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P325 A#1

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1.1

$$\vec{E}(r,t) = \vec{E}_0 e^{i(\vec{k} \cdot \vec{r} - \omega t)}$$

a) Prove:

i) $\nabla \cdot \vec{E} = i\vec{k} \cdot \vec{E}$

$$\begin{aligned} \text{LHS} = \nabla \cdot \vec{E} &= \frac{\partial E_x}{\partial x} + \frac{\partial E_y}{\partial y} + \frac{\partial E_z}{\partial z} = E_0 e^{i(\vec{k} \cdot \vec{r} - \omega t)} i k_x + E_0 e^{i(\vec{k} \cdot \vec{r} - \omega t)} i k_y + E_0 e^{i(\vec{k} \cdot \vec{r} - \omega t)} i k_z \\ &= i e^{i(\vec{k} \cdot \vec{r} - \omega t)} (E_0 k_x + E_0 k_y + E_0 k_z) = i e^{i(\vec{k} \cdot \vec{r} - \omega t)} \vec{E}_0 \cdot \vec{k} \\ &= \vec{E}_0 e^{i(\vec{k} \cdot \vec{r} - \omega t)} \cdot i\vec{k} = i\vec{k} \cdot \vec{E} = \text{RHS.} \end{aligned}$$

$$\therefore \nabla \cdot \vec{E} = i\vec{k} \cdot \vec{E}$$

ii) $\nabla \times \vec{E} = i\vec{k} \times \vec{E}$

$$\text{LHS} = \nabla \times \vec{E} = \left(\frac{\partial E_z}{\partial y} - \frac{\partial E_y}{\partial z} \right) \hat{i} + \left(\frac{\partial E_x}{\partial z} - \frac{\partial E_z}{\partial x} \right) \hat{j} + \left(\frac{\partial E_y}{\partial x} - \frac{\partial E_x}{\partial y} \right) \hat{k}$$

$$\nabla \times \vec{E} = (E_z i k_y - E_y i k_z) \hat{i} + (E_x i k_z - E_z i k_x) \hat{j} + (E_y i k_x - E_x i k_y) \hat{k}$$

$$\begin{aligned} \text{RHS} = i\vec{k} \times \vec{E} &= \begin{vmatrix} \hat{i} & \hat{j} & \hat{k} \\ k_x & k_y & k_z \\ E_x & E_y & E_z \end{vmatrix} = i \left[(k_y E_z - k_z E_y) \hat{i} + (k_z E_x - k_x E_z) \hat{j} + (k_x E_y - k_y E_x) \hat{k} \right] \end{aligned}$$

$$\text{LHS} = \text{RHS}, \therefore \nabla \times \vec{E} = i\vec{k} \times \vec{E}$$

iii) $\frac{d\vec{E}}{dt} = -i\omega \vec{E}$

$$\text{LHS} = \frac{d\vec{E}}{dt} = \frac{d}{dt} \vec{E}_0 e^{i(\vec{k} \cdot \vec{r} - \omega t)} = \vec{E}_0 e^{i(\vec{k} \cdot \vec{r} - \omega t)} (-i\omega) = -i\omega \vec{E} = \text{RHS}$$

$$\therefore \frac{d\vec{E}}{dt} = -i\omega \vec{E}$$

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b) Calculate $\nabla^2 \vec{E}$ and $\frac{\partial^2 \vec{E}}{\partial t^2}$. Show explicitly the wave equation is satisfied w/ $\omega = c|\vec{k}|$

$$\rightarrow \nabla^2 \vec{E} = \nabla^2 E_x \hat{i} + \nabla^2 E_y \hat{j} + \nabla^2 E_z \hat{k}$$

$$= \frac{\partial^2 E_x}{\partial x^2} \hat{i} + \frac{\partial^2 E_y}{\partial y^2} \hat{j} + \frac{\partial^2 E_z}{\partial z^2} \hat{k}$$

$$= \frac{\partial}{\partial x} i k_x E_x \hat{i} + \frac{\partial}{\partial y} i k_y E_y \hat{j} + \frac{\partial}{\partial z} i k_z E_z \hat{k}$$

$$= i k_x E_x i k_x \hat{i} + i^2 k_y^2 E_y \hat{j} + i^2 k_z^2 E_z \hat{k}$$

$$\nabla^2 \vec{E} = - (E_x k_x^2 \hat{i} + E_y k_y^2 \hat{j} + E_z k_z^2 \hat{k}) = - \begin{pmatrix} E_x k_x^2 \\ 0 \\ 0 \end{pmatrix} - \begin{pmatrix} 0 \\ E_y k_y^2 \\ 0 \end{pmatrix} - \begin{pmatrix} 0 \\ 0 \\ E_z k_z^2 \end{pmatrix} = \boxed{-\vec{k}(\vec{k} \cdot \vec{E})} \quad \textcircled{1}$$

$$\rightarrow \frac{\partial^2 \vec{E}}{\partial t^2} = \frac{\partial}{\partial t} -i\omega \vec{E} = -i^2(-\omega^2) \vec{E} = \boxed{-\omega^2 \vec{E} = \frac{\partial^2 \vec{E}}{\partial t^2}}$$

\rightarrow Consider $\omega = c|\vec{k}|$:

$$\vec{k} = k_x \hat{i} + k_y \hat{j} + k_z \hat{k}, |\vec{k}| = \sqrt{\vec{k} \cdot \vec{k}}, |\vec{k}|^2 = \vec{k} \cdot \vec{k}$$

$$\frac{\partial^2 \vec{E}}{\partial t^2} = -\omega^2 \vec{E}$$

$$= -c^2 (\vec{k} \cdot \vec{k}) (\vec{E})$$

$$= -c^2 (\vec{k}) (\vec{k} \cdot \vec{E})$$

$$= -c^2 (\nabla^2 \vec{E})$$

\rightarrow Since \vec{k} and \vec{E} are both in the $(\hat{i}, \hat{j}, \hat{k})$

$$\text{basis, } (\vec{k} \cdot \vec{k}) \vec{E} = \vec{k} (\vec{k} \cdot \vec{E}) = \begin{pmatrix} E_x k_x^2 \\ E_y k_y^2 \\ E_z k_z^2 \end{pmatrix}$$

\rightarrow From ①

$$\Rightarrow \boxed{\frac{\partial^2 \vec{E}}{\partial t^2} = c^2 \nabla^2 \vec{E}}, \therefore \text{THE WAVE EQ. IS SATISFIED WHEN } \omega = c|\vec{k}|$$

c) Use the property of waves that $\lambda v = c$ to show $|\vec{k}| = \frac{2\pi}{\lambda}$. Recall $\omega = 2\pi\nu$
 $\omega = c|\vec{k}|$

$$\lambda = \frac{c}{v}, v = \frac{\omega}{2\pi} \Rightarrow \lambda = \frac{c 2\pi}{\omega}, \omega = c|\vec{k}| \Rightarrow \lambda = \frac{c 2\pi}{c|\vec{k}|} \Rightarrow \boxed{|\vec{k}| = \frac{2\pi}{\lambda}}$$

d) Use results from (a) and Maxwell's eq. to prove plane waves are transverse. i.e. that $\vec{k} \cdot \vec{E} = 0$

$$\text{From (a): } \nabla \cdot \vec{E} = i\vec{k} \cdot \vec{E}$$

$$\text{From Maxwell's eq. } \nabla \cdot \vec{E} = 0$$

$$\rightarrow \text{So } i\vec{k} \cdot \vec{E} = 0 \Rightarrow \boxed{\vec{k} \cdot \vec{E} = 0}$$

\rightarrow This means the direction of propagation \perp electric field.

PLANE WAVES
 \therefore ARE
 TRANSVERSE

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1.2

a) Show that for an electromagnetic ^{plane} wave, the magnitudes of the electric and magnetic components are related through $|\vec{E}| = c|\vec{B}|$
 Assume $\vec{E}(\vec{r}, t) = \vec{E}_0 e^{i(\vec{k} \cdot \vec{r} - \omega t)}$, $\vec{B}(\vec{r}, t) = \vec{B}_0 e^{i(\vec{k} \cdot \vec{r} - \omega t)}$

Hint: use results from 1.1 and M.E. for Faraday's Law. Also $\vec{A} \times \vec{B} = |\vec{A}||\vec{B}| \sin \theta$

Faraday's Law: $\nabla \times \vec{E} = -\frac{\partial \vec{B}}{\partial t} = -\frac{\partial}{\partial t} \vec{B}_0 e^{i(\vec{k} \cdot \vec{r} - \omega t)} = i\omega \vec{B}_0 e^{i(\vec{k} \cdot \vec{r} - \omega t)} = i\omega \vec{B}$ ①

From 1.1(a)(ii): $\nabla \times \vec{E} = i\vec{k} \times \vec{E}$ ②

① = ② $\Rightarrow i\omega \vec{B} = i\vec{k} \times \vec{E} \Rightarrow \omega \vec{B} = \vec{k} \times \vec{E}$

$\Rightarrow \omega \vec{B} = |\vec{k}| |\vec{E}| \sin(\frac{\pi}{2})$, $\omega = c|\vec{k}|$

$\Rightarrow c|\vec{k}| \vec{B} = |\vec{k}| |\vec{E}|$

$\therefore c\vec{B} = \vec{E}$

b) Consider electromagnetic field passing through Hydrogen atom.

Angular momentum of e^- orbiting atom: $L_n = n\hbar$, $n=1,2,3,\dots$

Using $L_n = m v_n r_n$ and result that $r_n = 4\pi\epsilon_0 \frac{n^2 \hbar^2}{m_e e^2}$, show max velocity of e^- is $v = \frac{1}{4\pi\epsilon_0} \frac{e^2}{\hbar}$

$L_n = m v_n r_n \Rightarrow v_n = \frac{L_n}{m r_n} = \frac{n\hbar}{m 4\pi\epsilon_0 \frac{n^2 \hbar^2}{m_e e^2}} = \frac{n\hbar m_e e^2}{m_e 4\pi\epsilon_0 n^2 \hbar^2} = \frac{e^2}{4\pi\epsilon_0 n\hbar}$

v_n is largest when $n=1$, so: $V = \frac{1}{4\pi\epsilon_0} \frac{e^2}{\hbar}$

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$$c) \vec{F}_E = -e\vec{E}, \vec{F}_B = -e\vec{v} \times \vec{B}$$

What is max magnitude of magnetic component of force exerted on e^- ?

$$|\vec{F}_B| = e v_{\max} |\vec{B}|$$

$$\rightarrow \text{From (b), } v_{\max} = \frac{1}{4\pi\epsilon_0} \frac{e^2}{\hbar}$$

$$\rightarrow |\vec{B}| = \sqrt{B_0 e^{i(k \cdot r - \omega t)}} \cdot B_0 e^{-i(k \cdot r - \omega t)} = B_0$$

$$|\vec{F}_B| = e \frac{1}{4\pi\epsilon_0} \frac{e^2}{\hbar} B_0$$

$$\therefore \text{Max magnitude is } |\vec{F}_B| = \frac{B_0 e^3}{4\pi\epsilon_0 \hbar}$$

d) Show ratio of mag to electric force on atom given by light: $\frac{|\vec{F}_B|}{|\vec{F}_E|} = \alpha$
 $\alpha = \frac{1}{4\pi\epsilon_0} \frac{e^2}{\hbar c}$

$$\left. \begin{array}{l} |\vec{F}_E| = e|\vec{E}| \\ |\vec{F}_B| = ev|\vec{B}| \end{array} \right\} \frac{|\vec{F}_B|}{|\vec{F}_E|} = \frac{ev|\vec{B}|}{e|\vec{E}|} = v \frac{|\vec{B}|}{|\vec{E}|}$$

$$\text{From (a), } \frac{|\vec{B}|}{|\vec{E}|} = \frac{1}{c}; \text{ From (b), } v = \frac{e^2}{4\pi\epsilon_0 \hbar c}$$

$$\Rightarrow \frac{|\vec{F}_B|}{|\vec{F}_E|} = \frac{e^2}{4\pi\epsilon_0 \hbar c} = \alpha$$

1.3

a) Intensity of sunlight on earth, for surface \perp to sun's direction?

$P_{\text{sun}} = 3.8 \times 10^{26} \text{ W}$, spread out evenly over sphere w/ $r = 1.5 \times 10^{11} \text{ m}$

$$I_{\text{sun}} = \frac{P_{\text{sun}}}{A}, \quad A = 4\pi r^2$$

$$= \frac{3.8 \times 10^{26} \text{ W}}{4\pi (1.5 \times 10^{11} \text{ m})^2} = 1343.975 \text{ W/m}^2 \approx 1.3 \times 10^3 \text{ W/m}^2$$

b) Intensity of laser w/ $P = 1 \text{ mW}$, $\lambda = 650 \text{ nm}$, $d = .5 \text{ mm}$

$$I = \frac{P}{A}, \quad A = \pi r^2$$

$$= \frac{.001 \text{ W}}{\pi (0.25 \times 10^{-3} \text{ m})^2} = 1.27 \text{ W/m}^2$$