

# MADANALYSIS 5, a user-friendly framework for collider phenomenology

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## Abstract

We present MADANALYSIS 5, a new framework for phenomenological investigations at particle colliders. Based on a C++ kernel, this program allows us to efficiently perform, in a straightforward and user-friendly fashion, sophisticated physics analyses of event files such as those generated by a large class of Monte Carlo event generators. MADANALYSIS 5 comes with two modes of running. The first one, easier to handle, uses the strengths of a powerful PYTHON interface in order to implement physics analyses by means of a set of intuitive commands. The second one requires one to implement the analyses in the C++ programming language, directly within the core of the analysis framework. This opens unlimited possibilities concerning the level of complexity which can be reached, being only limited by the programming skills and the originality of the user.

*Keywords:* Particle physics phenomenology, Monte Carlo event generators, hadron colliders.

## PROGRAM SUMMARY

*Manuscript Title:* MADANALYSIS 5, a user-friendly framework for collider phenomenology.

*Authors:* Eric Conte, Benjamin Fuks, Guillaume Serret.

*Program Title:* MADANALYSIS 5

*Licensing provisions:* Permission to use, copy, modify and distribute this program is granted under the terms of the GNU General Public License.

*Programming language:* PYTHON, C++.

*Computer:* All platforms on which PYTHON version 2.7, ROOT version 5.27 and the G++ compiler are available. Compatibility with newer versions of these programs is also ensured. However, the PYTHON version must be below version 3.0.

*Operating system:* UNIX, LINUX and MAC OS operating systems on which the above-mentioned versions of PYTHON and ROOT, as well as G++, are available.

*Keywords:* Particle physics phenomenology, Monte Carlo event generators, hadron colliders.

*Classification:* 11.1 General, High Energy Physics and Computing.

*Nature of problem:* Implementing sophisticated phenomenological analyses in high-energy physics through a flexible, efficient and straightforward fashion, starting from event files as those produced by Monte Carlo event generators. The event files can have been matched or not to parton-showering and can have been processed or not by a (fast) simulation of a detector. According to the sophistication level of the event files (parton-level, hadron-level, reconstructed-level), one must note that several input formats are possible.

*Solution method:* We implement an interface allowing to produce predefined as well as user-defined histograms for a large class of kinematical distributions after applying a set of event selection cuts specified by the user. This therefore allows to devise robust and novel search strategies for collider experiments, such as those currently running at the Large Hadron Collider at CERN, in a very efficient way.

*Restrictions:* Unsupported event file format.

*Unusual features:* The code is fully based on object representations for events, particles, reconstructed objects and cuts, which facilitates the implementation of an analysis.

*Running time:* It depends on the purposes of the user and on the number of events to process. It varies from a few seconds to the order of the minute for several millions of events.

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## 1. Introduction

Among the key topics of the present experimental program of high-energy physics lies the quest for new physics and the identification of the fundamental building blocks of matter, together with their interactions. In these prospects, the Large Hadron Collider (LHC) is currently exploring the TeV scale and multi-purpose experiments, such as ATLAS or CMS, are currently pushing the limits on beyond the Standard Model physics to a further and further frontier. Discoveries from these experiments, together with their interpretation, are in general challenging and strongly rely on our ability to accurately simulate both the possible candidate signals and the backgrounds. This task is however rendered quite complicated due to the complexity of the typical final states to be produced at the LHC, which contain large numbers of light and heavy flavor jets, charged leptons and missing transverse energy. Consequently, the overwhelming sources of Standard Model background require the development of robust, and possibly novel, search strategies. In this context, tools allowing us to compute predictions for large classes of models are central.

This has triggered, during the last twenty years, a lot of efforts dedicated to the development of multi-purpose matrix-element based event generators such as ALPGEN [1], COMIX [2], COMPHEP/CALCHEP [3, 4, 5], HELAC [6], MADGRAPH/MADEVENT [7, 8, 9, 10, 11], SHERPA [12, 13] and WHIZARD [14, 15]. As a result, the problem of the generation of parton-level events, at the leading-order accuracy, for many renormalizable or non-renormalizable new physics theories has been solved. More recently, progress has also been achieved in the automation of next-to-leading-order computations. On the one hand, the generation of the real emission contributions with the appropriate subtraction terms has been achieved in an automatic way [16, 17, 18, 19, 20, 21]. On the other hand, several algorithms addressing the numerical calculation of loop amplitudes have been proposed [22, 23, 24, 25, 26] and successfully applied to the computation of Standard Model processes of physical interest [27, 28, 29, 30, 31].

Even if each of the above-mentioned tools is based on a different philosophy, uses a specific programming language and requires a well-defined input format for the physics models under consideration, programs such as FEYN-RULES [32, 33, 34, 35, 36] and LANHEP [37, 38] have alleviated the time-consuming and error-prone task of implementing new physics theories in the Monte Carlo tools. Furthermore, the introduction of the Universal FEYN-

RULES Output (UFO) format [39] and the development of an automated tool computing helicity amplitudes [40] have also streamlined the communication between the construction of a new physics theory and its implementation in the matrix-element generators through a standardized fashion.

Parton-level physics analyses based on event samples produced by matrix-element generators are far from describing the reality of what is observed in any existing detector. This kind of phenomenological work can however be useful in the prospects of investigating new types of signature and devising original search strategies, preliminary to more complete phenomenological investigations. In order to facilitate the information transfer from the matrix-element generators, a generic format for storing (parton-level) events and their properties has been proposed ten years ago, the so-called Les Houches Event (LHE) file format [41, 42]. Following the LHE standards, events are stored in a single file, following an XML-like structure, together with additional information related to the way in which the events have been generated.

Monte Carlo generator-based physics analyses at the parton-level then consist in first parsing LHE event files related to both the signal and the different sources of background, then implementing various selection cuts on the objects contained in the events, *i.e.*, quarks, leptons, neutrinos, new stable particles, *etc*, and finally in creating histograms representing several kinematical quantities. By means of signal over background ratios, optimization of the selection cuts can be achieved in order to maximize the chances to unveil the signal of interest. Of course, the only sensible conclusions which could be stated at this stage would be to motivate (or not) a more realistic analysis, including at least parton showering and hadronization.

An accurate simulation of the collision to be observed at hadron colliders indeed requires a proper modeling of the strong interaction, including parton showering, fragmentation and hadronization. This is efficiently provided by packages such as PYTHIA [43, 44] and HERWIG [45, 46] and several algorithms matching parton showering to hard scattering matrix elements have been recently developed [47, 48, 49, 50].

Even if not directly necessary for the analysis of the event samples, a large part of the information present in a parton-level LHE event file allows for a further matching procedure. Consequently, the generation of hadron-level events for a multitude of Standard Model and beyond the Standard Model processes can be (and has been) done in a systematic fashion. This starts from parton-level events stored in LHE format-compliant files, as generated

by matrix-element based event generators. These files are then further processed by a parton showering and hadronization code which allows us to match the strengths of both the description of the physics embedded into the matrix elements and the one modeled by the parton showering. Consequently, physics analyses at the hadron-level are far more sophisticated than their counterparts at the parton-level.

When analyzing hadron-level events, one must note that the key difference with respect to the parton-level case lies in the objects which are contained in the event files. In contrast to partonic events, hadronic events consist in general in a huge collection of hadrons given together with their four-momentum. The particle content of the final state is thus much richer and the storing of the information at the time of parsing the event file is hence rendered rather cumbersome.

In order to streamline this procedure, dedicated event class libraries have been developed and several common formats for outputting hadron-level event samples exist. Event files compliant with such formats can in general be straightforwardly read by fast detector simulation programs, necessary for a more advanced level of sophistication of the phenomenological analysis. As examples, one finds the STDHEP [51] or HEPMC [52] structure for event files, both widely used in the high-energy physics community. According to these conventions, an event contains, in addition to the final state particles and their properties, the whole event history as a set of mother-particle to daughter-particle relations. Let us also note that the task of parsing hadron-level event files can be highly simplified after clustering the hadrons into jets, using jet algorithms such as, *e.g.*, those included in the FASTJET program [53]. This simplified picture allows for the usage of a simpler event format, such as the LHE format introduced above, which subsequently renders an analysis easier to implement.

State-of-the-art phenomenological analyses require in general fast (and realistic) detector simulation in order to correctly estimate both the signals under consideration and the backgrounds. Starting from hadronized event samples, several frameworks, such as PGS 4 [54] and DELPHES [55] are dedicated to this task. They produce event samples containing reconstructed objects such as photons, jets, electrons, muons or missing energy, together with their properties. Whereas the description of a specific signature as provided by a fast detector simulation tool is still far from what could have been obtained using a full detector simulation, including among others the transport of the particles through the detector material, fast simulation of

collider experiments is often sufficient to study the feasibility of a specific analysis on realistic grounds. The results then motivate (or not) the associated analysis in the context of a full detector simulation, the latter being embedded in generally complex experimental software such as those used by large collaborations.

As stated above, after their processing by a fast detector simulation, the original sets of hadrons are replaced by high-level reconstructed objects stored together with properties such as their (smeared) four-momentum or if they have been tagged as a  $b$ -jet or not. The structure of the events and their properties can be very efficiently saved into a file following the LHCO conventions [56]. Let us note that this format has been designed in particular for that purpose.

As outlined above, the study of the hadron collider phenomenology of a given signature can be performed at several levels of sophistication. Sometimes, very preliminary works at the parton-level are necessary in order to motivate further investigations. Hadron-level analyses give a clearer hint about the expected sensitivity of colliders to the signal under consideration. This assumes a perfect and ideal detector. In contrast, state-of-the-art phenomenology includes a fast and semi-realistic simulation of the detectors. This allows us to derive conclusions closer to the expectations estimated with the help of a full simulation of the detector as performed by the large collaboration software. The major drawback of the latter is that such tools consist in very heavy, complex, time-consuming and non-public algorithms. This therefore motivates pioneering works under the parton-level, hadronic-level or reconstructed-level assumptions<sup>1</sup>. It is however important to keep in mind that the final word is always included in the data.

Performing a phenomenological analysis on the basis of the results provided by Monte Carlo generators always starts with the reading of several event samples. Since partonic or hadronic events (the latter including or not a fast detector simulation) are in general stored under different formats, this step demands to use appropriate routines capable of understanding the file structure. Selection cuts on the objects included in the events, *i.e.*, partons, hadrons or high-level objects, according to the level of sophistication, are

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<sup>1</sup>In this paper, we denote by *reconstructed-level* events which have already been processed by a fast detector simulation, in contrast to *hadron-level* events which have only been showered, fragmented and hadronized and *parton-level* events where the matching to a parton showering algorithm is fully absent.

then implemented and applied to both the signal and background samples. This allows for the creation of histograms of various quantities in order to be able to extract (or not, in the worst case scenario) information about a signal in general swamped by backgrounds.

This procedure is most of the time based on home-made and non-public programs, especially due to the lack of a dedicated framework. As a consequence, this can lead to various problems in the validation and the traceability of the analyses as well as possibly in the interpretation of the results. In this work, we are alleviating this issue by proposing a single efficient framework for phenomenological analyses at any level of sophistication. We introduce the package **MADANALYSIS** 5, an open source program based on a multi-purpose C++ kernel, denoted **SAMPLEANALYZER**, which uses the **ROOT** platform [57].

Using the strengths of a **PYTHON** interface, the user can define his own physics analysis in an efficient, flexible and straightforward way. Similarly to the older **FORTRAN** and **PERL** version of **MADANALYSIS** which was linked to version 4 of the **MADGRAPH** program, **MADANALYSIS** 5 can either be run within **MADGRAPH** 5 or as a standalone package. It includes a complete reorganization of the code and the implementation of many novel functionalities such as an efficient method to implement cuts as well as to generate histograms and cut-flow charts in an automated fashion. Moreover, care has been taken in developing a fast and optimized code.

Consequently, the procedure for performing a phenomenological analysis has been drastically simplified since the only task left to the user is to define the corresponding selection cuts and the distributions to be computed. Sometimes, these embedded features might however not be sufficient according to the needs of the user. In order to overcome this limitation, **MADANALYSIS** 5 offers an expert mode of running with unlimited possibilities. It is unlimited in the sense that the user directly implements, within the C++ kernel, his own analysis.

Finally, in order to release a single framework for the analysis of parton-level, hadron-level or reconstructed-level based event simulations, several event file formats are supported as input, from the **LHE** files which could describe partonic or hadronic events to the more complex **STDHEP**, **HEPMC** and **LHCO** file formats. Let us note that, according to the needs of the users, interfacing additional event formats, such as, *e.g.*, the **ExROOTANALYSIS** format [58], can be easily achieved.

This paper documents the fifth version of **MADANALYSIS** and consists in

its user guide. An up-to-date version of this document, together with the program can be found on the web page

<http://madanalysis.irmp.ucl.ac.be>

Moreover, an extended, more pedestrian, version of this manual is also available at the same Internet address. The outline of this paper is as follows. In Section 2, we give a general overview of the program, its philosophy and its structure. Section 3 illustrates the capabilities of MADANALYSIS 5 without entering into the details, the latter being given in the next sections. Hence, Section 4 provides the guidelines for implementing basic (but still professional) physics analyses, *i.e.*, implementing simple selection cuts and creating histograms for various kinematical distributions, whilst Section 5 is dedicated to a more expert usage of the program which allows us to design more sophisticated analyses. Our conclusions are presented in Section 6. Finally, a technical Appendix on the installation of the program and its dependencies follows.

## 2. Overview of MADANALYSIS 5

### 2.1. MADANALYSIS 5 *in a nutshell*

The MADANALYSIS 5 package is an open source program allowing us to perform physics analyses of Monte Carlo event samples in an efficient, flexible and straightforward way. It relies on a C++ kernel, named SAMPLEANALYZER, which uses the ROOT platform and interacts with the user by means of a PYTHON command line interface.

The distribution of the program includes a home-made reader of event files created by Monte Carlo event generators. The reader is compliant with Monte Carlo samples containing events either at the parton-level, at the hadron-level or at the reconstructed-level. Hence, a unique framework can be used for analyzing events embedded equivalently into simplified LHE files or into more complex STDHEP, HEPMC and LHCO files.

From the information included in the event files, the user can ask MADANALYSIS 5 to generate histograms illustrating various properties of the generated physics processes. These properties range from observables related to the whole content of the events, such as the multiplicity of a given particle species or the missing transverse energy distribution, to observables associated to a specific particle such as, for instance, the transverse-momentum distribution of the final-state muon with the highest energy or the angular

separation between the final-state jets. The user has also the possibility to compare event samples related to different physical processes in a straightforward fashion. On the one hand, the investigated distributions can be stacked on the same histogram, after automatic normalization of the samples to an integrated luminosity which the user can specify. On the other hand, the distributions can be normalized to unity in order to facilitate the design of event selection cuts allowing for an efficient background rejection based on the shape of the distribution. In this case, the histograms could be superimposed rather than stacked in the aim of improving readability.

Moreover, if available, the whole event history, *i.e.*, all the event information from the hard-scattering process to the final-state hadrons, including the mother-particle to daughter-particle relations among the different stages yielding the final state particle content, is stored. One can emphasize that this feature is important for sanity and consistency checks of the generated event samples.

Event selection can be easily performed by means of a series of intuitive PYTHON commands yielding the application of selection cuts to the samples under consideration. For instance, it is possible to only consider, in the current analysis, events which contain at least a certain amount of missing transverse energy or a given number of final-state leptons. In the same way, one can access a specific particle in the event, such as the jet with the hardest transverse momentum, and require that it fulfills a given geometrical or kinematical criterion. As a last example, the user could decide to only select events containing at least two muons whose invariant mass is compatible with the  $Z$ -boson mass. After applying a given cut, MADANALYSIS 5 proceeds with an automatic computation of the associated cut efficiency.

The ultimate goal of a physics analysis in particle physics is to extract some signal from a usually swamping background so that one could study its properties. The MADANALYSIS 5 framework offers the possibility to tag a specific event sample as a background or signal sample. This tagging allows for an automatic treatment of the signal over background ratio, or of any other similar observable which can be specified by the user, together with the associated uncertainty. This quantity is recomputed after each of the different selection cuts implemented by the user. With these pieces of information available, optimizing selection cuts consequently becomes easier, which allows us to investigate in a fast and efficient way whether a given signature could be observable at colliders.

To display the results in a human-readable form, MADANALYSIS 5 can

collect them either into a LATEX document to be further compiled or under the form of an HTML webpage.

## 2.2. Basic concepts

The MADANALYSIS 5 program provides to the user a platform including a wide class of functionalities allowing one to perform sophisticated physics analyses. Even if the existing possibilities are rather large, see Section 4 and Section 5, implementing an analysis within the MADANALYSIS 5 framework always follows the same steps, each of these steps being linked to one or several of the key features of the program. These key features are briefly described in this Section, whilst additional and more detailed information can be found in the rest of this manual.

### Sample declaration - datasets.

When implementing an analysis within the framework of MADANALYSIS 5, the first task which is asked of the user is to indicate the Monte Carlo event files to be processed on run-time. In general, a full Monte Carlo event generation requires the simulation of several physical processes, such as the signal under consideration and the associated sources of background. Furthermore, for the processes with the highest cross sections, more than one event sample may have to be generated in order to get enough statistical significance. The user then requires these samples, describing the same physics, to be treated on the same footing. In MADANALYSIS 5, this can be performed in a very natural way. The event files to be merged are gathered under the form of an object dubbed a *dataset*. When the analysis is effectively executed by the C++ core of the program, datasets, *i.e.*, collections of event files, are processed rather than the event files individually.

### The particle content of the events - particles and multiparticles.

According to the conventions of the different event formats introduced above, the particles described in the events are defined, in an unambiguous way, through their Particle Data Group identifier (PDG-id) [59]. Similarly, the algorithms generated by MADANALYSIS 5 are widely based on this concept to distinguish the various particle species. The drawback is obviously that this identifier is most of the time non-intuitive for the user and can even become unreadable when the set of particles included in the events becomes large. For instance, the particle content of hadron-level events consists in

hundreds of different particles, each of them being identified by a different PDG-id.

In the spirit of the MADGRAPH program, MADANALYSIS 5 alleviates this issue by allowing us to associate a *particle label* to a given PDG-id. Hence, instead of representing an electron and a positron by the integer numbers 11 and -11, one can create the (more intuitive) labels **e-** and **e+** and associate them to the PDG-ids 11 and -11, respectively. It is also possible to collect labels together through the concept of *multiparticles*. Hence, one could define a label **e** referring to both the electron and the positron.

In order to implement an analysis in an efficient way, it is recommended to the user to define, in a first step, a series of particle and multiparticle labels facilitating the readability (and then the validation) of his analysis.

### **Definition of the analysis - selections.**

Implementing the analysis is the core task of the user. It consists in defining the histograms that have to be generated and the selection cuts that need to be applied. These two types of objects, *i.e.*, histograms and cuts, are uniquely dubbed *selections*. It can be noted that even if asking for the generation of several histograms is a commutative operation, the ordering of the selection cuts is in contrast important. Applying the cuts in a different order indeed leads to the production of different (intermediate) histograms and efficiency tables.

### **Running the analysis - jobs.**

After having created a series of dataset objects and defined the analysis to be performed, *i.e.*, having implemented the histograms and selection cuts of interest, the user needs to execute the analysis on the datasets. Contrary to the previous steps which rely on the PYTHON module of MADANALYSIS 5, this task is built, for efficiency reasons, upon a C++ code which is generated by MADANALYSIS 5 and denoted as a *job*. After MADANALYSIS 5 creates the job, it has to be executed by the SAMPLEANALYZER kernel for each of the predefined datasets. Once the analysis has been performed, the results are subsequently imported in the PYTHON module and re-interpreted by MADANALYSIS 5.

### **Display of the results - reports.**

The generated results can be collected and displayed in a synthetic *report*. This report is given either under the HTML format, as a webpage, or as a

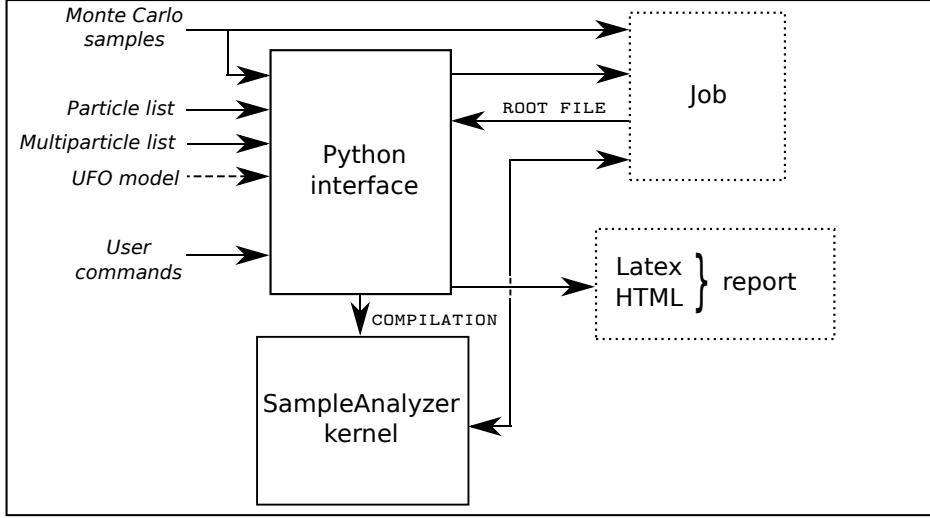


Figure 1: The MADANALYSIS 5 flowchart.

LATEX document that has to be further compiled. The report contains the histograms which have been generated and the efficiency tables related to the implemented selection cuts. Furthermore, a signal over background table is generated provided the datasets defined by the user have been tagged as signal and background samples.

### 2.3. Logical architecture of the program

The MADANALYSIS 5 program has two components, a PYTHON command line interface and a C++ kernel dubbed SAMPLEANALYZER. In the normal mode of running of the program, these two modules interact at different stages, as illustrated in Figure 1. The way to perform physics analyses in this mode is presented in detail in Section 3 and Section 4. For more sophisticated analyses, users with advanced skills in programming can bypass the PYTHON interface and directly implement their analysis within the SAMPLEANALYZER framework. This expert mode of running the code is described in detail in Section 5.

The PYTHON command line interface of MADANALYSIS 5 consists in a command prompt where the user accesses all the functionalities of the program through a set of commands. This allows to implement an analysis in a very user-friendly way. Each command entered by the user is first checked from the point of view of the syntax and if necessary, an error message is

printed to the screen. The issued command is then stored in the memory of the computer.

The interface can import different types of external information. Hence, a predefined list of particle and multiparticle labels could be imported if given as a text file or as a UFO model [39]. In a similar way, the list of the Monte Carlo event files to be analyzed can be easily loaded in the current session of MADANALYSIS 5.

After the user has typed in all the commands defining his analysis (definition of the datasets, histograms and selection cuts), the PYTHON interface creates a C++ code using the SAMPLEANALYZER framework, the previously introduced MADANALYSIS 5 job. This C++ code comes with a `Makefile` and is kept available to the user for further modifications or improvements of the analysis without having to be regenerated. After the compilation and execution of the job, the PYTHON interface loads the results and uses the ROOT library of functions [57] to normalize and draw the histograms according to the requirements of the user. The HTML and LATEX reports are eventually generated.

The SAMPLEANALYZER C++ kernel of MADANALYSIS 5 is, strictly speaking, the part of the program dedicated to the analysis of the Monte Carlo event samples itself. It consists in a framework built upon an adaptive data format common for all types of Monte Carlo event samples and which contains a series of well-suited functionalities. In addition, it includes a reader compliant with the LHE, STDHEP, HEPMC and LHCO event formats and a library of specific functions facilitating particle physics analyses. Among the latter, one finds, *e.g.*, methods allowing for boosting four-momenta or testing if a particle is a final-state particle or not. Let us finally note that the storage of information within the context of the SAMPLEANALYZER platform is widely based on the functions implemented within the ROOT library.

### 3. First steps with MADANALYSIS 5

In this Section, we present the philosophy, main features and the user-friendliness of MADANALYSIS 5 by means of a simple example. In contrast, the full list of capabilities of MADANALYSIS 5, which are much broader than what is shown in this Section, are described in Section 4.

For the sake of the illustration, we decide to perform a toy analysis at the parton-level. We consider several samples of 1000 events each describing various Standard Model hard-scattering processes at the LHC running at a

center-of-mass energy of 7 TeV. Events are generated with the Monte Carlo generator MADGRAPH 5 [11]. Neglecting all quark masses (but the top mass), we employ the leading order set of the CTEQ6 parton density fit [60] and identify both the renormalization and factorization scales as the transverse mass of the produced particles.

We consider four different event samples describing three different physical processes. They are stored into the directory `samples` of MADANALYSIS 5. If this directory is not present on the system of the user or if the event files are not there for any reason, one can type, once the command line interface of MADANALYSIS 5 has been started (see Section 3.1),

```
install samples
```

This leads to the creation of the directory `samples`, if relevant, and to the download from the Internet of the four samples

```
ttbar_sl_1.lhe.gz
ttbar_sl_2.lhe.gz
ttbar_fh.lhe.gz
zz.lhe.gz
```

The first two samples are related to the production of a semi-leptonically decaying top-antitop pair where the lepton is an electron or a muon, and the third one is related to the production of a fully hadronically decaying top-antitop pair. The last sample describes the production of a  $Z$ -boson pair, including also diagrams with virtual photons, where each of the bosons is decaying either to an electron pair or to a muon pair. At the time of event generation, we demand that the produced parton-level jets have a transverse momentum  $p_T > 20$  GeV, a pseudorapidity  $|\eta| < 2.5$  and a relative distance  $\Delta R > 0.4$ . Leptons are required to have a pseudorapidity  $|\eta| < 2.5$  and we ask the invariant mass of a pair of two leptons of the same flavor to be higher than 20 GeV.

### 3.1. Starting the command interface

Once downloaded from the web and unpacked, the MADANALYSIS 5 package does not require any compilation<sup>2</sup> or configuration and its command in-

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<sup>2</sup>The C++ core of MADANALYSIS 5 has in fact to be compiled but this task is performed automatically, behind the scenes, without requiring any interaction from the user.

terface, consisting in a command prompt `ma5>`, can immediately be launched by issuing

```
bin/ma5
```

from the directory where MADANALYSIS 5 has been installed. For the installation procedure of the program and all its dependencies, we refer to Appendix A. The user is now able to access all the functionalities of the tool and can start implementing an analysis.

When launched, the program firstly checks that all the required dependencies, such as the ROOT header files and libraries and a C++ compiler, are present on the system and correctly installed. If not, an error message is printed, so that the user has enough information to solve the issue, and the program exits. On its first run, MADANALYSIS 5 also compiles a static library which is stored into the directory `lib` of the distribution. This library is further used by the SAMPLEANALYZER kernel when analyses are executed.

Secondly, two lists of labels corresponding to standard definitions of particles and multiparticles are loaded into the memory of the current session of the program. By default, when MADANALYSIS 5 is used as a standalone package, they are imported from the corresponding files stored in the `input` directory of the distribution of the program. In contrast, if MADANALYSIS 5 has been installed in the directory where MADGRAPH 5 has been unpacked, the lists of particle and multiparticle labels are directly imported from MADGRAPH 5. Let us note that several multiparticles, such as `hadronic` or `invisible`, are essential for a correct running of MADANALYSIS 5. Therefore, if not included in the imported files, they are automatically created.

### 3.2. Particles and multiparticles

As soon as all the particle and multiparticle labels have been imported, MADANALYSIS 5 creates links pointing to lists containing all the declared labels. This allows us to access them in an easy way at any time of the analysis, by issuing in the command interface

```
display_particles  
display_multiparticles
```

New labels can be created on run-time with the command `define`, as, *e.g.*,

```
define mu = mu+ mu-
```

This command creates a multiparticle label `mu` which is associated to both the muon and the antimuon. Moreover, new labels are automatically added to the relevant list of (multi)particles.

The definition of a specific label can be retrieved with the help of the `display` command which outputs the PDG-id of the associated particle(s). For example, after invoking the two commands

```
display b  
display l+
```

the command interface outputs information about the (multi)particle labels `b` and `l+`,

```
The particle 'b' is defined by the PDG-id 5.  
The multiparticle 'l+' is defined by the PDG-ids -11 -13.
```

As can be seen from the screen output above, the two symbols `b` and `l+` define a  $b$ -quark and a positively charged lepton different from a tau, respectively.

### 3.3. Importing event samples

Preliminary to performing any analysis, event files under consideration must be parsed and loaded into the memory. The command `import` has been designed for that purpose. As a mandatory (single) argument, it requires the name of the Monte Carlo sample(s) to be parsed. In order to import several files at one time, the wildcard characters `*` and `?` are allowed when typing-in the argument of the function. Hence, the four samples introduced above can be loaded simultaneously by issuing the command

```
import samples/*.lhe.gz
```

which is equivalent to the set of four commands

```
import samples/ttbar_sl_1.lhe.gz  
import samples/ttbar_sl_2.lhe.gz  
import samples/ttbar_fh.lhe.gz  
import samples/zz.lhe.gz
```

The result is the creation of a unique event sample denoted by `defaultset` containing all the imported files. We are now ready to define selection cuts which the `SAMPLEANALYZER` kernel will apply to the events included in the dataset `defaultset`. The name as well as the way how to merge the samples can be tuned according to the needs of the user, but this goes beyond the scope of this Section and will be addressed in Section 4.

### 3.4. Selection cuts and creation of histograms

Creating a histogram representing a specific kinematical distribution has been made very efficient in the framework of **MADANALYSIS 5** through the command **plot**. This command requires one mandatory argument, the observable to be computed, and a set of optional arguments containing, among others, the number of bins of the histogram to be created and the lower and upper bounds of its  $x$ -axis. In the setup of the bounds of a histogram, one must note that the standard unit of energy used in **MADANALYSIS 5** is the GeV.

In the toy analysis which is implemented in the following, we focus on the missing transverse-energy distribution as well as on the transverse-momentum distribution of the final state muons. We recall that we are considering a unique event sample resulting from the merging of three  $t\bar{t}$  and one diboson event files, as explained in Section 3.3. In order to create the associated histograms, it is sufficient to issue the two commands

```
plot MET
plot PT(mu) 20 0 100
```

where we recall that the multiparticle **mu** has been defined in Section 3.2 and represents both the muon and the antimuon. The symbol **MET** is associated to the missing transverse energy whilst the function **PT** stands for the transverse momentum of a given (multi)particle provided as its argument. The next pieces of information to be passed to the command **plot** are optional and related to the binning of the histograms. By default, *i.e.*, in the case the binning information is not specified as in the first example above, **MADANALYSIS 5** uses hard-coded values which depend on the observable under consideration. In contrast, as in the second example above, the user can provide, at the time of typing-in the command, the number of bins (20 here) together with the values of the lowest and highest bins (which are chosen equal to 0 GeV and 100 GeV, respectively, in our example).

We now turn to illustrating the implementation of the two types of selection cuts which is possible to employ in **MADANALYSIS 5**. These cuts will be further applied, by the **SAMPLEANALYZER** kernel, to the events contained in the **defaultset** dataset. We recall that the program contains a set of possibilities much broader than what is presented in this Section and we refer to Section 4 for more information. In a first step, we decide to select events where the missing transverse energy ( $\cancel{E}_T$ ) is not too large, *i.e.*,  $\cancel{E}_T < 100$

GeV. Rejecting events not fulfilling this criterion can be very efficiently and easily performed by issuing the command

```
reject MET > 100
```

The function `reject` tells MADANALYSIS 5 to remove from the event selection any event where the missing energy (`MET`) is larger than 100 GeV.

As a second illustration about the implementation of selection cuts in MADANALYSIS 5, we focus on the kinematical properties of the muons contained in the selected events. In particular, a part of the events contain a muon–antimuon pair and we would like to represent the corresponding invariant-mass distribution through a histogram. However, many events do not exactly contain one muon and one antimuon. This is taken into account by MADANALYSIS 5 at the time of the generation of the histogram, where one entry is included for each different muon–antimuon pair which can be formed from the event particle content. For a given event, the number of entries can hence be zero, one or bigger than one according to the (anti)muon multiplicity of the final state.

Realistic detectors are in general not capable of correctly reconstructing too soft particles. Even at the parton-level, this effect can (and should) be implemented. In our case, we could consider as (anti)muon candidate only final state (anti)muons with, *e.g.*, a transverse momentum  $p_T > 20$  GeV. This cut can be implemented in MADANALYSIS 5 through the command `reject`, but following a syntax different from above,

```
reject (mu) PT < 20
```

The effect of this command is to consider, for each of the selected events, as muon and antimuon candidates only muons and antimuons with a transverse momentum (`PT`) harder than 20 GeV. For all the histograms which are generated after having typed the line above in the command interface, only the selected (anti)muon candidate's are considered. As stated above, we choose, as an example, to represent the muon-pair invariant mass distribution computed from the analyzed event samples included in the dataset `defaultset`. The command to create this histogram is similar to those presented in the beginning of this Section,

```
plot M(mu+ mu-) 20 0 100
```

where the arguments of the function  $M$ , associated to the invariant mass, are multiple. This denotes that we are summing the four-momenta of the muon and the antimuon to derive the invariant mass to be represented. As it can be seen from the optional arguments of the command `plot`, we once again ask for a histogram of 20 bins with its lower and upper bounds being fixed to 0 GeV and 100 GeV, respectively. We recall that for a specific event, the number of entries which are included in the histogram corresponds to the number of different muon–antimuon pairs that can be formed from the final-state particle content. This can then be any integer number.

Once the selection is defined, it has to be executed through a job to be run by `SAMPLEANALYZER`. This is done through the command `submit` which takes as an argument the name of a directory which will be created,

```
submit <dirname>
```

The created directory `<dirname>` contains a series of C++ source and header files that are necessary for `SAMPLEANALYZER` to properly run. The compilation, linking to the external static library of `MADANALYSIS 5` (see Section 3.1) and the execution of the resulting code is handled by `MADANALYSIS 5`. The screen output indicates the status of these different tasks and various information such as the detected format of the event samples or the number of processed events. In the case anything is not going as smoothly as it should, `MADANALYSIS 5` also prints warning and/or error messages to the user and the program exits in the worst case scenario.

### 3.5. Displaying the results

After having performed the analysis as indicated in the previous Section, `MADANALYSIS 5` offers several ways to present the results under a human-readable report. This report can be either generated under the HTML format or under a `TEX` format that could be compiled with a `LATEX` or a `PDFLATEX` compiler. The histograms are saved under a Portable Network Graphics file (`.png`) for HTML and `PDFLATEX` reports, and under an Encapsulated PostScript file (`.eps`) for `LATEX` reports. The `MADANALYSIS 5` commands generating the reports read

```
generate_latex <dirname>
generate_pdflatex <dirname>
generate_html <dirname>
```

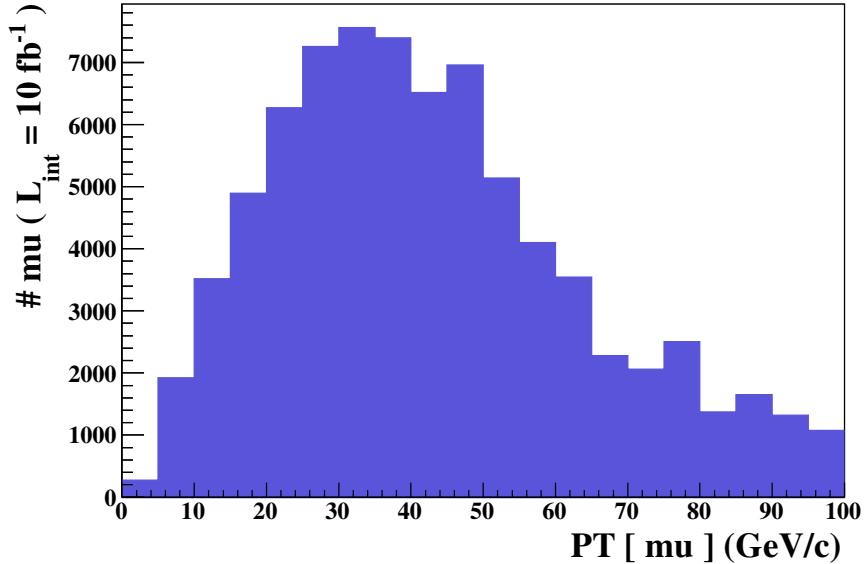


Figure 2: Transverse-momentum distribution of the muons for a dataset consisting of the four event samples introduced in Section 3.

where `<dirname>` stands for the directory in which the report is generated.

The structure of the report is similar whatever the adopted format is. It starts with a table of contents and firstly displays all the information necessary to ensure the reproducibility of the analysis. One hence finds the list of commands which have been issued, in the current session of the program, before the generation of the report, followed by the employed version of the code as well as the values of all the setup parameters. In this last category, one has, *e.g.*, the list of the event samples which have been analyzed, given together with the corresponding cross sections and the number of events contained in the event files.

The core of the report contains the results of all the commands related to histogram creation and to the application of a selection cut. These commands have been treated one by one by `SAMPLEANALYZER` and the report follows this pattern. By default, histograms are normalized to an integrated luminosity of  $10 \text{ fb}^{-1}$  and comes together with a summary table containing information on the mean value of the computed distribution and the associated root mean square (RMS), as well as on the presence of possible

Dataset	Integral	Entries / event	Mean	RMS	Underflow	Overflow
defaultset	82747	0.752	42.8177	21.36	0.0	6.181

Table 1: Statistics associated to the histogram of Figure 2.

Cuts	Signal (S)	Background (B)	S vs B
Initial (no cut)	109999 +/- 789		
cut 1	82994 +/- 375443		
cut 2	82994 +/- 612		

Table 2: Efficiencies of the cuts applied in Section 3.4.

underflow and overflow bins in the histogram. The latter are given as a ratio of the value of the integral of the represented distribution between its lower and upper bounds. A green–orange–red color code indicates their relative importance, an orange or red color suggesting more or less strongly to the user to modify the bounds of the histograms, if relevant. We refer to Section 4 for the description of all the properties of the histograms that can be modified by the user, such as the way to modify the luminosity or the binning of a given histogram.

In Figure 2, we take the example of the transverse-momentum distribution of the muons implemented in Section 3.4 and present the histogram generated by MADANALYSIS 5. As stated above, if several muons are contained in one specific event, they each correspond to a different entry in the histogram with the same weight. The summary table generated together with the histograms is given in Table 1. In the first column of the table, the name of the dataset is printed. In the second column, one finds the number of events normalized to an integrated luminosity of  $10 \text{ fb}^{-1}$ , since this is the default value which has not been modified. In the third column, the average number of histogram entries for each event is provided, with the mean and the root mean square of the distribution under consideration shown in the fourth and fifth columns. From the orange cells in the table, one immediately notices that the number of underflow and overflow entries given in the last two columns are only under a roughly reasonable control.

Let us finally note that for each selection cut, a summary table of the cut efficiency is also given, together with the corresponding effect on the signal over background ratio. In our toy example, we have not tagged any sample

as signal or background. Therefore, this feature is irrelevant here and we refer to Section 4 for more information. The generated summary of all cuts is provided in Table 2, where all the events are (by default) considered as signal events.

## 4. Implementing analyses in an efficient and user-friendly way

### 4.1. Starting a MADANALYSIS 5 session

The MADANALYSIS 5 package is able to handle several types of Monte Carlo samples. Supported event files can describe partonic, hadronic or reconstructed events. The key difference between these three types of events lies in the basic objects in which the event content is expressed in terms of. In the first case, an event consists in a set of partons (quark, gluons, charged leptons, neutrinos, new states, *etc*), whilst in the second case, it consists in hadrons (pions, kaons, baryons, *etc*). Finally, a reconstructed event contains a set of high-level reconstructed objects such as electrons, muons, jets, *etc*. For each of these classes of events, MADANALYSIS has a specific running mode which the user is required to specify when launching the code,

```
bin/ma5 [level]
```

Typing in a shell `bin/ma5 --partonlevel` or `bin/ma5 -P` allows us to run MADANALYSIS 5 in the mode required to analyze parton-level events, whilst issuing `bin/ma5 --hadronlevel` or `bin/ma5 -H` makes the program running in the hadron-level mode. In a similar fashion, the two shell commands `bin/ma5 --recolevel` and `bin/ma5 -R` allow us to start MADANALYSIS 5 in a ready way to analyze reconstructed-level events. In the case no argument is provided, as in the example of Section 3, the parton-level mode is automatically selected. Of course, the different modes cannot be combined, since the levels of sophistication are self-excluding. Once one of the previously introduced commands is issued, the command line interface of MADANALYSIS 5 is started. It consists in a command prompt `ma5>` where the user can access all the functionalities of the program and directly implement a physics analysis by issuing a set of commands. We refer to Section 4.2 for more information about the whole set of existing commands.

The list of commands to be typed can also be provided under the form of a script, *i.e.*, a simple text file. In order to execute the script, its path has to be provided as the last argument when typing the `bin/ma5` command in a shell. Hence,

```
bin/ma5 --recolevel <filename>
```

starts the command interface of MADANALYSIS 5 in the reconstructed-level mode. All the commands included in the file `<filename>` are then sequentially applied, one after the other. When having such a script executed by MADANALYSIS 5, it can be useful to bypass all the confirmation questions usually asked by the program. This can be done by issuing in a shell one of the two commands

```
bin/ma5 --script --recolevel <filename>
bin/ma5 -s --recolevel <filename>
```

This *script mode* also enforces the automatic exit of MADANALYSIS 5 after the completion of the script `<filename>`, in contrast to the *forced mode* described below. It is also possible to execute several scripts by providing various filenames,

```
bin/ma5 -s -R <filename1> <filename2> <filename3>
```

The different files are handled as concatenated in one single file, *i.e.*, all the commands included in `filename1` are processed sequentially, followed by the commands included in `filename2`, *etc.* Let us note that the character ‘#’ can be included when writing down script files. Everything standing to the right of this character is considered as a comment by MADANALYSIS 5 and ignored by the command line interpreter.

For highly sophisticated analyses, the features which are presented in the rest of this section might not be sufficient. Therefore, MADANALYSIS 5 can be run in an expert mode by issuing one of the three commands

```
bin/ma5 --expert      bin/ma5 -e      bin/ma5 -E
```

The possibilities to implement any analysis are here unlimited. More information about the expert mode is provided in Section 5. Let us however note that in this case, any other option which could be passed when starting the interface is ignored.

Among the other options which could be provided when starting the interpreter, one can note that the version of the distribution installed on the system of the user can be accessed through one of the following commands,

```
bin/ma5 -v      bin/ma5 --version      bin/ma5 --release
```

**Table 3: Syntax to launch the MADANALYSIS 5 program:**

`bin/ma5 -h or bin/ma5 --help`

This allows to display to in a shell a summary of the content of this Table.

`bin/ma5 -v , bin/ma5 --version or bin/ma5 --release`

This allows to display in a shell the number of the version of MADANALYSIS 5 present on the system.

`bin/ma5 -f , bin/ma5 --forced`

This allows to run MADANALYSIS 5 in a mode where confirmation questions are ignored.

`bin/ma5 -e , bin/ma5 -E or bin/ma5 --expert`

This indicates that MADANALYSIS 5 has to run in expert mode. In this case, any other option is ignored.

`bin/ma5 [level] [files]`

`[level]`

When provided, this optional argument allows to select the type of event files to analyze. If absent, parton-level events are assumed. The user can choose one of the following options:

- P or --partonlevel      Parton-level events.
- H or --hadronlevel      Hadron-level events.
- R or --recolevel      Reconstructed events.

`[files]`

When provided, this optional argument consists in a sequence of file-names, separated by blank characters, containing each a set of MADANALYSIS 5 commands. The files are handled as concatenated and the commands are applied sequentially. If the option pattern `-s` or `--script` is included, MADANALYSIS 5 exits after having executed the script and does not ask any confirmation question.

and that typing-in

```
bin/ma5 -f      bin/ma5 --forced
```

ensures a running mode of MADANALYSIS 5 where confirmation questions, such as, *e.g.*, those printed to the screen when a directory is about to be removed, are not asked of the user.

The different running modes of MADANALYSIS 5 and the way to cast them are summarized in Table 3. They can also be displayed by the program when the user types in a shell one of the following commands,

```
bin/ma5 --help      bin/ma5 -h
```

#### 4.2. The command line user interface of MADANALYSIS 5

The command line interface of MADANALYSIS 5 is built upon the PYTHON module `cmd` which allows for a flexible processing of commands issued by the user. It features, among others, text help, tab completion and shell access. In addition, the history of commands issued by the user can be obtained by means of the up and down keys of the keyboard. As presented in Section 4.1, the command interface can be started by issuing in a shell the command

```
bin/ma5 [options]
```

from the directory where MADANALYSIS 5 has been installed. However, MADANALYSIS 5 could also be started from any location on the computer of the user, the command above needing only to be modified accordingly.

The user interface consists in a command prompt `ma5>` where the user can implement a physics analysis very efficiently and easily by means of a series of (case sensitive) commands. The related syntax is inspired by the PYTHON programming language and a command therefore always starts with an *action*. Several actions can be typed together in a single command line after separating them with a semicolon ‘;’. The number of different implemented actions has been made as small as possible in order to guarantee a quick handling of the code by any physicist. Their list has been collected in Table 4, with basic instructions about how to use each of these actions. The information included in this table can also be displayed to the screen at run-time by issuing

```
help [<action>]
```

where `<action>` is the action under consideration. Moreover, if `help` is typed-in without any argument, the list of all the available actions is printed out. Let us emphasize that tab completion could also be used with the aim of obtaining that list. We now dedicate the rest of this Section 4 to a detailed description of each of those actions.

As stated above, the history of all the actions undertaken by the user can be accessed through the up and down keys of the keyboard, as for standard shells. They are read from the file `.ma5history` stored in the directory where MADANALYSIS 5 has been unpacked. This file is updated every time that the user types a command, and the list of the last 100 commands executed by the interpreter is saved. In addition, if one restricts ourselves to the current session of MADANALYSIS 5, the list of typed commands can be accessed by issuing in the command interface

### `history`

Besides actions, the language handled by the interpreter contains *objects* that actions are acting on. There exist various types of objects representing particles, multiparticles, datasets, selection cuts or even the configuration panel of the current session of the program, denoted by `main`. In addition, any object is described by several attributes related to its properties, which are generically denoted as *options*. These attributes can be accessed through the syntax `object.option`.

Finally, the language of MADANALYSIS 5 contains a set of reserved keywords, such as `all`, `or` or `as`, whose usage is described in the following. Moreover, it is important to keep in mind that both predefined and user-defined objects have a name which has to be unique. Therefore, in an analysis, any object which is created, such as, *e.g.*, a new instance of the multiparticle class, must have a name different from those of the existing actions (see Table 4) or those of the objects already defined and/or used in the current analysis.

The syntax introduced above has been developed with a focus on uniformization and intuition. For the sake of the example, we now focus on the three commands `display`, `set` and `remove`. To display to the screen an object denoted by `<object>` together with its properties, it is enough to type the command

### `display <object>`

Similarly, the display of a specific property, denoted by `<option>`, of the object under consideration can be performed by issuing

```
display <object>.<option>
```

On the same footing, the value of the attribute `<object>.<option>` can be easily modified via the action `set`

```
set <object>.<option> = x
```

where in the example above, the option under consideration is set to the value `x`. Finally, an object can be deleted from the computer memory with the help of the command `remove`,

```
remove <object>
```

However, three objects, the objects `main`, `hadronic` and `invisible`, as well as all the particles and multiparticles currently used in the analysis, *i.e.*, being related to a histogram or to a selection cut, cannot be deleted.

As already stated in the beginning of this Section, the `cmd` module of PYTHON allows for efficient access to shell commands directly from inside the command line interface of MADANALYSIS 5. This requires starting the command line either with an exclamation mark ‘!’ or with the command `shell`. Therefore, the syntax

```
!<command>
```

or equivalently

```
shell <command>
```

has to be followed in order to execute the command `<command>` from an external shell.

Finally, the configuration of the command line interpreter can be restored as for a new session of MADANALYSIS 5 by issuing the command

```
reset
```

Consequently, all the user-defined particles and multiparticles as well as all the implemented histograms and selection cuts are removed from the computer memory.

**Table 4: Actions available from the command line user interface of MADANALYSIS 5**

define <name> = <def>	This creates a new (multi)particle label <name> defined through the (list of) PDG-id(s) <def>.
display <var>	This displays to the screen either the properties of an object or the value of one of its options, both generically denoted by <var>.
display_datasets	This displays to the screen the list of the imported datasets.
display_multiparticles	This displays to the screen the list of defined multiparticle labels.
display_particles	This displays to the screen the list of defined particle labels.
EOF	This closes the current session of MADANALYSIS 5.
exit	This closes the current session of MADANALYSIS 5.
generate_html <dir>	The report of the current analysis is generated, under the HTML format, and stored in the directory <dir>.
generate_latex <dir>	The report of the current analysis is generated, under the $\text{\TeX}$ format, and stored in the directory <dir>. Compilation requires a LATEX compiler.
generate_pdflatex <dir>	The report of the current analysis is generated, under the $\text{\TeX}$ format, and stored in the directory <dir>. Compilation requires a PDFLATEX compiler.
help	This displays to the screen the list of all the actions implemented in MADANALYSIS 5.
history	This displays to the screen the history of the commands typed in the current session of the program.

**Table 4 (continued): Actions available from the command line user interface of MADANALYSIS 5**

<code>import &lt;obj&gt;</code>	This allows to import in the current session of MADANALYSIS 5 external information, generically denoted by <code>&lt;obj&gt;</code> , such as Monte Carlo samples, a configuration used by SAMPLEANALYZER in a previous analysis or a UFO model.
<code>install &lt;obj&gt;</code>	This allows to install the external object <code>&lt;obj&gt;</code> from the Internet. Currently, the only allowed choice for the value of the variable <code>&lt;obj&gt;</code> consists in the keyword <code>samples</code> which yields the download of the four example samples presented in Section 3.
<code>plot &lt;obs&gt; [options]</code>	This creates, in an analysis, an histogram associated to the distribution of the observable <code>&lt;obs&gt;</code> . We refer to Section 4.5 for the list of available options.
<code>open &lt;rep&gt;</code>	This opens the report <code>&lt;rep&gt;</code> of a given analysis, with the appropriate program.
<code>preview &lt;hist&gt;</code>	This displays an histogram <code>&lt;hist&gt;</code> in a graphical window.
<code>quit</code>	This closes the current session of MADANALYSIS 5.
<code>reject (&lt;prt&gt;) &lt;cond&gt;</code>	This adds a selection cut. In an analysis, a candidate to a given particle species <code>&lt;prt&gt;</code> is ignored if the condition <code>&lt;cond&gt;</code> is fulfilled.
<code>reject &lt;cond&gt;</code>	This adds a selection cut. An event is rejected if the condition <code>&lt;cond&gt;</code> is fulfilled.
<code>remove &lt;obj&gt;</code>	This removes the object <code>&lt;obj&gt;</code> from the memory.

**Table 4 (continued): Actions available from the command line user interface of MADANALYSIS 5**

<code>reset</code>	This reinitializes MADANALYSIS 5 as when a new session starts.
<code>resubmit</code>	This allows, for an analysis which has already been run by SAMPLEANALYZER and further modified, to be re-executed.
<code>select (&lt;prt&gt;) &lt;cond&gt;</code>	This adds a selection cut. In order to consider, in an analysis, a candidate to a given particle species <code>&lt;prt&gt;</code> , the condition <code>&lt;cond&gt;</code> has to be fulfilled.
<code>select &lt;cond&gt;</code>	This adds a selection cut. Events are selected only if the condition <code>&lt;cond&gt;</code> is fulfilled.
<code>set &lt;obj&gt;.&lt;opt&gt; = &lt;val&gt;</code>	This allows to set the attribute <code>&lt;opt&gt;</code> of the object <code>&lt;obj&gt;</code> to the value <code>&lt;val&gt;</code> .
<code>shell &lt;com&gt;</code>	This allows to run the shell command <code>&lt;com&gt;</code> from the command interface of MADANALYSIS 5. The action <code>shell</code> can be replaced by an exclamation mark ‘!’.
<code>submit &lt;dir&gt;</code>	This allows to execute an analysis as a job run by SAMPLEANALYZER. The C++ source and header files related to the job are created in the directory <code>&lt;dir&gt;</code> .
<code>swap &lt;sel1&gt; &lt;sel2&gt;</code>	This allows to permute the sequence of two instances of the class <code>selection</code> , <code>&lt;sel1&gt;</code> and <code>&lt;sel2&gt;</code> . We refer to Section 4.5 for more information.

Before moving on, some remarks on tab completion are in order. This feature allows for an easy typing of commands, attributes of the objects, *etc*. Typing on the ‘`tab`’ key has the effect of printing to the screen all the allowed possibilities for completing the command about to be written. It works equivalently for actions, objects and attributes. Let us emphasize that this makes tab completion in MADANALYSIS 5 much more advanced than its counterpart in standard command shells.

#### 4.3. Datasets

In Section 3, we have shown that four example Monte Carlo samples can be downloaded from the Internet by employing the command `install`,

```
install samples
```

and stored into a directory `samples` of the current distribution of MADANALYSIS 5. Actions, such as their gathering into datasets, the creation of histograms and the definition of selection cuts, have been subsequently executed on those samples. This Section is dedicated to the dataset class of objects, whilst Section 4.5 and Section 4.6 concerns histograms and cuts, respectively.

In order to load Monte Carlo samples into the computer memory, they have to be imported from outside the current session of MADANALYSIS 5. This requires the use of the command `import`, as briefly presented in both Section 3 and Table 4,

```
import <path-to-sample>
```

The syntax above allows us to import a Monte Carlo sample stored at a location `<path-to-sample>` on the computer of the user. Several samples can be imported at one time by employing the wildcard characters `*` and `?`. As a generic example, the command

```
import <directory>/*
```

imports all the samples stored in the directory `<directory>`. Moreover, the tilde character `~` can be used when typing the path to a sample. As in standard LINUX shells, it points to the location of the home directory of the user.

According to their event format, MADANALYSIS 5 handles differently the imported event samples. The key to the procedure lies in the extension of

the event files, up to a possible packing with GZIP. The formats currently supported are the LHE event file format, whose corresponding file extensions are `.lhe` or `.lhe.gz`, the STDHEP event file format, whose corresponding file extensions are `.hep` or `.hep.gz`, the HEPMC event file format, whose corresponding defining extensions are `.hepmc` or `.hepmc.gz` and the LHCO event file format, whose corresponding event file extensions are `.lhco` or `.lhco.gz`.

Of course, all the different formats are not appropriate for any level of sophistication of the analysis. For instance, using the parton-level mode of MADANALYSIS 5 to analyze an event file compliant with the LHCO format is clearly inadequate. In more detail, events stored in files compliant with the LHE format can be imported whatever is the level of sophistication of the analysis, *i.e.*, equally for parton-level, hadron-level or reconstructed-level analyses. A remark is in order here. In principle, the LHE format is not supposed to describe reconstructed events. However, there exist routines external to MADANALYSIS 5, such as the HEP2LHE converter included in MADGRAPH, which allow us to efficiently convert events as produced by hadronization algorithms or a fast detector simulation tool into reconstructed events. Those routines use a jet algorithm in order to reconstruct light and  $b$ -tagged jets, and the total missing transverse energy of the event is eventually computed. In contrast, STDHEP and HEPMC event files can only store parton-level and hadron-level events, whilst complementary LHCO files can only be used after high-level objects have been reconstructed, *i.e.* at the reconstructed-level after detector simulation.

The effect of the action `import` is to unify all the imported samples into one single object dubbed *dataset*. If nothing is specified by the user, events are gathered together into a dataset called `defaultset`. The possibility to import events and collect them into different datasets allows us, for instance, to differentiate background event samples from signal event samples, as with the following commands

```
import <path-to-signal-events>/* as signalset
import <path-to-background-events>/* as backgroundset
```

The first command imports all the event samples present in the directory `<path-to-signal-events>` and collects them into one single dataset labeled as `signalset`. Similarly, the second command loads into the current session of MADANALYSIS 5 all the event files located in the directory `<path-to-background-events>` and gathers them together into a dataset

labeled as `backgroundset`. On the same footing, this feature can be used to collect efficiently event files corresponding to different sources of background,

```
import <path-to-events>/ttbar* as ttbar
import <path-to-events>/zz* as zz
import <path-to-events>/drellyan* as drellyan
```

These three commands import into MADANALYSIS 5 three series of samples, related respectively to top-antitop pair,  $Z$ -boson pair and Drell-Yan events. Three different datasets, `ttbar`, `zz` and `drellyan`, are created so that these events can be treated separately when performing the analysis.

A dataset has several properties which are implemented as options of the `dataset` class and whose value can be modified with the command `set`. These attributes can be grouped into three categories of properties which are summarized in Table 5.

The first category of options consists in fact in one single property which is related to the attribute `type`. It allows us to tag the corresponding event sample(s) as a part of the signal or the background. Taking the example of a generic dataset `<dataset>`, the two commands

```
set <dataset>.type = background
set <dataset>.type = signal
```

tag it as a dataset belonging to the series of background or signal event samples, respectively. By default, a dataset is always considered as of the type `signal`. In an analysis, this distinction between signal and background is mandatory for a correct (automated) computation of the cut efficiencies by MADANALYSIS 5, as well as for the corresponding derivation of the signal over background ratio.

When creating histograms, their overall normalization is related to both the luminosity, which can be specified by the user (see Section 4.7) and the cross sections associated to the different datasets included in the histograms. These cross sections are included in LHE event files and are directly read and imported into the current session of the program, in the case LHE samples are analyzed. In contrast, the numerical value of the cross sections are absent from STDHEP, HEPMC and LHCO files. Therefore, in most of the cases, the user has to indicate the cross section manually. Moreover, in the case of LHE files, one could also want to modify the values which have been read. For instance, the cross section associated to an event sample which has been

generated with a leading order Monte Carlo tool could be modified by the user in order to account for next-to-leading-order normalization effects.

The second category of options related to the `dataset` class offers two ways to perform the above-mentioned task, either through the `weight` attribute or through the `xsection` attribute. A weight different from one can be assigned to each event of a dataset by modifying the value of the attribute `weight` of the `dataset` class,

```
set <dataset>.weight = <weight>
```

The command line above leads to the assignment of a weight `<weight>` to each event included in a generic dataset which has been defined as `<dataset>`. Consequently, the default value of one has been superseded by the value `<weight>`. In contrast, the user can also decide to leave the weight of each event unchanged and equal to unity, and modify instead the value of the cross section. The new value of the cross section has to be stored through the attribute `xsection` of the `dataset` class,

```
set <dataset>.xsection = <value>
```

If the value `<value>` is different from the default choice of zero, MADANALYSIS 5 uses it at the time of the creation of the histogram when calculating its normalization, ignoring the value of the cross section possibly included in the event file.

The last class of attributes associated to `dataset` objects concerns the layout of the histograms generated by MADANALYSIS 5, and in particular the style of the curves associated to the datasets which can be customized. On the one hand, styles and colors can be modified via the value of the attributes `linestyle`, `linewidth`, `linecolor`, `backstyle` and `backcolor`. On the other hand, the text used in histogram legends can be set by the user through the attribute `title`. The first three attributes above are associated to the type of lines (solid, dashed or dotted) used in the histograms, their width and their color (blue, green, none, purple, white, black, cyan, gray, orange, red or yellow). The two attributes `backstyle` and `backcolor` refer to the surface under a histogram and the style (solid, dotted, or hatched lines) and color (blue, green, none, purple, white, black, cyan, gray, orange, red or yellow) employed when it is drawn.

The default value for the two attributes related to colors in histograms, `linecolor` and `backcolor`, is `auto`. This means that MADANALYSIS 5 handles the color features automatically, assigning different colors to the datasets

**Table 5: List of the attributes of the dataset class**

```
set <dataset>.〈option〉 = <value>
```

<b>backcolor</b>	In an histogram, this changes the color filling the surface under the curve associated to the dataset under consideration. The allowed choices are <code>auto</code> (default), <code>blue</code> , <code>green</code> , <code>none</code> , <code>purple</code> , <code>white</code> , <code>black</code> , <code>cyan</code> , <code>gray</code> , <code>orange</code> , <code>red</code> and <code>yellow</code> . Integer numbers between one and four can be added or subtracted to these values (see Figure 3).
<b>backstyle</b>	In an histogram, this changes the style employed when filling the surface under the curve associated to the dataset under consideration. The allowed choices are <code>auto</code> , <code>solid</code> , <code>dotted</code> , <code>hline</code> , <code>dline</code> and <code>vline</code> (see Figure 4).
<b>linecolor</b>	In an histogram, this changes the color of the line of the curve associated to the dataset under consideration. For the allowed choices, we refer to the attribute <code>backcolor</code> .
<b>linestyle</b>	In an histogram, this changes the style of the line of the curve associated to the dataset under consideration. The allowed choices are <code>solid</code> , <code>dashed</code> , <code>dotted</code> and <code>dash-dotted</code> .
<b>linewidth</b>	In an histogram, this changes the width of the line of the curve associated to the dataset under consideration. It takes an integer value between one and ten.
<b>title</b>	This change the string used in histogram legends for the dataset under consideration. The possible choices are either <code>auto</code> (the name of the dataset) or any string under a valid T <sub>E</sub> X form.
<b>type</b>	This modifies the type of a dataset, associating it to the set of either signal ( <code>signal</code> ) or background ( <code>background</code> ) samples.
<b>weight</b>	This allows to change the weight of each event included in the dataset. The default value is one.
<b>xsection</b>	This allows to modify the total cross section associated to the events included in the dataset under consideration. The value can be any real number.

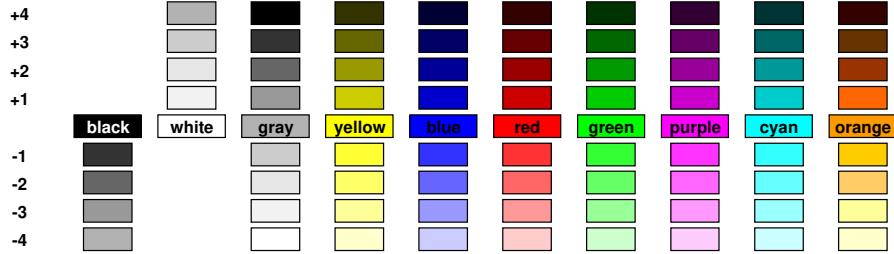


Figure 3: Complete list of allowed choices for the values taken by the attributes `backcolor` and `linecolor` of the `dataset` class.

created by the user. For a specific dataset object denoted by `<dataset>`, the values of the two color attributes can be, as usual, superseded by employing the command `set`,

```
set <dataset>.linecolor = <color>
set <dataset>.backcolor = <color>
```

The supported values for the variable `<color>` are intuitive and read `auto`, `blue`, `green`, `none`, `purple`, `white`, `black`, `cyan`, `gray`, `orange`, `red` or `yellow`. For histograms employing a large number of datasets, this panel of colors might however be insufficient. Shades of these basic colors can be used by adding or subtracting an integer number between one and four to those colors. For instance, the set of commands

```
set <dataset1>.backcolor = red+1
set <dataset2>.backcolor = red+2
set <dataset3>.backcolor = red+3
set <dataset4>.backcolor = red+4
```

allows us to assign four different shades of red to the four datasets `<dataset1>`, `<dataset2>`, `<dataset3>` and `<dataset4>`. Consequently, when a stacked histogram containing these four datasets is drawn, the associated surfaces will all be filled with different colors derived from the basic red. The complete list of colors integrated in MADANALYSIS 5 is illustrated in Figure 3.

The style of the lines of the curves drawn in the histograms can be modified through the attribute `linestyle` of the `dataset` class,

```
set <dataset>.linestyle = <value>
```

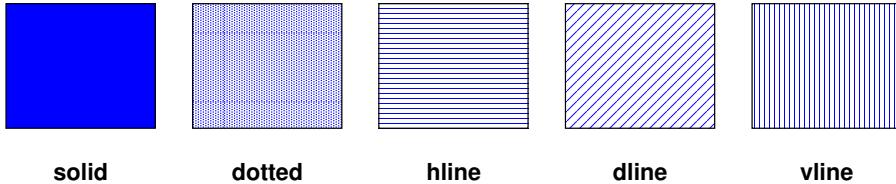


Figure 4: Complete list of styles allowed for the values possibly taken by the attribute `backstyle` of the `dataset` class.

This command changes the default employed solid style to the value `<value>`, which can be either `solid`, `dashed`, `dotted` or `dash-dotted`. Similarly, the attribute `linewidth` allows us to change the width of the drawn lines,

```
set <dataset>.linewidth = <value>
```

where `<value>` is an integer number between one and ten, the default value being unity. The option `backstyle` of the `dataset` class is linked to the style employed to fill the surface under a histogram and can be set to a new value by issuing in the command line interface

```
set <dataset>.backstyle = <value>
```

The allowed choices for the parameter `<value>` are either the default value `auto` or `solid`, `dotted`, `hline`, `dline` or `vline`, as presented in Figure 4. The last three choices, less intuitive, stand for horizontal, diagonal and vertical lines, respectively.

Finally, when creating a histogram where the curves related to several datasets are stacked or superimposed, MADANALYSIS 5 includes a legend with the explanation of the color/style code employed to distinguish the datasets. By default, the name of the dataset is used in this legend, but this can be modified by the user once setting the attribute `title` to a string consisting of a valid T<sub>E</sub>X expression,

```
set <dataset>.title = <string>
```

where `<string>` could stand, *e.g.*, in the case of a dataset describing events related to the production of a pair of  $W$ -bosons, for "`W^{+}` `W^{-}`" (including the quotation marks).

The effects of the customization of the histograms are only taken into account when the reports containing the results are generated. Therefore,

the user does not have to issue the command `submit` each time he modifies the style of a curve or how the area under a curve is filled inside a histogram.

Before moving on, let us recall that the command `display_datasets` introduced in Table 4 allows us to display to the screen the list of all the instances of the `dataset` class which have been created in the current session of MADANALYSIS 5. Furthermore, the action `display` can also be used on datasets. This leads to the printing to the screen of the current values of all the attributes of a given dataset. Hence, for a dataset labeled by `ttbar` where no attribute has been modified by the user, the effect of the command

```
display ttbar
```

is to print to the screen the following pieces of information,

```
Name of the dataset = ttbar (signal)
Title = 'ttbar'
User-imposed cross section = 0.0
User-imposed weight of the set = 1.0
Line color in histograms = auto
Line style in histograms = solid
Line width in histograms = 1
Background color in histograms = auto
Background style in histograms = solid
List of event files included in this dataset:
- ttbar.lhe.gz
```

This indeed summarizes the status of the instance of the `dataset` class labeled `ttbar` and displays the value of all the options.

Finally, a dataset can be deleted from the memory of the computer by employing the command `remove`. If several event files are included in a single dataset, it is however not possible to remove a particular file. In this case, the whole dataset has to be removed and then recreated without including this file.

#### 4.4. Particles and multiparticles

When starting a session of MADANALYSIS 5, lists of predefined particle and multiparticle labels are loaded into the memory, as it is mentioned in the example of Section 3. These lists are taken from the content of the directory `madanalysis/input` which contains several files with label definitions.

Not all the files present in this directory are loaded when a new session of MADANALYSIS 5 starts. Indeed, according to a running with the aim of analyzing parton-level, hadron-level or reconstructed-level events, only some of the files are appropriate, and only these are imported when MADANALYSIS 5 is launched.

For parton-level analyses, the particle labels are imported from the file `particles_name_default.txt` whilst the multiparticle labels are read from the file `multiparticles_default.txt`. These two files contains standard names for the Standard Model and the Minimal Supersymmetric Standard Model particles (*e.g.*, an antimuon is denoted by `mu+` and the lightest neutralino by `n1`) as well as appropriate definitions for the multiparticles `hadronic` and `invisible` (see below). For hadron-level analyses, only a single file with a list of 415 hadrons (see Ref. [59]), `hadron_default.txt`, is imported. Finally, for analyses at the reconstructed-level, the two lists `reco_default.txt` and `multiparticles_reco_default.txt` are imported. They contain labels for high-level reconstructed objects necessary for such a level of sophistication of the Monte Carlo simulations. The set of predefined labels is rather short and consists of `e+`, `e-`, `mu+`, `mu-`, `ta+`, `ta-` for charged leptons, and `j` and `b` and `nb` for jets, *b*-tagged jets and non-*b*-tagged jets. In addition, several (intuitive) multiparticle labels are included, `e`, `mu`, `ta`, `l+`, `l-` and `l` for various combinations of charged leptons, as well as two special labels dedicated to muons and antimuons, `mu-_isol` and `mu+_isol`, in case they are isolated.

Let us note that if MADANALYSIS 5 has been installed in the directory where MADGRAPH 5 is unpacked, the predefined particle and multiparticle labels are directly imported from MADGRAPH 5 rather than from the directory `madanalysis/input` of MADANALYSIS 5.

Two multiparticles, denoted by the labels `invisible` and `hadronic`, have a special role and are essential in any analysis. They are respectively related to the computation of observables related to the missing energy and to the hadronic activity. Therefore, if they are not included in the imported files, they are automatically created by MADANALYSIS 5 when a new session starts. Moreover, their special role prevents them from being deleted with the command `remove`.

Full support is also offered to load entire UFO models [39]. The UFO format is automatically detected by MADANALYSIS 5 so that, in order to load the model, it is enough to issue in the interpreter

```
import <path-to-UFO-files>
```

where `<path-to-UFO-files>` is the location of the UFO model under consideration. The list of particle labels is derived from the `particles.py` file containing the UFO implementation of the particles of the model. Moreover, any stable, electrically and color neutral particle, with the exception of the photon, is automatically added to the `invisible` multiparticle object.

As already presented in Section 3 and in Table 4, the command interpreter of MADANALYSIS 5 contains two actions for printing to the screen the list of all the particle and multiparticle labels which have been defined in the current session,

```
display_particles      display_multiparticles
```

Moreover, as for any instance of any class, the properties of a given particle or multiparticle object can be obtained with the action `display`,

```
display <label>
```

where `<label>` is the label of the (multi)particle under consideration. The effect of the command above is to print to the screen, in the case of a particle, the PDG-id linked to the label `<label>`. Similarly, for multiparticle labels, the whole list of associated PDG-ids is displayed.

As already shown in Section 3 where we have taken the example of muons and antimuons, new particle and multiparticle labels can be created, according to the needs of the user, by means of the command `define`,

```
define <label> = <identifiers>
```

The command above creates a label denoted by `<label>` and associates to it the content of the parameter `<identifiers>`. If the value of `<identifiers>` consists in one single integer number, a new particle label related to the corresponding PDG-id is created. In the case we have either several integer numbers separated by spaces or several other particle and multiparticle labels separated by spaces, a new multiparticle label is created and linked to the list of corresponding PDG-ids. The `define` command also allows for redefining existing labels or to add extra (multi)particles to a definition. In this last case, the user is allowed to issue commands such as, *e.g.*,

```
define <label> = <label> <identifiers>
```

The effect of the command above is to add the particle(s) contained in the variable `<identifiers>` to the definition of the label `<label>`.

Even if, in principle, the user is allowed to choose any name for the labels to be created, MADANALYSIS 5 contains a set of reserved keywords that cannot be used as labels, such as the symbols `all`, `or` and `and` (see below). In addition, the names of the different actions (see Table 4) cannot be used for (multi)particle labels.

At any time, a label not already used for the definition of a histogram or of a cut can be deleted from the memory of the computer by issuing

```
remove <label>
```

where the label to be deleted is denoted by `<label>`. Due to their particular nature, the labels `hadronic` and `invisible` however can never be removed.

#### 4.5. Creating histograms

There are two types of observables that can be represented as histograms by MADANALYSIS 5, global observables (see Table 7), related to the entire event content, and observables related to a given type of particle (see Table 8), thus specific to only a part of the event. In addition, a few additional observables are dedicated to analyses at the reconstructed-level and are collected in Table 9.

We first address the description of the global observables. Two of them are related to the missing energy. The corresponding symbols implemented in MADANALYSIS 5 are denoted by `MET` and `MHT` and correspond to the missing transverse energy  $\cancel{E}_T$  and the missing hadronic transverse energy  $\cancel{H}_T$ , respectively. The definitions of these two quantities are not unique in the literature, and we decide to adopt the choice

$$\cancel{E}_T = \left\| \sum_{\text{visible particles}} \vec{p}_T \right\| \quad \text{and} \quad \cancel{H}_T = \left\| \sum_{\text{hadronic particles}} \vec{p}_T \right\|, \quad (1)$$

where  $\vec{p}_T$  stands for the particle transverse momentum. At the parton-level, the generic name *hadronic particles* stands for gluons and light and  $b$ -quarks, but not for top quarks (decaying most of the time to a  $W$ -boson and a  $b$ -quark). At the hadronic-level and reconstructed-level, *hadronic particles* are trivially hadrons and jets, respectively. These definitions are related to the default values of the multiparticle labels `hadronic` and `invisible`. We

remind that the latter can be modified by the user by means of the command **define** (see Section 4.4).

Even if particle objects are not explicitly tagged as visible, any non-invisible object, *i.e.*, an object whose PDG-id is not included in the definition of the label **invisible**, is considered by MADANALYSIS 5 as visible. Similarly, any object which is not tagged as hadronic is considered as non-hadronic.

Concerning the missing transverse energy  $\cancel{E}_T$ , a remark is in order. Detector simulation tools are in general internally computing the missing energy, following a built-in definition which might be different from the one of Eq. (1). This value is stored under a special tag in the event files, compliant with the LHCO format relevant for analyses at the reconstructed-level. When this is read by MADANALYSIS 5, the value of the missing transverse energy is subsequently imported and always supersedes the one which would have been computed by employing Eq. (1).

Parallel to those observables related to the invisible sector of the events, two important global observables are related to the visible objects, the total transverse energy  $E_T$  and the total transverse hadronic energy  $H_T$ . In MADANALYSIS 5, these kinematical quantities are defined by

$$E_T = \sum_{\text{visible particles}} ||\vec{p}_T|| \quad \text{and} \quad H_T = \sum_{\text{hadronic particles}} ||\vec{p}_T|| , \quad (2)$$

and are associated to the symbols **TET** and **THT**, respectively.

Creating histograms associated to the four variables introduced above follows the standard syntax of the command **plot**,

```
plot <observable> <nbins> <min> <max>
```

where the value of the variable **<observable>** corresponds here either to **MET**, **MHT**, **TET** or **THT**. The number of bins, the value of the lowest bin and the one of the highest bin are optional information which can be passed through the symbols **<nbins>**, **<min>** and **<max>**, respectively. If left unspecified, MADANALYSIS 5 uses built-in values.

The effect of the command **plot** is to create a new instance of the class **selection** which has been designed to handle histograms (and cuts, as it is shown in Section 4.6). The labeling of the **selection** objects is internally handled by MADANALYSIS 5. The first time that the command **plot** is issued in the current session of MADANALYSIS 5, the label **selection[1]** is assigned

to the corresponding histogram. The second occurrence of the command `plot` leads to the creation of the object `selection[2]`, and so on. The full list of instances of the `selection` class which have been created can be displayed to the screen with the `display` command,

```
display selection
```

As for particle and multiparticle labels, typing in the interpreter

```
display selection[<i>]
```

allows us to display the properties of the  $<i>^{\text{th}}$  created `selection` object. This  $<i>^{\text{th}}$  selection can always be deleted from the memory of the computer with the action `remove`,

```
remove selection[<i>]
```

After the removal of a specific selection, the numbering of the different created selections is automatically adapted by MADANALYSIS 5 so that the sequence of the integer numbers is respected, without any hole. Even if handled automatically, the ordering of the selections can be modified by the user by means of the command `swap`. Issuing

```
swap selection[<i>] selection[<j>]
```

results then in the exchange of the  $<i>^{\text{th}}$  and  $<j>^{\text{th}}$  instances of the `selection` class.

Objects of the `selection` class have several attributes, as shown in Table 6. In this section, we however only focus on `selection` objects related to histograms. In the case of cuts, we refer to Section 4.6. The number of bins, the value of the lowest bin of a histogram and the one of its highest bin are stored in the attributes `nbins`, `xmin` and `xmax`, respectively. Their values are fixed at the time the command `plot` is issued in the interpreter. They can however be further modified by the use of the command `set`, as for any other attribute of an object, by typing in the command interface, *e.g.*,

```
set selection[<i>].xmax = 100
```

This command allows us to set the value of the highest bin of the histogram associated to the object `selection[<i>]` to 100.

Two of the other attributes of the `selection` class are related to the scales used when drawing the axes of the histograms. By default, a linear

**Table 6: List of the attributes of the selection class**

<code>logX</code>	If set to <code>true</code> , this enforces a logarithmic scale for the <i>x</i> -axis. If set to <code>false</code> (default), a linear scale is used.
<code>logY</code>	If set to <code>true</code> , this enforces a logarithmic scale for the <i>y</i> -axis. If set to <code>false</code> (default), a linear scale is used.
<code>nbins</code>	Number of bins of the histogram. The value taken by this attribute must be an integer.
<code>rank</code>	This refers to the observable to be used for ordering particles of the same type. The value of this option has to be anything among <code>ETAordering</code> (pseudorapidity ordering), <code>ETordering</code> (transverse-energy ordering), <code>Eordering</code> (energy ordering), <code>Pordering</code> (momentum ordering), <code>PTordering</code> (transverse-momentum ordering, the default choice), <code>PXordering</code> (ordering according to the <i>x</i> -component of the momentum), <code>PYordering</code> (ordering according to the <i>y</i> -component of the momentum) or <code>PZordering</code> (ordering according to the <i>z</i> -component of the momentum).
<code>stacking_method</code>	When several datasets are represented on an histogram, the different curves can be stacked ( <code>stack</code> ), superimposed ( <code>superimpose</code> ) or normalized to unity ( <code>normalize2one</code> ), including superimposing. The default value is <code>auto</code> and then refers to the properties of the attribute <code>stacking_method</code> of the object <code>main</code> (see Section 4.7).
<code>statuscode</code>	By default, only final state particles are considered in histograms. This attribute allows us to consider instead initial or intermediate particles, or all the particle content. The (intuitive) allowed values are <code>allstate</code> , <code>initialstate</code> , <code>finalstate</code> (default) and <code>interstate</code> .
<code>titleX</code>	The title of the <i>x</i> -axis of the histogram, given as a string.
<code>titleY</code>	The title of the <i>y</i> -axis of the histogram, given as a string.
<code>xmin</code>	Central value of the lowest bin of the histogram. This must be a real number.
<code>xmax</code>	Central value of the highest bin of the histogram. This must be a real number.

scale is employed. Logarithmic scales can be enforced through the boolean attributes `logX` and `logY` which can then take the values `true` or `false` (default choice). For instance, issuing

```
set selection[<i>].logY = true
set selection[<j>].logX = false
```

ensures the usage of a logarithmic scale for the  $y$ -axis of the histogram related to the object `selection[<i>]` and the one of a linear scale for the  $x$ -axis of the histogram associated to the object `selection[<j>]`.

Another important attribute related to the layout of the histograms concerns the stacking method used for the curves related to the different datasets created by the user. By default, the curves are drawn as stacked one above each other, but this can be modified through the attribute `stacking_method` of the class `selection`. For instance, focusing on the object `selection[<i>]`, the command

```
set selection[<i>].stacking_method = <value>
```

allows us to change the employed stacking method to the value `<value>`. By default, this attribute is set to the value `auto`. This means that the value of the attribute `stacking_method` of the object `main` (see Section 4.7) is employed. The other allowed choices are `stack`, `superimpose` and `normalize2one`. In the first case, the curves are all stacked, whilst in the second case, they are superimposed. The last possibility for the attribute `stacking_method`, *i.e.*, `normalize2one`, has been designed for comparing the shapes of the curves related to the different datasets. Here, the normalization of each curve is set to one and they are drawn superimposed. Let us note that this consists in a helpful feature when optimizing selection cuts.

To achieve the description of the functions allowing us to tune the layout of a histogram, the user is allowed to change the titles of the  $x$ -axis and  $y$ -axis by means of the attributes `titleX` and `titleY` of the `selection` class,

```
set selection[<i>].titleX = <string>
set selection[<i>].titleY = <string'>
```

The effect of the commands above leads to the use of the strings `<string>` and `<string'>` as titles for the  $x$ -axis and  $y$ -axis of the histogram represented by the object `selection[<i>]`, respectively.

The attributes of an instance of the `selection` class related to a histogram can also be directly set when issuing the command `plot`,

**Table 7: List of the global observables that can be represented by a histogram**

NAPID	The particle multiplicity of the events after mapping particles and antiparticles.
NPID	The particle multiplicity of the events.
MET	The missing transverse energy as defined in Eq. (1). At the reconstructed-level, the missing energy is however directly read from the LHCO file.
MHT	The missing transverse hadronic energy as defined in Eq. (1).
SQRTS	Partonic center of mass energy. Only available for parton-level and hadron-level event samples.
TET	The visible transverse energy as defined in Eq. (2).
THT	The visible transverse hadronic energy as defined in Eq. (2).

```
plot <observable> <nbins> <min> <max> [options]
```

The optional pattern [options] stands for `logX`, `logY`, `stack`, `superimpose`, `normalizeZone` or the value of any of the other attributes presented below, `statuscode` and `rank`. Multiple keywords are allowed (however only one for each attribute). In this way, the values of several options can be passed at one time, each separated by a space character. For instance, creating a histogram representing the visible transverse hadronic energy with a  $y$ -axis represented with a logarithmic scale and where all the curves are drawn superimposed can be done by issuing the command

```
plot THT [logY superimpose]
```

The binning is then automatically handled by MADANALYSIS 5, since it is not provided by the user.

In addition to the kinematical global observables  $E_T$ ,  $\cancel{E}_T$ ,  $H_T$  and  $\cancel{H}_T$  introduced in the beginning of this Section, the partonic center-of-mass energy can be represented by a histogram by typing the command

```
plot SQRTS
```

Finally, two other global observables are available. The latter are related to the particle multiplicity of the events. In order to draw the associated histograms, it is enough to enter in the command interface the two commands

```
plot NPID
plot NAPID
```

The first command, `plot NPID`, generates a histogram where one bin is associated to each possible type of final-state particle, the height of the bin being related to the multiplicity of the corresponding particle within the whole sample. Hence, if several particles of the same type are present in one specific event, they correspond to several entries in the histogram. The second command above, `plot NAPID`, produces a similar histogram after mapping antiparticles and particles. For both types of histograms, the labels of the *x*-axis correspond to the particle label imported in MADANALYSIS 5. If non-existing, the PDG-ids are used instead. We emphasize that these histograms are special histograms dedicated to the task of getting an idea about the particles present in the input sample. To compute the multiplicity of a given particle species, we refer to the observable `N` described below.

The second large class of observables that can be represented by histograms in MADANALYSIS 5 refers to the kinematical properties of the particles contained in the events. Hence, distributions such as the invariant mass or the transverse momentum of given particle species can be computed. The complete list of implemented observables can be found in Table 8.

Creating histograms associated to a given property `<observable>` of a specific particle represented by the label `<label>` is also based on the command `plot`. In this case, the syntax is however slightly different as for the global observables. The symbol `<observable>` has to be seen as a function which takes as the argument the label associated to the particle under consideration,

```
plot <observable>(<label>)
```

For instance, if `mu+` stands for the particle label related to antimuons, the command

```
plot PT(mu+)
```

results in representing by a histogram the transverse-momentum distribution of all the antimuons in the sample. As above, if several antimuons are included in one single event, they contribute to several entries in the histogram. In the case of the relative distance between two particles (denoted by `DELTAR`), two particle labels `<label1>` and `<label2>` are required,

**Table 8: List of the kinematical observables that can be represented by histograms**

BETA	Velocity $\beta = v/c$ .
DELTAR	Relative distance between two objects in the $\eta - \phi$ plane.
E	Energy.
ET	Transverse energy.
ETA	Pseudorapidity.
GAMMA	Lorentz factor.
M	Invariant mass.
N	Particle multiplicity.
MT	Transverse mass.
P	Norm of the momentum.
PHI	Azimuthal angle of the momentum.
PT	Norm of the transverse momentum.
PX	Projection of the momentum on the $x$ -axis.
PY	Projection of the momentum on the $y$ -axis.
PZ	Projection of the momentum on the $z$ -axis.
R	Position of the object in the $\eta - \phi$ plane.
THETA	Angle between the momentum and the beam axis.
Y	Rapidity.

```
plot DELTAR(<label1>, <label2>)
```

As illustrated above, when studying the kinematical properties of a given particle species, the normalization of the histograms reflects the total number of particles of this type included in the full sample. This can be different from the total number of events, since one single event could describe a final state with zero, one or several particles of the considered type, which then corresponds to zero, one or several entries in the produced histograms<sup>3</sup>.

The syntax detailed above is also valid for multiparticle objects. In this case, the label `<label>` is related to an instance of the multiparticle class. Histograms are created by treating on the same footing all the particles linked

---

<sup>3</sup>In this Section, we do not consider the luminosity which also affects the normalization of the histograms. This feature is described in more detail in Section 4.7.

to the label `<label>`. For example, the set of commands

```
define <multi> = <particle1> <particle2>
plot <observable>(<multi>)
```

defines, in a first step, a multiparticle label denoted by `<multi>` and linked to the particle species `<particle1>` and `<particle2>` (see Section 4.4). In a second step, a property represented by the observable `<observable>` is investigated. For a specific event, each particle of the type `<particle1>` or `<particle2>` is associated to one entry in the histogram.

As shown above, if several particles of the considered type appear in a specific event, they always correspond to several entries in the histograms. In phenomenological analyses, it is however often more relevant to only consider the leading particle, *i.e.*, the one which has the highest value of a kinematical variable such as the transverse momentum or the energy. The squared brackets ‘[ ]’ allow us to access leading, next-to-leading, *etc*, particles. For instance, the histogram resulting from the command

```
plot <observable>(<label>[<i>])
```

represents the distribution of the observable `<observable>` for the particle of the type `<label>` with the  $< i >^{\text{th}}$  largest transverse momentum, `<i>` being a positive integer. Similarly, a negative value of the parameter `<i>` corresponds to a pointer to the particle with the  $< i >^{\text{th}}$  smallest transverse momentum. Events where the number of particles of the type under consideration is smaller than the absolute value of `<i>` are ignored at the time of the creation of the histogram.

By default, the ordering variable employed in MADANALYSIS 5 is the transverse momentum. Other possible choices exist and the information is passed, for a given histogram, through the value of the attribute `rank` of the `selection` class. Hence, considering the instance of the class `selection[<i>]`, one can modify the ordering variable by issuing

```
set selection[<i>].rank = <value>
```

The parameter `<value>` can take any value among `ETAordering` (pseudorapidity ordering), `ETordering` (transverse-energy ordering), `Eordering` (energy ordering), `Pordering` (momentum ordering), `PTordering` (transverse-momentum ordering), `PXordering` (ordering according to the  $x$ -component of the momentum), `PYordering` (ordering according to the  $y$ -component of

the momentum) and `PZordering` (ordering according to the  $z$ -component of the momentum).

For all the histograms produced so far, `MADANALYSIS 5` only uses information related to the final-state particles. However, for phenomenological purposes, it is often interesting to investigate the properties of the intermediate particles or those of final-state particles issued from the decays of a specific type of intermediate particle. This requires, of course, that the relevant information is available. This is not the case for event samples at the reconstructed-level since these features are totally absent from event files under the LHCO format. In contrast, the user can access, for hadron-level or parton-level event files, to the whole particle history by a set of mother-to-daughter relations. To benefit from those relations, there exist two special functions in `MADANALYSIS 5`.

Firstly, the symbol ‘`<`’ links one particle or set of particles to their direct mother. The command line

```
plot <observable>(<type1> < <type2>)
```

allows us hence to study a given property, which is represented by the symbol `<observable>`, of the particles of type `<type1>` included in the final state of the events under consideration. However, in order to correspond to an entry in the histogram, the particles of type `type1` must be issued from the direct decay of a particle of type `<type2>`. Doubling the symbol ‘`<`’, *i.e.*, replacing it by ‘`<<`’, allows us to remove the restriction of a *direct* decay. Finally, the usage of these symbols recursively

```
plot <observable>(<type1> < <type2> << <type3>)
```

allows us to focus on entire decay chains. Here, we investigate the properties of the particle species `<type1>`, but only for particles of type `type1` issued from a direct decay of a particle of type `<type2>`, which is itself issued from a decay of a particle represented by the label `<type3>`, this last decay possibly occurring in several steps.

Secondly, kinematical properties of the intermediate particles can be directly investigated through the option `statuscode` of the `selection` class. As suggested above, by default, only final-state particles are considered. This corresponds to the value `finalstate` of the attribute `statuscode`. Issuing in the interpreter

```
set selection[<i>].statuscode = interstate
```

indicates that, for the selection `selection[<i>]`, we are not considering final-state particles anymore, but only intermediate states. On the same footing, setting `statuscode` to the value `initialstate` allows us to focus on the initial-state particles only, whilst setting it to `allstate` allows us to consider equivalently initial-state, final-state and intermediate-state particles.

Key observables to design efficient selection cuts in an analysis are in general related to more than one single particle. For instance, highlighting a new  $Z'$  gauge boson in Drell-Yan events and estimating its mass with a good accuracy rely on the invariant-mass distribution of the produced lepton pair. All the functions of Table 8, but the `DELTAR` observable which requires exactly two arguments, can take an arbitrary number of arguments. This allows us to combine particles before computing the kinematical distribution to be represented. Hence, the two equivalent command lines

```
plot <observable>(<prtc11> <prtc12>)
plot v<observable>(<prtc11> <prtc12>)
```

lead to the creation of a histogram showing the distribution of the observable `<observable>`. The observable is computed on the basis of the combined four-vector built from the sum of the four-momentum of the particles `<prtc11>` and `<prtc12>`. The optional prefix ‘v’ indicates that the four-momenta are combined vectorially (other options are shown below). If, for a given event, several pairs of particles `<prtc11>` and `<prtc12>` can be formed, each possible pair leads to one different entry in the histogram. Hence, the histogram describing the invariant mass of a muon pair can be created by issuing

```
plot M(mu+ mu-)
```

where, as before, the binning is automatically handled by MADANALYSIS 5. The observable that is computed corresponds to the norm of the sum of the four-momentum of the muon and the one of the antimuon. This syntax can straightforwardly be generalized to multiparticles or to combinations of more than two particles. A remark is however in order here. If `<multi>` denotes the label associated to an instance of the multiparticle class, issuing

```
plot <observable>(<multi> <multi>)
```

generates a histogram where each entry corresponds to one *different* combination of the particles represented by the multiparticle `<multi>`. MADANALYSIS 5 indeed forbids double-counting any combination.

By default, the particles are combined by adding vectorially their four-momentum, *i.e.*,

$$p_\mu = \sum_i p_\mu^i ,$$

where  $p_\mu$  is the resulting four-momentum and  $p_\mu^i$  are the four-momenta of the particles to be combined. The four-vector  $p^\mu$  is the one which is used when computing the value of the observable to be represented on the corresponding histogram. MADANALYSIS 5 offers additional ways to perform this combination. The sum could be, in contrast, done scalarly, *i.e.*, by firstly computing the considered observable for each of the particles to be combined and secondly adding the results. To select this option, the user must add a prefix ‘**s**’ in front of the name of the observable when typing the command in the interpreter,

```
plot s<observable>(<prtcl1> ...)
```

where the dots stand for the list of particles to be combined. In the case the user wants to combine all particles of a given type `<prtcl>` in an event, the reserved keyword `all` can be used in order to simplify the syntax,

```
plot <observable(all <prtcl>)
```

If two (and only two) particles are combined, differences can also be computed, rather than sums. To allow for this option, it is enough to include one of the the prefixes ‘**dv**’, ‘**vd**’, ‘**d**’, ‘**ds**’, ‘**sd**’ or ‘**r**’ in front of the symbol of the observable to be computed. For the options ‘**dv**’, ‘**vd**’ and ‘**d**’, vectorial differences are considered. Hence, the result of one of the equivalent commands

```
plot dv<observable>(<prtcl1> <prtcl2>)
plot vd<observable>(<prtcl1> <prtcl2>)
plot d<observable>(<prtcl1> <prtcl2>)
```

is to subtract the two four-momenta from each other,

$$p_\mu = p_\mu^1 - p_\mu^2 ,$$

**Table 9: List of the additional observables that can be represented by histograms for reconstructed events**

EE_HE	Ratio of the electromagnetic and hadronic energy for a given object.
HE_EE	Ratio of the hadronic and electromagnetic energy for a given object.
NTRACKS	Number of tracks in a jet. This returns zero for non-jet objects.

$p_\mu^1$  being the four-momentum of the particle associated to the label `<prtcl1>` and  $p_\mu^2$  the one of the particle associated to the label `<prtcl2>` and then compute the observable `<observable>` from the resulting four-vector  $p^\mu$ . In the case of the prefixes ‘ds’ and ‘sd’,

```
plot ds<observable>(<prtcl1> <prtcl2>)
plot sd<observable>(<prtcl1> <prtcl2>)
```

scalar differences are computed, *i.e.*, the observable `<observable>` is computed for each of the particles `<prtcl1>` and `<prtcl2>` and the results are subtracted from each other. Relative differences can also be computed by means of the prefix ‘r’,

```
plot r<observable>(<prtcl1> <prtcl2>)
```

In this case, the results consist in the scalar difference of the values of the observable computed for the two particles `<prtcl1>` and `<prtcl2>` taken individually, which is however given relative to the value of the observable for the first particle represented by the label `<prtcl1>`. This corresponds then to the quantity

$$\frac{\langle\text{observable}\rangle(\langle\text{prtcl1}\rangle) - \langle\text{observable}\rangle(\langle\text{prtcl2}\rangle)}{\langle\text{observable}\rangle(\langle\text{prtcl1}\rangle)}.$$

If the considered observable vanishes for the particle labeled by `<prtcl1>`, the value zero is returned.

A final way for combining observables is related to the reserved word `and`. When the considered observable has to be evaluated for several possible

combinations of particles, the user can use the keyword `and` to efficiently implement the corresponding command in MADANALYSIS 5,

```
plot <observable>(<p1> <p2> and <p3> <p4>)
```

The command line above computes the observable represented by the symbol `<observable>`, firstly, for the vectorial combination of the particles `<p1>` and `<p2>` and secondly, for the vectorial combination of the particles `<p3>` and `<p4>`. The two distributions are eventually summed.

Three additional observables can be represented by histograms in the case of fully-reconstructed objects. To allow for generating histograms for these observables, MADANALYSIS 5 must be run in the reconstructed-level mode. These observables consist of the ratio between the hadronic and the electromagnetic energy for a given object (the ratio between the energy deposited in the electromagnetic and hadronic calorimeters of a detector), the inverse ratio and the number of tracks within a (reconstructed) jet. For objects different from a jet, this last observable always returns the number zero. The associated symbols are `HE_EE`, `EE_HE` and `NTRACKS` and their definitions are collected in Table 9. The corresponding histograms can be created by following the usual syntax,

```
plot <observable>(<label>)
```

Let us note that for these three observables, combining objects is not supported, *i.e.*, only one single (multi)particle label can be passed as an argument.

#### 4.6. Selection cuts

In MADANALYSIS 5, the process of event selection is based on two equivalent classes of kinematical cuts which can be applied to the imported datasets. The program offers to the user the two choices of either selecting or rejecting events in the case a certain condition is fulfilled. This task can be performed at the level of the command line interpreter of the program by means of the two actions `select` and `reject`. The associated syntax is very intuitive and reads

```
select <condition> [<options>]
reject <condition> [<options>]
```

For the first (second) command, events are selected (rejected) if the condition `<condition>` is satisfied. Hence, the command

```
reject PT(mu) > 50
```

leads to the rejection of all the events where at least one muon with a transverse momentum  $p_T > 50$  GeV is found, whilst issuing

```
select M(e+ e-) > 100
```

allows for the selection of all the events where we have a electron-positron pair with an invariant mass larger than 100 GeV. Internally, the effects of the commands above are to create instances of the **selection** class with special properties. Contrary to the command **plot** which is related to the creation of histograms, the commands **select** and **reject** lead to the production of tables of cut efficiencies. Consequently, the only attributes of the **selection** class which are relevant are **rank** and **statuscode** (see Table 6). They can be either passed directly at the time of typing-in the commands, by including the desired values as the optional parameter **<options>** above, or at a later stage by means of the command **set** (see Section 4.5).

Modifying the **rank** attribute only plays a role if the ordering of the particles is necessary information for a good application of the condition, as *e.g.*, if we are constraining some observable related to the leading or next-to-leading particles. Furthermore, setting the option **statuscode** to a non-default value allows us to apply, if needed by the user, cuts on initial or intermediate states rather than on final states only.

The condition **<condition>** must be given according to the pattern

```
<observable> <logical-operator> <value>
```

where the observable **<observable>** can be any observable from Tables 7, 8 and 9, computed for a given particle or for any combination of particles, with the exception of the global variables related to the symbols **NPID** and **NAPID**<sup>4</sup>. The supported logical operators are ‘>’ (greater than), ‘ $\geq$ ’ (greater than or equal to), ‘<’ (smaller than), ‘ $\leq$ ’ (smaller than or equal to), ‘==’ (equal to) and ‘!=’ (different from).

Conditions can also be combined by using one or several of the connecting keywords **and** (logical *and*) and **or** (logical *inclusive or*), such as in

```
<cond1> <connector1> <cond2> <connector2> <cond3>
```

---

<sup>4</sup>To implement selection cuts on the multiplicity of a given particle species, cuts on the observable **N** have to be performed, rather than on the global observable **NPID** and **NAPID**.

The condition above consists of the combination of the three conditions `<cond1>`, `<cond2>` and `<cond3>` by employing the two connecting keywords `<connector1>` and `<connector2>` being `and` or `or`. Moreover, brackets are also authorized by the syntax for handling more complex conditions. In the special case both upper and lower limits are imposed on a given observable, the condition can be easily implemented by means of the keyword `and`. However, there exists a more compact syntax

```
<value1> <logical-operator> <obs> <logical-operator> <value2>
```

which could be employed. The logical operator `<logical-operator>` is here used twice. The parameter `<obs>` denotes the observable and `<value1>` and `<value2>` the imposed bounds.

So far, we have supposed that all the objects present in each event can be used for applying selection cuts. However, this is barely the case in most realistic phenomenological analyses, since, for example, too soft particles are in general omitted. This feature can also be implemented in analyses performed with MADANALYSIS 5 by means of the commands `select` and `reject`, but following a slightly different syntax as the one shown before,

```
select (<particle>) <condition> [<options>]
reject (<particle>) <condition> [<options>]
```

The commands above allow us to respectively select and reject any particle associated to the label `<particle>`, if the condition `<condition>` is fulfilled. The syntax for typing-in the condition is similar to the one employed for implementing conditions associated to the selection or rejection of events, with the difference that the observable entering the condition cannot here take any argument and must simply be one of the symbols presented in Tables 8 and 9.

In the case the particle candidate is rejected, the event is considered as without containing this particle. Let us note that whilst selecting and rejecting events have a direct influence on the signal over background ratio, selection or rejecting candidates to a particle type only affects the number of entries for one event in the histograms.

#### *4.7. Executing an analysis and displaying the results*

Once an analysis has been implemented, it must be passed to SAMPLE-ANALYZER, the C++ kernel of MADANALYSIS 5, for execution by means of the command

```
submit <dirname>
```

A directory named `<dirname>` is created and all the files necessary for SAMPLEANALYZER to properly run are generated and included in this directory. The code is further compiled and linked to the external static library of MADANALYSIS (see Section 3.1), and the execution of the resulting program is eventually managed by MADANALYSIS 5.

This execution starts with the reading of the event samples under consideration and their storing in the memory of the computer according to a format internal to MADANALYSIS 5. All the histograms required by the user are then sequentially created, including the application to all the events of the defined selection cuts. As an output, a ROOT file [57] is generated so that the analysis can be accessed later, as, *e.g.*, directly in the ROOT framework or in a new session of MADANALYSIS 5. In this last case, the SAMPLEANALYZER (executed) job can be imported by means of the command `import`

```
import <dirname>
```

where the directory `<dirname>` contains the analysis previously performed.

After modifications such as asking for the creation of a new histogram or the application of a new selection cut, the user does not have to submit the SAMPLEANALYZER job entirely again. A much faster option, saving a sensible amount of computing time, consists in the command

```
resubmit
```

This updates the already existing SAMPLEANALYZER directory and only the differences with respect to the original analysis are executed.

Once the ROOT file has been created by SAMPLEANALYZER, the command `preview` allows for the display of a single histogram,

```
preview selection[<i>]
```

where in the generic example above, the `<i>th` histogram is asked to be displayed to the screen, in a ROOT popup window. The command `preview` only works for displaying histograms. Consequently, previewing the efficiency table associated to a selection cut is not possible.

A more complete report, with all the selection cuts, efficiency tables and histograms can be generated by issuing in the command line interface one of the three commands

```
generate_html <html dirname>
generate_latex <tex dirname>
generate_pdflatex <pdflatex dirname>
```

The first command, `generate_html`, generates the report under the HTML format and stores the files in the directory `<html dirname>`. The second and third commands, `generate_latex` and `generate_pdflatex`, are related to the creation of `TEX` files to be compiled with the help of the shell commands `latex` and `pdflatex`, respectively. In these two cases, the `TEX` files are already compiled by MADANALYSIS 5, if the `latex` and `pdflatex` commands are available on the system of the user.

Histograms are exported to (non-ROOT) figures under either the Encapsulated PostScript (`.eps`) format, required for a proper compilation with `latex`, or to the Portable Network Graphics (`.png`) format for HTML files or `TEX` files to be compiled with `pdflatex`. In addition, each figure is also saved as a ROOT macro which can be modified by the user for further processing. Once generated, the report can be immediately displayed by typing-in the command

```
open <report dirname>
```

which opens the report in a web browser.

The analysis, and thus the generation of the histograms and the efficiency tables, is so far based on the default configuration of MADANALYSIS 5. This configuration can be modified by superseding the default values of the attributes of the object `main`, which are listed in Table 10 and Table 11.

Two options allow us to control the normalization of the histogram, `lumi` and `normalize`. This last attribute defines the way the histograms are normalized. The allowed values are `none`, `lumi` and `lumi_weight` and can be set as for any other attribute of any class,

```
set main.normalize = <value>
```

For the first choice (`<value> = none`), the total number of events included in each dataset is kept. The two other options imply that the histograms are normalized with respect to the integrated luminosity. The difference lies in the weight which can be possibly associated to each event (see Section 4.3), which can be ignored (`<value> = lumi`) or included (`<value> = lumi_weight`). This last possibility consists in the default choice.

**Table 10: List of the attributes of the object `main`**  
`set main.<option> = <value>`

<code>currentdir</code>	Current directory in which any directory created by MADANALYSIS 5 is stored.
<code>lumi</code>	This allows us to modify the value of the integrated luminosity, in $\text{fb}^{-1}$ , used for the normalization of the histograms. The default value is $10 \text{ fb}^{-1}$ .
<code>normalize</code>	This defines the way the histograms are normalized. The allowed choices are the number of events included in the different datasets ( <code>none</code> ), the integrated luminosity without ( <code>lumi</code> ) or accounting for the event weights associated to each dataset (the default choice, <code>lumi_weight</code> ).
<code>SBerror</code>	This fixes the way the uncertainty on the signal over background ratios is computed by MADANALYSIS 5. The attribute <code>&lt;value&gt;</code> is the corresponding analytical formula passed as a valid PYTHON expression given as a string. It has to depend on <code>S</code> (number of signal events), <code>B</code> (number of background events), <code>ES</code> (uncertainty on the number of signal events) and <code>EB</code> (uncertainty on the number of background events). It is automatically handled for the most simple cases (see Eq. (3)).
<code>SBratio</code>	This fixes the way signal over background ratios are computed by MADANALYSIS 5. The attribute <code>&lt;value&gt;</code> is the corresponding analytical formula passed as a valid PYTHON expression given as a string. It has to depend on <code>S</code> (number of signal events) and <code>B</code> (number of background events).
<code>stacking_method</code>	When several datasets are represented on histograms, the different curves can be stacked ( <code>stack</code> , default), superimposed ( <code>superimpose</code> ) or normalized to unity ( <code>normalize2one</code> ), including superimposing. When the <code>stacking_method</code> attribute of an instance of the class <code>selection</code> is set to <code>auto</code> , the value of <code>main.stacking_method</code> is employed.

By default, all the histograms are normalized to an integrated luminosity of  $10 \text{ fb}^{-1}$ . This value can be updated by modifying the attribute `lumi` of the object `main`,

```
set main.lumi = <new-value>
```

where the new value of the integrated luminosity, `<new-value>`, is given in  $\text{fb}^{-1}$ .

Once all the datasets have been defined as part of the signal or background samples, MADANALYSIS 5 can compute automatically the signal ( $S$ ) over background ( $B$ ) ratio, and thus the efficiency of each selection cut. This feature is related to an attribute of the object `main` denoted by `SBratio`. It refers to an analytical formula, expressed as a valid PYTHON expression, which indicates how the signal over background ratio must be calculated<sup>5</sup>. When implementing this formula, the symbols related to the signal and background number of events are `S` and `B`, respectively. By default, the signal over background ratio is computed according to  $S/B$ . This can be modified through the command `set`, by issuing in the interpreter,

```
set main.SBratio = '<formula>'
```

where `<formula>` is, as sketched above, a valid PYTHON expression depending on the two variables `S` and `B`. For instance, the command

```
set main.SBratio = 'S/sqrt(S+B)'
```

enforces the signal over background ratio  $r$  to be computed according to

$$r = \frac{S}{\sqrt{S+B}} .$$

In the case the signal over background ratio is undefined due, *e.g.*, to the evaluation of the squared root of a negative number or to a division by zero, the value zero is returned. Let us emphasize that it is safer for the user to use real numbers when typing-in the analytical expression `<formula>`, rather than integer numbers. We indeed recall that `1/2` is evaluated as 0 whilst `1./2` returns 0.5.

---

<sup>5</sup>The formula is internally stored as an instance of the `TFormula` class. This structure is defined in the ROOT library linked to MADANALYSIS 5, and all the associated attributes can therefore be used. We refer to the ROOT manual for more information [57].

It is fundamental to associate an uncertainty to the signal over background ratio. The way to compute this quantity is related to the **SBerror** attribute of the object **main**. As for **SBratio**, it refers to an analytical formula, given as a valid PYTHON expression, which indicates how the uncertainty on the signal over background ratio must be computed. For the three choices

$$r_1 = \frac{S}{B}, \quad r_2 = \frac{S}{S+B} \quad \text{and} \quad r_3 = \frac{S}{\sqrt{S+B}}, \quad (3)$$

as well as for the three additional cases obtained when  $S$  and  $B$  are exchanged, MADANALYSIS 5 automatically detects the formula stored in the **SBratio** attribute. It then updates accordingly the uncertainty, setting the **SBerror** attribute to the values  $\Delta r_i$  given by

$$\begin{aligned}\Delta r_1 &= \frac{\sqrt{B^2(\Delta S)^2 + S^2(\Delta B)^2}}{B^2}, \\ \Delta r_2 &= \frac{\sqrt{B^2(\Delta S)^2 + S^2(\Delta B)^2}}{(S+B)^2}, \\ \Delta r_3 &= \frac{\sqrt{(S+2B)^2(\Delta S)^2 + S^2(\Delta B)^2}}{2(S+B)^{3/2}},\end{aligned}$$

where  $\Delta S$  and  $\Delta B$  are the uncertainties on the signal and on the background number of events. If the user wants to use another formula or to set himself the **SBerror** attribute, the command **set** has to be employed,

```
set main.SBerror = '<formula>'
```

where **<formula>** depends this time on the number of signal and background events **S** and **B** as well as on the associated uncertainties **ES** ( $\equiv \Delta S$ ) and **EB** ( $\equiv \Delta B$ ). For instance, implementing by hand the error  $\Delta r_1$  above would give

```
set main.SBerror = 'sqrt(B**2*ES**2+S**2*EB**2)/B**2'
```

Four specific attributes of the object **main** concern muon isolation when MADANALYSIS 5 is run in order to analyze reconstructed-level events. The user has the possibility to choose the algorithm which is employed by MADANALYSIS 5 when defining an isolated muon. Two choices are implemented and can be adopted by issuing one of the two commands

**Table 11: List of the attributes of the object main associated to the isolation of the muons.**

```
set main.isolation.<option> = <value>
```

<b>algo</b>	This specifies the algorithm to be employed for muon isolation. The allowed choices are DELTAR (default) and SUMPT. In the first case, we require that no tracks lies inside a cone around the muon. In the second case, the sum of the transverse-momentum of all the tracks inside the cone and the ratio of the summed transverse energy of these tracks over their summed transverse momentum must be lower than values specified by the user. For the second algorithm, the size of the cone is fixed by the detector simulation tool.
<b>deltaR</b>	This specifies the radius of the isolation cone to be used by the DELTAR isolation algorithm.
<b>ET_PT</b>	This specifies the value of the ratio of the sum of the transverse energy of the tracks lying in the isolation cone over the sum of their transverse-momentum to be used by the SUMPT algorithm.
<b>sumPT</b>	This specifies the transverse-momentum threshold to be used by the SUMPT algorithm.

```
set main.isolation.algo = DELTAR  
set main.isolation.algo = SUMPT
```

In the first case, a muon is tagged as isolated when no track lies inside a cone around the muon. The size of this cone can be specified by the user by typing in the command interface

```
set main.isolation.deltaR = <value>
```

where `<value>` is a floating-point number. For the second algorithm, a muon is considered as isolated when the sum of the transverse momentum of all the tracks lying in a cone around the muon is lower than a value `<value>`. The latter can be specified by typing

```
set main.isolation.sumPT = <value>
```

In addition, the ratio of the sum of the transverse energy of these tracks over the sum of their transverse momentum has also to be lower than a value `<value'>` to be specified. This value is provided by means of the command

```
set main.isolation.ET_PT = <value'>
```

On a fairly different line, the last attribute of the object `main` which can be modified by the user is denoted by `currentdir`

```
set main.currentdir = <dirname>
```

It fixes the path to the directory in which any file and/or directory created by MADANALYSIS 5 is stored.

## 5. MADANALYSIS 5 for expert users

Besides its user-friendliness, the way of using MADANALYSIS 5 described in Section 4 is (obviously) restricted by the set of functionalities that have been implemented. The latter allow, in general, to perform rather traditional and standard analyses but might not be sufficient for more sophisticated and/or exotic investigations. For example, the normal running mode of the program does not allow us to create a histogram related to either the distribution of a new observable or to the one of an existing observable that needs to be computed in a reference frame different from the laboratory reference frame. Along the same lines, selection cuts must match the pattern

presented in Section 4.6, which forbids any other selection than cutting on the implemented kinematical variables by means of assigning a lower and/or upper bound on the result of the computation of the associated observable. Finally, as a last example, two- or three-dimensional histograms, necessary for correlation studies, are not included.

In order to overcome the above-mentioned restrictions as well as any type of, even unforeseen, limitations, MADANALYSIS 5 comes with an expert mode of running. In this case, the possibilities are only limited by the programming skills of the user and his originality in designing the analysis. The user is asked to implement the entire analysis himself, knowing that he is able to benefit from all the strengths of the SAMPLEANALYZER framework. The latter comes indeed with its own set of reading routines for the event samples, its own data format and a large class of functions and methods ready to be employed. In addition, an automated compilation of the analysis C++ files as well as a programming error management system are included.

As already presented in Section 4.1, the expert mode of MADANALYSIS 5 can be set by issuing in a shell one of the three commands

```
bin/ma5 --expert      bin/ma5 -e      bin/ma5 -E
```

### 5.1. The SAMPLEANALYZER framework

In the expert mode, there are two possibilities in order to create an analysis, *i.e.*, to implement the C++ source and header files which are automatically generated by MADANALYSIS 5 in its normal mode of running. Either one can start from and extend an existing analysis already generated by MADANALYSIS 5, or one can design a new analysis from scratch.

In the first case, a working directory has already been created through the issue, in the command line interface of MADANALYSIS 5, of the command **submit** (see Section 4.7). Consequently, at least one dataset has been imported (see Section 4.3) and at least one histogram has been created (see Section 4.5). The created working directory contains three sub-directories, **SampleAnalyzer**, **lists** and **root**. The first of these directories, **SampleAnalyzer**, includes all the files necessary for having the C++ kernel of MADANALYSIS 5 properly running and executing the analysis implemented by the user. When issuing the command **submit** in the command interface, all the PYTHON commands entered by the user are translated into a set of three C++ files, **user.cpp**, **user.h** and **analysisList.cpp**, stored in the sub-directory **Analysis** of **SampleAnalyzer**.

The two other sub-directories included in the working directory, `lists` and `root`, are dedicated to the input and output files, respectively. As mentioned in Section 4.3, the imported events are gathered into different datasets. For each dataset, a single list, containing the paths to the various associated event files, is stored in the directory `lists`. After being executed, SAMPLE-ANALYZER creates a ROOT file for each of the defined datasets. These files, stored into the directory `root`, contain the necessary information to create the histograms requested by the user.

When the user starts an analysis from scratch, the situation is similar to the one described in the paragraphs above, with the exception that the working directory is initialized directly from the shell, when MADANALYSIS 5 is launched. When typing in a shell the command

```
bin/ma5 --expert
```

or equivalently any of the three commands recalled in the introduction of this section, MADANALYSIS 5 indeed asks a series of questions to the user such as the name of the working directory to be created or the chosen label for the analysis to be created (see below).

As a result, a working directory is created, together with the sub-directory `SampleAnalyzer` which contains a blank analysis. The implementation of this analysis, as in any analysis, is divided into the implementation of three core functions. The latter have to be provided by the user in the file `user.cpp` which comes together with the corresponding header file `user.h`. Both files are stored in the sub-directory `SampleAnalyzer/Analysis`, together with an additional file, `analysisList.cpp`, that contains the list of all the analyses which are/will be implemented in the working directory under consideration. After the creation of a fresh working directory, this list only contains a single analysis, the blank analysis pre-implemented in the file `user.cpp`. However, there is no limitation on the number of analyses which can be included and nothing prevents this list from becoming very large.

Once the three C++ files introduced above have been created, MADANALYSIS 5 exits and the user can then start implementing his analysis by modifying these files. In this Section 5.1, we do not address the way in which an analysis has to be implemented, since this is the scope of Section 5.3, Section 5.4 and Section 5.5, but we only describe the steps leading to the execution of the analysis and the generation of the output files.

In order to properly compile and execute the implemented analysis, the linking to the external dependencies, such as the ROOT header files and li-

braries, must be performed appropriately. This is allowed by a correct setting of the environment variables `LD_LIBRARY_PATH` (or rather `DYLD_LIBRARY_PATH` for MACOS systems), `LIBRARY_PATH` and `CPLUS_INCLUDE_PATH` prior to the compilation. This task has been rendered automatic by means of the scripts `setup.sh` and `setup.csh` included in the `SampleAnalyzer` sub-directory. All the above-mentioned environment variables can be appropriately set at once by issuing

```
source setup.sh
```

in a `bash` shell or

```
source setup.csh
```

in a `tcsh` shell before compiling and/or executing `SAMPLEANALYZER`.

After this step, the analysis can be compiled with the help of the `Makefile` present in the `SampleAnalyzer` directory, assuming that the GNU MAKE package has been installed on the computer of the user. As for any program to be compiled with GNU MAKE, it is enough to issue in a shell the command

```
make
```

which also takes care of linking the external libraries to the `SAMPLEANALYZER` program. In addition, the command

```
make clean
```

has also been implemented and leads to the cleaning of the various sub-directories included in the current working directory.

Once compiled, the `SAMPLEANALYZER` core, containing the user's analysis, is ready to be executed. The only input left to be provided consists in the location paths of the event samples. They have, for a given dataset, to be collected into a single text file, denoted in the following by `<datasetlist>`. The `SAMPLEANALYZER` package is then simply run by issuing in a shell

```
SampleAnalyzer [ options ] <datasetlist>
```

If the event samples are collected into several datasets, `SAMPLEANALYZER` has to be run once for each of the datasets to be included in the analysis, with a different text file with the paths to the relevant event samples.

The only option supported by `SAMPLEANALYZER` consists in specifying the label of the analysis to be performed. In particular, this allows us to

include several analyses, each of them specified by a unique label, in one single working directory of **SAMPLEANALYZER**. The user can hence indicate which analysis to perform on run-time. If the option pattern is not provided, **SAMPLEANALYZER** lists to the screen all the analyses included in the file **analysisList.cpp**, together with the corresponding labels, and asks to the user to make his choice.

The label of an analysis is defined in the corresponding header file. In the case of a fresh working directory just created by **MADANALYSIS 5**, the chosen name for the label is asked by the program and exported to the file **user.h**. We refer to Section 5.2 for more information about the way to declare labels independently from **MADANALYSIS 5**. Assuming that the label under consideration is denoted by **<label>**, **SAMPLEANALYZER** is launched by issuing in a shell

```
SampleAnalyzer --analysis=<label> <datasetlist>
```

As a result, the analysis defined by the label **<label>** is executed by **SAMPLEANALYZER** and the output files are generated according to what the user has implemented in the C++ files related to the corresponding analysis.

### 5.2. Implementing new analyses using the analysis template

As mentioned in Section 5.1, the implementation of an analysis within the **SAMPLEANALYZER** framework consists in the writing of three files. Two of them are related to the analysis itself, *i.e.*, one C++ source file together with the associated header file. Since a given working directory of **SAMPLEANALYZER** can include many analyses, their list must be provided. This is the aim of the third file, **analysisList.cpp**, which is common to all the present analyses.

As presented in the previous Section, a pair of such analysis source/header files is automatically created when **MADANALYSIS 5** is run in expert mode. These files are denoted by **user.cpp** and **user.h**. However, we adopt, in the following, the generic names **name.cpp** and **name.h** since the user has in fact the freedom to choose the name of the files. The only requirement is that all the filenames are different. Moreover, all these files have to be stored, together with the list of the implemented analyses included in the file **analysisList.cpp**, in the sub-directory **SampleAnalyzer/Analysis**.

The pair of (generic) header and source analysis files **name.cpp** and **name.h** contains the declaration of a class denoted by **name**, *i.e.*, having the same

name as the files. The class `name` is a daughter class inheriting from the base class `AnalysisBase` that contains (empty) analysis methods. These methods are then specified at the level of the definition of the daughter classes.

The structure of the header file `name.h` follows, for any of the analyses included in the working directory,

```
#ifndef analysis_name_h
#define analysis_name_h

#include "Core/AnalysisBase.h"

class name: public AnalysisBase
{
    INIT_ANALYSIS(name,label)

public:
    virtual void Initialize();
    virtual void Finalize(const SampleFormat& summary,
                          const std::vector<SampleFormat>& files);
    virtual void Execute(const SampleFormat& sample,
                         const EventFormat& event);

private:
};

#endif
```

The only pieces among these predefined lines to be modified by the user are the name tag `name`, related to the filename, and the label of the analysis `label` which is a string. This label is the one that can be provided as an option when running the `SAMPLEANALYZER` code (see Section 5.1).

In the C++ code above, the `INIT_ANALYSIS` macro automatically creates the constructor and destructor methods associated to the class `name` and consistently links the filename to the name of the analysis, which must be the same. As for any class derived from the mother class `AnalysisBase`, the class `name` must contain the three methods `Initialize`, `Execute` and `Finalize`. Apart from this, the user is free to include his own set of additional functions and variables as required to perform his analysis.

The role of the three functions above is intuitive and related to their names. When `SAMPLEANALYZER` is executed in a shell, the program starts

by calling (once) the function `Initialize`

```
void Initialize()
```

Its aim is to initialize all the internal variables, as well as the entire set of user-defined variables such as the histograms that have been requested (and that must be properly declared in the header file `name.h`).

After initialization, `SAMPLEANALYZER` loops over all the events which are passed to the code. As shown in Section 5.1, the list of event files is collected into a single file which is passed as an argument when executing the program from a shell. For each event, the function `Execute` is called in order to perform the analysis. Among others, the histograms are filled event by event and the cuts are applied. The method,

```
void Execute(const SampleFormat& sample,
             const EventFormat& event)
```

takes as arguments, on the one hand, the event `event` under consideration, which is passed as a set of particles with specific properties, and on the other hand, general information on the entire event sample `sample` currently analyzed, such as the corresponding integrated cross section necessary for a proper normalization of the histograms to be produced.

Once all the events have been processed, the function `Finalize` is eventually called (once) by `SAMPLEANALYZER`

```
void Finalize(const SampleFormat& summary,
               const std::vector<SampleFormat>& files)
```

Its aim is twofold. Firstly, it allows for the creation of the histograms and cut-flow charts requested by the user. To implement the latter in an easy way, the user can employ all the functionalities included in the ROOT library and we therefore refer to the ROOT manual for more information [57]. Secondly, the method `Finalize` stores the results of the analysis as an output ROOT files which can be further processed or modified. The method takes as arguments general information related to the event samples, `summary`, such as the total cross section and the associated uncertainty. However, this time, instead of having this information linked to a specific event sample, an average over the whole list of samples included in the dataset under consideration has been performed. In the case the user wants to compute additional quantities when implementing the method `Finalize`, the same information, related to

each of the samples individually, is passed as the second argument which is denoted, in the example above, as `files`.

These three predefined functions have to be implemented by the user in the corresponding source file `name.cpp`. To this aim, we refer to the description of the internal data format used by SAMPLEANALYZER for event processing (see Section 5.3 and Section 5.4) and to the one of the methods included in the mother class `AnalysisBase` (see Section 5.5).

As mentioned above, several analyses can be included in the same working directory. The only requirement consists in implementing them in different files and classes since each analysis is unambiguously defined by its (file)name, which must therefore be unique. In order for SAMPLEANALYZER to correctly handle them, the user must refer to the various classes and files in the C++ source file `analysisList.cpp`, located in the sub-directory `SampleAnalyzer/Analysis`. This file contains a link to each of the header files associated to the analyses to be included. In addition, one instance of each analysis class is created. The architecture of this file follows the structure

```
#include "Analysis/name1.h"
#include "Analysis/name2.h"
...
#include "Core/AnalysisManager.h"
#include "Core/logger.h"

void AnalysisManager::BuildTable()
{
    Add(new name1);
    Add(new name2);
    ....
}
```

At least two analyses, denoted by `name1` and `name2` are implemented and the corresponding header files have been included in SAMPLEANALYZER. In the example above, the dots stand for possible additional analyses that the user might want to embed too. In the case the user has only implemented one single analysis, the file `analysisList.cpp` must still be present. It however then only includes the header file of this analysis and creates an instance of the related class.

In order to facilitate the implementation of new analyses, SAMPLEANALYZER comes with a PYTHON script `newAnalysis.py` located in the directory `SampleAnalyzer`. As shown above, creating a new analysis with the name `newname` requires the implementation of the two C++ files `newname.h` and `newname.cpp` and then update the file `analysisList.cpp` in order to include the new analysis. The analysis-independent part of this task has been automated through this PYTHON script. The user can use it from a shell, by typing

```
newAnalysis.py newname
```

The script starts by asking the user to type-in the label of the new analysis. As a result, the two files containing the declaration of the new analysis class are created (with a blank analysis included) and the list of the existing analyses in `analysisList.cpp` is updated. The user must now start implementing the analysis itself.

To this aim, he has to declare the variables necessary for the analysis and implement the three methods `Initialize`, `Execute` and `Finalize`, together with possible additional user-defined functions. This step requires a knowledge of the methods already included in the base class `AnalysisBase` and the one of the data format used internally by SAMPLEANALYZER. This is the scope of Section 5.3, Section 5.4 and Section 5.5.

### 5.3. The data format used by SAMPLEANALYZER

The function `Execute` introduced in the previous Section takes as a second argument an object of the type `EventFormat`, which points to the current event being analyzed denoted by `event` in the following. We dedicate this section to the description of the `EventFormat` class, which can also be found as a DOXYGEN documentation on the MADANALYSIS 5 website,

<http://madanalysis.irmp.ucl.ac.be>

In order to implement his analysis and to apply it to the event under consideration, the user generally needs to access various pieces of information related to this event. For example, the momentum of a specific type of object could be needed to implement some sophisticated cuts. To this aim, the SAMPLEANALYZER framework contains many built-in methods. They are split into two categories, the first one being related to hadron-level and parton-level events, which are generically named, in the following, as Monte Carlo level events, and the second one to reconstructed events.

When processing the event `event`, SAMPLEANALYZER automatically creates, when reading the event, an object with a structure appropriate to store the information included in the event under consideration,

```
event.mc()           event.rec()
```

These two objects are C++ pointers to all the methods implemented to facilitate the design of an analysis at the Monte Carlo level and at the reconstructed-level, respectively. These methods are extensively described in the rest of this Section.

### 5.3.1. The data format for parton-level or hadron-level events

The pointer `event.mc()` introduced above allows us to access general information related to a (Monte Carlo) event denoted by `event` such as the weight of the event or the employed value of the strong coupling constant. In addition, the whole set of initial-, intermediate- and final-state particles, together with their kinematical properties, is available. These properties form the so-called data format of SAMPLEANALYZER for Monte Carlo level events and are summarized in Table 12.

For Monte Carlo samples which include events related to several physical processes, matrix-element generators usually assign to each event a tag, *i.e.*, an unsigned integer number, that allows us to identify the physical process which the event is originating from. If the user needs to access this tag in the analysis, it is available through the function

```
event.mc()->processId()
```

which returns an unsigned integer. With a similar syntax, the weight associated to the event `event` can be used in the C++ source file of the analysis through,

```
event.mc()->weight()
```

which returns a floating-point number. If the implementation of the analysis requires the evaluation of the factorization scale, it can be obtained from the method

```
event.mc()->scale()
```

which gives the results as a floating-point number. In contrast, the renormalization scale is not available but the values of the strong and electromagnetic coupling constants (including a possible running) can also be employed in the analysis,

**Table 12: Methods related to the parton-level  
and hadron-level event format**

Let `ev` be an `EventFormat` object, *i.e.*, an instance of the class related to events issued from a partonic or hadronic Monte Carlo sample.

<code>ev.mc()&gt;alphaQCD()</code>	This returns, as a floating-point number, the value of the strong coupling constant used in the event.
<code>ev.mc()&gt;alphaQED()</code>	This returns, as a floating-point number, the value of the electromagnetic coupling constant used in the event.
<code>ev.mc()&gt;particles()</code>	This returns a vector of <code>MCParticleFormat</code> objects whose each entry consists in one of the particles present in the event, together with its properties. All initial, intermediate and final state particles are included.
<code>ev.mc()&gt;processId()</code>	This returns an unsigned integer number related to the tag of the physical process the event is originating from. It is especially useful when several physical processes are merged into one event sample.
<code>ev.mc()&gt;scale()</code>	This returns, as a floating-point number, the factorization scale.
<code>ev.mc()&gt;weight()</code>	This returns, as a floating-point number, the event weight.

```
event.mc()->alphaQCD()
event.mc()->alphaQED()
```

These last two methods also return floating-point numbers.

More importantly, implementing histograms or selection cuts requires us to investigate the particle content of the event, as well as the properties of one or several of these particles. All the initial-, intermediate- and final-state particles included in the event `event` are stored in a vector of `MCParticleFormat` objects, which can be called in the analysis through

```
event.mc()->particles()
```

The syntax above returns a vector that each entry consists in a particle present in the event, together with its properties, given as an instance of the `MCParticleFormat` class.

In order to implement a loop over all the particle content of the event `event`, in, *e.g.*, the function `Execute` of the analysis source file, it is sufficient to program

```
unsigned int n = event.mc()->particles().size();
for (unsigned int i=0; i<n; i++) { ... }
```

where we recall that at this stage, all the initial-, intermediate- and final-state particles are considered equivalently. We will show in Section 5.5 how to implement loops over, *e.g.*, the final-state particles only. Similarly, denoting by the object `prt` an instance of the `MCParticleFormat` class, the  $i^{\text{th}}$  particle is given by

```
MCParticleFormat* prt = &event.mc()->particles()[i];
```

In the rest of this Section, we focus on the attributes of the object `prt`.

The class `MCParticleFormat` contains seven methods which allow to extract and use the properties of a given particle within the analysis. These methods are summarized in Table 13.

The PDG-id of the particle `prt` can be accessed through

```
prt->pdgid()
```

which returns the corresponding signed integer number. The `statuscode` of the particle, *i.e.*, its tag as an initial-state, intermediate-state or final-state particle, can be obtained through the function

**Table 13: Methods related to the MCParticleFormat class**

Let `prt` be a `MCParticleFormat` object.

<code>prt.ctau()</code>	This returns, as a floating-point number, the decay length of the particle, assuming that it moves at the speed of light.
<code>prt.momentum()</code>	This returns a <code>T LorentzVector</code> containing the four-momentum of the particle <code>prt</code> .
<code>prt.mother1()</code>	This returns a <code>MCParticleFormat</code> object pointing to the mother particle of the particle <code>prt</code> , <i>i.e.</i> , either the particle which decays into <code>prt</code> (plus other particles) or a particle which interacts with another particle (see <code>mother2()</code> ) to produce the particle <code>prt</code> .
<code>prt.mother2()</code>	This returns a <code>MCParticleFormat</code> object pointing to the mother particle of the particle <code>prt</code> , but only in the case <code>prt</code> is produced from the interaction of two particles. These particles can be accessed through the two methods <code>mother1()</code> and <code>mother2()</code> .
<code>prt.pdgid()</code>	This returns an integer number standing for the PDG-id of the particle.
<code>prt.spin()</code>	This returns a floating-point number, the cosine of the angle between the momentum of the particle <code>prt</code> and its spin vector, computed in the laboratory reference frame.
<code>prt.statuscode()</code>	This returns an integer depending on the initial-state, intermediate-state or final-state nature of the particle. The conventions on this integer number are taken from Ref. [41].

```
prt->statuscode()
```

which returns an integer associated to the initial-state, intermediate-state or final-state nature of the particle. For the conventions on this integer number we refer to Ref. [41].

For non-initial-state particles, information on mother particle(s) can be extracted (and used in the analysis) through the two methods

```
prt->mother1()  
prt->mother2()
```

which return the `MCParticleFormat` objects related to the particle(s) from which the current particle `prt` is issued. In the case of a decay chain, only the `mother1` method has to be used. In contrast, when the particle `prt` is produced from the interaction of two particles, both the `mother1` and `mother2` give results, each of them pointing to one of the initial particles.

The most important property of the `MCParticleFormat` class to be used in physics analyses consists in the particle four-momentum, accessible from

```
prt->momentum()
```

The result of this function is given as a `T LorentzVector`, a ROOT class appropriate to store four-momentum. In addition, this class contains various methods to compute a large set of observables which can be derived from the knowledge of the four-momentum, such as the energy or the transverse momentum. The complete list of these observables is given, together with the associated syntax, in Table 14.

An additional observable related to the four-momentum, and more in particular to the three-momentum component of the four-momentum, which can be employed in an analysis reads

```
prt->spin()
```

This returns a floating-point number which stands for the cosine of the angle between the momentum of the particle `prt` and its spin vector, evaluated in the laboratory reference frame.

Finally, the decay length of the particle (assuming that the particle is moving at the speed of light) can also be easily obtained through the function

```
prt->ctau()
```

which returns a floating-point number too.

**Table 14: Common methods related to the MCParticleFormat and RecParticleFormat classes**

Let  $P$  be a MCParticleFormat or RecParticleFormat object.

$P.\text{angle}(P2)$	Angle between the momenta of the objects $P$ and $P2$ , where $P2$ is an instance of the MCParticleFormat or RecParticleFormat class.
$P.\text{beta}()$	Velocity $\beta = v/c$ .
$P.\text{dr}(P2)$	Relative distance between the objects $P$ and $P2$ in the $\eta - \phi$ plane, where $P2$ is an instance of the MCParticleFormat or RecParticleFormat class.
$P.\text{e}()$	Energy.
$P.\text{et}()$	Transverse energy.
$P.\text{eta}()$	Pseudorapidity.
$P.\text{gamma}()$	Lorentz factor.
$P.\text{m}()$	Invariant mass.
$P.\text{mt}()$	Transverse mass.
$P.\text{p}()$	Norm of the momentum.
$P.\text{phi}()$	Azimuthal angle of the momentum.
$P.\text{pt}()$	Norm of the transverse momentum.
$P.\text{px}()$	Projection of the momentum on the $x$ -axis.
$P.\text{py}()$	Projection of the momentum on the $y$ -axis.
$P.\text{pz}()$	Projection of the momentum on the $z$ -axis.
$P.\text{r}()$	Position of the object in the $\eta - \phi$ plane.
$P.\text{theta}()$	Angle between the momentum and the beam axis.
$P.\text{y}()$	Rapidity.

**Table 15: Methods related to the reconstructed event format**

Let `ev` be an `EventFormat` object, *i.e.*, an instance of the class related to events issued from a reconstructed Monte Carlo sample.

<code>ev.rec()&gt;electrons()</code>	This returns a vector with all the electrons of the event, encoded as <code>RecLeptonFormat</code> objects.
<code>ev.rec()&gt;jets()</code>	This returns a vector with all the jets of the event, given as <code>RecJetFormat</code> objects.
<code>ev.rec()&gt;met()</code>	This points to the missing energy of the event, encoded as a <code>RecMETFormat</code> object.
<code>ev.rec()&gt;muons()</code>	This returns a vector with all the muons of the event, encoded as <code>RecLeptonFormat</code> objects.
<code>ev.rec()&gt;taus()</code>	This returns a vector with all the taus of the event, given as <code>RecTauFormat</code> objects.

### 5.3.2. The data format for reconstructed events

At the beginning of this Section, we have introduced two types of data format which are used for event processing by `SAMPLEANALYZER`. They consist in the two sides of the more general `EventFormat` class. In the previous Section, we have focused on the description of an event object `event` read from a parton-level or hadron-level sample. We have shown that its properties are embedded, in the `SAMPLEANALYZER` framework, into a structure which can be browsed from the C++ pointer `event.mc()`.

Since the properties of reconstructed events are in general very different from those of Monte Carlo events, `SAMPLEANALYZER` embeds the latter into another structure which is, this time, linked to the C++ pointer `event.rec()`. It consists in five methods, collected in Table 15, associated to the five types of objects that can be reconstructed, *i.e.*, electrons, muons, taus, jets and missing energy<sup>6</sup>. In the `SAMPLEANALYZER` framework, each

---

<sup>6</sup>By electrons, muons and taus, we are considering both the corresponding particles

reconstructed objects is associated to a specific class derived from the generic `RecParticleFormat` mother class and depending on its species. In the following, we adopt the choice of describing the daughter classes, more relevant for the user, rather than the generic class.

Two methods are associated to first and second generation charged leptons, *i.e.*, to the reconstructed electrons and muons, present in the event `event`,

```
event.rec()->electrons()  
event.rec()->muons()
```

They return vectors that the entries are the different reconstructed electrons and muons of the event, respectively. Each reconstructed electron or muon is encoded as a `RecLeptonFormat` object. This last class has six associated methods, gathered in Table 16, related to the attributes of the reconstructed leptonic electron and muon objects.

As for Monte Carlo event samples, the four-momentum of the reconstructed objects is one of the most incontrovertible variables to be employed in analyses. For a reconstructed lepton `lep`, its four-momentum can be included and used in the analysis source file by calling the function

```
lep.momentum()
```

the syntax being similar to the one employed for particles included in parton-level and hadron-level events. This method returns a `TLorentzVector` object containing the corresponding four-momentum. Therefore, the methods associated to all the observables which can be derived from the knowledge of the four-momentum, collected in Table 14, are also methods of the `RecLeptonFormat` class.

Reconstructed leptons are considered as electrons or muons regardless of their charge. In the case the user wants to select them according to the electric charge, the method

```
lep.charge()
```

of the `RecLeptonFormat` class can be used. It returns, as a floating-point number, the electric charge of the reconstructed lepton which can be equal to either  $-1$  or  $+1$  according to its antiparticle or particle nature.

---

and antiparticles.

**Table 16: Methods related to the `RecLeptonFormat` class**

Let `lep` be an `RecLeptonFormat` object, *i.e.*, an instance of the class describing the reconstructed electrons and muons.

<code>lep.charge()</code>	This returns the electric charge of the reconstructed object <code>lep</code> as a floating-point number. The returned value consists in $-1$ or $+1$ .
<code>lep.EEoverHE()</code>	This returns the ratio of the electromagnetic and hadronic energy for the reconstructed object <code>lep</code> , given as a floating-point number.
<code>lep.ET_PT_isol()</code>	This returns the ratio of the values of the functions <code>sumET_isol()</code> and <code>sumPT_isol()</code> .
<code>lep.HEoverEE()</code>	This returns the ratio of the hadronic and electromagnetic energy for the reconstructed object <code>lep</code> , given as a floating-point number.
<code>lep.momentum()</code>	This returns a <code>TLorentzVector</code> containing the four-momentum of the reconstructed object <code>lep</code> .
<code>lep.sumET_isol()</code>	This returns, if the object <code>lep</code> is a muon, the sum of the transverse energy of all the tracks lying in a cone around the muon. The size of the cone is fixed by the detector simulation tool. If <code>lep</code> is an electron, this function returns zero.
<code>lep.sumPT_isol()</code>	This returns, if the object <code>lep</code> is a muon, the sum of the transverse momentum of all the tracks lying in a cone around the muon. The size of the cone is fixed by the detector simulation tool. If <code>lep</code> is an electron, this function returns zero.

All the methods presented in Table 14 can also be used.

Two methods allow us to get information on the splitting of the reconstructed electron or muon energy between the electromagnetic and the hadronic parts of the detector,

```
lep.EOverHE()
lep.HOverEE()
```

These methods compute the ratio between the energy deposited in the electromagnetic calorimeter and the one deposited in the hadronic calorimeter, and *vice-versa*, and return them as floating-point numbers.

Finally, the `RecLeptonFormat` class contains three specific methods related to muon isolation (see also Section 5.5). The algorithms to be employed for deciding if a muon is considered as isolated or not require in general the evaluation of two quantities, the sum of the transverse momentum of all tracks lying in a cone around the muon candidate and the sum of their transverse energy. These two observables can be accessed by typing

```
lep.sumPT_isol()
lep.sumET_isol()
```

which return a zero value in the case the lepton `lep` is an electron. In contrast, for muons, the value read from the event file is employed. The size of the cone is fixed by the fast detector simulation tool and is not available in the LHCO event format. The ratio of the values of these two functions can be obtained via the function

```
lep.ET_PT_isol()
```

Tau leptons being unstable, they always decay either into a narrow jet, into a muon or into an electron, each time in association with missing energy. Therefore, a specific class, different from the `RecLeptonFormat` class, exists in order to embed reconstructed taus. This class is denoted by `RecTauFormat`. In the internal data format used by `SAMPLEANALYZER`, the pointer to the reconstructed event, `event.rec()`, contains a specific method to access all the taus present in the event

```
event.rec()->taus()
```

This returns, as a vector of `RecTauFormat` objects, all the reconstructed tau leptons of the event (regardless of their electric charge).

**Table 17: Methods related to the `RecTauFormat` class**

Let `tau` be an instance of the `RecTauFormat` class, *i.e.*, the class describing the reconstructed taus.

<code>tau.charge()</code>	This returns the electric charge of the object <code>tau</code> as a floating-point number. The returned value consists in $-1$ or $+1$ .
<code>tau.EEoverHE()</code>	This returns the ratio of the electromagnetic and hadronic energy for the object <code>tau</code> , given as a floating-point number.
<code>tau.HEoverEE()</code>	This returns the ratio of the hadronic and electromagnetic energy for the object <code>tau</code> , given as a floating-point number.
<code>tau.momentum()</code>	This returns a <code>T LorentzVector</code> containing the four-momentum of the object <code>tau</code> .
<code>tau.ntracks()</code>	This returns, as a short integer, the number of charged tracks contained in the object <code>tau</code> .

All the methods presented in Table 14 can also be used.

In addition to the four methods `momentum()`, `charge()`, `EEoverHE()`, `HEoverEE()` and the methods given in Table 14 already present for charged leptons of the first and second generations, the `RecTauFormat` class includes an additional method extracting the number of charged tracks issued from the decaying tau and included in the reconstructed object. Denoting by `tau` the reconstructed object, this number, given as a short integer, can be obtained and further used in the implementation of the analysis by

```
tau.ntracks()
```

The entire list of methods associated to the `RecTauFormat` class can be found in Table 17.

As for electrons, muons and taus, the entire set of reconstructed jets

**Table 18: Methods related to the `RecJetFormat` class**

Let `j` be an `RecJetFormat` object, *i.e.*, an instance of the class describing the reconstructed jets.

<code>j.btag()</code>	This returns <code>true</code> or <code>false</code> according to the <i>b</i> -tagging (or not) of the object <code>j</code> .
<code>j.EEoverHE()</code>	This returns the ratio of the electromagnetic and hadronic energy for the object <code>j</code> , given as a floating-point number.
<code>j.HEoverEE()</code>	This returns the ratio of the hadronic and electromagnetic energy for the object <code>j</code> , given as a floating-point number.
<code>j.momentum()</code>	This returns a <code>TLorentzVector</code> containing the four-momentum of the object <code>j</code> .
<code>j.ntracks()</code>	This returns, as a short integer, the number of charged tracks contained in the object <code>j</code> .

All the methods presented in Table 14 can also be used.

included in the event `event` can be obtained through the intuitive method of the `event.rec()` structure,

```
event.rec()->jets()
```

This returns a vector of instances of the `RecJetFormat` class whose properties are collected in Table 18.

Reconstructed jets are characterized by their momentum, the number of charged tracks induced by the parton showering, fragmentation and hadronization of the initial partons and the ratio of the hadronic and electromagnetic parts of the jet energy. These properties are related to the methods `momentum()`, `ntracks()`, `EEoverHE()` and `HEoverEE()` as well as all those included in Table 14, of the `RecJetFormat` class. In addition, an extra method returns `true` or `false` according to the fact that the jet has been tagged as a *b*-jet or not,

**Table 19: Methods related to the RecMetFormat class**

Let `miss` be a `RecMetFormat` object, *i.e.*, the variable containing the reconstructed missing energy associated to an event.

<code>miss.mag()</code>	This returns the magnitude of the missing energy represented by the object <code>miss</code> as a floating-point number.
<code>miss.phi()</code>	This returns the azimuthal angle of the missing energy represented by the object <code>miss</code> as a floating-point number.
<code>miss.x()</code>	This returns the $x$ -component of the missing energy represented by the object <code>miss</code> as a floating-point number.
<code>miss.y()</code>	This returns the $y$ -component of the missing energy represented by the object <code>miss</code> as a floating-point number.

`j.btag()`

where `j` denotes an instance of the `RecJetFormat` class.

The last type of objects which are included in reconstructed events consists in the associated missing transverse energy. The structure `event.rec()` comes with the related method

`event.rec() -> met()`

which returns a two-dimensional vector implemented as a `TVector2` object. It contains then the direction and magnitude of the missing energy in the transverse plane as shown in Table 19.

The  $x$ -component and  $y$ -component of the missing transverse energy can be used when implementing an analysis through the two methods

`miss.x()`  
`miss.y()`

which return floating-point numbers associated to the two components of the two-dimensional vector describing the missing energy, the object `miss` being an instance of the `RecMETFormat` class. In the case the user prefers to use polar coordinates instead of Cartesian coordinates when implementing his analysis, he can use the two methods

```
miss.mag()
miss.phi()
```

which return the magnitude and the azimuthal angle of the two-dimensional transverse-momentum vector as floating-point numbers.

#### *5.4. The sample format used by SAMPLEANALYZER*

In Section 5.2, we have introduced the function `Execute` as the main function of the analysis class. We have shown that it requires two arguments, the event being currently analyzed, stored as an `EventFormat` object, and general information about the sample which this event belongs to, passed as a `SampleFormat` object. This format is also the one to be used for the arguments of the function `Finalize`, dedicated to the creation of the output files containing the results of the analysis. In this case, the user must pass two arguments, an instance of the `SampleFormat` class associated to all the event samples included in the dataset under consideration, *i.e.*, information averaged over all the samples, and one vector of `SampleFormat` objects with one entry for each of the samples included in the dataset, *i.e.*, the same information, but given for each individual sample.

The `SampleFormat` class contains several methods allowing us to access general information about the sample under consideration, provided the information is available. If not, the associated entries in the `SampleFormat` structure are left non-initialized and the corresponding quantity cannot consequently be used in an analysis. A full DOXYGEN documentation is available on the MADANALYSIS 5 website,

<http://madanalysis.irmp.ucl.ac.be>

Since information supplementing the events is only available within LHE and STDHEP files, the `SampleFormat` structure is then only relevant for the analysis of such event files. In particular, LHCO files with reconstructed events do not include anything but the events. In this case, it is however still possible to pass additional information such as cross sections to SAMPLEANALYZER through the attributes of the `dataset` class (see Section 4.3).

In Section 5.3, we have shown that `SAMPLEANALYZER` is creating an `EventFormat` structure each time an event is processed, resulting in a C++ pointer pointing to the whole methods available for the `EventFormat` structure. Similarly, when processing a new sample generically denoted by `sample`, `SAMPLEANALYZER` creates a C++ pointer

```
sample.mc()
```

which points to a structure containing all the methods allowing us to access the global information of the event file. These methods are collected in Table 20 and Table 21.

As shown by its name, `sample.mc()` is related to parton-level or hadron-level events. The counterpart of this object in the case of a sample containing reconstructed events, `sample.rec()`, has been implemented within the `SAMPLEANALYZER` framework. However, the only event file format appropriate for reconstructed events, the LHCO format, does not leave a possibility for including additional information to the events. Therefore, the pointer `sample.rec()` points to a set of null information. If in the future, a new format for reconstructed events is designed, with the room for global information about the event sample, the related structure in the `SAMPLEANALYZER` framework will be ready for the new format.

The first series of methods available within the `SampleFormat` structure, collected in Table 20, are related to the description of the initial colliding beams. The two members of the `SampleFormat` class

```
sample.mc() -> beamPDGID().first
sample.mc() -> beamPDGID().second
```

return, as integer numbers, the PDG-id of the first and second beams, respectively, whilst the four class members

```
sample.mc() -> beamPDFauthor().first
sample.mc() -> beamPDFID().first
sample.mc() -> beamPDFauthor().second
sample.mc() -> beamPDFID().second
```

return, as four unsigned integer numbers, information with respect to the set of parton density functions used for the beams. These identifying numbers are exported from the event samples (if available) and correspond to the numbering scheme employed by matrix element generators to identify

**Table 20: Methods related to the SampleFormat class giving information on the colliding beams**

Let `sample` be a `SampleFormat` object.

`sample.mc()>beamE().first`

This returns, as a floating-point number, the energy of the first of the colliding beams.

`sample.mc()>beamPDFauthor().first`

This returns the author group of the parton density set used for the first of the colliding beams, as an unsigned integer number.

`sample.mc()>beamPDFID().first`

This returns the identifier, as an unsigned integer number, of the parton density set used for the first of the colliding beams, within a given author group of parton densities.

`sample.mc()>beamPDGID().first`

This returns the PDG-id of the first of the colliding beams, as an integer number.

`sample.mc()>beamE().second`

This returns, as a floating-point number, the energy of the second of the colliding beams.

`sample.mc()>beamPDFauthor().second`

This returns the author group of the parton density set used for the second of the colliding beams, as an unsigned integer number.

`sample.mc()>beamPDFID().second`

This returns the identifier, as an unsigned integer number, of the parton density set used for the second of the colliding beams, within a given author group of parton densities.

`sample.mc()>beamPDGID().second`

This returns the PDG-id of the second of the colliding beams, as an integer number.

**Table 21: Other methods related to the SampleFormat class**

Let `sample` be a `SampleFormat` object.

`sample.mc()>weightingmode()`

This returns, as an integer number, information on the type of events present in the sample (*e.g.*, weighted versus unweighted events).

`sample.mc()>processes()`

This returns a vector where the entries are `ProcessFormat` objects containing basic information about each of the physical processes included in the sample `sample`.

`sample.mc()>xsection()`

This returns the cross section associated to the event sample `sample`, as a floating-point number.

`sample.mc()>xsection_error()`

This returns the Monte Carlo uncertainty related to the cross section associated to the event sample `sample`, as a floating-point number.

the set of parton densities employed. According to the Les Houches conventions, this scheme is based on the PDFLIB [61] and LHAPDF [62] packages. Hence, the method `beamPDFauthor()` is related to the author group which has released the parton density relevant for the sample under consideration and the method `beamPDFID()` indicates which specific set of parton densities of the corresponding group has been used.

The last pieces of information available with respect to the beams which can be stored in the event files, and thus exported to the SAMPLEANALYZER framework, consist in their energy. This can be used in the implementation of an analysis by means of the two methods

```
sample.mc()> beamE().first
sample.mc()> beamE().second
```

for the first and second beams, respectively.

The second series of methods available from the `SampleFormat` class are related to the sample itself. When available, information about the weighting of the events is passed to `SAMPLEANALYZER` and can be accessed through

```
sample.mc() -> weightingmode()
```

According to the LHE format [41, 42], this type of information is stored as an integer number which indicates how event weights and cross sections have to be interpreted, and the `weightingmode()` method returns this integer number. Even if presently, `SAMPLEANALYZER` does not fully support weighted events, the architecture has already been implemented with respect to future developments of the program. Finally, the most important quantities included in this second series of methods concern the cross section associated to the sample `sample`, together with the corresponding uncertainty. Both can be called at the level of the analysis by using the methods

```
sample.mc() ->xsection()  
sample.mc() ->xsection_error()
```

which return floating-point numbers.

As mentioned in Section 5.3.1, several physical processes can be included within the same event sample. In this case, each of the processes is associated with an identifying tag (see the `processId` method introduced in Section 5.3.1). The `SampleFormat` structure allows us to easily access detailed information about the processes individually. This information is stored into a new structure, the `ProcessFormat` class, all of whose associated methods are summarized in Table 22. The method

```
sample.mc() ->processes()
```

returns a vector that each entry is a `ProcessFormat` object. For an instance of this class denoted by `process`, information on the identifier of the process can be obtained through

```
process.processId()
```

the corresponding cross section together with its associated uncertainty through the methods

```
process.xsection()  
process.xsection_error()
```

**Table 22: Methods related to the ProcessFormat class**

Let `proc` be an instance of the `ProcessFormat` class.

`proc.processId()`

This returns an unsigned integer number related to the tag of the physical process the event is originating from.

`proc.xsection()`

This returns the cross section associated to the process `proc`, as a floating-point number.

`proc.xsection_error()`

This returns the Monte Carlo uncertainty related to the cross section associated to the process `proc`, as a floating-point number.

`proc.maxweight()`

This returns the maximum weight carried by an event associated to the process `proc`, as a floating-point number.

and finally, the maximum weight carried by a single event originating from the subprocess `process` through

`process.maxweight()`

### 5.5. Framework services

In this Section, we describe the various services included in the SAMPLEANALYZER framework. These services are components of the program that are initialized (internally or by the user) at the beginning of the execution of SAMPLEANALYZER (within the `Initialize` method) and can then be further called within the analysis as many times as necessary. Two series of services are currently available, message services and physics services and are described in the rest of this Section. We recall that full DOXYGEN documentation is available on the MADANALYSIS 5 website,

<http://madanalysis.irmp.ucl.ac.be>

### 5.5.1. Message services

This class of functions has been implemented and can be used by the users in order to print text to the screen during an on-going analysis in a rather sophisticated fashion.

As for the implementation of any C++ program, the user can, when designing his analysis, use the standard C++ streamers `std::cout` and `std::cerr` in order to implement messages to be printed to the screen. However, these methods only include two levels of messages, *i.e.*, normal and error messages. It is hence rather cumbersome to handle multi-level messages and give information both on the reason leading to the printing of the message and on its location in the analysis code in a clear and useful way.

Therefore, `SAMPLEANALYZER` includes its own message services with four different levels of printing, `DEBUG`, `INFO`, `WARNING` and `ERROR`. The way to use these four modes mimics the one of the C++ commands `std::cout` and `std::cerr`,

```
DEBUG    << "Debug message." << std::endl;
INFO     << "Information message." << std::endl;
WARNING << "Warning message." << std::endl;
ERROR    << "Error message." << std::endl;
```

Each message level is associated to a different color. Debugging messages are printed in yellow, information messages in white, warning messages in purple and error messages in red. However, if the user is not interested in the color of the messages, this message can be switched off by including, in the source code of the analysis, the line

```
LEVEL->DisableColor()
```

The color can be enabled again through the command

```
LEVEL->EnableColor()
```

In the command lines above, the keyword `LEVEL` stands for any of the four levels of message, `DEBUG`, `INFO`, `WARNING` or `ERROR`.

Warning and error messages have a special role concerning the debugging of the analysis code. In addition to the message, the line number of the code having generated the message is also printed in order to facilitate the debugging.

Finally, let us note that messages associated to a given level can be fully switched off, if this is needed by the user, by including in the analysis code the command

```
LEVEL->Mute()
```

Messages can be restored by implementing

```
LEVEL->UnMute()
```

where the keyword LEVEL again stands for any of the four levels of message.

### 5.5.2. Physics services

Under the name *physics services*, we collect a series of methods and functions aiming to facilitate the writing of an analysis by the user. It includes, among others, functions to compute global observables related to the entire event, such as the transverse missing energy, and this for any type of event (parton-level, hadron-level and reconstructed-level). All these functions can be called through the C++ pointer PHYSICS.

Before moving on to the description of these functions, one must note that in the case the user wants to use, within his analysis, one or several of the methods related to the invisible or hadronizing particles, the definition of these invisible and hadronic particles must first be provided. This task is performed within the `Initialize` function introduced in Section 5.2 with the help of two different functions, one being associated to the invisible particles and another one being associated to the particles taking part in the hadronic activity. However, before declaring a new particle as invisible or hadronizing, the physics services class must be initialized by means of one of the commands

```
PHYSICS->mcConfig()->Reset()
PHYSICS->recConfig()->Reset()
```

In the case the event samples which are analyzed consist in parton-level and hadron-level event files, the method `mcConfig` is used whilst for reconstructed-level event files, `recConfig` is instead used. Then, a particle whose PDG-id is PDGID can be declared as an invisible particle by means of

```
PHYSICS->mcConfig()->AddInvisibleId(PDGID)
```

In the non-expert mode of MADANALYSIS 5, the definition of the multiparticle **invisible** (see Section 4.4) corresponds to several calls of the function above behind the scenes. Taking the example of an analysis at the parton-level in the context of the Minimal Supersymmetric Standard Model, the invisible particles consist of the neutrinos and the lightest superpartner assumed to be the lightest neutralino. The corresponding declaration within the SAMPLEANALYZER framework reads

```
PHYSICS->mcConfig().AddInvisibleId(-16);
PHYSICS->mcConfig().AddInvisibleId(-14);
PHYSICS->mcConfig().AddInvisibleId(-12);
PHYSICS->mcConfig().AddInvisibleId(12);
PHYSICS->mcConfig().AddInvisibleId(14);
PHYSICS->mcConfig().AddInvisibleId(16);
PHYSICS->mcConfig().AddInvisibleId(1000022);
```

Similarly, the method

```
PHYSICS->mcConfig()->AddHadronicId(PDGID)
```

declares the particle whose PDG-id is given by PDGID as a particle related to the hadronic activity in an event. Again, in the non-expert mode of MADANALYSIS 5, the definition of the multiparticle **hadronic** (see Section 4.4) corresponds to a series of calls to this last function. For the example of a parton-level analysis within the Standard Model, the related declaration would be given by

```
PHYSICS->mcConfig().AddHadronicId(-5);
PHYSICS->mcConfig().AddHadronicId(-4);
PHYSICS->mcConfig().AddHadronicId(-3);
PHYSICS->mcConfig().AddHadronicId(-2);
PHYSICS->mcConfig().AddHadronicId(-1);
PHYSICS->mcConfig().AddHadronicId(1);
PHYSICS->mcConfig().AddHadronicId(2);
PHYSICS->mcConfig().AddHadronicId(3);
PHYSICS->mcConfig().AddHadronicId(4);
PHYSICS->mcConfig().AddHadronicId(5);
PHYSICS->mcConfig().AddHadronicId(21);
```

**Table 23: Methods related to the initialization  
of the physics services**

`PHYSICS->mcConfig()->AddHadronic(PDGID)`

This adds the particle whose PDG-id is PDGID to the list of the invisible particles.

`PHYSICS->mcConfig()->AddInvisible(PDGID)`

This adds the particle whose PDG-id is PDGID to the list of the particles taking part of the hadronic activity of an event.

`PHYSICS->mcConfig()->Reset()`

This initializes physics services when analyzing parton-level or hadron-level events.

`PHYSICS->recConfig()->Reset()`

This initializes physics services when analyzing reconstructed-level events.

In the case of events at the reconstructed-level, the user may need to specify the algorithm to be employed when testing the isolation of a muon. This can be done through the two (self-excluding) methods

```
PHYSICS->recConfig().UseDeltaRIsolation(deltaR=0.5);
PHYSICS->recConfig().UseSumPTIsolation(sumPT, ET_PT);
```

In the first case, a muon is tagged as isolated when no track lies inside a cone of size `deltaR` around the muon. The default size of this cone is set to 0.5. In the second case, the muon is tagged as isolated when the sum of the transverse momentum of all the tracks lying in a cone around the muon is lower than `sumPT` and in addition, when the sum of the transverse energy of these tracks over the sum of their transverse momentum is lower than `ET_PT`. By default, the first algorithm is adopted with `deltaR` being set to 0.5.

The usage of all the methods associated to the initialization of the physics services, *i.e.*, to the `PHYSICS->mcConfig()` and `PHYSICS->recConfig()` objects, is summarized in Table 23. All the other methods included in the physics services are in general called within the function `Execute`, at the core of the analysis and are summarized in Table 24 and Table 25.

Three boolean functions exist in order to test if a particle `prt` is invisible or visible as well as if this particle takes part in the hadronic activity of the event or not,

```
PHYSICS->IsHadronic(prt)
PHYSICS->IsInvisible(prt)
PHYSICS->IsVisible(prt)
```

When analyzing partonic or hadronic event samples, the particle `prt` is passed as a `MCParticleFormat` object. In contrast, for analyses at the reconstructed-level, `prt` is an instance of the `RecParticleFormat` class<sup>7</sup>. In order for these three methods to correctly work, it is necessary to declare which particle is invisible and which particle takes part in the hadronic activity as shown above. In the reconstructed-level case, an additional method exists to test whether a muon is isolated within an event `event`,

```
PHYSICS->IsIsolatedMuon(prt, event)
```

which returns always `false` for particles which are not muon candidates. In the case of muons, this method applies the isolation algorithm chosen by the user when setting the `PHYSICS->recConfig()` properties.

Another set of three methods checks whether a given particle `prt` is a final-state, initial-state or intermediate-state particle,

```
PHYSICS->IsFinalState(prt)
PHYSICS->IsInitialState(prt)
PHYSICS->IsInterState(prt)
```

These functions all return a boolean value according to the final-state, initial-state or intermediate-state nature of the particle under consideration. As above, those methods work equivalently for analyses at the hadronic, partonic and reconstructed levels, the particle `prt` being hence either an instance of the `MCParticleFormat` or of the `RecParticleFormat` classes. Since status codes are defined in a different fashion according to the event file format, we have adopted the choice to include these features within the physics services rather than the data format itself, which allows us to have a unified way to probe the initial-, intermediate- or final-state nature of the particles included

---

<sup>7</sup>The `RecParticleFormat` class is the mother class of all the classes defining reconstructed objects, *i.e.*, the `RecLeptonFormat`, `RecJetFormat` and `RecTauFormat` classes.

**Table 24: Boolean methods included in the physics services**

Let `prt` be an instance of either the `MCParticleFormat` or the `RecParticleFormat` classes.

`PHYSICS->IsFinalState(prt)`

This checks if the particle `prt` is a final-state particle.

`PHYSICS->IsHadronic(prt)`

This checks if the particle `prt` takes part to the hadronic activity in an event.

`PHYSICS->IsInitinalState(prt)`

This checks if the particle `prt` is an initial-state particle.

`PHYSICS->IsInterState(prt)`

This checks if the particle `prt` is an intermediate-state particle.

`PHYSICS->IsInvisible(prt)`

This checks if the particle `prt` is an invisible particle.

`IsIsolatedMuon(prt, evt)`

This checks if the particle `prt` is a muon isolated from the other particles of the event `evt`.

`PHYSICS->IsVisible(prt)`

This checks if the particle `prt` is a visible particle.

in an event. These last series of functions allows us to implement efficiently a loop over, *e.g.*, the final-state particles of the event, and not on all the particles, in contrast to the example given in Section 5.3.1,

```
unsigned int n = event.mc()->particles().size();
for (unsigned int i=0; i<n; i++)
{
    MCParticleFormat* prt = &event.mc()->particles()[i];
    if(PHYSICS->IsFinalState(prt))
        { ... }
}
```

Five methods included in the physics services are dedicated to the computation of global observables related to the entire event content: the missing

transverse energy  $\cancel{E}_T$ , the missing hadronic energy  $\cancel{H}_T$ , the total transverse energy  $E_T$  and the total hadronic energy  $H_T$  as defined in Eq. (1) and Eq. (2), as well as the partonic center-of-mass energy. The associated functions read, respectively,

<code>PHYSICS-&gt;EventMET(evt-&gt;mc())</code>	<code>PHYSICS-&gt;EventMET(evt-&gt;rec())</code>
<code>PHYSICS-&gt;EventMHT(evt-&gt;mc())</code>	<code>PHYSICS-&gt;EventMHT(evt-&gt;rec())</code>
<code>PHYSICS-&gt;EventTET(evt-&gt;mc())</code>	<code>PHYSICS-&gt;EventTET(evt-&gt;rec())</code>
<code>PHYSICS-&gt;EventTHT(evt-&gt;mc())</code>	<code>PHYSICS-&gt;EventTHT(evt-&gt;rec())</code>
<code>PHYSICS-&gt;SqrtS(event-&gt;mc())</code>	<code>PHYSICS-&gt;SqrtS(event-&gt;rec())</code>

where `evt` is an instance of the `EventFormat` class. These five methods work for any level of analysis (partonic, hadronic and reconstructed) and then take accordingly as argument either an `event->mc()` or an `event->rec()` object. The value, in GeV, of the corresponding observable is returned as a floating-point number.

For most analyses, it is important to be able to order the particles according to a specific variable. Therefore, in order to avoid requiring the user to implement his own sorting algorithm, physics services include one dedicated function,

```
PHYSICS->sort(prtVect, OrderingObs)
```

The command above allows us to sort a vector of particles `prtVect`, that each entry is either a `MCParticleFormat` or a `RecParticleFormat` object, according to the ordering observable `OrderingObs`. The implemented choices for the latter are `ETAordering` (pseudorapidity ordering), `ETordering` (transverse-energy ordering), `Eordering` (energy ordering), `Pordering` (momentum ordering), `PTordering` (transverse-momentum ordering), `PXordering` (ordering according to the  $x$ -component of the momentum), `PYordering` (ordering according to the  $y$ -component of the momentum) and `PZordering` (ordering according to the  $z$ -component of the momentum).

The last method included in the physics services that can be useful when implementing an analysis is related to the reference frame in which the four-momenta of the particles included in the event are computed. When reading the event files, all the four-momenta are, by convention, given in the laboratory reference frame. However, for specific observables, it is necessary to recalculate one or several of the four-momenta in the rest frame of a given particle, which is equivalent to applying a Lorentz boost to these four-vectors. This task can be done automatically by means of the method

**Table 25: Other methods included in the physics services**

Let `evt` be an instance of the `EventFormat` class assuming to point to an `evt->mc()` or a `evt->rec()` structure.

`PHYSICS->ToRestFrame(prt,prt1)`

This recalculates (and modifies) the four-momentum of the particle `prt` after a Lorentz boost to the rest frame of the particle `prt1`. The objects `prt` and `prt1` are two instances of the `MCParticleFormat` or `RecParticleFormat` classes.

`PHYSICS->EventMET(evt)`

This computes the missing transverse energy  $\cancel{E}_T$  associated to the event `evt`.

`PHYSICS->EventMHT(evt)`

This computes the missing transverse hadronic energy  $\cancel{H}_T$  associated to the event `evt`.

`PHYSICS->EventTET(evt)`

This computes the total transverse energy  $E_T$  associated to the event `evt`.

`PHYSICS->EventTHT(evt)`

This computes the total transverse hadronic energy  $H_T$  associated to the event `evt`.

`PHYSICS->sort(prtvector,orderobs)`

This sorts a vector of `MCParticleFormat` or `MCRecFormat` objects according to the ordering observable `orderobs`. The allowed choices for the ordering variable are `ETAordering` (pseudorapidity ordering), `ETordering` (transverse-energy ordering), `Eordering` (energy ordering), `Pordering` (momentum ordering), `PTordering` (transverse-momentum ordering), `PXordering` (ordering according to the  $x$ -component of the momentum), `PYordering` (ordering according to the  $y$ -component of the momentum) and `PZordering` (ordering according to the  $z$ -component of the momentum).

`PHYSICS->SqrtS(evt)`

This returns the partonic center-of-mass energy of the event in GeV.

```
PHYSICS->ToRestFrame(prt,prt1)
```

where the objects `prt` and `prt1` are two instances of the `MCParticleFormat` or `RecParticleFormat` classes. The four-momentum of the particle `prt` is boosted to the rest frame of the particle `prt1` and overwritten. Let us note that the Lorentz transformation is chosen in such a way that the tri-dimensional ( $x, y, z$ ) axes are kept unchanged.

### 5.6. A detailed example

In this Section, we give a detailed example about how to implement an analysis within the expert mode of MADANALYSIS 5. We choose to investigate the polarization of the  $W$ -boson issued from a top quark decaying leptonically,

$$t \rightarrow W^+ b \rightarrow \ell^+ \nu_l b . \quad (4)$$

This property of the  $W$ -boson is usually studied through the shape of a particular angular distribution,  $d\sigma/d\cos\theta^*$ . The  $\theta^*$  angle is defined as the angle between the momentum of the  $W$ -boson evaluated in the rest frame of the originating top quark and the one of the lepton  $\ell^+$  evaluated in the rest frame of the  $W$ -boson.

To investigate this observable, we focus on the production of a top-antitop pair where one of the final state top quarks undergoes a leptonic decay and the other one decays hadronically,

$$pp \rightarrow t\bar{t} \rightarrow (b \ j \ j)(\bar{b} \ \ell^- \ \bar{\nu}_l) \quad \text{or} \quad pp \rightarrow t\bar{t} \rightarrow (b \ \ell^+ \ \nu_l)(\bar{b} \ j \ j) ,$$

where  $\ell$  stands for an electron or a muon and  $j$  for a light jet. From the parton-level samples stored in the directory `samples` of MADANALYSIS 5<sup>8</sup>, `ttbar_s1_1.lhe.gz` and `ttbar_s1_2.lhe.gz`, we illustrate the implementation, within the SAMPLEANALYZER framework, of an analysis code leading to the creation of a histogram representing the  $d\sigma/d\cos\theta^*$  distribution introduced above. We refer to Section 3 concerning details on the generation of these two Monte Carlo samples.

When launching MADANALYSIS 5 in expert mode by issuing in a shell

```
bin/ma5 -e
```

---

<sup>8</sup>If the `samples` directory is absent, we recall that it can be created by issuing, in the command line interface of MADANALYSIS 5, the command `install samples`.

the user is asked to enter the name of the directory to be created and the label of the analysis. For the sake of the example, the name of the directory is chosen to be `Wpol` whilst the string `W polarization from a top decay` is entered as the label, or title, of the analysis. The three files `analysisList.cpp`, `user.cpp` and `user.h` are automatically created and stored in the directory `Wpol/SampleAnalyzer/Analysis`, as mentioned in Section 5.1. Whilst the first of these files does not have to be modified by the user, the two others must be updated to include the analysis which has to be implemented.

The file `user.h` strictly follows the structure presented in Section 5.2,

```
#ifndef analysis_user_h
#define analysis_user_h

#include "Core/AnalysisBase.h"

class user: public AnalysisBase
{
    INIT_ANALYSIS(user,"W polarization from a top decay")

public:
    virtual void Initialize();
    virtual void Finalize(const SampleFormat& summary,
                          const std::vector<SampleFormat>& files);
    virtual void Execute(const SampleFormat& sample,
                        const EventFormat& event);

private:
    TH1F* myHisto;
};

#endif
```

The label of the analysis has been automatically included into the arguments of the `INIT_ANALYSIS` function at the execution of MADANALYSIS 5, together with the declaration of the three main functions `Initialize`, `Execute` and `Finalize`. Since we are interested in the creation of a single histogram, we therefore need to declare, as a private member of the `user` class, an instance

of the `TH1F` class<sup>9</sup> which we denote by `myHisto`.

The C++ source file `user.cpp` contains the implementation of the three core functions `Initialize`, `Execute` and `Finalize`. Its structure reads thus

```
#include "Analysis/user.h"

void user::Initialize()
{ ... }

void user::Execute(const SampleFormat& sample,
                   const EventFormat& event)
{ ... }

void user::Finalize(const SampleFormat& summary,
                    const std::vector<SampleFormat>& files)
{ ... }
```

where the dots are specified in the remainder of this Section. The function `Initialize` contains on the one hand the initialization of the physics services as well as, on the other hand, the one of the histogram to be drawn, following the ROOT syntax,

```
void user::Initialize()
{
    // Initializing Physics services
    PHYSICS->mcConfig().Reset();

    // Initializing the histogram
    myHisto = new TH1F("myHisto","cos#theta^{*}",15,-1.,1.);
}
```

We here ask for the creation of a histogram containing 15 bins ranging from  $-1$  to  $+1$ , the minimal and maximal values for the cosine of the  $\theta^*$ -angle.

The implementation of the function `Execute` is a bit more complicated. First of all, the observable of interest can only be computed once the three relevant particles have been identified, *i.e.*, the final-state lepton  $\ell$ , the intermediate  $W$ -boson decaying into this lepton  $\ell$  and the originating top quark.

---

<sup>9</sup>We refer to the ROOT manual [57] for the definition of the `TH1F` class and its attributes.

Then, the four-momentum of the  $W$ -boson has to be boosted into the rest frame of the top quark and the one of the lepton into the rest frame of the  $W$ -boson. Once this is done, the cosine of the angle  $\theta^*$  can be computed and the value filled into the histogram. The skeleton of the function `Execute` reads thus

```
void user::Execute(const SampleFormat& sample,
                   const EventFormat& event)
{
    // Initialization of three pointers to the lepton,
    // W and top quark
    { ... }

    // Identification of the three particles of interest
    { ... }

    // Computing the observable; filling the histogram
    { ... }
}
```

The first series of dots concerns the declaration of the three objects which contain the three particles of interest, once identified, and which are necessary to compute the observable  $\cos \theta^*$ ,

```
const MCParticleFormat* top      = 0;
const MCParticleFormat* w       = 0;
const MCParticleFormat* lepton = 0;
```

In the lines above, three pointers to a `MCParticleFormat` structure are created and initialized to the zero value.

The second series of dots are related to the identification of the top quark, the  $W$ -boson and the lepton. In particular, the analysis code needs to check that the mother-to-daughter relations given by the decay chain of Eq. (4) are fulfilled. Otherwise, the event must be rejected. In practice, this task is performed through a loop over the particle content of the event. The lepton is firstly selected as a final-state particle whose PDG-id reads  $\pm 11$  (electron) or  $\pm 13$  (muon). The  $W$ -boson is then further identified by requiring that the mother particle of this lepton has a PDG-id equal to  $\pm 24$ . The top quark is finally identified by requiring that the ‘grandmother’ of the lepton has a PDG-id equal to  $\pm 6$ . The implemented code is given by

```

for (unsigned int i=0;i<event.mc()->particles().size();i++)
{
    const MCParticleFormat* prt=&(event.mc()->particles()[i]);

    // The lepton is a final state particle
    if (!PHYSICS->IsFinalState(prt)) continue;

    // Lepton selection based on the PDG-id
    if ( std::abs(prt->pdgid())!=11 &&
        std::abs(prt->pdgid())!=13 ) continue;

    // Getting the mother of the lepton and
    // checking if it is a W-boson
    const MCParticleFormat* mother = prt->mother1();
    if (mother==0) continue;
    if (std::abs(mother->pdgid())!=24) continue;

    // Getting the grand-mother of the lepton and
    // checking if it is a top quark
    const MCParticleFormat* grandmother = mother->mother1();
    if (grandmother==0) continue;
    if (std::abs(grandmother->pdgid())!=6) continue;

    // Saving the selected particles
    lepton = prt;
    w      = mother;
    top    = grandmother;

    // Particles are found: breaking the loop
    break;
}

```

In the case the three particles are not identified, the event must be rejected, which is done by implementing

```

// Rejection of the event if the decay chain is not found
if (lepton==0)
{
    WARNING << "(t > b W > b l nu) decay chain not found!"

```

```

        << std::endl;
    return;
}

```

where we have used the message services included in the `SAMPLEANALYZER` framework to print a warning message to the screen.

In the third and last series of dots included in the skeleton of the `user.cpp` file, the observable  $\cos \theta^*$  is computed and the histogram filled, this last task being performed by means of standard ROOT commands,

```

// Boosting the lepton four-momentum to the W rest frame
MCParticleFormat lepton_new = *lepton;
PHYSICS->ToRestFrame(lepton_new,w);

// Boosting the W four-momentum to the top rest frame
MCParticleFormat w_new = *w;
PHYSICS->ToRestFrame(w_new,top);

// Computing the observable; filling the histogram
myHisto -> Fill( cos(lepton_new.angle(w_new)) );

```

The histogram must now be created and exported to a human-readable format. This is done at the level of the function `Finalize` by means of a series of ROOT commands. For the sake of the example, we have decided to normalize the histogram to the number of events  $\mathcal{N}$  expected for an integrated luminosity of  $10 \text{ fb}^{-1}$ ,

$$\mathcal{N} = \sigma \mathcal{L}_{\text{int}} / N ,$$

where  $\sigma$  is the cross section, in pb, corresponding to the process under consideration,  $\mathcal{L}_{\text{int}}$  the integrated luminosity (thus set to  $10000 \text{ pb}^{-1}$ ) and  $N$  the number of events included in the samples. As stated above, the value of the cross section read from the LHE files is stored in the `SampleFormat` object `summary` (see Section 5.4) where the average over both input samples has been performed. The corresponding C++ implementation of the function `Finalize` reads

```

void user::Finalize(const SampleFormat& summary,
                    const std::vector<SampleFormat>& files)
{

```

```

// Color of the canvas background
gStyle->SetCanvasColor(0);

// Turning off the border lines of the canvas
gStyle->SetCanvasBorderMode(0);

// Configuring statistics printing
gStyle->SetOptStat(111110);

// Creating the output root file
TCanvas* myCanvas = new TCanvas("myCanvas","");
myCanvas->cd();

// Setting background color
myHisto->SetFillColor(kBlue);

// Normalization of the histogram: L = 10 fb-1
double nrm = summary.mc()->xsection() * 10000. /
    static_cast<float>(summary.nevents());
myHisto->Scale(nrm);

// Setting axis title
myHisto->GetXaxis()->SetTitle("cos#theta^{*}");

// Drawing histogram
myHisto->Draw();

// Saving plot
myCanvas->SaveAs((outputName_+".eps").c_str());
}

```

After compiling the analysis and linking it to the external libraries with the help of the provided `Makefile` located in the `SampleAnalyzer` directory, the analysis can be executed,

```
SampleAnalyzer --analysis="W polarization from a top decay" \
list.txt
```

where `list.txt` is a text file containing the absolute paths to the two event samples under consideration. A figure, named `list.eps`, is generated in the

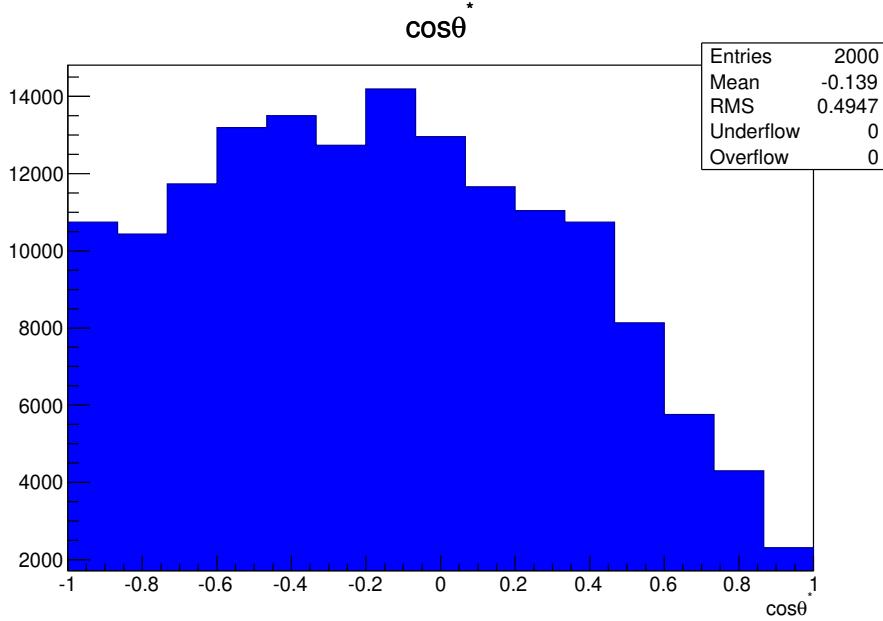


Figure 5:  $\cos\theta^*$  distribution for the production of a top-antitop pair semi-leptonically decaying. The  $\theta^*$  angle is defined as in the text.

directory where the executable is stored. The results, shown in Figure 5, agree, for instance, with Ref. [63].

## 6. Conclusions

The task of performing a phenomenological analysis based on event files such as those generated by Monte Carlo generators can be divided into three stages. Firstly, the event samples must be read and loaded into the memory of the computer. The format of these files depends in general on the level of sophistication of the analysis (at the parton-level, hadron-level or reconstructed-level). Secondly, the analysis itself must be performed, *i.e.*, selection cuts are applied on the signal and background event samples with the aim of being able to extract information on the signal from the often overwhelming background. Finally, the results are outputted as histograms and/or cut-flow charts to improve their readability and interpretation.

In this work, we have presented MADANALYSIS 5, a new user-friendly and efficient framework aiming to facilitate the implementation of phenomenological analyses such as the one described above. We have explained how to

implement and run an analysis within this framework in a straightforward way and have given several detailed examples addressing the different facets of the program.

**MADANALYSIS** 5 is based on a multi-purpose C++ kernel, **SAMPLEANALYZER**, which uses the **ROOT** platform. Compatible with most of the event file formats commonly used for analyzing parton-level, hadron-level and reconstructed events, **MADANALYSIS** 5 offers two modes of running according to the needs and the expertise of the user.

A highly-user-friendly mode, the normal running mode of the program, uses the strengths of a **PYTHON** interface to reduce the implementation of an analysis to a set of intuitive commands whose syntax has been inspired by the **PYTHON** programming language. Therefore, rather complex analyses can be achieved without too much effort.

For users with more advanced C++ and **ROOT** programming skills, **MADANALYSIS** 5 can also be run in its expert mode. This overcomes the limitations of the normal mode of running in which the scope is restricted by the set of functionalities that have been implemented. The user is in this case required to directly implement his analysis in C++ within the **SAMPLEANALYZER** framework, rendering the possibilities, at the analysis level, only limited by the imagination and the skills of the user. However, even if this mode of running is in principle more complicated to handle, the existence of many built-in functions and methods renders the task of implementing the analysis easier and more straightforward.

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## Appendix A. Installation of the program

### Appendix A.1. Requirements

For a proper running, MADANALYSIS 5 requires several mandatory external libraries. In addition, the full set of functionalities of the program can be made available by installing optional external dependencies on the computer of the user. If these optional libraries are absent, several of the functionalities of MADANALYSIS 5 are deactivated but analyses of event files can still be performed. In contrast, the absence of one of the mandatory external libraries simply does not allow use of the program.

In order to run MADANALYSIS 5 locally, PYTHON 2.6 or a more recent version (however not from the 3.x series) must be installed on the computer of the user [64]. We recall that in order to check the version of PYTHON present on a system, it is enough to type in a shell

```
python --version
```

The installation of the PYTHON package is one of the three mandatory requirements without which the program cannot run. The two other external packages that have to be installed are related to C++ and ROOT.

The SAMPLEANALYZER kernel requires the installation of a C++ compiler together with the associated Standard Template Libraries (STL). Since the validation procedure of MADANALYSIS 5 has only been achieved within the context of the GNU GCC compiler and since this compiler is available on most operating systems [65], we have adopted the choice of requiring the installation of GCC. The program has been validated with the versions 4.3.x and 4.4.x. We recall that the version of the GCC installed on a system can be obtained by issuing in a shell

```
g++ --version
```

Let us however stress that compatibility with any other C++ compiler is in principle ensured, but requires the modification of several core files of MADANALYSIS 5. Therefore, it is currently not supported.

Concerning the ROOT package, a version more recent than version 5.27 has to be installed [66], and the user has to check that the PYTHON functionalities of the ROOT library are available. We remind that in order to install a version of ROOT including its PYTHON library, the LINUX package PYTHON-DEVEL has to be present on the system of the user and the Root configuration script must be run as

```
./configure --with-python
```

Very importantly, the version of PYTHON employed to start MADANALYSIS 5 must match the one used to generate the PYTHON library of ROOT. Finally, the version of ROOT present on a computer, together with the compatibility with the PYTHON language, can be checked by issuing in a shell the commands

```
root-config --version  
root-config --has-python
```

The three above-mentioned programs, *i.e.*, PYTHON, the GCC compiler and ROOT, must be available from any location of the computer. If this is not the case, the user has to modify the system variable \$PATH. If these programs have been installed at standard location on the system, the necessary environment variables are set automatically by MADANALYSIS 5. Contrary, it is left to the user to check that the paths to the associated header and library files are included in the environment variables of the GCC compiler, \$CPLUS\_INCLUDE\_PATH and \$LIBRARY\_PATH. Finally, for a proper linking of the external libraries, the variable \$LD\_LIBRARY\_PATH (or, on MACOS operating systems, \$DYLD\_LIBRARY\_PATH) must also contain the paths to the libraries to be linked.

We now turn to the optional libraries that can be linked to MADANALYSIS 5. In order to handle zipped Monte Carlo event samples, the ZLIB library has to be installed on the system [67]. However, if ZLIB is not locally installed, the possibility of analyzing zipped event samples is simply deactivated, *i.e.*, the user will have to unzip the event files manually before running the code.

### *Appendix A.2. Downloading the program*

It is recommended to always use the latest stable version of the MADANALYSIS 5 package, which contains, when downloaded from the web, both the PYTHON command line interface and the SAMPLEANALYZER framework. It can always be found together with an up-to-date manual on the webpage  
<http://madanalysis.irmp.ucl.ac.be>

The package does not require any compilation or configuration. After having downloaded the tar-ball from the website, it can be unpacked either in the directory where MADGRAPH 5 is installed,

```
cd <path-to-madgraph>  
tar -xvf ma5_xxxx.tgz
```

or in any location on the computer of the user,

```
mkdir <ma5 dirname>
cd <ma5 dirname>
tar -xvf ma5_xxxx.tgz
```

where `xxxx` stands for the version number of the MADANALYSIS 5 release which has been downloaded. When MADANALYSIS 5 is installed as a dependency of MADGRAPH 5, the list of predefined particle and multiparticle labels is directly updated from the MADGRAPH 5 standards. Moreover, this allows to perform analyses at the time of the event generation by MADGRAPH 5 which, in this case, pilots MADANALYSIS 5. We emphasize that several predefined analyses exist, but the user has the freedom to tune the desired analysis to be performed according to his own aims.

Once installed and unpacked, MADANALYSIS 5 can be immediately launched by issuing in a shell

```
bin/ma5
```

as presented in Section 4.1.

### *Appendix A.3. Running MADANALYSIS 5*

When started, MADANALYSIS 5 first checks that all the dependencies (GCC, PYTHON, ROOT, ZLIB) are present on the system and that compatibility is ensured with the installed versions. In the case of any problem, a message is printed to the screen and the code exists if it is found that it cannot properly run. On the first session of MADANALYSIS 5, the SAMPLEANALYZER core is compiled behind the scene as a static library stored in the directory `lib`. For the next sessions, the kernel is only recompiled if the configuration of the system has changed (new version of the dependencies or of the main program).

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