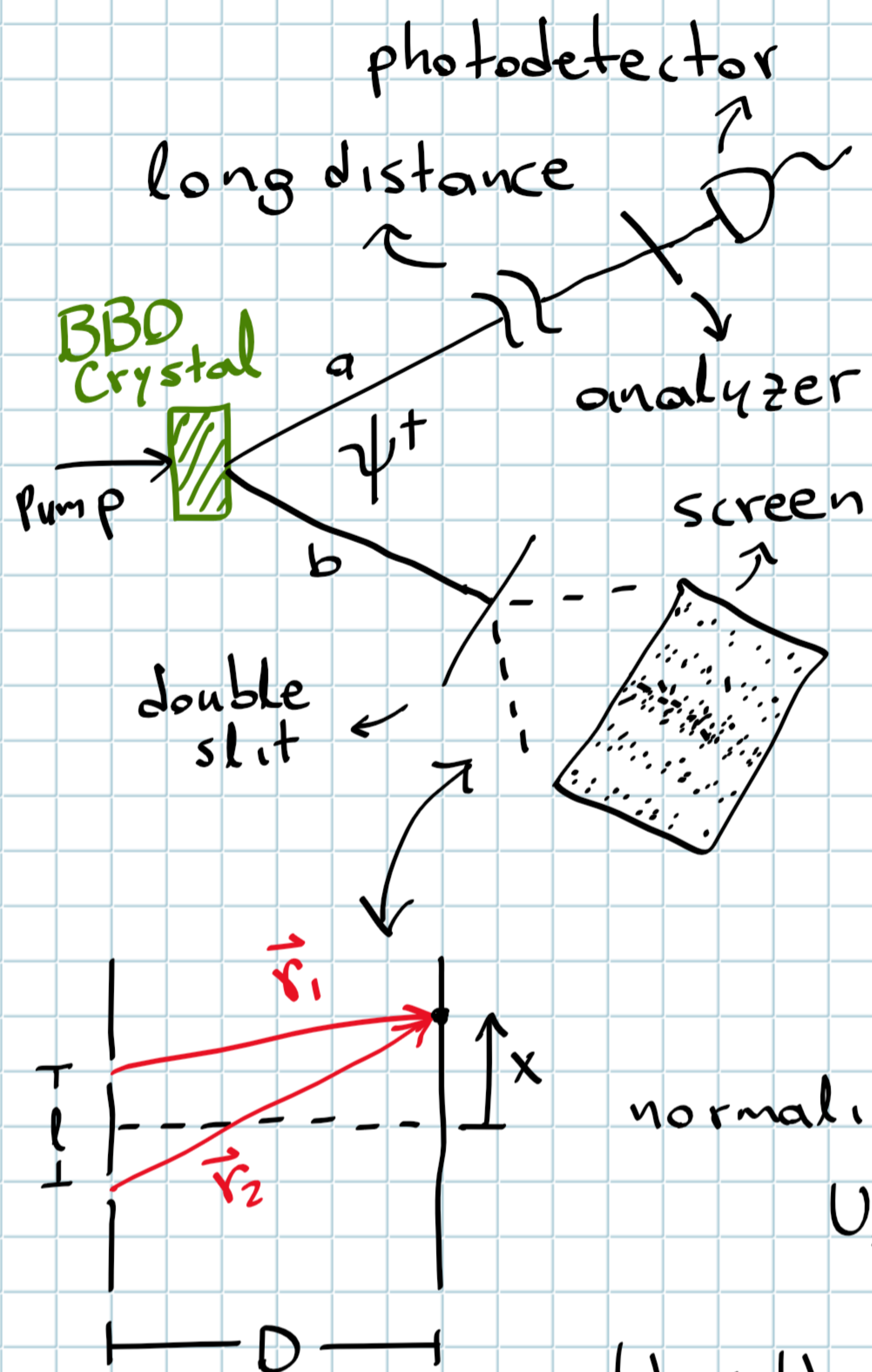


# Quantum Eraser

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two-photon state generated by the BBO crystal

$$|\psi^+\rangle_{ab} = \frac{1}{\sqrt{2}}(|H\rangle_a |V\rangle_b + |V\rangle_a |H\rangle_b)$$

state immediately after the double slit (no QWPs)

$$|\psi^+\rangle_{ab} \otimes \frac{1}{\sqrt{2}}(|S_1\rangle_b + |S_2\rangle_b) \quad \text{which-path states}$$

the effect of free-space propagation is

$$|S_j\rangle_b \rightarrow \int_{-\infty}^{\infty} dx \frac{e^{ikr_j}}{r_j} |x\rangle_b, \quad \text{screen states}$$

normalizing...

$$U_{fsp}|S_j\rangle = \sqrt{\frac{D}{\pi}} \int_{-\infty}^{\infty} dx \frac{e^{ikr_j}}{r_j} |x\rangle_b$$

thus the final state is

$$|\psi_f\rangle = \frac{1}{\sqrt{2}}(|H\rangle_a |V\rangle_b + |V\rangle_a |H\rangle_b) \otimes \sqrt{\frac{D}{2\pi}} \int dx \left( \frac{e^{ikr_1}}{r_1} + \frac{e^{ikr_2}}{r_2} \right) |x\rangle_b$$

In this state, polarization and position degrees of freedom are separable (not entangled).

Since we do not measure polarization of the b photon, we can trace out this degree of freedom, which gives us

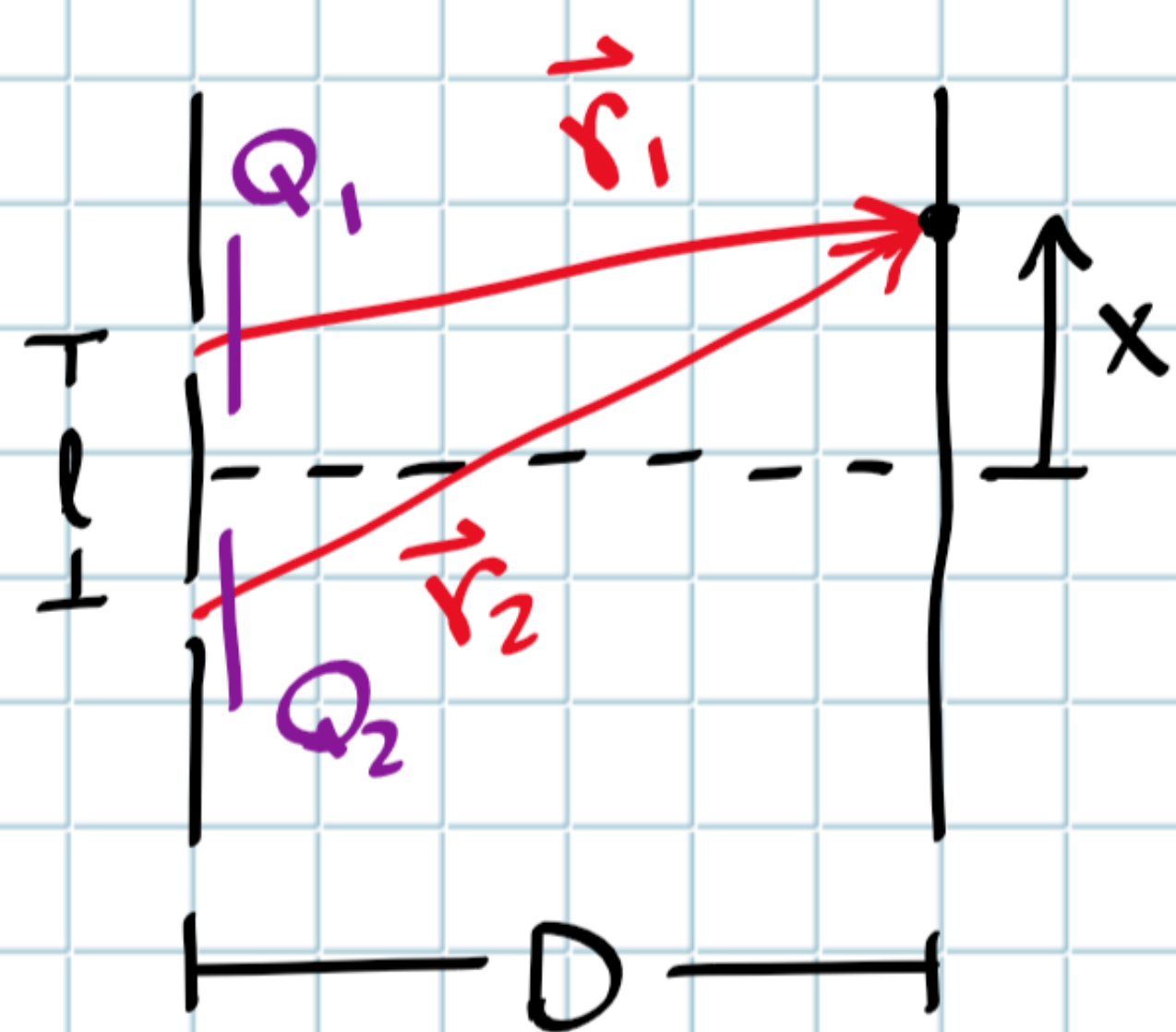
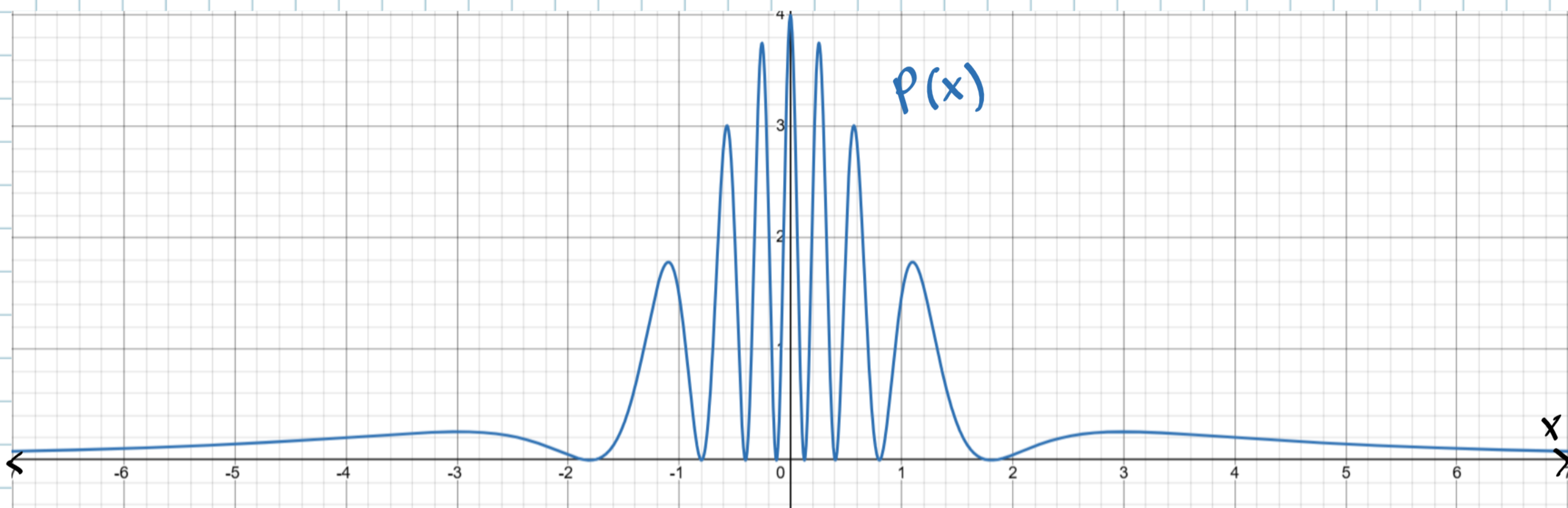
$$\rho_f = \left( \frac{1}{2} |H\rangle\langle H|_a + \frac{1}{2} |V\rangle\langle V|_a \right) \otimes \frac{D}{2\pi} \int dx \int dx' \left[ \frac{e^{ik(r_1-r'_1)}}{r_1 r'_1} + \frac{e^{ik(r_1-r'_2)}}{r_1 r'_2} + \text{same } 1 \leftrightarrow 2 \right] |x\rangle\langle x'|_b$$

thus, the photon detection probability density at the screen is

$$P(x) = \langle x | \text{Tr}_a(\rho_f) | x \rangle = \frac{D}{2\pi} \left( \frac{1}{r_1^2} + \frac{1}{r_2^2} + \frac{2}{r_1 r_2} \cos[k(r_1 - r_2)] \right)$$

where  $r_j = \sqrt{D^2 + (x \pm l/2)^2}$ , regardless of what happens to photon a.





What if we introduced two QWPs  $Q_1$  and  $Q_2$  with their fast axes at  $\pm 45^\circ$  with respect to the horizontal? That is, such that

$$Q_1 |H\rangle = |R\rangle$$

$$Q_1 |V\rangle = i |L\rangle$$

$$Q_2 |H\rangle = |L\rangle$$

$$Q_2 |V\rangle = -i |R\rangle$$

Then, the state immediately after the double slit would be

$$Q_1 = \frac{1}{2} \begin{pmatrix} 1 & i \\ i & 1 \end{pmatrix}$$

$$Q_2 = \frac{1}{2} \begin{pmatrix} 1 & -i \\ -i & 1 \end{pmatrix}$$

$$\frac{1}{2} (i |H\rangle_a \otimes (|L, s_1\rangle_b - |R, s_2\rangle_b) + |V\rangle_a \otimes (|R, s_1\rangle_b + |L, s_2\rangle_b))$$

After free-space propagation we get

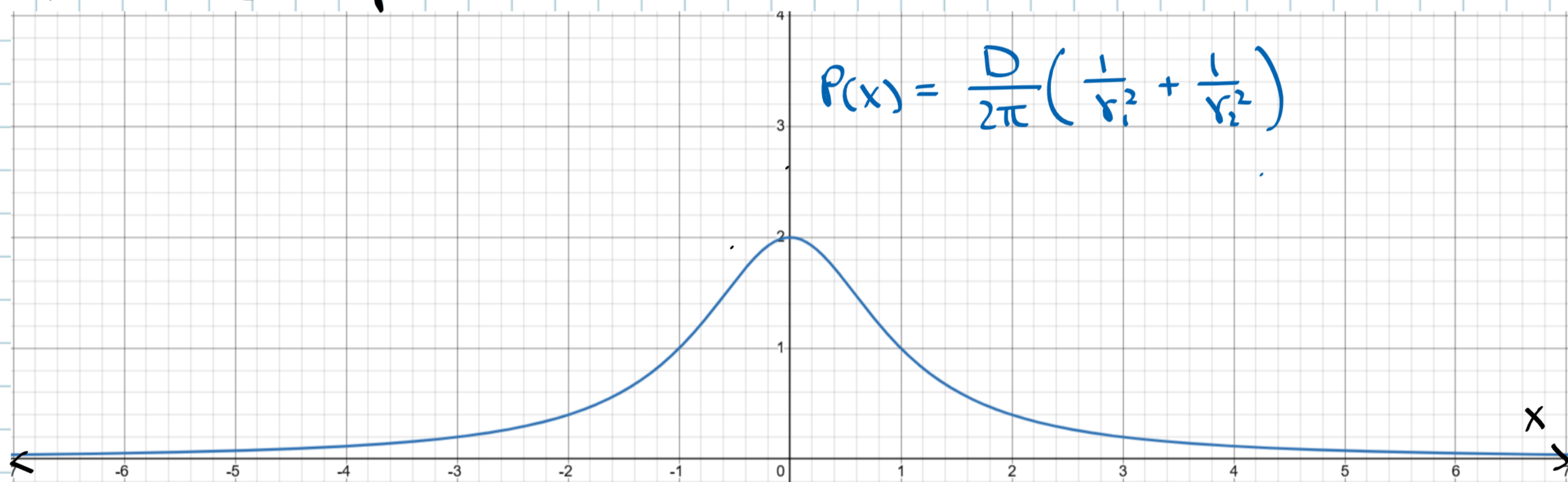
$$|\psi_f\rangle = \frac{i}{2} |L\rangle_b \otimes \sqrt{\frac{D}{\pi}} \int dx \left( \frac{e^{ikr_1}}{r_1} |H\rangle_a - i \frac{e^{ikr_2}}{r_2} |V\rangle_a \right) |x\rangle_b \\ + \frac{1}{2} |R\rangle_b \otimes \sqrt{\frac{D}{\pi}} \int dx \left( \frac{e^{ikr_1}}{r_1} |V\rangle_a - i \frac{e^{ikr_2}}{r_2} |H\rangle_a \right) |x\rangle_b$$

Tracing out unmeasured degrees of freedom...

$$P_f = \frac{D}{4\pi} \int dx \int dx' \left[ \left( \frac{e^{ik(r_1-r'_1)}}{r_1 r'_1} + \frac{e^{ik(r_2-r'_2)}}{r_2 r'_2} \right) \overbrace{( |H\rangle\langle H|_a + |V\rangle\langle V|_a )}^{1_a = \begin{pmatrix} 1 & 0 \\ 0 & 1 \end{pmatrix}} \right. \\ \left. + i \left( \frac{e^{ik(r_1-r'_2)}}{r_1 r'_2} - \frac{e^{ik(r_2-r'_1)}}{r'_1 r_2} \right) \overbrace{( |H\rangle\langle V|_a + |V\rangle\langle H|_a )}^{X_a = \begin{pmatrix} 0 & 1 \\ 1 & 0 \end{pmatrix}} \right] |x\rangle\langle x'|_b$$

Here, regardless of what happens to photon a, we see no interference pattern

$$P(x) = \frac{D}{2\pi} \left( \frac{1}{r_1^2} + \frac{1}{r_2^2} \right)$$





If photon a is measured in the HV-basis then we can post-select the outcomes of the measurements of photon b based on the results. This leads to

$$P_H(x) = \langle x |_b \langle H |_a |\rho_+ |H \rangle_a |x \rangle_b = \frac{D}{4\pi} \left( \frac{1}{r_1^2} + \frac{1}{r_2^2} \right)$$

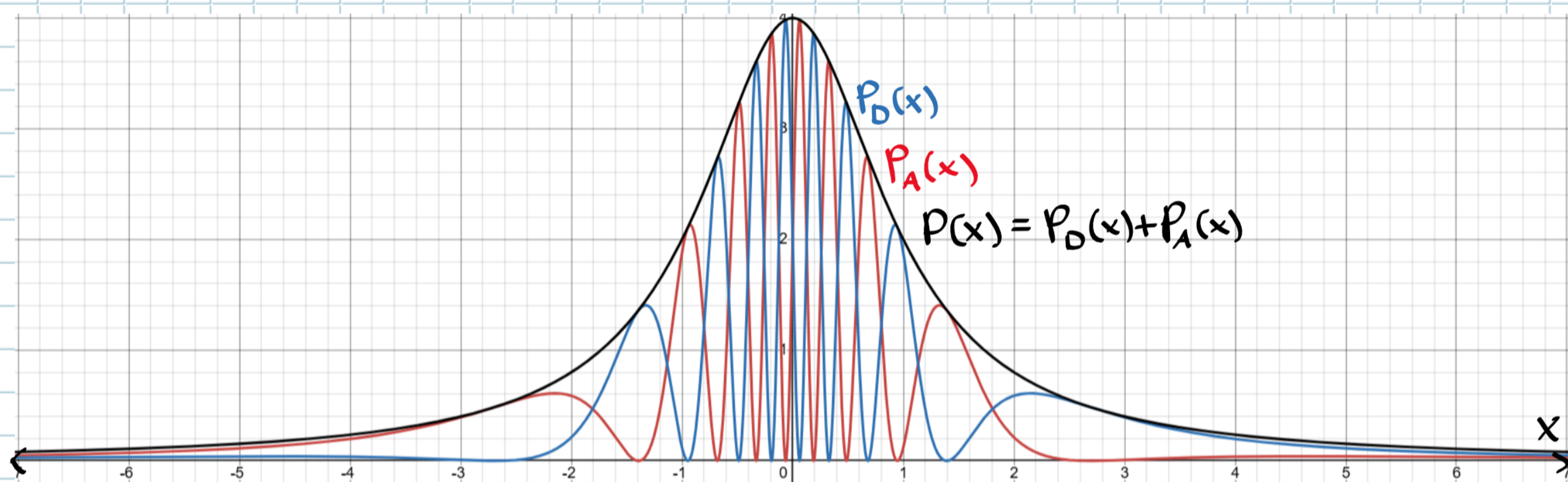
$$P_V(x) = \langle x |_b \langle V |_a |\rho_+ |V \rangle_a |x \rangle_b = \frac{D}{4\pi} \left( \frac{1}{r_1^2} + \frac{1}{r_2^2} \right)$$

these are  $P(x, H)$  and  $P(x, V)$

Again, no interference is observed. What if Bob measured in the DA-basis instead? Then, post-selection allow us to "recover" the interference pattern

$$P_D(x) = \frac{D}{4\pi} \left( \frac{1}{r_1^2} + \frac{1}{r_2^2} - \frac{2}{r_1 r_2} \sin[k(r_1 - r_2)] \right)$$

$$P_A(x) = \frac{D}{4\pi} \left( \frac{1}{r_1^2} + \frac{1}{r_2^2} + \frac{2}{r_1 r_2} \sin[k(r_1 - r_2)] \right)$$



Is "Quantum erasing" simply post-selection?

What would happen if Bob measured in the LR-basis?