

03.20.2023

PHYS 427 A#7

$$10^{-20} \text{ m}^3 \left(\frac{\text{kg}}{1000 \text{ m}^3} \right) = 10^{-23} \text{ kg}$$

● #1 Apply the Adams-Williamson equation in integral form to compute the density $\rho(r)$ at the base of a 30km thick outer-most shell of the Earth (i.e. $r_0 = r_E = 6400 \text{ km}$ and $r = 6370 \text{ km}$), assuming constant seismic velocities of $\alpha = 7.00 \text{ km/s}$ and $\beta = 4.50 \text{ km/s}$ over the shell, and a density at the top of the shell of $\rho(r_0) = 2800 \text{ kg/m}^3$. Take the mass of the Earth to be $6.0 \times 10^{24} \text{ kg}$.

→ Want density at r , i.e. base of shell



A-W Eq: $\frac{d\rho}{dr} = \frac{-GM_r \rho(r)}{r^2(\alpha^2(r) - \frac{4}{3}\beta^2(r))} \Rightarrow \ln\left(\frac{\rho(r)}{\rho(r_0)}\right) = -GM_r \int_{r_0}^r \frac{du}{u^2(\alpha^2(u) - \frac{4}{3}\beta^2(u))}$

close to surface,

- Since thin shell relative to the Earth's radius, can assume mass of shell is negligible relative to Earth's mass, i.e. $M_r = M_E$.

$$\ln(\rho(r)) = -GM_E \int_{r_0}^r \frac{du}{u^2(\alpha^2 - \frac{4}{3}\beta^2)} + \ln(\rho(r_0)) = \frac{-GM_E}{(\alpha^2 - \frac{4}{3}\beta^2)} \int_{r_0}^r \frac{1}{u^2} du + \ln(\rho(r_0))$$

$$\ln(\rho(r)) = \frac{GM_E}{(\alpha^2 - \frac{4}{3}\beta^2)} \left[\frac{1}{r} - \frac{1}{r_0} \right] + \ln(\rho(r_0))$$

$$\ln(\rho(6370)) = \frac{(6.67 \times 10^{-20} \frac{\text{km}^3}{\text{kg s}^2})(6.00 \times 10^{24} \text{ kg})}{(7.00 \text{ km/s})^2 - \frac{4}{3}(4.50 \text{ km/s})^2} \left[\frac{1}{6370 \text{ km}} - \frac{1}{6400 \text{ km}} \right] + \ln(2800 \text{ kg/m}^3)$$

$$\rho(r) = e^{(0.013386167 + \ln(2800 \text{ kg/m}^3))}$$

$$\rho(r=6370 \text{ km}) = 2837.733257 \text{ kg/m}^3$$

$$\rho(r=6370 \text{ km}) \approx 2838 \text{ kg/m}^3$$