 \end{bmatrix} + \begin{bmatrix} 0 & -\omega_3 & \omega_2 \\ \omega_3 & 0 & -\omega_1 \\ -\omega_2 & \omega_1 & 0 \end{bmatrix} \begin{bmatrix} I_{TR1} \omega_1 \\ I_{TR2} \omega_2 \\ I_{spin} \omega_3 \end{bmatrix} = \begin{bmatrix} 0 \\ 0 \\ 0 \end{bmatrix}$$

$$I_{TR1} \dot{\omega}_1 + \omega_2 \omega_3 (I_{spin} - I_{TR2}) = 0$$

$$I_{TR2} \dot{\omega}_2 + \omega_3 \omega_1 (I_{TR1} - I_{spin}) = 0$$

$$I_{spin} \dot{\omega}_3 + \omega_1 \omega_2 (I_{TR2} - I_{TR1}) = 0$$

$$\dot{\omega}_1 = \frac{(I_{TR2} - I_{spin})}{I_{TR1}} \omega_2 \omega_3$$

$$\dot{\omega}_2 = \frac{(I_{spin} - I_{TR1})}{I_{TR2}} \omega_1 \omega_3$$

$$\dot{\omega}_3 = \frac{(I_{TR1} - I_{TR2})}{I_{spin}} \omega_1 \omega_2$$

If ω_1 & ω_2 are very small, then $\omega_1, \omega_2 \ll 0$

$$\therefore \dot{\omega}_3 = 0, \text{ so } \omega_3(t) = \text{const} = \omega_{3,avg}$$

We substitute $\omega_3(t) = \omega_{3avg}$ in $\dot{\omega}_1$ & $\dot{\omega}_2$ Eq to get

$$\dot{\omega}_2 = \left(\frac{I_{TR2} - I_{spin}}{I_{TR1}} \right) \omega_{3avg} \omega_2$$

differentiation once w.r.t to time

$$\ddot{\omega}_2 = \left(\frac{I_{TR2} - I_{spin}}{I_{TR1}} \right) \omega_{3avg} \dot{\omega}_2$$

Substituting the Eq for $\dot{\omega}_2 = \left(\frac{I_{TR1} - I_{spin}}{I_{TR2}} \right) \omega_{3avg} \omega_1$
we get

$$\ddot{\omega}_1 = \frac{(I_{TR2} - I_{spin})(I_{TR1} - I_{spin})}{I_{TR1} I_{TR2}} \omega_{3avg}^2 \omega_1$$

give the 2nd order ODE as follows.

$$\ddot{\omega}_1 + \underbrace{\left(\frac{(I_{TR2} - I_{spin})(I_{TR1} - I_{spin})}{I_{TR1} I_{TR2}} \omega_{3avg}^2 \right)}_{\omega_{nut}^2} \omega_1 = 0$$

$\omega_{nut}^2 \rightarrow$ the nutation frequency

$$\therefore \omega_{nut} = \sqrt{\frac{(I_{spin} - I_{TR1})(I_{spin} - I_{TR2})}{I_{TR1} I_{TR2}}} \omega_{3avg} \quad \blacksquare$$

the assumption here is $I_{spin} > I_{TR1}$ & $I_{spin} > I_{TR2}$
(or) $I_{spin} < I_{TR1}$ & $I_{spin} < I_{TR2}$

The general solutions to the 2nd order PDEs are

$$\omega_1(t) = A \cos(\omega_{\text{nut}} t) + B \sin(\omega_{\text{nut}} t)$$

$$\omega_2(t) = C \cos(\omega_{\text{nut}} t) + D \sin(\omega_{\text{nut}} t)$$

$$\dot{\omega}_1(t) = -A \omega_{\text{nut}} \sin(\omega_{\text{nut}} t) + B \omega_{\text{nut}} \cos(\omega_{\text{nut}} t)$$

using the initial dynamics Eq

$$-A \omega_{\text{nut}} \sin(\omega_{\text{nut}} t) + B \omega_{\text{nut}} \cos(\omega_{\text{nut}} t)$$

$$= \left(\frac{I_{\text{TR2}} - I_{\text{spin}}}{I_{\text{TR1}}} \right) \omega_{\text{3avg}} \omega_2$$

$$= \left(\frac{I_{\text{TR2}} - I_{\text{spin}}}{I_{\text{TR1}}} \right) \omega_{\text{3avg}} \left[C \cos(\omega_{\text{nut}} t) + D \sin(\omega_{\text{nut}} t) \right]$$

Equating coeff of $\sin(\omega_{\text{nut}} t)$ & $\cos(\omega_{\text{nut}} t)$, as this relation need to hold for all t , we get

$$B \omega_{\text{nut}} = C \left(\frac{I_{\text{TR2}} - I_{\text{spin}}}{I_{\text{TR1}}} \right) \omega_{\text{3avg}}$$

$$-A \omega_{\text{nut}} = D \left(\frac{I_{\text{TR2}} - I_{\text{spin}}}{I_{\text{TR1}}} \right) \omega_{\text{3avg}}$$

$$B \sqrt{\frac{(I_{sp} - I_{TR1})(I_{sp} - I_{TR2})}{I_{TR1} I_{TR2}}} = C \text{Sign}(I_{TR2} - I_{spin}) \sqrt{\frac{(I_{TR2} - I_{spin})^2}{I_{TR1}^2}}$$

$$-A \sqrt{\frac{(I_{sp} - I_{TR1})(I_{sp} - I_{TR2})}{I_{TR1} I_{TR2}}} = D \text{Sign}(I_{TR2} - I_{spin}) \sqrt{\frac{(I_{TR2} - I_{spin})^2}{I_{TR1}^2}}$$

$$B = C \text{Sign}(I_{TR2} - I_{spin}) \sqrt{\frac{I_{TR2}(I_{spin} - I_{TR2})}{I_{TR1}(I_{spin} - I_{TR1})}}$$

$$D = -A \text{Sign}(I_{TR2} - I_{spin}) \sqrt{\frac{I_{TR1}(I_{spin} - I_{TR1})}{I_{TR2}(I_{spin} - I_{TR2})}}$$

We realize that when $t=0$

$$\omega_1(0) = A \cos(0) = A \text{ \& \> } \omega_2(0) = C \cos(0) = C$$

$$\therefore A = \omega_1(0) \text{ \& \> } C = \omega_2(0)$$

So the final dynamics' solution look like

$$\omega_1(t) = \omega_1(0) \cos(\omega_{und} t) + \omega_2(0) \text{Sign}(I_{TR2} - I_{spin}) \sqrt{\frac{I_{TR2}(I_{spin} - I_{TR2})}{I_{TR1}(I_{spin} - I_{TR1})}} \sin(\omega_{und} t)$$

$$\omega_2(t) = \omega_2(0) \cos(\omega_{und} t) - \omega_1(0) \text{Sign}(I_{TR2} - I_{spin}) \sqrt{\frac{I_{TR1}(I_{spin} - I_{TR1})}{I_{TR2}(I_{spin} - I_{TR2})}} \sin(\omega_{und} t)$$

Problem 5

```
%script_simgravgradsc12.m
%
% Copyright (c) 2019 Mark L. Psiaki. All rights reserved.
%
% This Matlab script simulates the torque-free motion of
% a non-axi-symmetric spinning satellite.
%
% This script makes a plot of the
% angular momentum time history in body-fixed
% coordinates and in inertial
% coordinates. It also makes plots
% of the time histories of the 3 body-axis spin-
% rate vector elements.
%
% Clear the Matlab workspace.
%
clear;clc;close all;
%
% Set up the simulation parameters. Load the body-axes moment-of-
% inertia matrix and the initial body-axes angular velocity
% from simgravgradsc12_data.mat.
%
load simgravgradsc12_data
%
% Set up the orbital angular rate.
%
norbit = 0; % eliminate gravity-gradient torque and
            % rotation of the reference frame relative to which
            % x(1:4,1) defines the attitude quaternion so that
            % it becomes an inertial reference frame rather
            % than a non-inertial orbit-following local-level
            % reference frame.
q0 = [0;0;0;1];
x0 = [q0;omegabody0];
%
% Define the aircraft dynamics function handle
% in a form that is suitable for input to ode45.m.
%
ffunctode45 = @(tdum,xlum) ...
    ffunctgravgradsc02(tdum,xlum,IMoIbody,norbit);
%
% Define the time span of the simulation, computing outputs
% every 0.5. This time span should be large enough
% to see several spins periods and several nutation periods.
%
tspan = ((0:900)')*0.5;
%
% Set up numerical integration options for ode45.m
% in a way that uses a tighter relative tolerance than
% is normally used.
```

```

%
optionsode45 = odeset('RelTol',1.e-10);
%
% Call ode45.m in order to perform numerical integration.
%
tic
[thist,xhist] = ode45(ffunctode45,tspan,x0,optionsode45);
timetosim = toc;
%
% Compute the angular momentum vector time history in
% inertial coordinates.
%
tic
N = size(thist,1);
hvecbodyhist = zeros(N,3);
hvecinertialhist = zeros(N,3);
for k = 1:N
    xk = xhist(k,:)';
    hvecbodyk = IMoIbody*xk(5:7,1);
    hvecbodyhist(k,:) = hvecbodyk';
    qk = xk(1:4,1);
    qknorm = qk*(1/sqrt(sum(qk.^2)));
    Rk = rotmatquaternion(qknorm);
    hvecinertialk = (Rk')*hvecbodyk;
    hvecinertialhist(k,:) = hvecinertialk';
end
timetohvecinertial = toc
clear k xk hvecbodyk qk qknorm Rk hvecinertialk
%
% Transform to principal axes and assume that the
% principal axis whose moment of inertia is the most
% different from the other two is the spin-axis
% inertia. This is a nearly axially-symmetric
% spacecraft whose principal axes do not
% exactly align with the body axes in which
% the simulation has been conducted.
%
% This eigenvalue decomposition computes the 3-by-3
% matrix Roldnew and IMoIbodynew such that
% Roldnew*IMoIbodynew*inv(Roldnew) = IMoIbody
% with IMoIbodynew being a diagonal matrix. The symmetry
% of IMoIbody should ensure that inv(Roldnew) = Roldnew'
% so that Roldnew*IMoIbodynew*(Roldnew') = IMoIbody.
% Symmetry should also ensure that IMoIbody is truly
% diagonalizable (Some matrices can only be put into a
% form known as Jordan form that is not completely
% diagonal if there are repeated eigenvalues.)
%
[Roldnew,IMoIbodynew] = eig(IMoIbody);
%
% Check that Roldnew is orthonormal.
%
errdum = norm(Roldnew*(Roldnew') - eye(3));
if errdum > 1.e-12

```

```

        disp('Warning in script_simgravgradsc12.m: IMoIbody')
        disp(' does not appear to have orthonormal eigenvectors.')
        disp(' maybe it is not exactly diagonal.')
        disp(' ')
    end
    clear errdum
%
% Extract the eigenvalues and arrange them in ascending order.
%
    Iprsvec = diag(IMoIbodynew);
    [Iprsvec,idumsortvec] = sort(Iprsvec);
    Roldnew = Roldnew(:,idumsortvec);
    if det(Roldnew) < 0
        Roldnew(:,3) = - Roldnew(:,3);
    end
%
% Check that the physical constraint on the maximum
% eigenvalue is respected.
%
    if Iprsvec(3,1) > (Iprsvec(1,1) + Iprsvec(2,1))
        disp('Warning in script_simgravgradsc12.m: IMoIbody's')
        disp(' largest eigenvalue is more than the sum of it's')
        disp(' other two eigenvalues.')
        disp(' ')
    end
    if Iprsvec(1,1) <= 0
        disp('Warning in script_simgravgradsc12.m: IMoIbody')
        disp(' has one or more non-positive eigenvalues.')
        disp(' ')
    end
%
% Transform the body-axes spin rate vector time history
% into the principal axis coordinate system.
% Note that omegabodynewk = (Roldnew')*omegabodyk,
% as should be the case, where omegabodyk = xhist(k,5:7)'
% and omegabodynewk = omegabodynewhist(k,:)'
%
    omegabodynewhist = xhist(:,5:7)*Roldnew;
%
% Determine whether the spacecraft is spinning
% primarily about its minor axis or about its
% major axis. Make the third axis be the
% spin axis in either case.
%
    if mean(abs(omegabodynewhist(:,1))) > ...
        mean(abs(omegabodynewhist(:,3)))
        idumsortvec = [2;3;1];
        Iprsvec = Iprsvec(idumsortvec,1);
        Roldnew = Roldnew(:,idumsortvec);
        omegabodynewhist = omegabodynewhist(:,idumsortvec);
        clear idumsortvec
    end
%
% Make change the sign definitions on the first

```

```

% and third axes, if necessary, in order
% to ensure that the spin is about the positive
% third axis.
%
if mean(omegabodynewhist(:,3)) < 0
    Roldnew(:,1) = - Roldnew(:,1);
    Roldnew(:,3) = - Roldnew(:,3);
    omegabodynewhist(:,1) = - omegabodynewhist(:,1);
    omegabodynewhist(:,3) = - omegabodynewhist(:,3);
end
%
% Check that the new diagonal moment-of-inertia matrix
% and the rotation matrix agree with the original
% moment-of-inertia matrix.
%
IMoIbodynew = diag(Iprsvec);
IMoIbody_re = Roldnew*IMoIbodynew*(Roldnew');
errdum = norm(IMoIbody_re - IMoIbody)/norm(IMoIbody);
if errdum > 1.e-12
    disp('Warning in script_simgravgradsc12.m: There is')
    disp(' an inaccuracy in the principal axes model')
    disp(' ')
end
clear errdum
%
% Compute an approximate nutation frequency based
% on an approximate model that assumes axial symmetry
% even though this assumption is not quite right.
% The calculation of the mean omegaspin should
% average over an integer number of nutation periods.
% The length of the simulation has been chosen to
% make this likely.
%
omegaspinavg = mean(omegabodynewhist(:,3));
Itr1 = Iprsvec(1,1);
Itr2 = Iprsvec(2,1);
Ispin = Iprsvec(3,1);
omeganut = omegaspinavg*sqrt( (Ispin - Itr1)*(Ispin - Itr2)/Itr2/
Itr1);
%
% Compute the theoretical body-axis spin vector component
% time histories that are valid for this axially-symmetric
% spacecraft. This analysis assumes that
%
omegabodynewapprox1(t) = A*cos(omeganut*t) + B*sin(omeganut*t)
%
omegabodynewapprox2(t) = C*cos(omeganut*t) + D*sin(omeganut*t)
%
% The constant DoverA = D/A and the constant BoverC = B/C
% prove helpful in writing these approximation theoretical
% solutions in terms of the initial values omegabodynewhist(1,1)
% and omegabodynewhist(1,2).
%

```

```

    DoverA = -sign(Itr2-Ispin)*sqrt(Itr1*(Ispin - Itr1)/Itr2/(Ispin -
Itr2));
    BoverC = - (1/DoverA);
    omegabodynewapprox1hist = omegabodynewhist(1,1)*cos(omeganut*thist)
+ omegabodynewhist(1,2)*BoverC*sin(omeganut*thist);
    omegabodynewapprox2hist = omegabodynewhist(1,2)*cos(omeganut*thist)
+ omegabodynewhist(1,1)*DoverA*sin(omeganut*thist);
    omegabodynewapprox3hist = omegaspinavg*ones(N,1);
    clear omegatRmaghist omegatRmagmean omegatrmagratio0 ...
           omeganewbody10approx omeganewbody20approx
%
% Plot the body-axes angular momentum time history.
%
figure(1)
hold off
plot(thist,hvecbodyhist,'LineWidth',2)
set(get(gcf,'CurrentAxes'),'FontSize',10)
grid
xlabel('Time (sec)')
ylabel('Angular Momentum (N-m-sec)')
title(['Original Body-Axes Angular Momentum','...
' script\_simgravgradsc12.m'])
legend('h1 body','h2 body','h3 body')
%
% Plot the inertial angular momentum time history.
%
figure(2)
hold off
plot(thist,hvecinertialhist,'LineWidth',2)
set(get(gcf,'CurrentAxes'),'FontSize',10)
grid
xlabel('Time (sec)')
ylabel('Angular Momentum (N-m-sec)')
title(['Inertial Angular Momentum','...
' script\_simgravgradsc12.m'])
legend('h1 ECIF','h2 ECIF','h3 ECIF')
%
% Plot the body-axis spin vector component time histories,
% both from the numerical integration and the theoretical
% values.
%
figure(3)
hold off
plot(thist,omegabodynewhist,'-','LineWidth',2)
set(get(gcf,'CurrentAxes'),'FontSize',8)
hold on
plot(thist,[omegabodynewapprox1hist,omegabodynewapprox2hist,...
omegabodynewapprox3hist],'.','MarkerSize',10)
hold off
grid
xlabel('Time (sec)')
ylabel('Angular Velocity (radians/sec)')
title(['Principal Axes Angular Velocity','...
' script\_simgravgradsc12.m'])

```

```

    legend('omegabodynew1 sim','omegabodynew2 sim','omegabodynew3
sim',...
          'omegabody1 approx. theory','omegabody2 approx. theory',...
          'omegabody3 approx. theory')
%
% Plot the difference between the nonlinear simulation
% of the velocity vector components and the theoretical
% models in order to focus in on the errors.
%
figure(4)
hold off
omegaapproxerrhist = ...
    [omegabodynewapprox1hist,omegabodynewapprox2hist,...
    omegabodynewapprox3hist] - omegabodynewhist;
plot(thist,omegaapproxerrhist,'-','LineWidth',2)
set(get(gcf,'CurrentAxes'),'FontSize',10)
grid
xlabel('Time (sec)')
ylabel('Angular Velocity Approx. Errors (radians/sec)')
title(['Principal-Axes Ang. Vel. Errors','...
' script\_simgravgradsc12.m'])
legend('omegabody1 approx.-true','omegabody2 approx.-true',...
'omegabody3 approx.-true')
%
% Save the results.
%
textcommands = ['These data have been generated by the',...
' commands in script_simgravgradsc12.m'];
save simgravgradsc12
format long
xfinal = xhist(end,:)

disp('The following are the differences compared to
script_simgravgradsc11.m:')
disp('1) The max amplitudes for angular velocities and angular
momentums are different between w1 and w2. This is because of the
difference I_tr1 and I_tr2');
disp('2) The w3 is not constant in the ode45, true to the dynamics.
But the theoretical w3 is constant because of the assumption of small
w1 and w2');

timetohvecinertial =

    0.0104469000000000

xfinal =

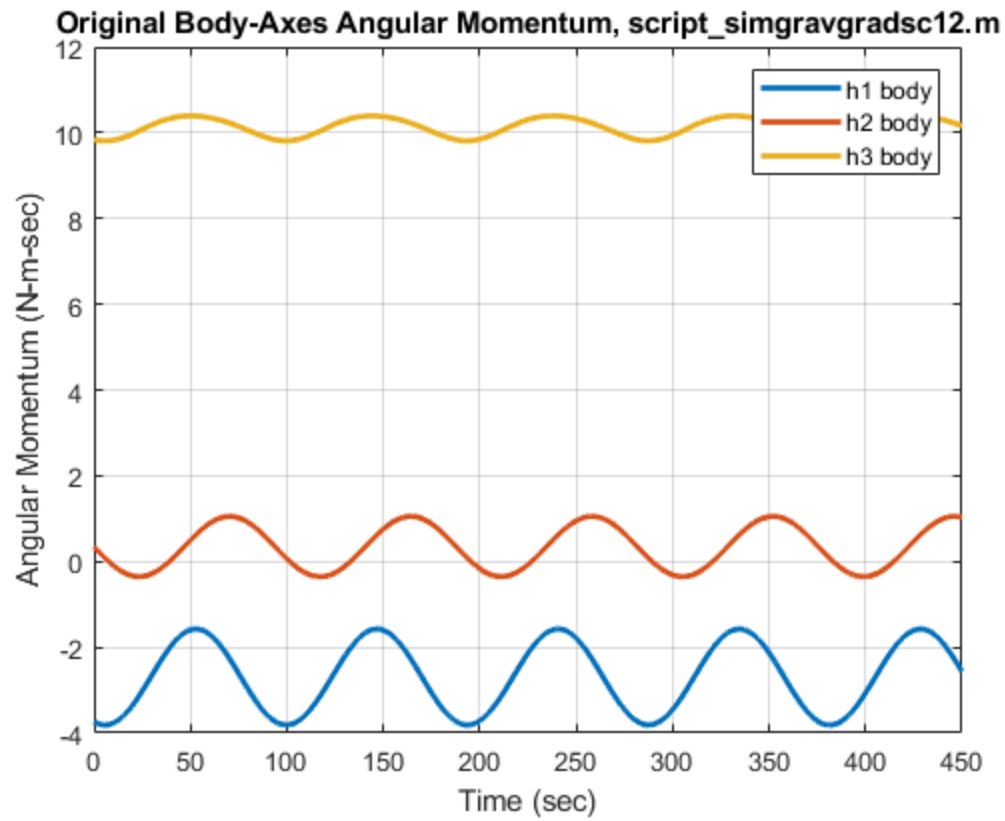
    0.300577539317991
   -0.074095730362301
   -0.935199318657734
    0.171928828717826
   -0.026762379581941

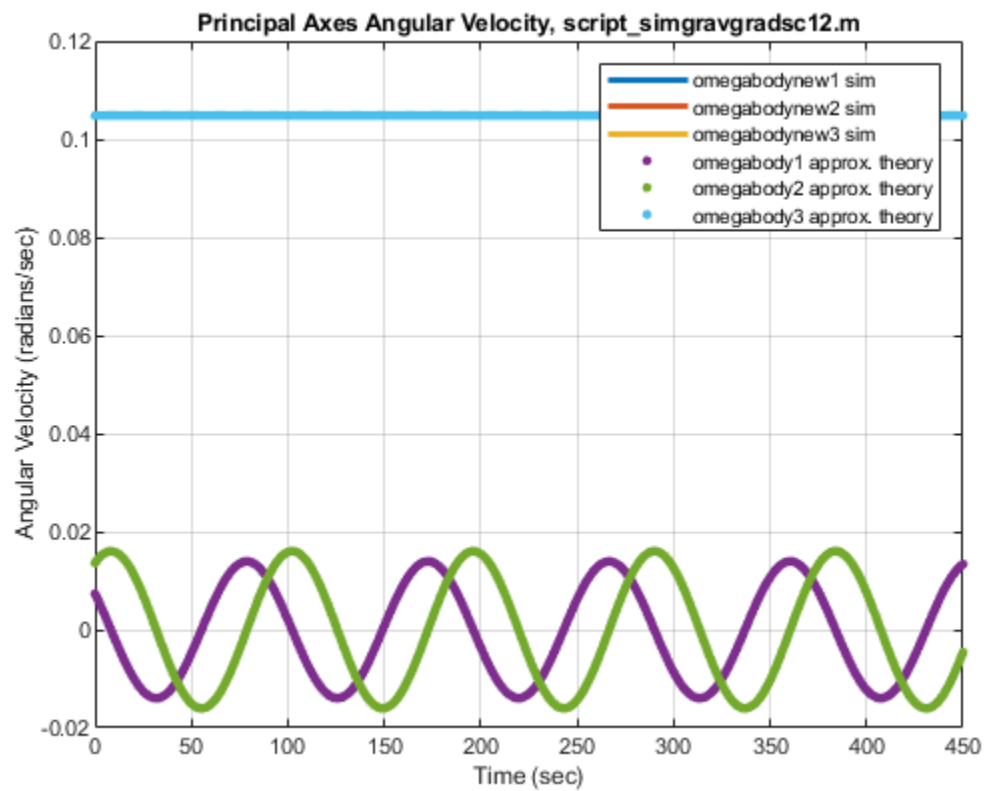
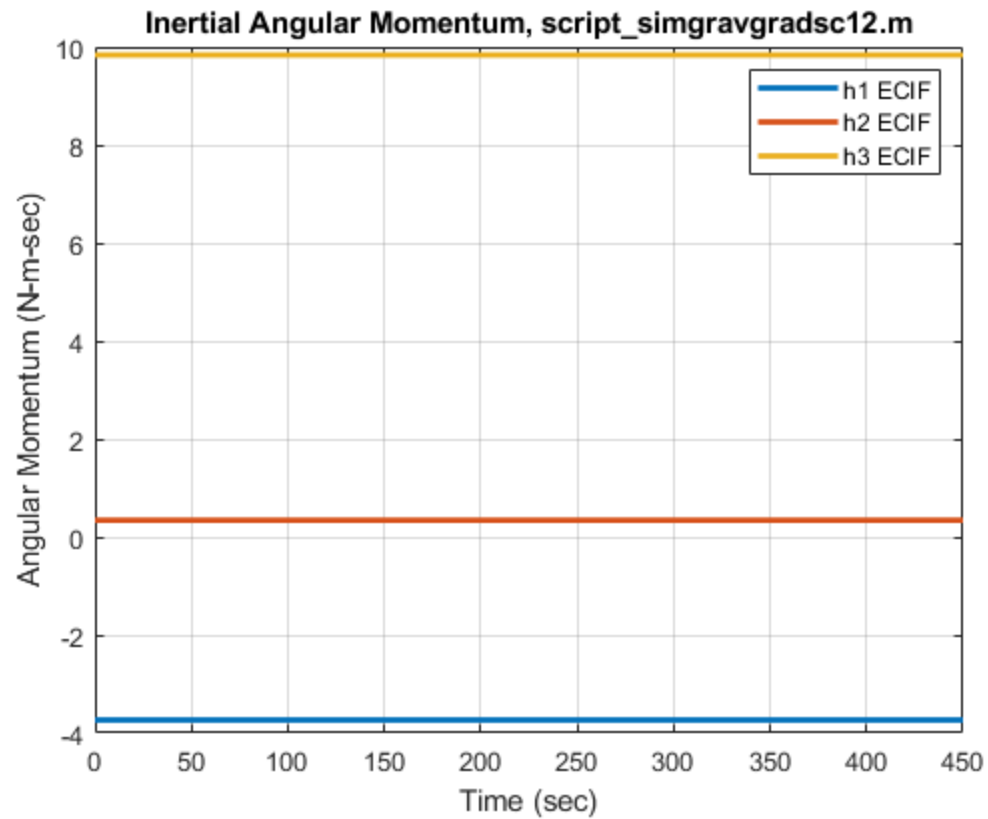
```

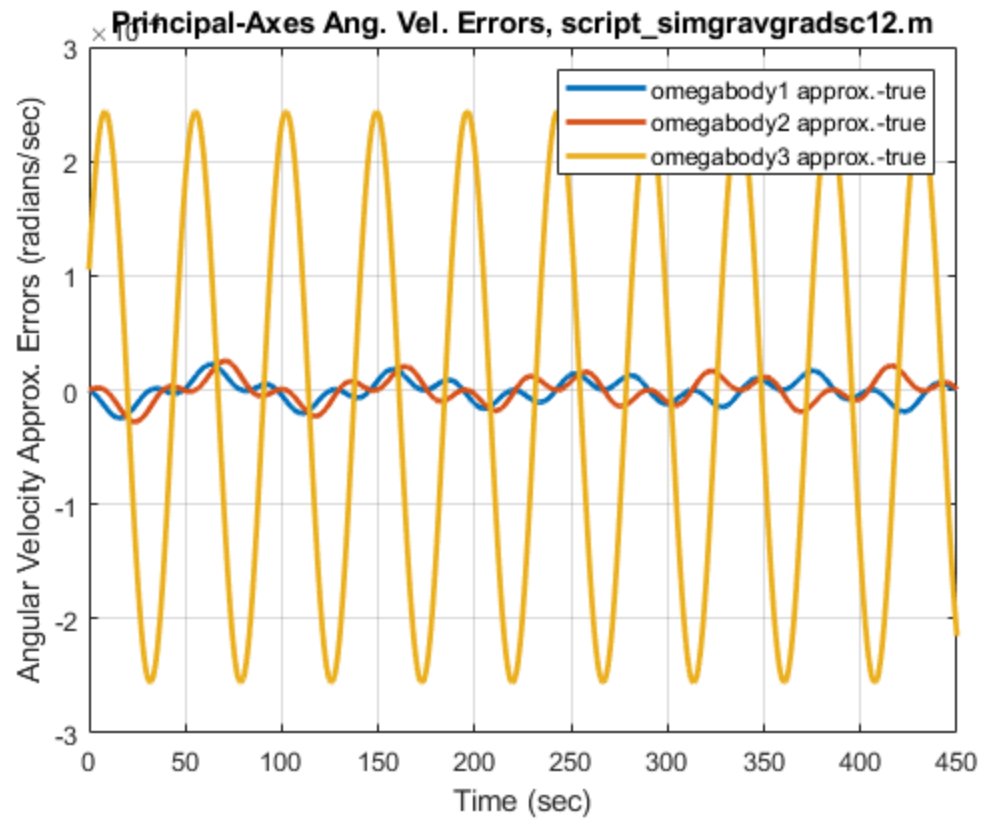
0.017598485086061
0.101028409577155

The following are the differences compared to
script_simgravgradsc11.m:

- 1) The max amplitudes for angular velocities and angular momentums are different between w_1 and w_2 . This is because of the difference I_{tr1} and I_{tr2}
- 2) The w_3 is not constant in the ode45, true to the dynamics. But the theoretical w_3 is constant because of the assumption of small w_1 and w_2







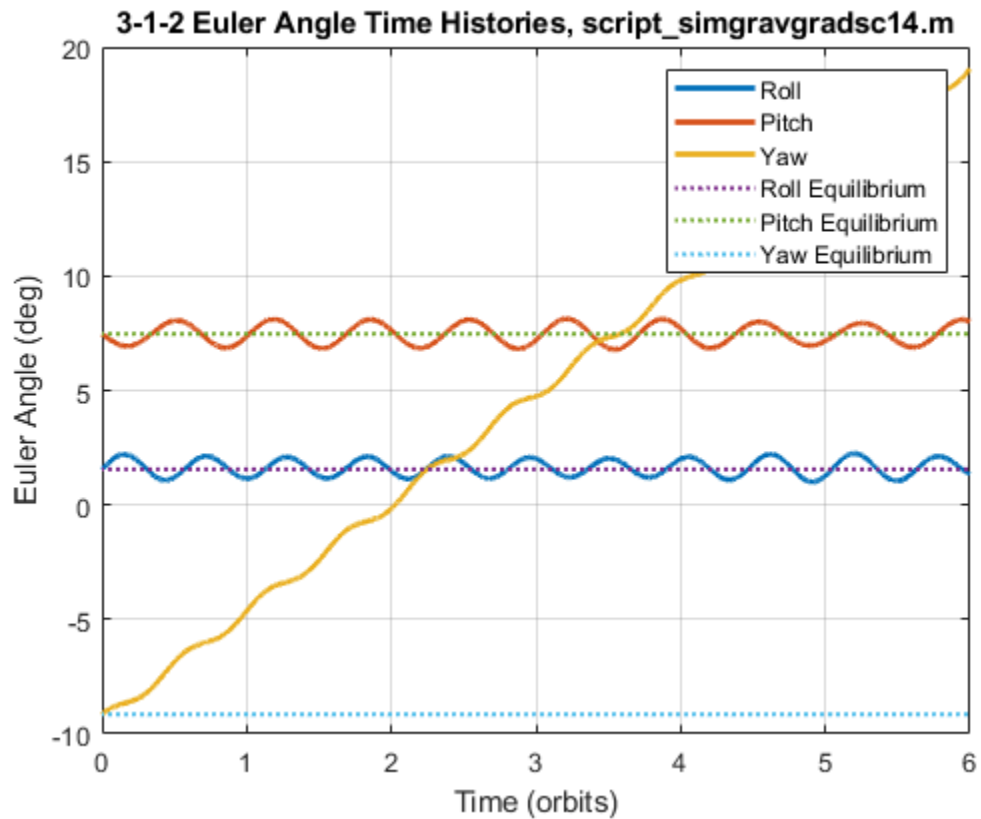
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Problem 6

$x_{final} =$

```
-0.000024674605850  
0.070758815234848  
0.165798947726374  
0.983617801330580  
-0.000353720202840  
-0.000998579313148  
-0.000003710776134
```

The system is unstable, evident from diverging yaw angle

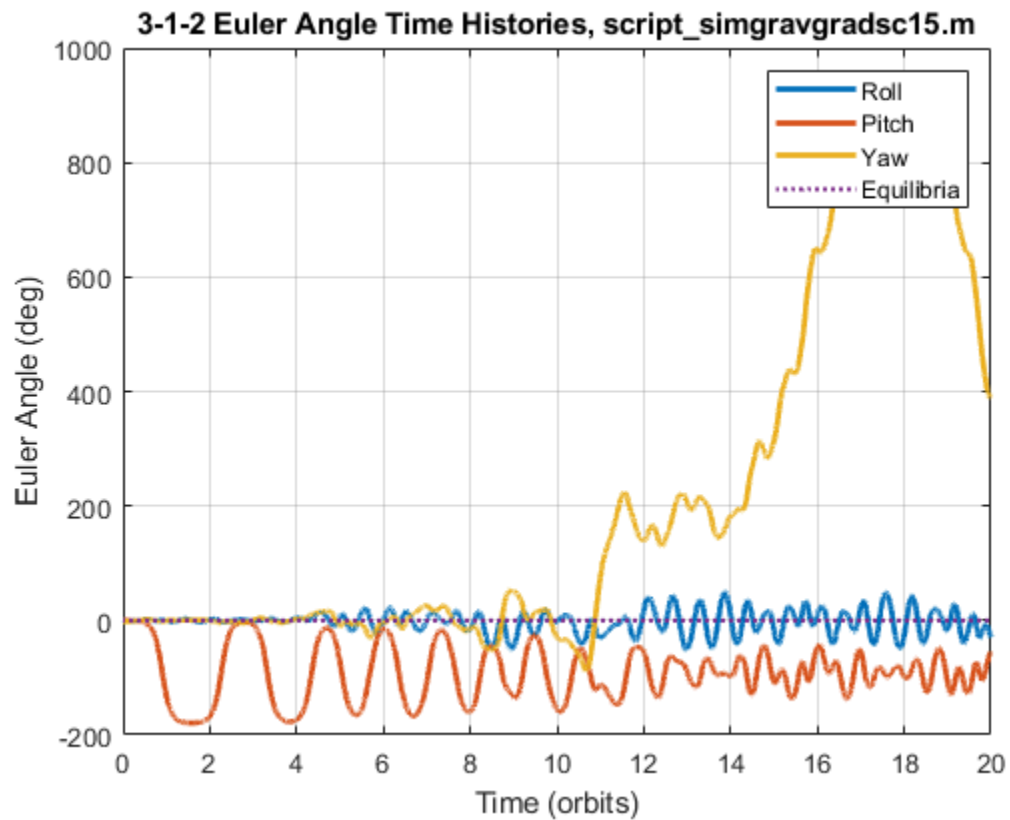


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Problem 7

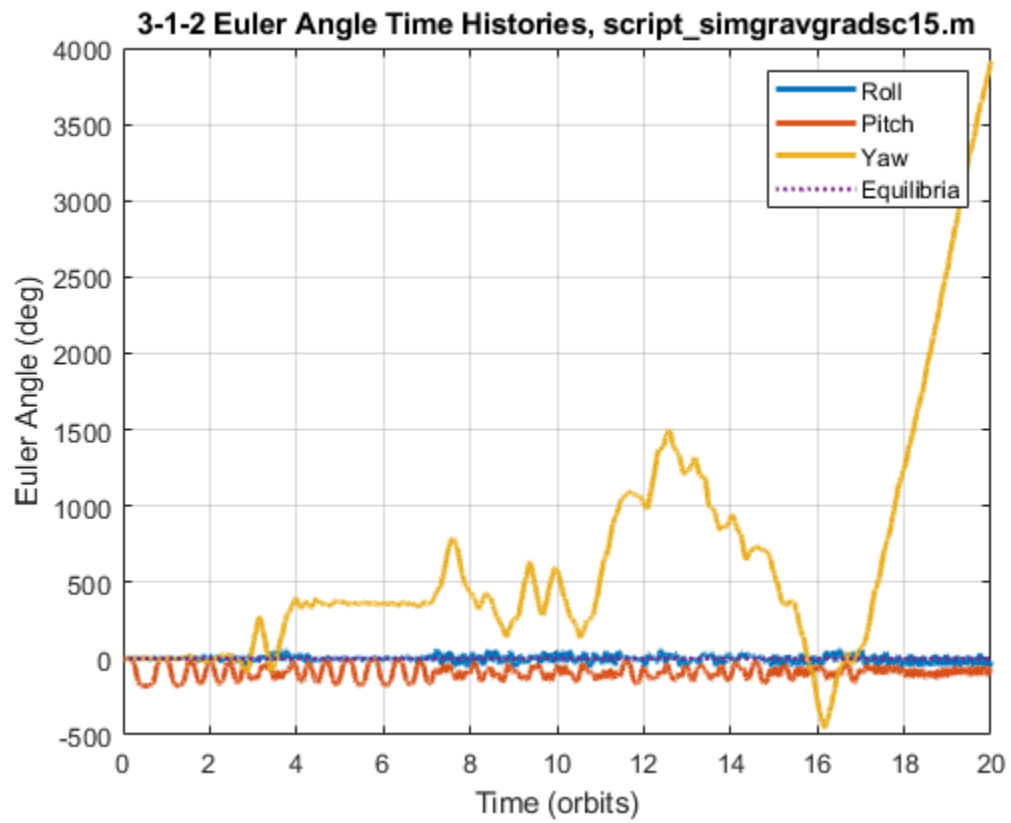
```
xfinal =  
  
    0.117496010185310  
    0.480952819413493  
   -0.323683410192931  
   -0.806297823411958  
   -0.001205763144471  
   -0.000264064182243  
    0.000547611740791
```

The system is unstable, evident from diverging yaw angle. Though the time history shows behavior of marginally stable system, it reveals an chaotic behavior is larger time scales



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Problem 7 (longer simulation)



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