

## Solutions for Tutorial #5

1. A wire bent as a parabola  $y = ax^2$  is located in a uniform magnetic field  $B$  oriented perpendicular to the plane  $x-y$ . At the moment  $t = 0$  a connector parallel to  $x$  axis starts sliding from the apex of the parabola with a constant acceleration  $w$ , see Figure 1. Find the emf of the electromagnetic induction in the loop as a function of  $y$ .

According to Lenz's law induced current & induced emf will be anticlockwise in the loop

We can obtain emf from Faraday's law  $\epsilon_0 = \left| \frac{d\phi}{dt} \right| - (1)$

where,  $d\phi = \vec{B} \cdot d\vec{S} - (2)$

As the connector slides parallel to x-axis, area of the parabola

For small displacement  $dy$  is  $2x dy$

Hence,  $d\phi = 2Bx dy - (3)$

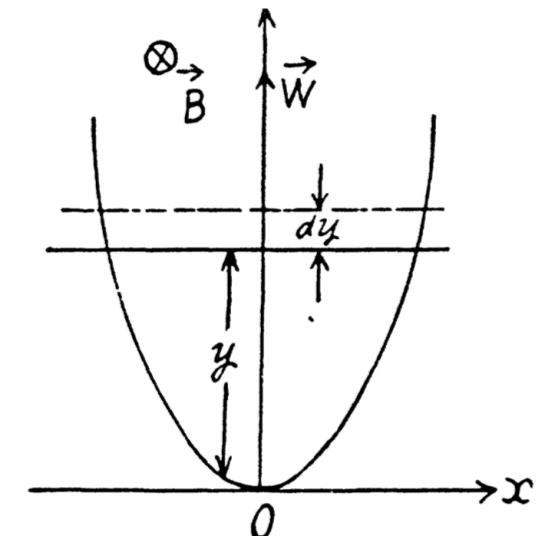
Wire is bent like a parabola  $y = ax^2$  Therefore,  $x = \sqrt{\frac{y}{a}} - (4)$

Also as the connector moves with constant acceleration,  $y = \frac{1}{2}wt^2 - (5)$  So,  $t = \sqrt{\frac{2y}{w}} - (6)$

From (5),  $\frac{dy}{dt} = wt = \sqrt{2yw} - (7)$

Hence From (1),  $\epsilon_0 = \left| \frac{d\phi}{dt} \right| = 2B \sqrt{\frac{y}{a}} \frac{dy}{dt}$  [Using (3) & (4)]

$$\epsilon_0 = By \sqrt{\frac{8w}{a}} \quad \text{[Using (7)]}$$



2. Plot the vector field  $\frac{\hat{r}}{r^2}$  and compute the divergence

$$\vec{\nabla} \cdot \frac{\hat{r}}{r^2}$$

where  $\vec{r} = x \hat{i} + y \hat{j} + z \hat{k}$ . Explain the result. Is this calculation valid everywhere?

$$\text{Let, } V = \frac{\hat{r}}{r^2}$$

$$\vec{\nabla} \cdot \vec{V} = \frac{1}{r^2} \frac{\partial}{\partial r} (r^2 V_r) + \frac{1}{r \sin \theta} \frac{\partial}{\partial \theta} (\sin \theta V_\theta) + \frac{1}{r \sin \theta} \frac{\partial V_\phi}{\partial \phi}$$

$$\vec{\nabla} \cdot \vec{V} = \frac{1}{r^2} \frac{\partial}{\partial r} \left( r^2 \frac{1}{r^2} \right) = 0$$

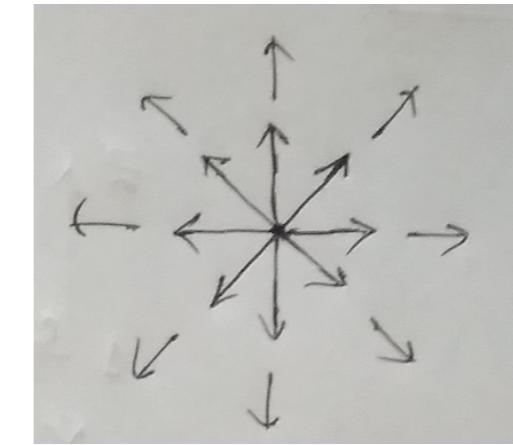
So  $\vec{\nabla} \cdot \vec{V}$  will be zero everywhere except at the origin  $r=0$

At origin consider a sphere of radius  $R$

$$\int_V (\vec{\nabla} \cdot \vec{V}) dV = \oint_S \vec{V} \cdot d\vec{S} = \int \left( \frac{1}{R^2} \hat{r} \right) (R^2 \sin \theta d\theta d\phi) \hat{r}$$

$$\Rightarrow \int_V (\vec{\nabla} \cdot \vec{V}) dV = \int_0^\pi \sin \theta d\theta \int_0^{2\pi} d\phi = 4\pi$$

The integral value of  $\vec{\nabla} \cdot \vec{V}$  over a volume containing origin gives a non-zero value. So it behaves like a Dirac Delta function.

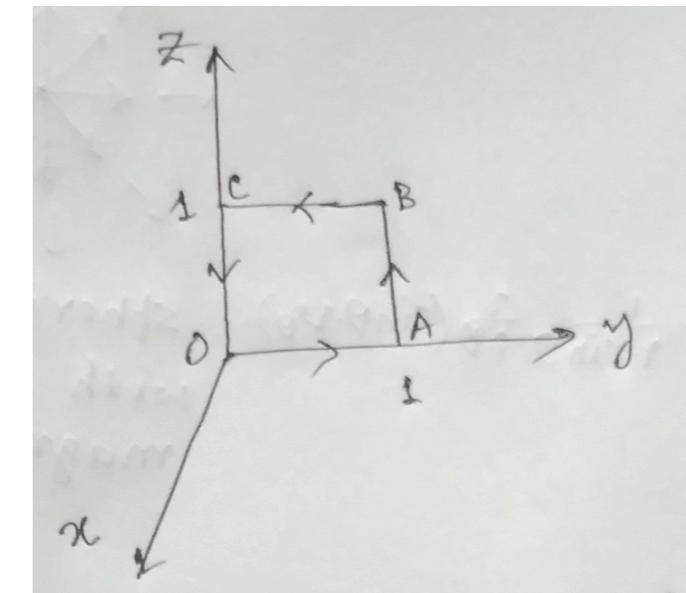


Spherically symmetric with decreasing magnitude

3. Check Stokes theorem for the vector field  $\vec{W}(x, y, z) = (2xz + 3y^2)\hat{j} + (4yz^2)\hat{k}$  over a square surface of unit length in the  $y - z$  plane positive quadrant with one edge at the origin.

$$\vec{W}(x, y, z) = (2xz + 3y^2)\hat{j} + 4yz^2\hat{k}$$

$$d\vec{S} = dydz \cdot \hat{i}$$



$$\vec{\nabla} \times \vec{W} = \begin{vmatrix} \hat{i} & \hat{j} & \hat{k} \\ \frac{\partial}{\partial x} & \frac{\partial}{\partial y} & \frac{\partial}{\partial z} \\ 0 & 2xz + 3y^2 & 4yz^2 \end{vmatrix}$$

$$\Rightarrow \vec{\nabla} \times \vec{W} = (4z^2 - 2x)\hat{i} + 2z\hat{z}$$

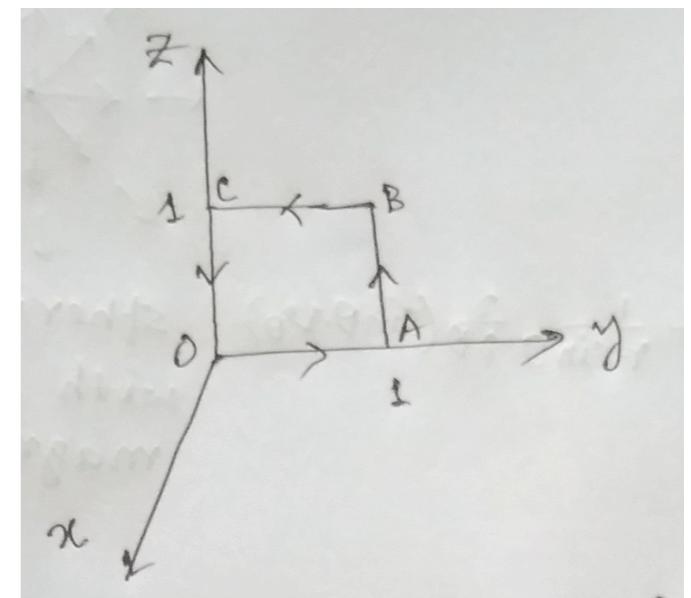
Therefore,  $\int_S (\vec{\nabla} \times \vec{W}) \cdot d\vec{S} = \int (4z^2 - 2x) dy dz$

$$\Rightarrow \int_S (\vec{\nabla} \times \vec{W}) \cdot d\vec{S} = \int_0^1 dy \int_0^1 4z^2 dz = \frac{4}{3} \quad \text{Since } [x=0]$$

### Question 3 continued ..

Line integral,

$$d\vec{l} = dx \cdot \hat{i} + dy \cdot \hat{j} + dz \cdot \hat{k}$$



$$\oint_{OABC} \vec{W} \cdot d\vec{l} = \int_{OA} \vec{W} \cdot d\vec{l} + \int_{AB} \vec{W} \cdot d\vec{l} + \int_{BC} \vec{W} \cdot d\vec{l} + \int_{CO} \vec{W} \cdot d\vec{l}$$

$$\Rightarrow \oint_{OABC} \vec{W} \cdot d\vec{l} = \int_0^1 3y^2 dy + \int_0^1 4z^2 dz + \int_1^0 3y^2 dy + \int_1^0 0 dz$$

$$\Rightarrow \oint_{OABC} \vec{W} \cdot d\vec{l} = 1 + \frac{4}{3} - 1 + 0 = \frac{4}{3}$$

$$\text{Therefore, } \oint_{OABC} \vec{W} \cdot d\vec{l} = \int_S (\vec{\nabla} \times \vec{W}) \cdot d\vec{S}$$

So Stokes theorem is satisfied

4. Given the electric flux law:

$$\Phi_E = \oint \vec{E} \cdot d\vec{A} = \frac{q}{\epsilon_0}$$

Use the Gauss's divergence theorem to show

$$\vec{\nabla} \cdot \vec{E} = \frac{\rho}{\epsilon_0}$$

where  $\rho$  is the volume charge density.

$$Given, \phi_E = \oint \vec{E} \cdot d\vec{A} = \frac{q}{\epsilon_0}$$

According to Gauss's divergence theorem : Volume integral of the divergence  
Is equal to the surface integral over the boundary that encloses the volume

$$\oint \vec{E} \cdot d\vec{A} = \int_V (\vec{\nabla} \cdot \vec{E}) dV$$

$$Again, \oint \vec{E} \cdot d\vec{A} = \frac{q}{\epsilon_0} = \frac{\int_V \rho dV}{\epsilon_0} \quad \text{Where } \rho \text{ is volume charge density}$$

$$Thus, \oint \vec{E} \cdot d\vec{A} = \int_V (\vec{\nabla} \cdot \vec{E}) dV = \frac{\int_V \rho dV}{\epsilon_0}$$

$$Hence, \vec{\nabla} \cdot \vec{E} = \frac{\rho}{\epsilon_0}$$

5. Given  $\vec{E} = \rho_0 \epsilon_0 x \hat{i}$  find the charge density responsible for this electric field.

$$\vec{E} = \rho_0 \epsilon_0 x \cdot \hat{i}$$

$$\vec{\nabla} \cdot \vec{E} = \frac{\partial}{\partial x} (\rho_0 \epsilon_0 x) = \rho_0 \epsilon_0$$

According to differential form of Gauss's law,

$$\vec{\nabla} \cdot \vec{E} = \frac{\rho}{\epsilon_0} \quad \text{Where } \rho \text{ is charge density}$$

$$So, \rho = \epsilon_0 (\vec{\nabla} \cdot \vec{E}) = \rho_0 \epsilon_0^2$$