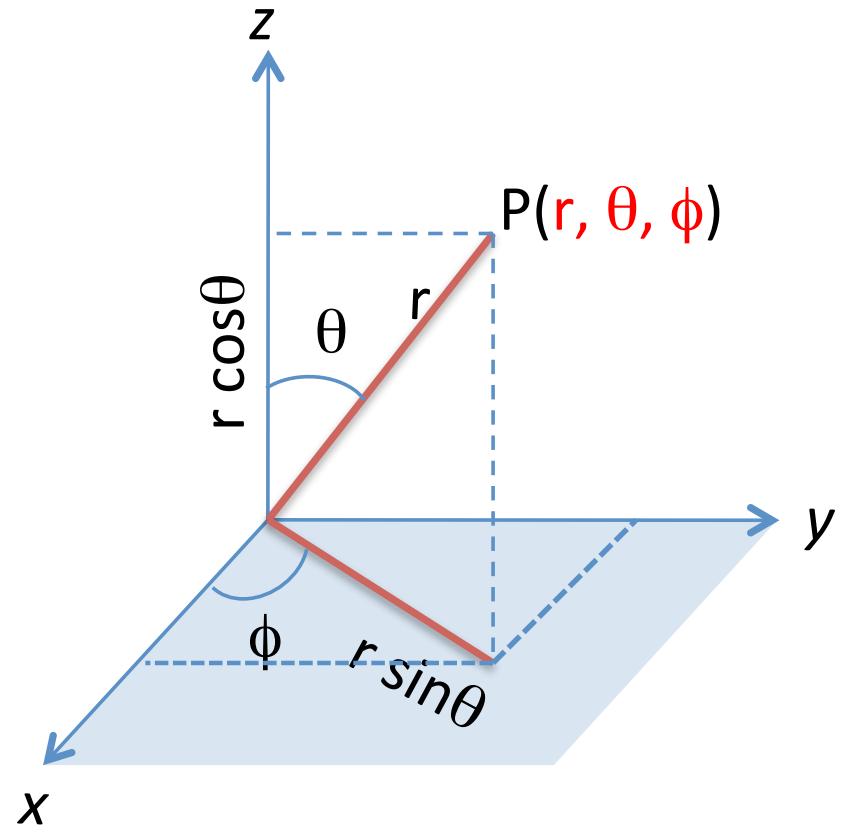
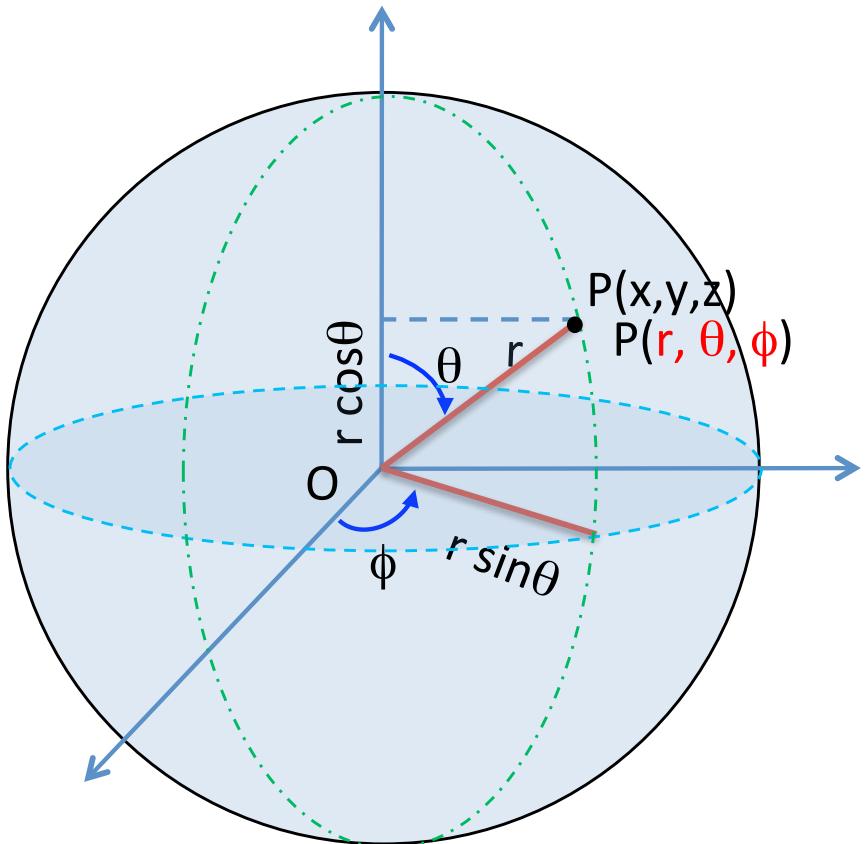


# Spherical Polar Coordinates



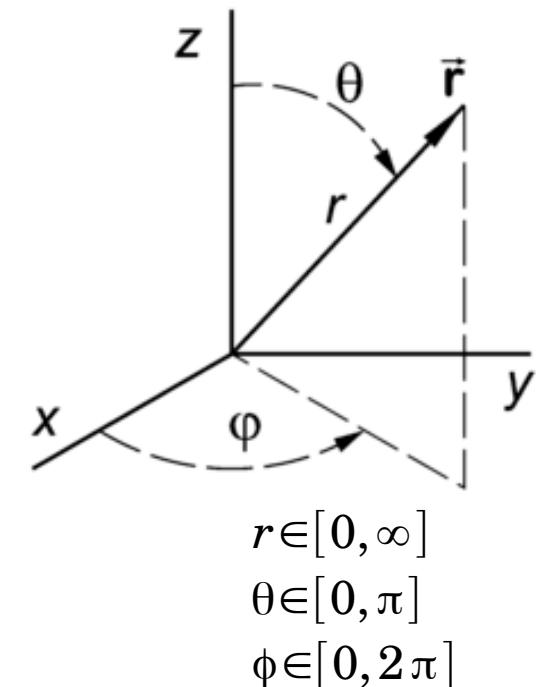
# SPHERICAL POLAR COORDINATES

$$x = r \sin \theta \cos \phi$$

$$y = r \sin \theta \sin \phi$$

$$z = r \cos \theta$$

$$\begin{pmatrix} \delta x \\ \delta y \\ \delta z \end{pmatrix} = \begin{pmatrix} \sin \theta \cos \phi & r \cos \theta \cos \phi & -r \sin \theta \sin \phi \\ \sin \theta \sin \phi & r \cos \theta \sin \phi & r \sin \theta \cos \phi \\ \cos \theta & -r \sin \theta & 0 \end{pmatrix} \begin{pmatrix} \delta r \\ \delta \theta \\ \delta \phi \end{pmatrix}$$



$$\begin{pmatrix} \hat{r} \\ \hat{\theta} \\ \hat{\phi} \end{pmatrix} = \begin{pmatrix} \sin \theta \cos \phi & \sin \theta \sin \phi & \cos \theta \\ \cos \theta \cos \phi & \cos \theta \sin \phi & -\sin \theta \\ -\sin \phi & \cos \phi & 0 \end{pmatrix} \begin{pmatrix} \hat{x} \\ \hat{y} \\ \hat{z} \end{pmatrix}$$

$$\hat{r} \cdot \hat{\theta} = \hat{r} \cdot \hat{\phi} = \hat{\theta} \cdot \hat{\phi} = 0$$

# SPHERICAL POLAR COORDINATES

$$dl_r = \delta r \quad dl_\theta = r \delta \theta \quad dl_\phi = r \sin \theta \delta \phi$$

Infinitesimal displacement  $d\vec{l} = \delta r \hat{r} + r \delta \theta \hat{\theta} + r \sin \theta \delta \phi \hat{\phi}$

Arc length  $ds^2 = d\vec{l} \cdot d\vec{l} = \delta r^2 + r^2 \delta \theta^2 + r^2 \sin^2 \theta \delta \phi^2$

$$dV = dl_r dl_\theta dl_\phi = r^2 \sin \theta \delta r \delta \theta \delta \phi$$

Area element:

If  $r = \text{constant}$  (surface of a sphere)  $\delta r = 0$

$$dA = dl_\theta dl_\phi = r^2 \sin \theta d\theta d\phi$$

If  $\theta = \text{constant}$ ,  $\delta \theta = 0$

$$dA = dl_r dl_\phi = r \sin \theta dr d\phi$$

If  $\phi = \text{constant}$ ,  $\delta \phi = 0$

$$dA = dl_r dl_\theta = r dr d\theta$$

# SPHERICAL POLAR COORDINATES

Gradient of a scalar field

If we have a function  $T(r, \theta, \phi)$  then we want

$$\begin{aligned} dT &= \frac{\partial T}{\partial r} \delta r + \frac{\partial T}{\partial \theta} \delta \theta + \frac{\partial T}{\partial \phi} \delta \phi \\ &= \nabla T \cdot \delta \vec{l} \end{aligned}$$

$$\text{Now, } \delta \vec{l} = \hat{r} \delta r + \hat{\theta} r \delta \theta + \hat{\phi} r \sin \theta \delta \phi$$

The gradient in the spherical polar system can be identified as

$$\nabla T = \hat{r} \frac{\partial T}{\partial r} + \hat{\theta} \frac{1}{r} \frac{\partial T}{\partial \theta} + \hat{\phi} \frac{1}{r \sin \theta} \frac{\partial T}{\partial \phi}$$

# SPHERICAL POLAR COORDINATES

Velocity and acceleration

$$\begin{pmatrix} \hat{r} \\ \hat{\theta} \\ \hat{\phi} \end{pmatrix} = \mathbf{M} \begin{pmatrix} \hat{x} \\ \hat{y} \\ \hat{z} \end{pmatrix}$$

$$\mathbf{M} = \begin{pmatrix} \sin \theta \cos \phi & \sin \theta \sin \phi & \cos \theta \\ \cos \theta \cos \phi & \cos \theta \sin \phi & -\sin \theta \\ -\sin \phi & \cos \phi & 0 \end{pmatrix}$$

$$\begin{pmatrix} \hat{x} \\ \hat{y} \\ \hat{z} \end{pmatrix} = \mathbf{M}^T \begin{pmatrix} \hat{r} \\ \hat{\theta} \\ \hat{\phi} \end{pmatrix}$$

$$\begin{pmatrix} \dot{\hat{r}} \\ \dot{\hat{\theta}} \\ \dot{\hat{\phi}} \end{pmatrix} = \dot{\mathbf{M}} \begin{pmatrix} \hat{x} \\ \hat{y} \\ \hat{z} \end{pmatrix}$$

$$\mathbf{M}^T = \begin{pmatrix} \sin \theta \cos \phi & \cos \theta \cos \phi & -\sin \phi \\ \sin \theta \sin \phi & \cos \theta \sin \phi & \cos \phi \\ \cos \theta & -\sin \theta & 0 \end{pmatrix}$$

$$\dot{\mathbf{M}} = \begin{pmatrix} \cos \theta \cos \phi \dot{\theta} - \sin \theta \sin \phi \dot{\phi} & \cos \theta \sin \phi \dot{\theta} + \sin \theta \cos \phi \dot{\phi} & -\sin \theta \dot{\phi} \\ -\sin \theta \cos \phi \dot{\theta} - \cos \theta \sin \phi \dot{\phi} & -\sin \theta \sin \phi \dot{\theta} + \cos \theta \cos \phi \dot{\phi} & -\cos \theta \dot{\phi} \\ -\cos \phi \dot{\theta} & -\sin \phi \dot{\phi} & 0 \end{pmatrix}$$

# SPHERICAL POLAR COORDINATES

Velocity and acceleration

$$\begin{pmatrix} \dot{\hat{r}} \\ \dot{\hat{\theta}} \\ \dot{\hat{\phi}} \end{pmatrix} = \dot{\mathbf{M}} \mathbf{M}^T \begin{pmatrix} \hat{r} \\ \hat{\theta} \\ \hat{\phi} \end{pmatrix}$$
$$= \begin{pmatrix} 0 & \dot{\theta} & \sin \theta \dot{\phi} \\ -\dot{\theta} & 0 & \cos \theta \dot{\phi} \\ -\sin \theta \dot{\phi} & -\cos \theta \dot{\phi} & 0 \end{pmatrix} \begin{pmatrix} \hat{r} \\ \hat{\theta} \\ \hat{\phi} \end{pmatrix}$$

Why are the diagonal terms zero? Can you see the physical implication?

Notice that the matrix connecting the two vectors is anti-symmetric.

This was also the case in the plane polar co-ordinates. But we didn't mention it there.

The problem for velocity and acceleration components can now be completed...

# SPHERICAL POLAR COORDINATES

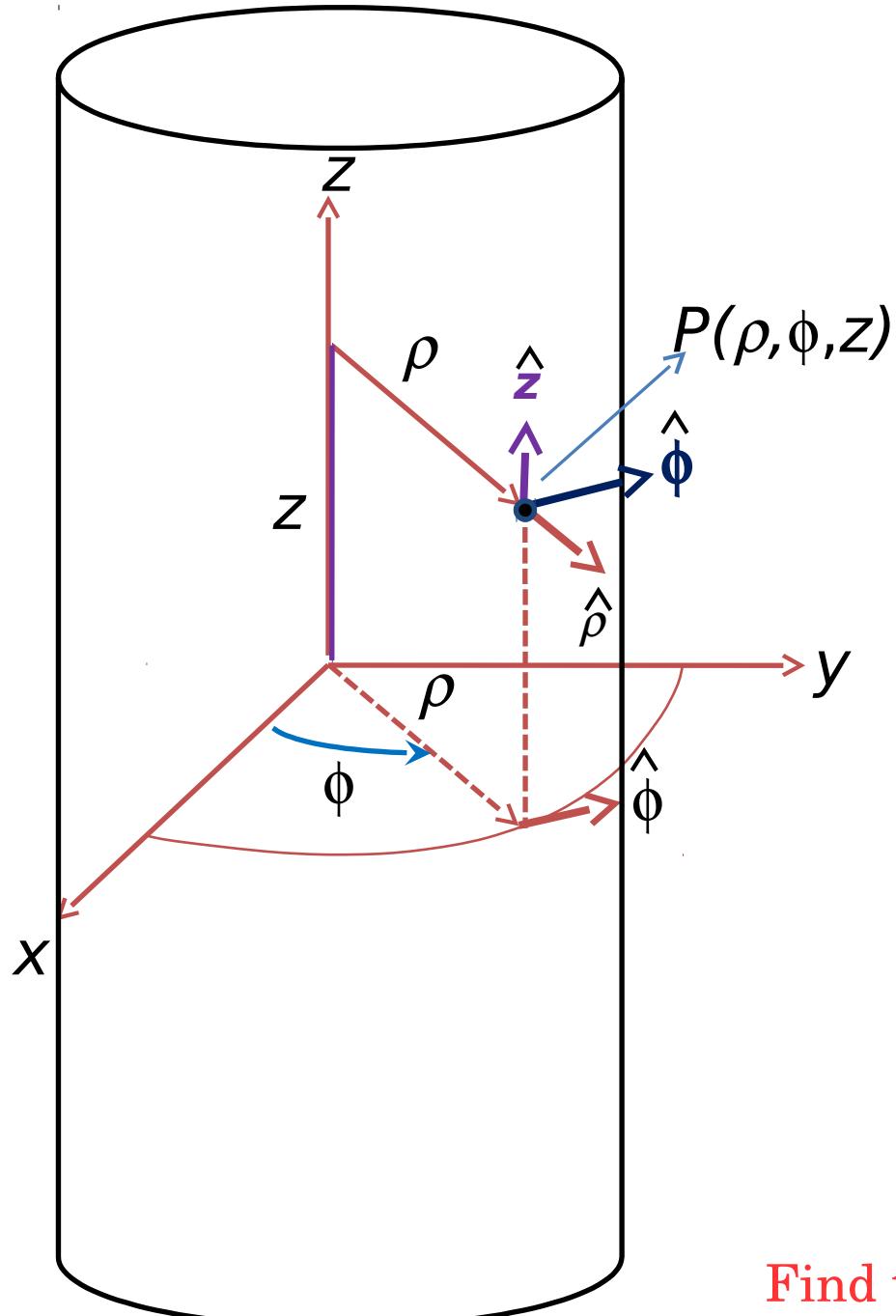
$$\vec{r} = r \hat{r}$$

$$\vec{v} = \dot{r} \hat{r} + r \dot{\hat{r}} = \dot{r} \hat{r} + r \dot{\theta} \hat{\theta} + r \sin \theta \dot{\phi} \hat{\phi}$$

$$\vec{a} = \hat{r} \left( \ddot{r} - r \dot{\theta}^2 - r \dot{\phi}^2 \sin^2 \theta \right) + \hat{\theta} \left( r \ddot{\theta} + 2 \dot{r} \dot{\theta} - r \dot{\phi}^2 \sin \theta \cos \theta \right) + \hat{\phi} \left( r \ddot{\phi} \sin \theta + 2 \dot{r} \dot{\phi} \sin \theta + 2 r \dot{\theta} \dot{\phi} \cos \theta \right)$$

We now have all ingredients to solve dynamical coordinates in this system.

# CYLINDRICAL POLAR COORDINATES



$$\begin{aligned}x &= \rho \cos \phi & 0 \leq \rho \leq \infty \\y &= \rho \sin \phi & 0 \leq \phi \leq 2\pi \\z &= z & 0 \leq z \leq \infty\end{aligned}$$

$$\begin{pmatrix} \hat{r} \\ \hat{\phi} \\ \hat{z} \end{pmatrix} = \begin{pmatrix} \cos \phi & \sin \phi & 0 \\ -\sin \phi & \cos \phi & 0 \\ 0 & 0 & 1 \end{pmatrix} \begin{pmatrix} \hat{x} \\ \hat{y} \\ \hat{z} \end{pmatrix}$$

$$\delta \vec{l} = \delta \rho \hat{r} + \rho \delta \phi \hat{\phi} + \delta z \hat{z}$$

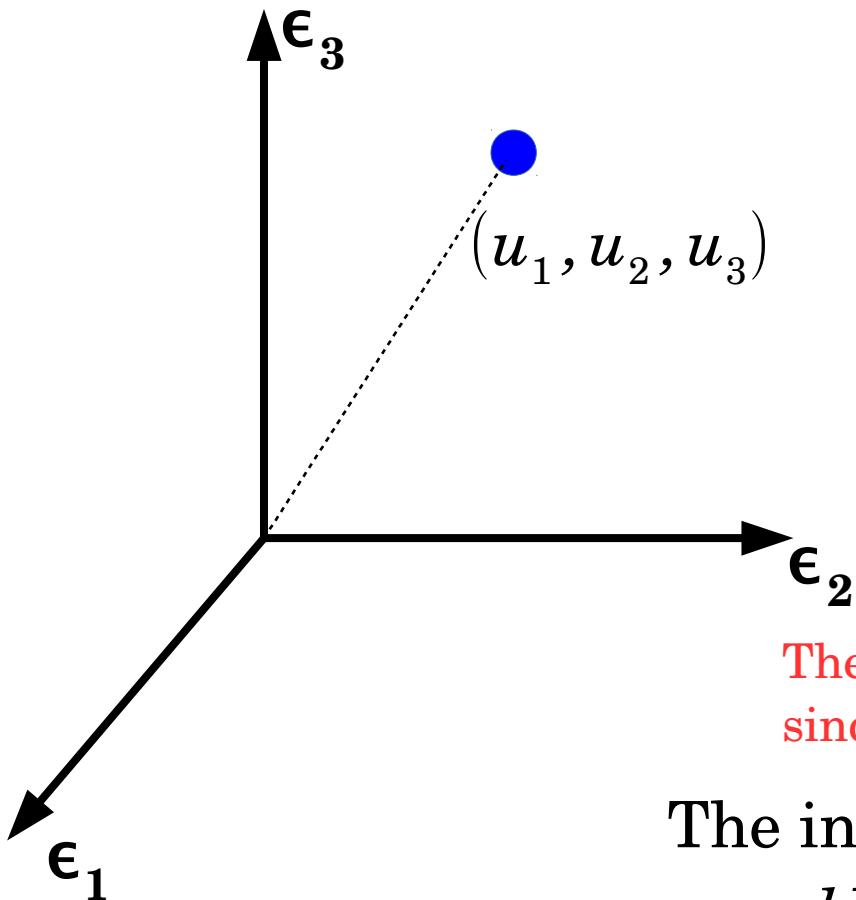
$$ds^2 = d\rho^2 + \rho^2 d\phi^2 + dz^2$$

$$dV = \rho d\rho d\phi dz$$

$$\nabla T = \hat{r} \frac{\partial T}{\partial \rho} + \hat{\phi} \frac{1}{\rho} \frac{\partial T}{\partial \phi} + \hat{z} \frac{\partial T}{\partial z}$$

Find the velocity and acceleration components

# GENERAL ORTHOGONAL CURVILINEAR COORDINATES



Identify a point in space by its three coordinates  $(u_1, u_2, u_3)$

If the system is orthogonal,

$$\epsilon_i \cdot \epsilon_j = 0 \quad \forall i \neq j$$

$$\epsilon_i \cdot \epsilon_j = \delta_{ij}$$

The unit vectors in general are functions of position, since their direction can vary from point to point.

The infinitesimal displacement vector is

$$d\mathbf{l} = \epsilon_1 h_1 du_1 + \epsilon_2 h_2 du_2 + \epsilon_3 h_3 du_3$$

$$h_i = \sqrt{\left(\frac{\partial x}{\partial u_i}\right)^2 + \left(\frac{\partial y}{\partial u_i}\right)^2 + \left(\frac{\partial z}{\partial u_i}\right)^2}$$

$$ds^2 = h_i^2 du_i^2$$

$$dV = \prod_i h_i du_i$$

$$= \sum_i \epsilon_i h_i du_i = \epsilon_i h_i du_i$$

**REPEATED INDEX IMPLIES SUMMATION**

# GENERAL ORTHOGONAL CURVILINEAR COORDINATES

System	$u_1$	$u_2$	$u_3$	$h_1$	$h_2$	$h_3$
Cartesian	$x$	$y$	$z$	1	1	1
Spherical	$r$	$\theta$	$\phi$	1	$r$	$r \sin \theta$
Cylindrical	$\rho$	$\phi$	$z$	1	$r$	1

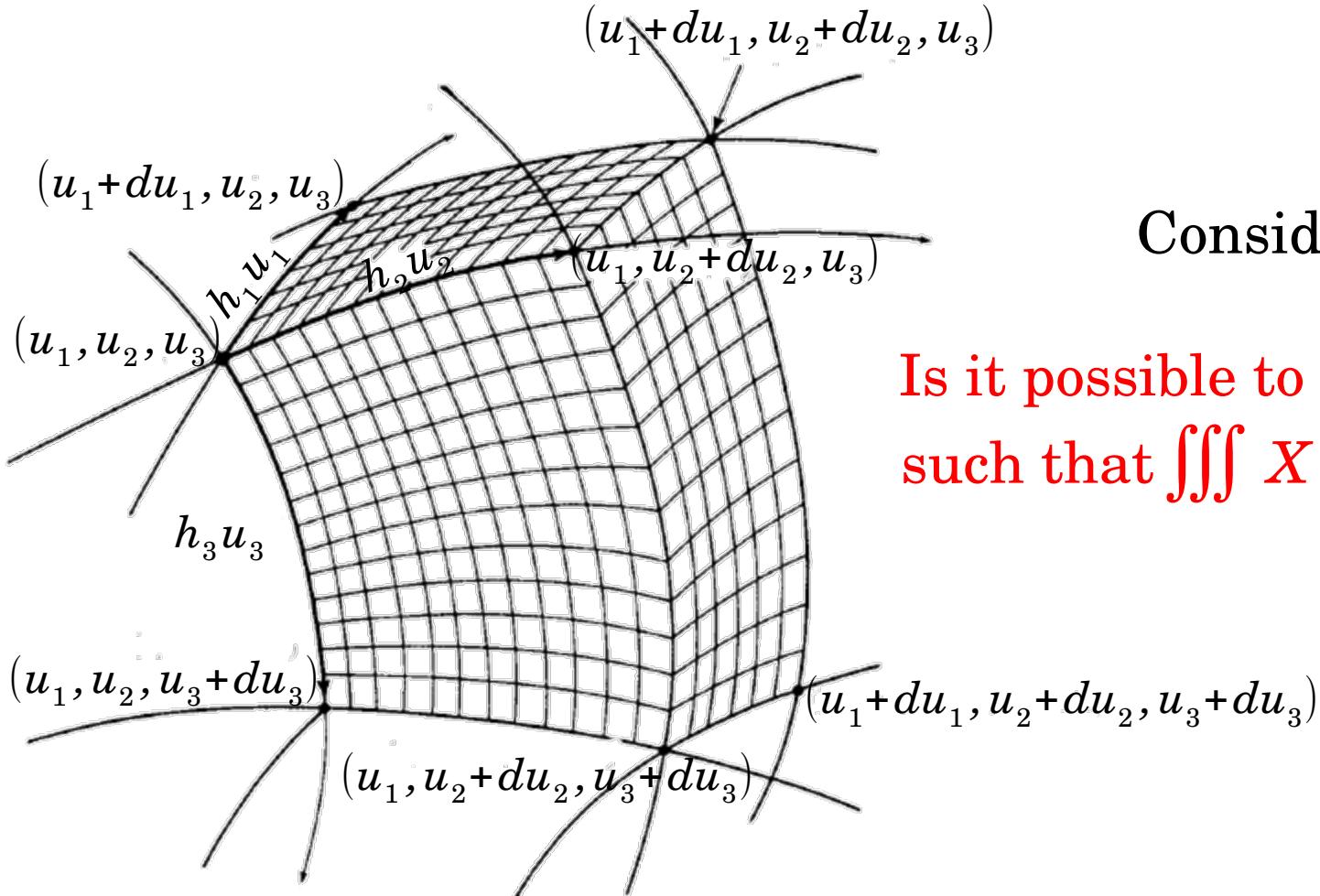
## Gradient of a scalar field

$$dT = \frac{\partial T}{\partial u_i} \delta u_i = \nabla T \cdot d\mathbf{l} = [(\nabla T)_i \epsilon_i] (h_i \delta u_i \epsilon_i)$$

$$\nabla T = \frac{1}{h_1} \frac{\partial T}{\partial u_1} \epsilon_1 + \frac{1}{h_2} \frac{\partial T}{\partial u_2} \epsilon_2 + \frac{1}{h_3} \frac{\partial T}{\partial u_3} \epsilon_3 = \frac{1}{h_i} \frac{\partial T}{\partial u_i} \epsilon_i$$

Does it match with the gradient for the three coordinate systems derived previously?

# DIVERGENCE IN CURVILINEAR COORDINATES



Consider a vector  $\vec{F}$ .

Is it possible to have a function  $X(\vec{F})$  such that  $\iiint X(\vec{F}) dV = \oint \vec{F} \cdot d\vec{S}$ ?

Consider the infinitesimal volume generated by starting at the point  $(u_1, u_2, u_3)$  and increasing each of the coordinates in succession by an infinitesimal amount.

# DIVERGENCE IN CURVILINEAR COORDINATES

Volume element  $dV = (h_1 h_2 h_3) du_1 du_2 du_3$

For the front surface  $d\mathbf{S} = -(h_2 h_3) du_2 du_3 \epsilon_1$

Flux through the front surface  $[\mathbf{F} \cdot d\mathbf{S}]_F = -(h_2 h_3 F_1) du_2 du_3$

Flux through the back surface  $[\mathbf{F} \cdot d\mathbf{S}]_B = (h_2 h_3 F_1) \Big|_{u_1+du_1} du_2 du_3$   
 $= (h_2 h_3 F_1) du_2 du_3 + \frac{\partial}{\partial u_1} (h_2 h_3 F_1) du_1 du_2 du_3$

Front and back together give a contribution

$$[\mathbf{F} \cdot d\mathbf{S}]_F + [\mathbf{F} \cdot d\mathbf{S}]_B = \frac{\partial}{\partial u_1} (h_2 h_3 F_1) du_1 du_2 du_3 = \frac{1}{h_1 h_2 h_3} \frac{\partial}{\partial u_1} (h_2 h_3 F_1) dV$$

Similarly, the left and right sides together give a contribution

$$[\mathbf{F} \cdot d\mathbf{S}]_L + [\mathbf{F} \cdot d\mathbf{S}]_R = \frac{1}{h_1 h_2 h_3} \frac{\partial}{\partial u_2} (h_1 h_3 F_2) dV$$

and  $[\mathbf{F} \cdot d\mathbf{S}]_T + [\mathbf{F} \cdot d\mathbf{S}]_B = \frac{1}{h_1 h_2 h_3} \frac{\partial}{\partial u_3} (h_1 h_2 F_3) dV$

# DIVERGENCE IN CURVILINEAR COORDINATES

Adding the flux from the individual sides, the total flux is

$$\mathbf{F} \cdot d\mathbf{S} = \frac{1}{h_1 h_2 h_3} \left[ \frac{\partial}{\partial u_1} (h_2 h_3 F_1) + \frac{\partial}{\partial u_2} (h_1 h_3 F_2) + \frac{\partial}{\partial u_3} (h_1 h_2 F_3) \right] dV$$

The coefficient of  $dV$  defines the  $\nabla \cdot \mathbf{F}$  in curvilinear coordinates

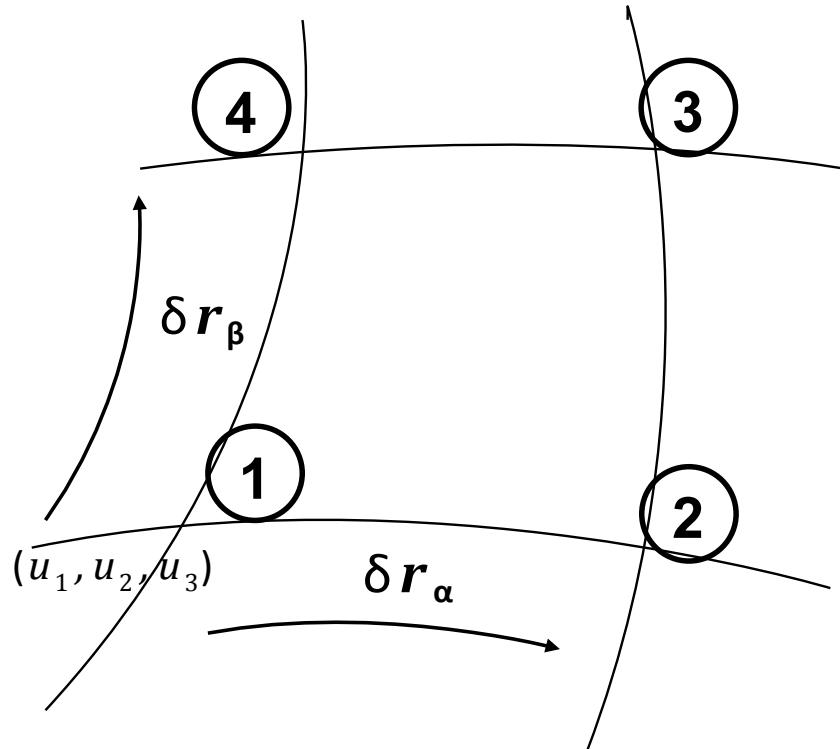
$$\nabla \cdot \mathbf{F} = \frac{1}{h_1 h_2 h_3} \left[ \frac{\partial}{\partial u_1} (h_2 h_3 F_1) + \frac{\partial}{\partial u_2} (h_1 h_3 F_2) + \frac{\partial}{\partial u_3} (h_1 h_2 F_3) \right]$$

Cartesian:  $\nabla \cdot \mathbf{F} = \frac{\partial F_x}{\partial x} + \frac{\partial F_y}{\partial y} + \frac{\partial F_z}{\partial z}$

Spherical:  $\nabla \cdot \mathbf{F} = \frac{1}{r^2} \frac{\partial}{\partial r} (r^2 F_r) + \frac{1}{r \sin \theta} \frac{\partial}{\partial \theta} (\sin \theta F_\theta) + \frac{1}{r \sin \theta} \frac{\partial}{\partial \phi} (F_\phi)$

Cylindrical:  $\nabla \cdot \mathbf{F} = \frac{1}{\rho} \frac{\partial}{\partial \rho} (\rho F_\rho) + \frac{1}{\rho} \frac{\partial F_\theta}{\partial \theta} + \frac{\partial F_z}{\partial \phi}$

# CURL OF VECTOR FIELDS



Consider a vector field  $\mathbf{F}$ .  
Given two arbitrary displacements  
 $\delta \mathbf{r}_\alpha$  and  $\delta \mathbf{r}_\beta$  around the point  $(u_1, u_2, u_3)$ ,

Can we define a function  $\mathbf{X}(\mathbf{F})$  such that  
$$\iint \mathbf{X}(\mathbf{F}) \cdot d\mathbf{S} = \oint \mathbf{F} \cdot d\mathbf{l}$$

$$\delta \mathbf{r}^\alpha = \epsilon_1 h_1 \delta u_1^\alpha + \epsilon_2 h_2 \delta u_2^\alpha + \epsilon_3 h_3 \delta u_3^\alpha$$

$$\delta \mathbf{r}^\beta = \epsilon_1 h_1 \delta u_1^\beta + \epsilon_2 h_2 \delta u_2^\beta + \epsilon_3 h_3 \delta u_3^\beta$$

# CURL OF VECTOR FIELDS

What is the area element?

$$d\mathbf{S} = \delta \mathbf{r}^\alpha \times \delta \mathbf{r}^\beta = \begin{vmatrix} \epsilon_1 & \epsilon_2 & \epsilon_3 \\ h_1 \delta u_1^\alpha & h_2 \delta u_2^\alpha & h_3 \delta u_3^\alpha \\ h_1 \delta u_1^\beta & h_2 \delta u_2^\beta & h_3 \delta u_3^\beta \end{vmatrix}$$

What is the flux of the vector  $\mathbf{X}(F)$  through this area?

$$\begin{aligned} \mathbf{X}(F) \cdot d\mathbf{S} = & X_1 h_2 h_3 [\delta u_2^\alpha \delta u_3^\beta - \delta u_3^\alpha \delta u_2^\beta] \\ & - X_2 h_1 h_3 [\delta u_1^\alpha \delta u_3^\beta - \delta u_3^\alpha \delta u_1^\beta] \\ & + X_3 h_1 h_2 [\delta u_1^\alpha \delta u_2^\beta - \delta u_2^\alpha \delta u_1^\beta] \end{aligned}$$

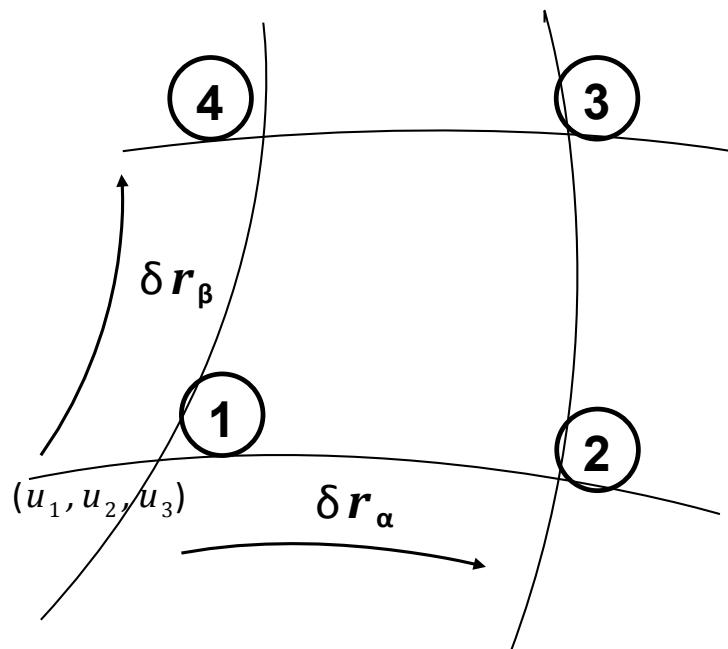
# CURL OF VECTOR FIELDS

Consider the pair of paths  $(1 \rightarrow 2)$  and  $(3 \rightarrow 4)$

$$\begin{aligned}\mathbf{F} \cdot \delta \mathbf{l}|_{1 \rightarrow 2} &= F_1 h_1 \delta u_1^\alpha + F_2 h_2 \delta u_2^\alpha + F_3 h_3 \delta u_3^\alpha \\ \mathbf{F} \cdot \delta \mathbf{l}|_{3 \rightarrow 4} &= [F_i h_i + (\nabla F_i h_i) \cdot \delta \mathbf{r}^\beta] (-\delta u_i^\alpha) \quad (i=1,2,3)\end{aligned}$$

Write contributions from  $\mathbf{F} \cdot \delta \mathbf{l}|_{2 \rightarrow 3}$  &  $\mathbf{F} \cdot \delta \mathbf{l}|_{4 \rightarrow 1}$  similarly

Full path gives:



$$\begin{aligned}& (\nabla \mathbf{F} \cdot \delta \mathbf{r}^\beta) \cdot \delta \mathbf{r}^\alpha - (\nabla \mathbf{F} \cdot \delta \mathbf{r}^\alpha) \cdot \delta \mathbf{r}^\beta \\ &= \sum_{k,i} \left[ \frac{1}{h_k} \frac{\partial F_i h_i}{\partial u_k} \delta u_i^\beta \right] h_k \delta u_k^\alpha \\ &\quad - \sum_{k,i} \left[ \frac{1}{h_k} \frac{\partial F_i h_i}{\partial u_k} \delta u_i^\alpha \right] h_k \delta u_k^\beta \\ &= \sum_{k,i} \left[ \frac{\partial F_i h_i}{\partial u_k} - \frac{\partial F_k h_k}{\partial u_i} \right] \delta u_i^\beta \delta u_k^\alpha\end{aligned}$$

# CURL OF VECTOR FIELDS

Comparing coefficients,

$$X_1 h_2 h_3 = \left[ \frac{\partial F_3 h_3}{\partial u_2} - \frac{\partial F_2 h_2}{\partial u_3} \right]$$

The function  $X(F)$  is then,

$$\text{So } X(\mathbf{F}) = \frac{1}{h_1 h_2 h_3} \begin{vmatrix} h_1 \epsilon_1 & h_2 \epsilon_2 & h_3 \epsilon_3 \\ \frac{\partial}{\partial u_1} & \frac{\partial}{\partial u_2} & \frac{\partial}{\partial u_3} \\ h_1 F_1 & h_2 F_2 & h_3 F_3 \end{vmatrix} \equiv \begin{cases} \nabla \times \mathbf{F} \\ \text{curl } \mathbf{F} \\ \text{rot } \mathbf{F} \end{cases}$$

We have  $\oint \nabla \times \mathbf{F} \cdot d\mathbf{S} = \oint \mathbf{F} \cdot d\mathbf{l}$  (Stoke's theorem)

# GRADIENT, DIVERGENCE, CURL - GENERALIZED COORDINATES

$$\nabla T = \frac{1}{h_1} \frac{\partial T}{\partial u_1} \epsilon_1 + \frac{1}{h_2} \frac{\partial T}{\partial u_2} \epsilon_2 + \frac{1}{h_3} \frac{\partial T}{\partial u_3} \epsilon_3 = \frac{1}{h_i} \frac{\partial T}{\partial u_i} \epsilon_i$$

$$\nabla \cdot \mathbf{F} = \frac{1}{h_1 h_2 h_3} \left[ \frac{\partial}{\partial u_1} (h_2 h_3 F_1) + \frac{\partial}{\partial u_2} (h_1 h_3 F_2) + \frac{\partial}{\partial u_3} (h_1 h_2 F_3) \right]$$

$$\begin{aligned} \nabla \times \mathbf{F} = & \frac{1}{h_1 h_2 h_3} \left[ h_1 \epsilon_1 \left\{ \frac{\partial}{\partial u_2} (h_3 F_3) - \frac{\partial}{\partial u_3} (h_2 F_2) \right\} \right. \\ & + h_2 \epsilon_2 \left\{ \frac{\partial}{\partial u_3} (h_1 F_1) - \frac{\partial}{\partial u_1} (h_3 F_3) \right\} \\ & \left. + h_3 \epsilon_3 \left\{ \frac{\partial}{\partial u_1} (h_2 F_2) - \frac{\partial}{\partial u_2} (h_1 F_1) \right\} \right] \end{aligned}$$

# IS EVERYTHING CONSISTENT? - GRADIENT

$$f = \rho^3 \sin \phi$$

$$f = x^2 y + y^3$$

$$\nabla f = 3\rho^2 \sin \phi \hat{\rho} + \rho^2 \cos \phi \hat{\phi}$$

$$\nabla f = 2x y \hat{x} + (3y^2 + x^2) \hat{y}$$

$$\begin{aligned}\nabla f &= 3\rho^2 \sin \phi (\cos \phi \hat{x} + \sin \phi \hat{y}) + \rho^2 \cos \phi (-\sin \phi \hat{x} + \cos \phi \hat{y}) \\ &= [2\rho^2 \sin \phi \cos \phi] \hat{x} + [3\rho^2 \sin^2 \phi + \rho^2 \cos^2 \phi] \hat{y} \\ &= 2x y \hat{x} + (3y^2 + x^2) \hat{y}\end{aligned}$$


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$$f = \frac{1}{r}$$

$$f = \frac{1}{\sqrt{x^2 + y^2 + z^2}}$$

$$\nabla f = -\frac{1}{r^2} \hat{r}$$

$$\nabla f = -\frac{1}{r^2} (\sin \theta \cos \phi \hat{x} + \sin \theta \sin \phi \hat{y} + \cos \theta \hat{z}) = -\frac{1}{r^2} \left( \frac{x}{r} \hat{x} + \frac{y}{r} \hat{y} + \frac{z}{r} \hat{z} \right)$$

$$= -\left( \frac{x}{(x^2 + y^2 + z^2)^{3/2}} \hat{x} + \frac{y}{(x^2 + y^2 + z^2)^{3/2}} \hat{y} + \frac{z}{(x^2 + y^2 + z^2)^{3/2}} \hat{z} \right)$$

# IS EVERYTHING CONSISTENT? - DIVERGENCE

$$\vec{F} = \rho^3 \hat{\rho} + \rho z \hat{\phi} + \rho z \sin \phi \hat{z}$$

$$\nabla \cdot \vec{F} = \frac{1}{\rho} \left\{ \frac{\partial}{\partial \rho} (\rho F_\rho) + \frac{\partial}{\partial \phi} (F_\phi) + \frac{\partial}{\partial z} (\rho F_z) \right\} = 4\rho^2 + \rho \sin \phi$$

$$\vec{F} = (x^3 + xy^2 - yz) \hat{x} + (x^2y + y^3 + xz) \hat{y} + yz \hat{z}$$

$$\nabla \cdot \vec{F} = 4(x^2 + y^2) + y$$

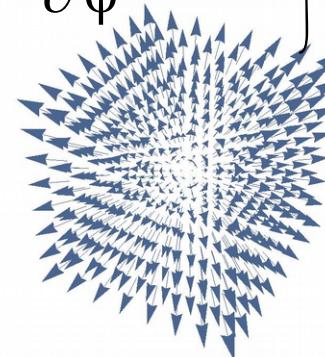
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$$\vec{F} = r \hat{r}$$

$$\nabla \cdot \vec{F} = \frac{1}{r^2 \sin \theta} \left\{ \frac{\partial}{\partial r} (r^2 \sin \theta F_r) + \frac{\partial}{\partial \theta} (r \sin \theta F_\theta) + \frac{\partial}{\partial \phi} (r F_\phi) \right\} = 3$$

$$\vec{F} = x \hat{x} + y \hat{y} + z \hat{z}$$

$$\nabla \cdot \vec{F} = 3$$



# IS EVERYTHING CONSISTENT? - CURL

$$\mathbf{F} = \rho^3 \hat{\rho} + \rho z \hat{\phi} + \rho z \sin \phi \hat{z}$$

$$\nabla \times \mathbf{F} = \frac{1}{\rho} \begin{vmatrix} \hat{\rho} & \hat{\phi} & \hat{z} \\ \frac{\partial}{\partial \rho} & \frac{\partial}{\partial \phi} & \frac{\partial}{\partial z} \\ \rho^3 & \rho^2 z & \rho z \sin \phi \end{vmatrix} = (z \cos \phi - \rho) \hat{\rho} + z \sin \phi \hat{\phi} + 2z \hat{z}$$

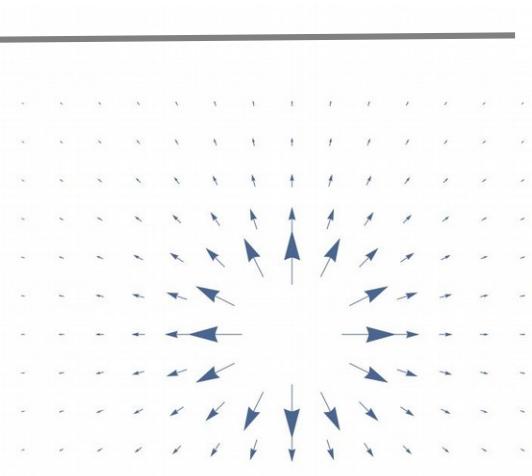
$$\mathbf{F} = (x^3 + xy^2 - yz) \hat{x} + (x^2y + y^3 + xz) \hat{y} + yz \hat{z}$$


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$$\mathbf{F} = r^k \hat{r}$$

$$\nabla \times \mathbf{F} = \frac{1}{r^2 \sin \theta} \begin{vmatrix} \hat{r} & \hat{\theta} & \hat{\phi} \\ \frac{\partial}{\partial r} & \frac{\partial}{\partial \theta} & \frac{\partial}{\partial \phi} \\ r^k & r \cdot 0 & r \sin \theta \cdot 0 \end{vmatrix} = 0$$

$$\mathbf{F} = \left( \frac{x}{(x^2 + y^2 + z^2)^{(k+1)/2}} \hat{x} + \frac{y}{(x^2 + y^2 + z^2)^{(k+1)/2}} \hat{y} + \frac{z}{(x^2 + y^2 + z^2)^{(k+1)/2}} \hat{z} \right)$$



# DIVERGENCE OF $1/R^2$

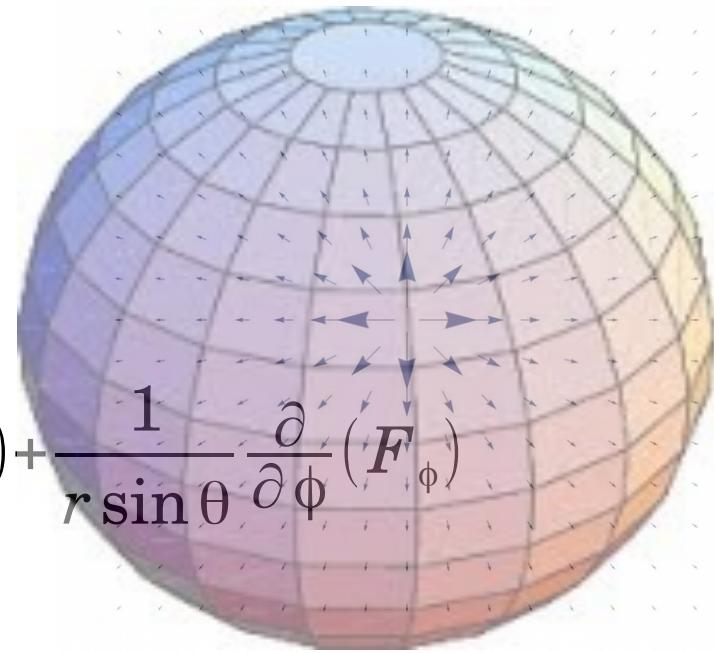
Consider the vector function  $\mathbf{v} = \frac{\hat{\mathbf{r}}}{r^2}$

What is  $\nabla \cdot \mathbf{v}$ ?

$$\nabla \cdot \mathbf{F} = \frac{1}{r^2} \frac{\partial}{\partial r} (r^2 F_r) + \frac{1}{r \sin \theta} \frac{\partial}{\partial \theta} (\sin \theta F_\theta) + \frac{1}{r \sin \theta} \frac{\partial}{\partial \phi} (F_\phi)$$

$$\nabla \cdot \mathbf{v} = \frac{1}{r^2} \frac{\partial}{\partial r} \left( r^2 \frac{1}{r^2} \right) = 0$$

$$\iiint \nabla \cdot \mathbf{v} dV = 0$$

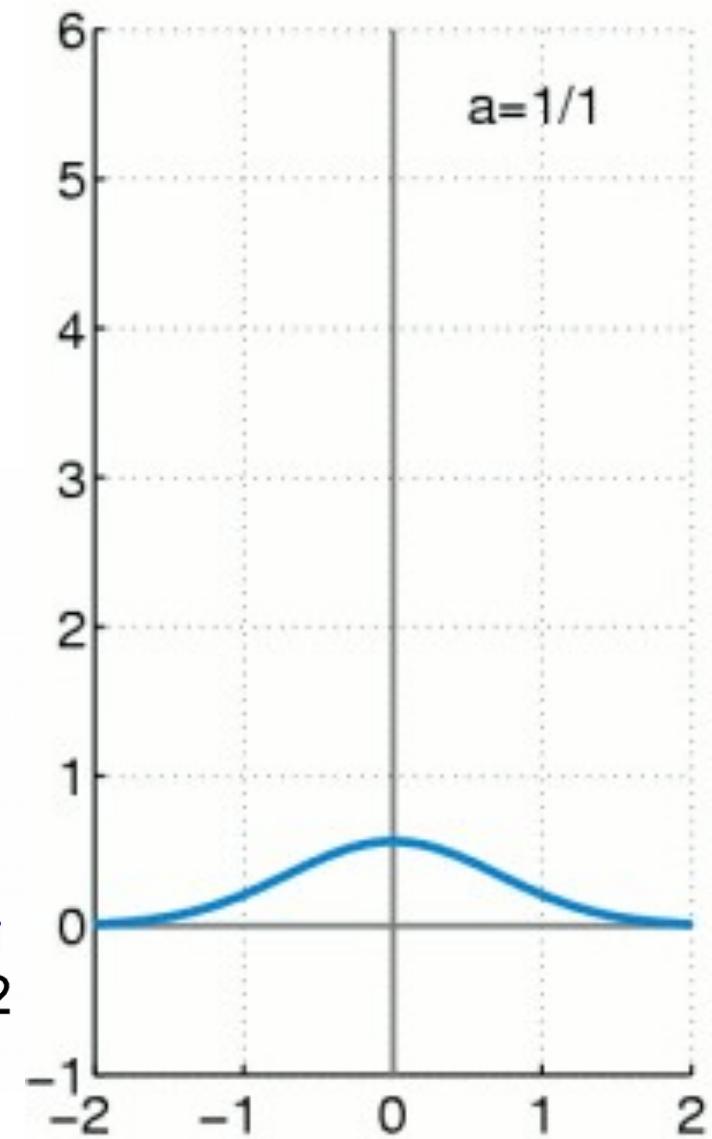
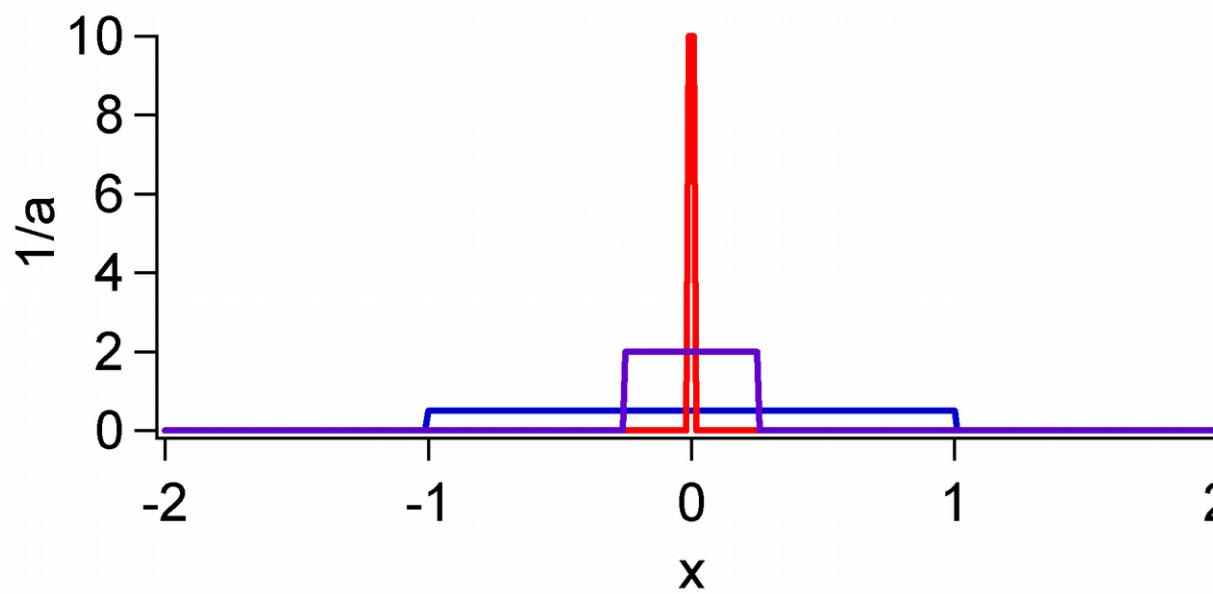


Something is wrong!

$$\oint \mathbf{v} \cdot d\mathbf{a} = \iint \left( \frac{1}{R^2} \hat{\mathbf{r}} \right) \cdot (R^2 \sin \theta d\theta d\phi \hat{\mathbf{r}}) = \left( \int_0^\pi \sin \theta d\theta \right) \left( \int_0^{2\pi} d\phi \right) = 4\pi$$

# THE DIRAC DELTA FUNCTION

$$\int_{-\infty}^{\infty} \delta(x - x_0) f(x) dx = f(x_0)$$



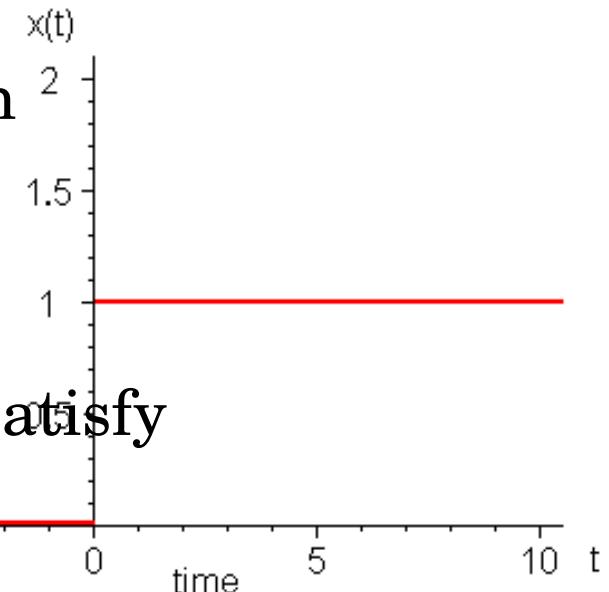
# THE DIRAC DELTA FUNCTION

Consider a simpler example: the step function

$$x(t) = \begin{cases} 0 & t < 0 \\ 1 & t > 0 \end{cases}$$

$\frac{dx}{dt} = ?$  looks like zero everywhere but must satisfy

$$\int_{0-|\epsilon|}^{0+|\epsilon|} \left( \frac{dx}{dt} \right) dt = x(|\epsilon|) - x(-|\epsilon|) = 1 \quad \forall \epsilon \neq 0$$

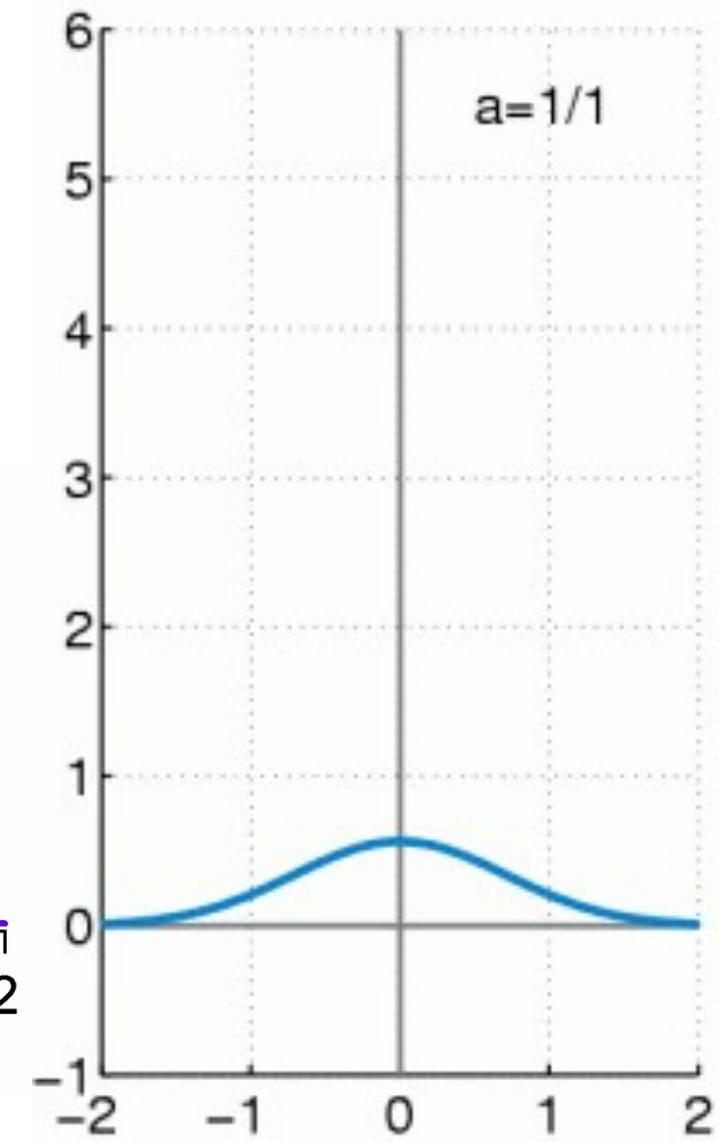
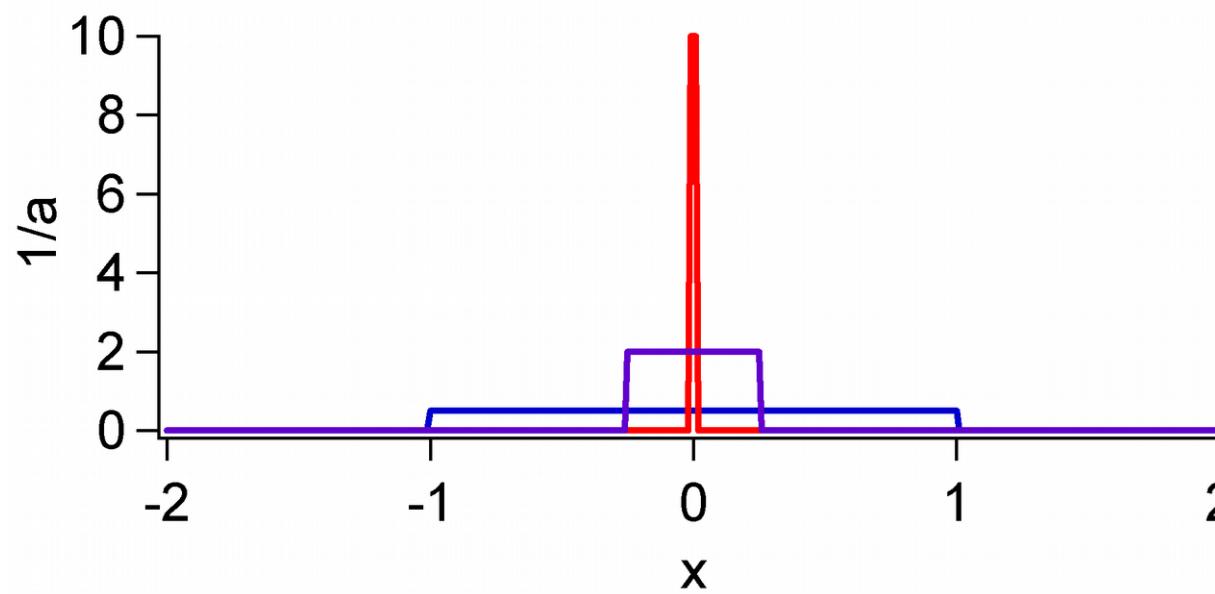


Such integrable singularities are treated by **defining** the  $\delta$  function

$$\int_a^b \delta(x - x_0) f(x) dx = \begin{cases} f(x_0), & \text{if } x_0 \text{ is within the limits} \\ 0 & \text{otherwise} \end{cases}$$

The Delta function is not a proper function, but rather the limit of a sequence of functions. It's called a generalized function or distribution, and is properly defined only within an integral.

$$\int_{-\infty}^{\infty} \delta(x - x_0) f(x) dx = f(x_0)$$



# THE 3D DIRAC DELTA FUNCTION

$$\delta(\mathbf{r}) = \delta(x)\delta(y)\delta(z)$$

$$\int_{\text{all space}} \delta(\mathbf{r}) dV = \int_{-\infty}^{\infty} \int_{-\infty}^{\infty} \int_{-\infty}^{\infty} \delta(x)\delta(y)\delta(z) dx dy dz = 1$$

$$\int_{\text{all space}} f(\mathbf{r}) \delta(\mathbf{r} - \mathbf{a}) dV = f(\mathbf{a})$$

$$\iiint \nabla \cdot \frac{\hat{\mathbf{r}}}{\mathbf{r}^2} dV = 4\pi$$

$$\nabla \cdot \frac{\hat{\mathbf{r}}}{\mathbf{r}^2} = 4\pi \delta^3(\mathbf{r})$$

Prove:

$$\delta(-x) = \delta(x)$$

$$\delta(\alpha x) = \frac{\delta(x)}{|\alpha|}$$

$$x \frac{d}{dx}(\delta(x)) = -\delta(x)$$

# UNIQUE SPECIFICATION OF A VECTOR FIELD

If we specify the curl of a function  $\nabla \times \mathbf{F} = \mathbf{C}$ ,  
can we determine the function  $\mathbf{F}$ ?

NO.  $\mathbf{F}' = \mathbf{F} + \nabla \phi$  will give same result!

If we specify the divergence of a function  $\nabla \cdot \mathbf{F} = D$ ,  
can we determine the function  $\mathbf{F}$ ?

NO.  $\mathbf{F}' = \mathbf{F} + \nabla \times \mathbf{A}$  will give same result!

If we specify the divergence **AND** the curl of a function  
 $\nabla \times \mathbf{F} = \mathbf{C}$  **AND**  $\nabla \cdot \mathbf{F} = D$ ,  
can we determine the function  $\mathbf{F}$ ?

YES! (Provided the boundary conditions are specified)

**HELMHOLTZ'S THEOREM**

# HELMHOLTZ'S THEOREM

If  $\nabla \times \mathbf{F} = \mathbf{C}$  AND  $\nabla \cdot \mathbf{F} = D$ ,

Note that, necessarily  $\nabla \cdot \mathbf{C} = 0$

We claim

$$\mathbf{F} = -\nabla U + \nabla \times \mathbf{W}$$

where

$$U(\mathbf{r}) = \frac{1}{4\pi} \int \frac{D(\mathbf{r}')}{|\mathbf{r} - \mathbf{r}'|} dV'$$

and

$$\mathbf{W}(\mathbf{r}) = \frac{1}{4\pi} \int \frac{\mathbf{C}(\mathbf{r}')}{|\mathbf{r} - \mathbf{r}'|} dV'$$

$$\begin{aligned}\nabla \cdot \mathbf{F} &= -\nabla^2 U - \frac{1}{4\pi} \int D(\mathbf{r}') \nabla^2 \left( \frac{1}{|\mathbf{r} - \mathbf{r}'|} \right) dV' = \int D(\mathbf{r}') \delta^3(\mathbf{r} - \mathbf{r}') dV' \\ &= D(\mathbf{r})\end{aligned}$$

# HELMHOLTZ'S THEOREM

$$\nabla \times \mathbf{F} = \nabla \times (\nabla \times \mathbf{W}) = -\nabla^2 \mathbf{W} + \nabla(\nabla \cdot \mathbf{W})$$

$$-\nabla^2 \mathbf{W} = -\frac{1}{4\pi} \int \mathbf{C}(\mathbf{r}') \nabla^2 \left( \frac{1}{|\mathbf{r} - \mathbf{r}'|} \right) dV' = \int \mathbf{C}(\mathbf{r}') \delta^3(\mathbf{r} - \mathbf{r}') dV' = \mathbf{C}(\mathbf{r})$$

$$\begin{aligned} 4\pi \nabla \cdot \mathbf{W} &= \int \mathbf{C}(\mathbf{r}') \cdot \nabla \left( \frac{1}{|\mathbf{r} - \mathbf{r}'|} \right) dV' = - \int \mathbf{C}(\mathbf{r}') \cdot \nabla' \left( \frac{1}{|\mathbf{r} - \mathbf{r}'|} \right) dV' \\ &= \int \frac{1}{|\mathbf{r} - \mathbf{r}'|} \nabla' \cdot \mathbf{C}(\mathbf{r}') dV' - \oint \frac{1}{|\mathbf{r} - \mathbf{r}'|} \mathbf{C} \cdot da \\ &= 0 \quad (\text{if } \mathbf{C} \text{ goes to zero sufficiently rapidly}) \end{aligned}$$

$$\nabla \cdot (f \vec{A}) = f(\nabla \cdot \vec{A}) + \vec{A} \cdot (\nabla f)$$

How rapidly?  $\mathbf{C}$  and  $\mathbf{D}$  must go to zero more rapidly than  $1/r^2$

# HELMHOLTZ'S THEOREM

Is this solution unique?

YES, as long as the vector field  $\mathbf{F}(\mathbf{r})$  itself goes to zero at infinity

If the divergence  $D(\mathbf{r})$  and the curl  $\mathbf{C}(\mathbf{r})$  of a vector function  $\mathbf{F}(\mathbf{r})$  are specified, and if they both go to zero faster than  $1/r^2$  as  $r$  goes to infinity, and if  $\mathbf{F}(\mathbf{r})$  itself goes to zero as  $r$  goes to infinity, then  $\mathbf{F}(\mathbf{r})$  is uniquely given by

$$\mathbf{F} = -\nabla U + \nabla \times \mathbf{W}$$

# MULTIPLE VECTOR PRODUCTS - $\epsilon$ $\delta$ (LEVI-CIVITA)

Write the dot product as  $\mathbf{A} \cdot \mathbf{B} = \delta_{ij} A_i B_j$  where  $\delta_{ij} = \begin{cases} 1 & \text{if } i=j \\ 0 & \text{if } i \neq j \end{cases}$

Write the cross product as

$\mathbf{A} \times \mathbf{B} |_i = \epsilon_{ijk} A_j B_k$  where  $\epsilon_{ijk} = \begin{cases} 1 & \text{if } ijk \text{ is an even permutation of } 123 \\ -1 & \text{if } ijk \text{ is an odd permutation of } 123 \\ 0 & \text{otherwise} \end{cases}$

Convince yourself that  $\epsilon_{ijk} = \epsilon_i \cdot \epsilon_j \times \epsilon_k$

This works with operators also: with  $x_i$  for  $x, y, z$

$$\nabla \cdot \mathbf{A} = \frac{\partial A_i}{\partial x_i}$$

$$\nabla \times \mathbf{A} |_i = \epsilon_{ijk} \frac{\partial A_k}{\partial x_j}$$

# MULTIPLE VECTOR PRODUCTS - $\epsilon \delta$ (LEVI-CIVITA)

Consider a vector triple product  $\mathbf{A} \times (\mathbf{B} \times \mathbf{C})$

$$\begin{aligned}\mathbf{A} \times (\mathbf{B} \times \mathbf{C})|_i &= \epsilon_{ijk} A_j (\mathbf{B} \times \mathbf{C})_k \\ &= \epsilon_{ijk} A_j (\epsilon_{kpq} B_p C_q) \\ &= \epsilon_{kij} \epsilon_{kpq} A_j B_p C_q\end{aligned}$$

What to do with a product like  $\epsilon_{ijk} \epsilon_{lpq}$  ?

This is either -1 or 0 or 1

$$\epsilon_{ijk} \epsilon_{lpq} = \begin{vmatrix} \delta_{il} & \delta_{ip} & \delta_{iq} \\ \delta_{jl} & \delta_{jp} & \delta_{jq} \\ \delta_{kl} & \delta_{kp} & \delta_{kq} \end{vmatrix}$$

# MULTIPLE VECTOR PRODUCTS - $\epsilon \delta$ (LEVI-CIVITA)

$$\begin{aligned}\epsilon_{kij} \epsilon_{kpq} &= \delta_{kk} \delta_{ip} \delta_{jq} + \delta_{kp} \delta_{iq} \delta_{jk} + \delta_{kq} \delta_{ik} \delta_{jp} \\&\quad - \delta_{kk} \delta_{iq} \delta_{jp} - \delta_{kp} \delta_{ik} \delta_{jq} - \delta_{kq} \delta_{ip} \delta_{jk} \\&= \delta_{kk} (\delta_{ip} \delta_{jq} - \delta_{iq} \delta_{jp}) + \delta_{kp} (\delta_{iq} \delta_{jk} - \delta_{ik} \delta_{jq}) + \delta_{kq} (\delta_{ik} \delta_{jp} - \delta_{ip} \delta_{jk}) \\&= 3(\delta_{ip} \delta_{jq} - \delta_{iq} \delta_{jp}) + (\delta_{iq} \delta_{jp} - \delta_{ip} \delta_{jq}) + (\delta_{iq} \delta_{jp} - \delta_{ip} \delta_{jq}) \\&= \delta_{ip} \delta_{jq} - \delta_{iq} \delta_{jp}\end{aligned}$$

Using the last result (with  $i = p$ )

$$\begin{aligned}\epsilon_{kij} \epsilon_{kiq} &= \delta_{ii} \delta_{jq} - \delta_{iq} \delta_{ji} \\&= 3\delta_{jq} - \delta_{jq} \\&= 2\delta_{jq}\end{aligned}$$

Using the last result (with  $j = q$ )

$$\begin{aligned}\epsilon_{kij} \epsilon_{kij} &= 2\delta_{jj} \\&= 6\end{aligned}$$

Successive summation over indices.

The first sum is most frequently encountered. It allows you to write a cross product in terms of dot product like terms .

# PROOF OF VECTOR IDENTITIES

$$\begin{aligned} [\nabla \times (\nabla \times \mathbf{A})]_i &= \epsilon_{ijk} \partial_j \epsilon_{kmn} \partial_m A_n = \epsilon_{kij} \epsilon_{kmn} \partial_j \partial_m A_n \\ &= (\delta_{im} \delta_{jn} - \delta_{in} \delta_{jm}) \partial_j \partial_m A_n = \partial_j \partial_i A_j - \partial_j \partial_j A_i \\ &= [\nabla (\nabla \cdot \mathbf{A}) - \nabla^2 \mathbf{A}]_i \end{aligned}$$

$$\begin{aligned} [\nabla \times (\mathbf{A} \times \mathbf{B})]_i &= \epsilon_{ijk} \partial_j \epsilon_{kmn} A_m B_n = \epsilon_{kij} \partial_j \epsilon_{kmn} \partial_j A_m B_n \\ &= (\delta_{im} \delta_{jn} - \delta_{in} \delta_{jm}) \partial_j (A_m B_n) \\ &= \partial_j (A_i B_j) - \partial_j (A_j B_i) \\ &= B_j \partial_j A_i + A_i \partial_j B_j - B_i \partial_j A_j - A_j \partial_j B_i \\ &= (\vec{B} \cdot \nabla) \vec{A} + \vec{A} (\nabla \cdot \vec{B}) - \vec{B} (\nabla \cdot \vec{A}) - (\vec{A} \cdot \nabla) \vec{B} \end{aligned}$$