

PH423 Assignment 2

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1. Your question here.

[Sankalp: I got this one.]

We start with the expansion of the operators \hat{J}_x and \hat{J}_y in terms of the ladder operators

$$\hat{J}_x = \frac{1}{2} \cdot (\hat{J}_+ + \hat{J}_-) \quad (1)$$

and

$$\hat{J}_y = \frac{1}{2i} \cdot (\hat{J}_+ - \hat{J}_-) . \quad (2)$$

The application of the ladder operators on a state $|j, m\rangle$ changes it to a state of the form $c \cdot |j, m \pm 1\rangle$ for some $c \in \mathbb{C}$. So, given the orthogonality of the $|j, m\rangle$ states, we get that

$$\langle j, m | \hat{J}_x | j, m \rangle = \langle j, m | \hat{J}_y | j, m \rangle = 0 \quad \forall |j, m\rangle . \quad (3)$$

Squaring Equation 1 and 2, we get the operators \hat{J}_x^2 and \hat{J}_y^2 in terms of the ladder operators. With the same argument as before, we see that only terms with equal powers of the two ladder operators will contribute, and using

$$\hat{J}_{\pm} |j, m\rangle = \hbar \sqrt{(j \mp m)(j \pm m + 1)} |j, m \pm 1\rangle , \quad (4)$$

we get

$$\langle j, m | \hat{J}_y^2 | j, m \rangle = \langle j, m | \hat{J}_x^2 | j, m \rangle \quad (5)$$

$$= \langle j, m | \frac{1}{4} \cdot (\hat{J}_+^2 + \hat{J}_+ \hat{J}_- + \hat{J}_- \hat{J}_+ + \hat{J}_-^2) | j, m \rangle \quad (6)$$

$$= \langle j, m | \frac{1}{4} \cdot (\hat{J}_+ \hat{J}_- + \hat{J}_- \hat{J}_+) | j, m \rangle \quad (7)$$

$$= \langle j, m | \frac{1}{2} \cdot \left(\sqrt{(j+m+1)(j-m)} \sqrt{(j-m)(j+m+1)} + \sqrt{(j-m)(j+m+1)} \sqrt{(j+m+1)(j-m)} \right) \cdot | j, m \rangle \quad (8)$$

$$= (j + m + 1)(j - m) \quad (9)$$

15 The values for x and y are not separately calculated as a trivial calculation shows they're equal. The same
16 is easily argued using symmetry in the x-y plane.

17 **2. Determine the eigenvalues and eigenvectors of the 2 x 2 matrix $\sigma \cdot \hat{n}$, where \hat{n} is a unit vector along the (θ, ϕ) direction and σ are the three Pauli matrices. This is basically the projection of the spin 1/2 operator (apart from $\frac{\hbar}{2}$) along the direction of the unit vector \hat{n} . Do this in two ways:**

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19 [Parth: Doing question 2, might have issues with part (b) make sure that it's correct]

20 (a) First by explicitly diagonalizing the matrix $\sigma \cdot \hat{n}$.

21 The vector $\sigma = (\sigma_1, \sigma_2, \sigma_3)$, where the σ_i matrices are -

$$\sigma_1 = \begin{pmatrix} 0 & 1 \\ 1 & 0 \end{pmatrix}, \sigma_2 = \begin{pmatrix} 0 & -i \\ i & 0 \end{pmatrix}, \sigma_3 = \begin{pmatrix} 1 & 0 \\ 0 & -1 \end{pmatrix}$$

Now we need to figure out what \hat{n} is. The unit vector points along the (θ, ϕ) direction. This is nothing but the unit vector \hat{r} in Polar co-ordinates.

$$\hat{n} = \hat{r} = \cos(\phi)\sin(\theta)\hat{i} + \sin(\phi)\sin(\theta)\hat{j} + \cos(\theta)\hat{k}$$

Thus, $\hat{n} = (\cos(\phi)\sin(\theta), \sin(\phi)\sin(\theta), \cos(\theta))$. We know that $\mathbf{a} \cdot \mathbf{b} = a_i b_i$ (implicit summation over i)

Thus, $\sigma \cdot \hat{n} = \sigma_i n_i$.

$$\begin{aligned} \sigma \cdot \hat{n} &= \cos(\phi)\sin(\theta) \begin{pmatrix} 0 & 1 \\ 1 & 0 \end{pmatrix} + \sin(\phi)\sin(\theta) \begin{pmatrix} 0 & -i \\ i & 0 \end{pmatrix} + \cos(\theta) \begin{pmatrix} 1 & 0 \\ 0 & -1 \end{pmatrix} \\ \therefore \sigma \cdot \hat{n} &= \sin(\theta) \begin{pmatrix} 0 & \cos(\phi) - i * \sin(\phi) \\ \cos(\phi) + i * \sin(\phi) & 0 \end{pmatrix} + \cos(\theta) \begin{pmatrix} 1 & 0 \\ 0 & -1 \end{pmatrix} \\ &= \sin(\theta) \begin{pmatrix} 0 & e^{-i\phi} \\ e^{i\phi} & 0 \end{pmatrix} + \cos(\theta) \begin{pmatrix} 1 & 0 \\ 0 & -1 \end{pmatrix} = \begin{pmatrix} \cos(\theta) & \sin(\theta)e^{-i\phi} \\ \sin(\theta)e^{i\phi} & -\cos(\theta) \end{pmatrix} \end{aligned}$$

To find the eigenvalues and eigenvectors, we now need to diagonalize this matrix. Let the eigenvalues be represented by λ . The characteristic polynomial takes the following form.

$$\begin{aligned} (\cos(\theta) - \lambda)(-\cos(\theta) - \lambda) - \sin(\theta)e^{-i\phi} * \sin(\theta)e^{i\phi} &= 0 \\ \therefore -\cos^2(\theta) + \lambda^2 - \sin^2(\theta) &= 0 \Rightarrow \lambda^2 - 1 = 0 \\ \therefore \lambda &= \pm 1 \end{aligned}$$

22 for $\lambda = 1$, let the eigenvector be $\mathbf{v}_1 = (v_{1,1}, v_{1,2})$, thus

$$\begin{pmatrix} \cos(\theta) & \sin(\theta)e^{-i\phi} \\ \sin(\theta)e^{i\phi} & -\cos(\theta) \end{pmatrix} \begin{pmatrix} v_{1,1} \\ v_{1,2} \end{pmatrix} = \begin{pmatrix} v_{1,1} \\ v_{1,2} \end{pmatrix}$$

$$\begin{aligned}\therefore \cos(\theta) * v_{1,1} + \sin(\theta)e^{-i\phi} * v_{1,2} &= v_{1,1}, \quad \sin(\theta)e^{i\phi} * v_{1,1} - \cos(\theta) * v_{1,2} = v_{1,2} \\ v_{1,2} &= e^{i\phi} \frac{\sin(\theta)}{(\cos(\theta) + 1)} * v_{1,1}\end{aligned}$$

23 Thus, for eigenvalue $\lambda = 1$, the eigenvector $\mathbf{v}_1 = (v_{1,1}, e^{i\phi} \frac{\sin(\theta)}{(\cos(\theta)+1)} * v_{1,1})$

24 Likewise, for $\lambda = -1$, let the eigenvector be $\mathbf{v}_2 = (v_{2,1}, v_{2,2})$, thus

$$\begin{aligned}\begin{pmatrix} \cos(\theta) & \sin(\theta)e^{-i\phi} \\ \sin(\theta)e^{i\phi} & -\cos(\theta) \end{pmatrix} \begin{pmatrix} v_{2,1} \\ v_{2,2} \end{pmatrix} &= \begin{pmatrix} -v_{2,1} \\ -v_{2,2} \end{pmatrix} \\ \therefore \cos(\theta) * v_{2,1} + \sin(\theta)e^{-i\phi} * v_{2,2} &= -v_{2,1}, \quad \sin(\theta)e^{i\phi} * v_{2,1} - \cos(\theta) * v_{2,2} = -v_{2,2} \\ v_{2,2} &= e^{i\phi} \frac{\sin(\theta)}{(1 - \cos(\theta))} * v_{2,1}\end{aligned}$$

25 Thus, for eigenvalue $\lambda = -1$, the eigenvector $\mathbf{v}_2 = (v_{2,1}, e^{i\phi} \frac{\sin(\theta)}{(1-\cos(\theta))} * v_{2,1})$.

26 We thus have our two eigenvalues (± 1) and our two eigenvectors (\mathbf{v}_1 and \mathbf{v}_2)

27

28 **(b)** By rotating the spinor pointing initially along the $+\hat{z}$ axis direction by appropriate angles, using the
29 appropriate rotation operator. Convince yourself that one has to rotate by an angle θ counterclock-
30 wise around the y -axis and then by ϕ around the z -axis. Apart from overall phases, is the resultant
31 spinor the same as the spin up eigenvector obtained in part **(a)**?

32 Let's start with the spinor pointing in the $+z$ -direction.

$$\left| s_z = +\frac{\hbar}{2} \right\rangle = \begin{bmatrix} 1 \\ 0 \end{bmatrix}, \quad \text{s.t. } S_z \left| s_z = +\frac{\hbar}{2} \right\rangle = +\frac{\hbar}{2} \left| s_z = +\frac{\hbar}{2} \right\rangle$$

33 If we apply consecutive rotation operators, we should be able to rotate this spinor into a general
34 state, pointing in an arbitrary direction $\hat{\mathbf{n}}$, where $\hat{\mathbf{n}}$ points in the (θ, ϕ) direction.

35 We first rotate this spinor by θ around the y -axis, and then by ϕ around the z -axis. The axis of spin
36 now points in the direction $\hat{\mathbf{n}}$. Thus -

$$|\hat{n}+\rangle = U[R(\phi\hat{z})]U[R(\theta\hat{y})] \begin{bmatrix} 1 \\ 0 \end{bmatrix}$$

37 To find the explicit form of $|\hat{n}+\rangle$, we'll need the forms of the unitary matrices $U[R(\phi\hat{z})]$ and
38 $U[R(\theta\hat{y})]$. We'll use the result given in Shankar -

$$U[R(\theta)] = \cos\frac{\theta}{2}I - i\sin\frac{\theta}{2}(\hat{\theta} \cdot \boldsymbol{\sigma})$$

39 Looking at the particular case of rotation around y -axis by amount θ and then subsequently around
40 z -axis by amount ϕ -

$$\begin{aligned}
U[R(\theta\hat{y})] \begin{bmatrix} 1 \\ 0 \end{bmatrix} &= \begin{bmatrix} \cos\frac{\theta}{2}I - i\sin\frac{\theta}{2}\sigma_y \end{bmatrix} \begin{bmatrix} 1 \\ 0 \end{bmatrix} \\
&= \begin{bmatrix} \cos\frac{\theta}{2} \\ 0 \end{bmatrix} - i\sin\frac{\theta}{2} \begin{bmatrix} 0 & -i \\ i & 0 \end{bmatrix} \begin{bmatrix} 1 \\ 0 \end{bmatrix} \\
&= \begin{bmatrix} \cos\frac{\theta}{2} \\ \sin\frac{\theta}{2} \end{bmatrix}
\end{aligned}$$

Applying rotation around z -axis by amount ϕ now, we get

$$\begin{aligned}
U[R(\phi\hat{z})] \begin{bmatrix} \cos\frac{\theta}{2} \\ \sin\frac{\theta}{2} \end{bmatrix} &= \begin{bmatrix} \cos\frac{\phi}{2}I - i\sin\frac{\phi}{2}\sigma_z \end{bmatrix} \begin{bmatrix} \cos\frac{\theta}{2} \\ \sin\frac{\theta}{2} \end{bmatrix} \\
&= \begin{bmatrix} \cos\frac{\phi}{2}\cos\frac{\theta}{2} \\ \cos\frac{\phi}{2}\sin\frac{\theta}{2} \end{bmatrix} - i\sin\frac{\phi}{2} \begin{bmatrix} 1 & 0 \\ 0 & -1 \end{bmatrix} \begin{bmatrix} \cos\frac{\theta}{2} \\ \sin\frac{\theta}{2} \end{bmatrix} \\
&= \begin{bmatrix} \cos\frac{\theta}{2} \left(\cos\frac{\phi}{2} - i\sin\frac{\phi}{2} \right) \\ \sin\frac{\theta}{2} \left(\cos\frac{\phi}{2} + i\sin\frac{\phi}{2} \right) \end{bmatrix} \\
&= \begin{bmatrix} \cos\frac{\theta}{2}e^{-i\frac{\phi}{2}} \\ \sin\frac{\theta}{2}e^{i\frac{\phi}{2}} \end{bmatrix}
\end{aligned}$$

41 This gives us a spinor $s_n = (s_{n1}, s_{n2}) = (\cos\frac{\theta}{2}e^{-i\frac{\phi}{2}}, \sin\frac{\theta}{2}e^{i\frac{\phi}{2}})$. If we recall our $\mathbf{v}_1 = (v_{1,1}, v_{1,2})$ from
42 part (a), we recall the relation we obtained at the end.

$$v_{1,2} = e^{i\phi} \frac{\sin(\theta)}{(\cos(\theta) + 1)} * v_{1,1}$$

43 Substituting $v_{1,1} = s_{n1} = \cos\frac{\theta}{2}e^{-i\frac{\phi}{2}}$ (as our final spinor seems to suggest), we get -

$$\begin{aligned}
v_{1,2} &= e^{i\phi} \frac{\sin(\theta)}{(\cos(\theta) + 1)} * v_{1,1} \\
&= e^{i\phi} \frac{\sin(\theta)}{(\cos(\theta) + 1)} * \cos\frac{\theta}{2}e^{-i\frac{\phi}{2}}
\end{aligned}$$

44 Recall $1 + \cos(A) = 2 * \cos^2(\frac{A}{2})$ and $\sin(A) = 2 * \sin(\frac{A}{2})\cos(\frac{A}{2})$

$$\begin{aligned}
e^{i\phi} \frac{\sin(\theta)}{(\cos(\theta) + 1)} * \cos\frac{\theta}{2}e^{-i\frac{\phi}{2}} &= e^{i\frac{\phi}{2}} \frac{\sin(\theta)}{2\cos^2(\frac{\theta}{2})} * \cos\frac{\theta}{2} \\
&= e^{i\frac{\phi}{2}} \frac{2\sin(\frac{\theta}{2})\cos(\frac{\theta}{2})}{2\cos^2(\frac{\theta}{2})} * \cos\frac{\theta}{2} \\
&= e^{i\frac{\phi}{2}} \sin(\frac{\theta}{2}) = s_{n2}
\end{aligned}$$

Therefore, apart from phase factors, the resultant spinor is the same as the spin up eigenvector we got in part (a).

3. Your question here.

[Sahas: I got this one.]

4. Your question here.

[Sankalp: I got this one.]

5. Prove that any function of the radial coordinate $f(r)$ where $r = |\mathbf{r}|$ and $\mathbf{X} \cdot \mathbf{P}$, where \mathbf{X} and \mathbf{P} are the position and momentum operators, are both scalar operators.

[Parth: Doing question 5, I'm not spending as much time on this as question 2]

Under a symmetry operator U , operators change as $\mathcal{O}' = U^\dagger \mathcal{O} U$. A scalar operator being one which is invariant under rotations, i.e

$$S' = U^\dagger [R] S U [R] = S$$

where $U(R(\boldsymbol{\alpha}) = e^{-\frac{i}{\hbar} \boldsymbol{\alpha} \cdot \mathbf{J}}$.

By considering infinitesimal rotations $\boldsymbol{\alpha} = \boldsymbol{\epsilon}$, we have

$$U[R(\boldsymbol{\alpha})] = \left(1 - \frac{i}{\hbar} \boldsymbol{\epsilon}_i J_i\right)$$

Thus, our definition for a scalar operator becomes -

$$S' = \left(1 + \frac{i}{\hbar} \boldsymbol{\epsilon}_i J_i\right) S \left(1 - \frac{i}{\hbar} \boldsymbol{\epsilon}_i J_i\right) = S$$

which gives us $\frac{i}{\hbar} \boldsymbol{\epsilon}_i [J_i, S] = 0$. Since $\boldsymbol{\epsilon}$ was an arbitrary choice, we have

$$[J_i, S] = 0$$

as our definition of a scalar operator.

Considering $f(r)$, where $r = |\mathbf{r}|$ as our operator.

$$[J_i, f(r)] = [J_i, r] * f'(r)$$

$r = \sqrt{\sum_{i=1}^3 X_i^2}$, Thus

$$[J_i, r] = [J_i, X_1] * \frac{X_1}{r} + [J_i, X_2] * \frac{X_2}{r} + [J_i, X_3] * \frac{X_3}{r}$$

we know that $[J_i, X_j] = i\hbar\epsilon_{ijl}X_l$. Thus

$$[J_i, r] = [J_i, X_j] * \frac{X_j}{r} = \frac{1}{r} (i\hbar\epsilon_{ijl}X_lX_j)$$

$$\epsilon_{ijl}X_lX_j = [X_l, X_j] = 0 (l \neq j) \Rightarrow [J_i, r] = 0$$

Thus, since $[J_i, r] = 0$, we have $[J_i, f(r)] = [J_i, r] * f'(r) = 0 * f'(r) = 0$.
Thus, $f(r)$ is a scalar operator.

Now considering $O = \mathbf{X} \cdot \mathbf{P}$ as our operator, we need to show $[J_i, O] = 0$

$$\mathbf{X} \cdot \mathbf{P} = X_i P_i \quad \text{implicit summation}$$

$$\begin{aligned} \therefore [J_i, O] &= [J_i, X_j P_j] \\ &= [J_i, X_j] P_j + X_j [J_i, P_j] \\ &= i\hbar\epsilon_{ijl}(X_l P_j + X_j P_l) \end{aligned}$$

Now, $\epsilon_{ijl}X_l P_j = [X_l, P_j]$ for $l \neq j$, but $[X_l, P_j] = 0, l \neq j$. Thus

$$i\hbar\epsilon_{ijl}(X_l P_j + X_j P_l) = 0 \Rightarrow [J_i, O] = 0$$

Since $[J_i, O] = 0$, we can say that the operator O is a scalar operator.
Thus, $\mathbf{X} \cdot \mathbf{P}$ is a scalar operator

6. Your question here.

[Sahas: I got this one.]