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_{
m s}^{} table _{
m 2}003, takahashi _{
m r} evising _{
m 2}012 or machine learning estimators like the cosmological emulator (Cosmic Emu) heitmann _{
m c}
\label{eq:cosmological} its discontinuous and typically permitane xploration of the parameter space just around a fiducial cosmology or cosmology with linear regime like baryonic effects such as in [?] or to parametrize the cosmological dependence of non-linear clustering, beyond the Zeldovich approximation (see [?]). These fitting functions have, however, an intrinsicer ror that linear; there for eone has to be aware of the range in scales and redshifts they were designed to work on, in order to keep the error, renormalized 2006 and all other resummation methods derived from it (for example [?]), the time renormalisation group (TRO) is a superconstant of the stream representation does not be blacked as a superconstant of the superc
   linear scales, when the single stream approximation does not hold anylonger.\\
  ^{\textstyle \Lambda}_{1995, Amendola_2000, Amendola_2004, pettorino_baccigalupi_2008, involve an extra degree of freedom, associated to a scalar fine the properties of 
   analyticalnon
  linear analysis \cite{N-1} and cosmological N-1
  body simulations with incoupled \bar{D}ark Energy have been performed by many different groups baldiet al_2010, Li_Barrow_2011, respectively. The properties of the properties 
  {\it mass}_2006, maccio_c oncentration {\it 2008}, baldic larifying {\it 2011}, cui_h alo {\it 2012}, baldie ffect {\it 2011}, sutter_observability {\it 2014} and therefore {\it 2011}, baldie free {\it 2011}, baldie {\it 2012}, baldie {\it 2012}
   scale structure creminellising le-field {\it 2014}, CMB morris {\it cosmic microwave} {\it 2014} or laborator y experiments hamilton {\it atom}
     _{c}ode cs_{2}012 has shown that CDE models have characteristic and measurable features in the morphology and history of non-
   \overline{l} in ear \overline{s} tructures, such a shalos, subhalos and voids, black and therefore in the non-linear power spectrum.
     \dot{2}000, Amendola_2004, pettorino_baccigalupi_2008, baldi_etal_2010.
\begin{split} \ddot{\phi} + 3H\dot{\phi} + \frac{dV}{d\phi} &= \sqrt{\frac{2}{3}}\beta(\phi)\frac{\rho_c}{M_{Pl}} \ , \\ \dot{\rho}_c + 3H\rho_c &= -\sqrt{\frac{2}{3}}\beta(\phi)\frac{\rho_c\dot{\phi}}{M_{Pl}} \ , \end{split}
                                                                                                                                                                                                                                                                           (1)
                                                                                                                                                                                                                                                                            (2)
                                \dot{\rho}_b + 3H\rho_b = 0 ,
\dot{\rho}_r + 4H\rho_r = 0 ,
                                                                        4H\rho_{r}=0, 		(4)
3H^{2}=\frac{1}{M_{Pl}^{2}}(\rho_{b}+\rho_{c}+\rho_{r}+\rho_{\phi})
   \widetilde{M} at arrese_1984, Wetterich_1988. The coupling function <math>\beta(\phi)

\frac{1005}{4}, Bertotti_1 ess_T ortora_2 003) roughly to \bar{\gamma}| \leq 4 \cdot 10^{-5} G_{eff} = 6 \cdot 10^{-5}

 G_N(1-\bar{\gamma}/2)
G_{eff} = G_N(1-\bar{\gamma}/2)
 G_N(1+4\beta^2/3)
\beta^2 =
      -3\bar{\gamma}/8
\beta_{baryons}^2
  m_c(z) = m_{c,0}e^{-\beta(\phi(z)-\phi(0))}.
  ParameterExplanation
                                                                                                                                                                                                                                                      Value
                                                                                                                                                                                                                                                                                                                                                                 Reference \LambdaCDM
                                                                                                                                                                                                                 \{0.05, 0.10, 0.15\}
         w_{\phi}(z=0)
                                                                                                                                                                                  \{-0.997, -0.995, -0.992\}
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-1

0.809

 $\{0.825, 0.875, 0.967\}$ 

 $\sigma_8(z=0)$ 

 $\Omega_{CDM}^{H_0}$   $\Omega_{DE}$ 

 $2.42 \times 10^{-9}$ 

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k_{Ny}
    \begin{array}{l} k_{Ny} \\ P(k) \\ P(k) \\ P_{top} \\ P_{fold} \\ P_{shot} \\ P_{s
         P_{Exp}(k; \beta, z) \equiv P_{Exp}(k; \beta, z) / P_{LCDM}(k; \beta, z)
                  z = \{0, 0.55, 1.0, 1.61\}

\begin{array}{l}
\rho = \\
\{0.05, 0.10, 0.15\} \\
g_8 = \\
0
\end{array}

         {}^0_codecs_2012. Therefore, in order to see the net effect of the coupling at non-linear scales, each power spectrum ratio has been re-linear scales. \\
         linear scales, each power spectrum ration as been re-normalized by dividing each model by its respective all the ratios are unity. The net effect of the fifth force is a "bump" at non-linear scales, whose amplitude increases with higher couplings and whose maximum is shifted into higher wavenumbers for linear power spectrum is what we want to use to improve the estimation of parameters using future surveys. To achieve this, we have a survey of the property of the property
          \begin{array}{l} k \\ R(k;\beta,z) \equiv \\ P_{EXP}(k;\beta,z) / P_{LCDM}(k;\beta,z) \end{array} 
         k
R(k; \beta, z)
R(k; \beta, z)
         \tilde{a}_i =
             a_0, a_1, c, b, k_0
         0, ..., 5
                  k_0
\begin{array}{c} k_0 \\ b \\ i \\ \beta \\ \tilde{z} 
         f(k) = 1 + a_0 + a_1 \cdot k + c \cdot k \cdot \arctan((k - k_0) \cdot b)
f(k) = 1 + a_0 + a_1 \cdot k + c \cdot k \cdot \frac{b \cdot (k - k_0)}{\sqrt{1 + b^2 \cdot (k - k_0)^2}}
    a_{0}^{2}a_{0}^{1}k_{0}^{2}k_{0}^{2}k_{0}^{2}k_{0}^{2}k_{0}^{2}k_{0}^{2}k_{0}^{2}k_{0}^{2}k_{0}^{2}k_{0}^{2}k_{0}^{2}k_{0}^{2}k_{0}^{2}k_{0}^{2}k_{0}^{2}k_{0}^{2}k_{0}^{2}k_{0}^{2}k_{0}^{2}k_{0}^{2}k_{0}^{2}k_{0}^{2}k_{0}^{2}k_{0}^{2}k_{0}^{2}k_{0}^{2}k_{0}^{2}k_{0}^{2}k_{0}^{2}k_{0}^{2}k_{0}^{2}k_{0}^{2}k_{0}^{2}k_{0}^{2}k_{0}^{2}k_{0}^{2}k_{0}^{2}k_{0}^{2}k_{0}^{2}k_{0}^{2}k_{0}^{2}k_{0}^{2}k_{0}^{2}k_{0}^{2}k_{0}^{2}k_{0}^{2}k_{0}^{2}k_{0}^{2}k_{0}^{2}k_{0}^{2}k_{0}^{2}k_{0}^{2}k_{0}^{2}k_{0}^{2}k_{0}^{2}k_{0}^{2}k_{0}^{2}k_{0}^{2}k_{0}^{2}k_{0}^{2}k_{0}^{2}k_{0}^{2}k_{0}^{2}k_{0}^{2}k_{0}^{2}k_{0}^{2}k_{0}^{2}k_{0}^{2}k_{0}^{2}k_{0}^{2}k_{0}^{2}k_{0}^{2}k_{0}^{2}k_{0}^{2}k_{0}^{2}k_{0}^{2}k_{0}^{2}k_{0}^{2}k_{0}^{2}k_{0}^{2}k_{0}^{2}k_{0}^{2}k_{0}^{2}k_{0}^{2}k_{0}^{2}k_{0}^{2}k_{0}^{2}k_{0}^{2}k_{0}^{2}k_{0}^{2}k_{0}^{2}k_{0}^{2}k_{0}^{2}k_{0}^{2}k_{0}^{2}k_{0}^{2}k_{0}^{2}k_{0}^{2}k_{0}^{2}k_{0}^{2}k_{0}^{2}k_{0}^{2}k_{0}^{2}k_{0}^{2}k_{0}^{2}k_{0}^{2}k_{0}^{2}k_{0}^{2}k_{0}^{2}k_{0}^{2}k_{0}^{2}k_{0}^{2}k_{0}^{2}k_{0}^{2}k_{0}^{2}k_{0}^{2}k_{0}^{2}k_{0}^{2}k_{0}^{2}k_{0}^{2}k_{0}^{2}k_{0}^{2}k_{0}^{2}k_{0}^{2}k_{0}^{2}k_{0}^{2}k_{0}^{2}k_{0}^{2}k_{0}^{2}k_{0}^{2}k_{0}^{2}k_{0}^{2}k_{0}^{2}k_{0}^{2}k_{0}^{2}k_{0}^{2}k_{0}^{2}k_{0}^{2}k_{0}^{2}k_{0}^{2}k_{0}^{2}k_{0}^{2}k_{0}^{2}k_{0}^{2}k_{0}^{2}k_{0}^{2}k_{0}^{2}k_{0}^{2}k_{0}^{2}k_{0}^{2}k_{0}^{2}k_{0}^{2}k_{0}^{2}k_{0}^{2}k_{0}^{2}k_{0}^{2}k_{0}^{2}k_{0}^{2}k_{0}^{2}k_{0}^{2}k_{0}^{2}k_{0}^{2}k_{0}^{2}k_{0}^{2}k_{0}^{2}k_{0}^{2}k_{0}^{2}k_{0}^{2}k_{0}^{2}k_{0}^{2}k_{0}^{2}k_{0}^{2}k_{0}^{2}k_{0}^{2}k_{0}^{2}k_{0}^{2}k_{0}^{2}k_{0}^{2}k_{0}^{2}k_{0}^{2}k_{0}^{2}k_{0}^{2}k_{0}^{2}k_{0}^{2}k_{0}^{2}k_{0}^{2}k_{0}^{2}k_{0}^{2}k_{0}^{2}k_{0}^{2}k_{0}^{2}k_{0}^{2}k_{0}^{2}k_{0}^{2}k_{0}^{2}k_{0}^{2}k_{0}^{2}k_{0}^{2}k_{0}^{2}k_{0}^{2}k_{0}^{2}k_{0}^{2}k_{0}^{2}k_{0}^{2}k_{0}^{2}k_{0}^{2}k_{0}^{2}k_{0}^{2}k_{0}^{2}k_{0}^{2}k_{0}^{2}k_{0}^{2}k_{0}^{2}k_{0}^{2}k_{0}^{2}k_{0}^{2}k_{0}^{2}k_{0}^{2}k_{0}^{2}k_{0}^{2}k_{0}^{2}k_{0}^{2}k_{0}^{2}k_{0}^{2}k_{0}^{2}k
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